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Holographic entanglement beyond classical gravity

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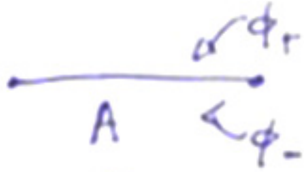
HOLOGRAPHIC ENTANGLEMENT BEYOND CLASSICAL GRAVITY

w/ T. Barrella, X. Dong, V.L. Martin
@ 1306.4682.

Reduced density matrix:



$$\rho[\phi_+,\phi_-] = \frac{1}{Z}\int \mathcal{D}\phi\, e^{-S[\phi]}\, \delta(\phi(A,t^+) = \phi_+) \times \delta(\phi(A,t^-) = \phi_-)$$

Clearly then:

$$\mathrm{Tr}\rho^n = \frac{Z_n}{Z_1^n}$$ ← partition function on n-sheeted cover:

Renyi entropy:

$$S_n = -\frac{1}{n-1}\log \mathrm{Tr}\rho^n$$

Entanglement entropy

$$S = \lim_{n\to 1} S_n = -\mathrm{tr}(\rho\log\rho)$$

small x expansion known from OPEs.

$$S_n = -N \frac{n}{2(n-1)} \left(\frac{x}{4n^2}\right)^{2h} \sum_{k=1}^{n-1} \frac{1}{\left(\sin\frac{\pi k}{n}\right)^{4h}}$$

multiplicity

dimension

$$\to S = -N \left(\frac{x}{4}\right)^{2h} \frac{\sqrt{\pi}}{4} \frac{\Gamma(2h+1)}{\Gamma(2h+\frac{3}{2})}$$

[Cardy-Calabrese]

We managed to implement the strategy exactly in an expansion in x.

<u>One key step</u>: To each order in x, only a finite number of words contribute to the sum.

<u>leading order</u>: <u>CDWs</u> "consecutively decreasing words"

$$\gamma_{k,m} = L_{k+m} L_{k+m-1} \cdots L_{m+1}$$

$\to$ can find eigenvalues explicitly

<u>higher orders</u>: products of <u>CDWs</u>.

$$\boxed{S = -\left(\frac{x^4}{630} + \underbrace{\frac{2x^5}{693} + \frac{15x^6}{4004} + \frac{x^7}{234}}_{\text{only CDWs}} + \frac{167x^8}{36936} + O(x^9)\right)}$$

leading

matches CC result.

Giombi-Maloney-Yin derived a nice formula for functional determinants on AdS_3/Γ:

$$\log Z\big|_{\text{one-loop}} = -\sum_{\gamma \in P}\sum_{m=2}^{\infty} \log\left|1-q_\gamma^m\right|$$

"virasoro descendants of ground state" on torus.

for metric fluctuations.

γ has eigenvalues: $q_\gamma^{\pm 1/2}$, $|q_\gamma|<1$.

conjugacy classes of primitive elements of Γ (i.e. γ s.t. $\gamma \neq \gamma_q^n$, $n>1$)

strategy:
- find Γ for all n (Schottky uniformize)
- generate P for all $n \to$ primitive words.
- find eigenvalues of words
- do sum
- analytically continue to $n \to 1$.

Cross ratio: $x = \dfrac{(z_3-z_2)(z_4-z_1)}{(z_3-z_1)(z_4-z_2)}$

Mutual information only depends on this. (4 pt fn of twist operators)

In well-separated phase, mutual information vanishes.

$$I_n(L_1,L_2) \equiv S_n(L_1) + S_n(L_2) - S_n(L_1 \cup L_2)$$

However, general results ((Cardy-Calabrese), this vanishing cannot be exact

↳ will recover leading mutual information in 1-loop correction.

Similarly in torus (us):

- above HP: $S = \frac{c}{6} \log\left(\sinh^2 \frac{T\Delta\phi}{2}\right)$ (1)
- below HP: $S = \frac{c}{6} \log \sin^2 \frac{\pi\Delta\phi}{R}$ (2)

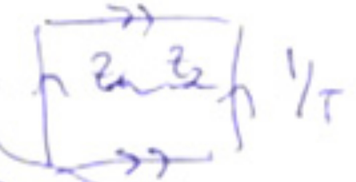
These are known exact answers for CFTs at ∞ and 0 temperature!

→ (1) is missing finite size corrections.

(2) is missing mixed nature of density matrix $S(L) \neq S(R \setminus L)$.

→ find these at 1-loop.

Similar for torus



$$\psi''(z) + \frac{1}{2}\sum_{i=1}^{2}\Big(\Delta\,\wp(z-z_i) + \gamma(-1)^{i+1}\zeta(z-z_i) + \delta\Big)\psi(z) = 0$$

R

Faulkner (following Zograf/Takhtadzhyan and Krasnov) showed that on-shell bulk action of AdS_3/Γ satisfied:

$$\frac{\partial S_E}{\partial z_i} = -\frac{cn}{6}\gamma_i$$

c → central charge

γ_i → found by trivializing monodromy

⇒ entanglement entropy

$$\frac{\partial S}{\partial z_i} = -\lim_{n\to 1}\frac{cn}{6(n-1)}\gamma_i$$

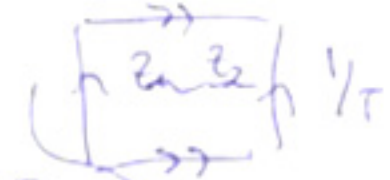
by solving for γ_i analytically as $n\to 1$, rederived the RT answer for two intervals on plane.

Interesting feature (known from RT)

small separation: large separation

phase transition → change of dominant uniformization

Similar for torus



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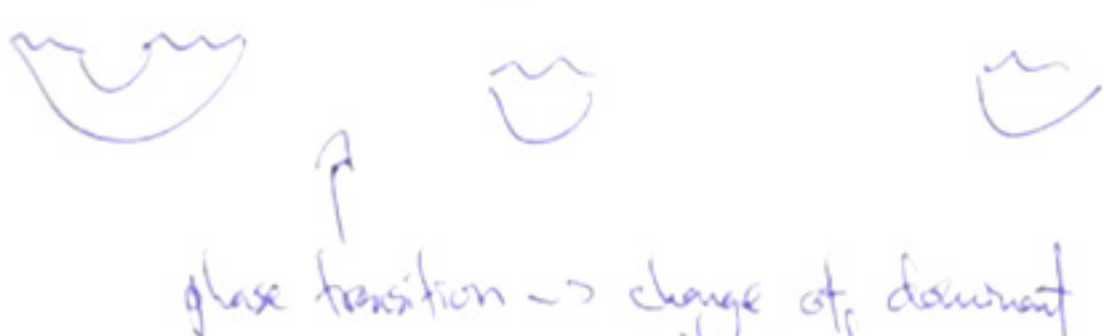


phase transition → change of dominant uniformization.

choose accessory parameters γ_i such that:

(I) ~~no~~ monodromy at ∞

↳ fixes 3 of 4 γ_is.

(II) one of the two cycles has trivial monodromy:

→ choice of Schottky uniformization

→ corresponds to different bulk AdS_3/Γ → lowest action dominates.

→ can have phase transitions (**)

Having trivialized all but one cycle, have generators of Schottky gp: M_1 M_2

$$\left.\begin{aligned} L_1 &= M_2 M_1 \\ L_i &= M_2^{i-1} L_1 M_2^{-(i-1)} \qquad i = 2, n-1 \end{aligned}\right\} \begin{array}{l} n-1 = g \\ \text{generators.} \end{array}$$

[Assumed $\mathbb{Z}_n$ symmetric uniformization]

Equ has 2 independent solns: ψ_1, ψ_2.
<u>Define</u>:

$$\omega = \frac{\psi_1(z)}{\psi_2(z)}.$$

near z_i: $\psi \sim (z-z_i)^{(1\pm 1/n)/2}$

$$\Rightarrow \omega \sim (z-z_i)^{1/n} \rightarrow \text{single valued on } n\text{-th cover.}$$

If we integrate pair (ψ_1, ψ_2) around contour C enclosing some $\{z_i\}$, get <u>monodromy</u>:

$$\begin{pmatrix}\psi_1\\ \psi_2\end{pmatrix} \mapsto M(C)\begin{pmatrix}\psi_1\\ \psi_2\end{pmatrix} \qquad M(C) = \begin{pmatrix}a & b\\ c & d\end{pmatrix} \in PSL(2,\mathbb{C})$$

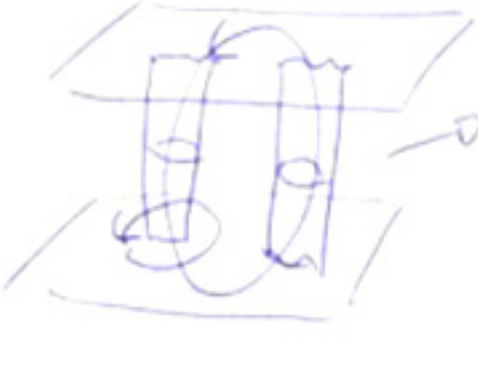
e.g.

This tells us that for ω must be identified:

$$\omega \sim \frac{a\omega+b}{c\omega+d} \rightarrow \text{quotient } \mathbb{C}/\Gamma\,!$$

Final step: too many quotients! $\qquad g = n-1$

eg. $\underline{n=2}$:



$\rightarrow$ two nontrivial cycles

$\Rightarrow$ Schottky group should be generated by <u>one</u> <u>not</u> <u>two</u> generators.

under $\Sigma = \mathbb{C}/\Gamma$, ~~g~~ circles give g generators of $\pi_1(\Sigma)$. Paths connecting circles give remaining g generators of genus g Riemann surface

How to find Γ?

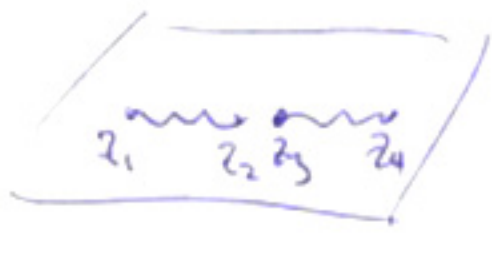
- Given multi-valued coordinate z on branched cover, want to find w s.t.
  1. single valued
  2. Γ acts by Möbius xfm on w.

- ~~[illegible]~~ Simplest nontrivial cases: • two intervals on plane.
  • one interval on torus. (CFT on circle, finite T)

- Define:

$$\psi''(z) + \frac{1}{2}\sum_{i=1}^{4}\left(\frac{\Delta}{(z-z_i)^2} + \frac{\gamma_i}{z-z_i}\right)\psi = 0$$

$$\left.\begin{array}{l}\Delta = \frac{1}{2}\left(1-\frac{1}{n^2}\right)\\ \gamma_i \text{ fixed shortly.}\end{array}\right\}$$



AdS_3: $ds^2 = \dfrac{d\xi^2 + dw\,d\bar{w}}{\xi^2}$

near $\xi \to 0$ (conformal bdy.) $PSL(2,\mathbb{C})$ acts via Möbius xfms:

$$w \mapsto L(w) = \frac{aw+b}{cw+d} \qquad ad-bc=1.$$

$$\xi \longmapsto |L'(w)|\,\xi$$

$\Rightarrow$ need to find Γ s.t. $\Sigma = \mathbb{C}/\Gamma$.

with Γ subgp of Möbius xfms.

$\Rightarrow$ called Schottky uniformization

Why reasonable?

Möbius maps circles to circles:

~~Sup~~ The Schottky gp will be freely generated by g elements of $PSL(2,\mathbb{C})$, Consider $2g$ disjoint circles s.t. $C_i' = L_i(C_i)$

S in general difficult to compute.
"outrageous" Ryu-Takayanagi proposal for theories with classical gravity duals.

$$S_A = \frac{\mathrm{Area}(\Sigma)}{4G_N}$$

It's true! $\underbrace{\text{Hartman; Faulkner;}}_{1+1\ \text{CFTs.}}$ Lewkowycz + Maldacena

Connection spacetime ⟷ entanglement?
→ Go beyond classical gravity in simplest setting, 1+1 CFTs.
↳ directly compute the bulk partition function for these Z_n at one loop.

We will obtain one-loop corrections to the RT formula in 1+1 dimensions.

Step I: Schottky uniformization

Einstein's equations in 2+1 bulk: geometry is a quotient AdS_3/Γ, with $\Gamma \subset PSL(2,\mathbb{C})$

holography: $Z_{CFT}(\Sigma) = Z_{bulk}(\partial(AdS_3/\Gamma) = \Sigma)$

for us: Σ: n^{th} branched cover of $\mathbb{C}$ over region A.
therefore need to: (i) find Γ

(ii) evaluate $Z(AdS_3/\Gamma)$