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Unwinding Inflation

New Lights in Cosmology from the CMB
ICTP Trieste, Summer 2013

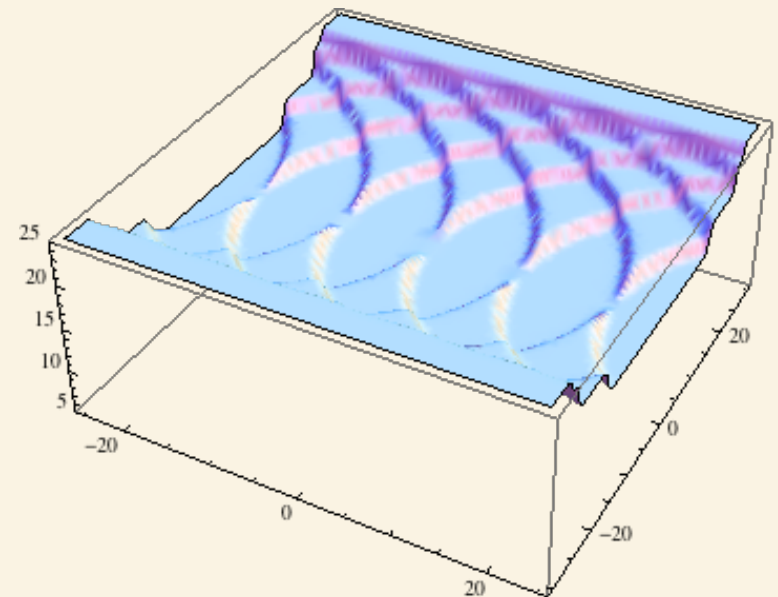
with Roberto Gobbetti, Matthew Kleban, Marjorie Schillo
1211.3416 (short), 1211.4589 (long)

What our model is

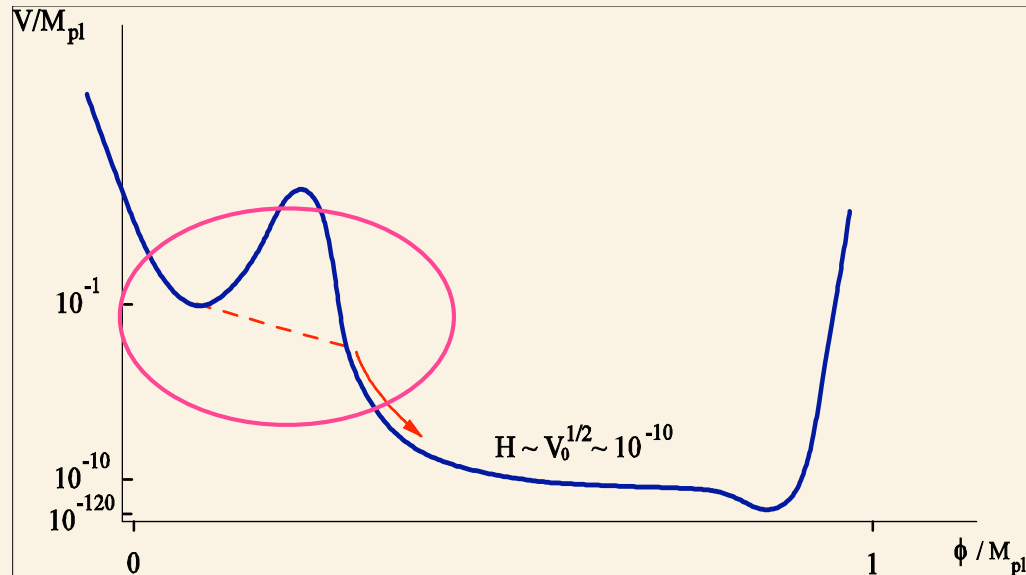
Take:

- Boom and bust inflation
- Monodromy inflation
- Chain inflation
- $m^2 \varphi^2$
- DBI inflation
- Old inflation
- Cascade inflation
- Trapped inflation

in no particular order.
Shake, do not stir...
Our model is ready!



False vacuum eternal inflation



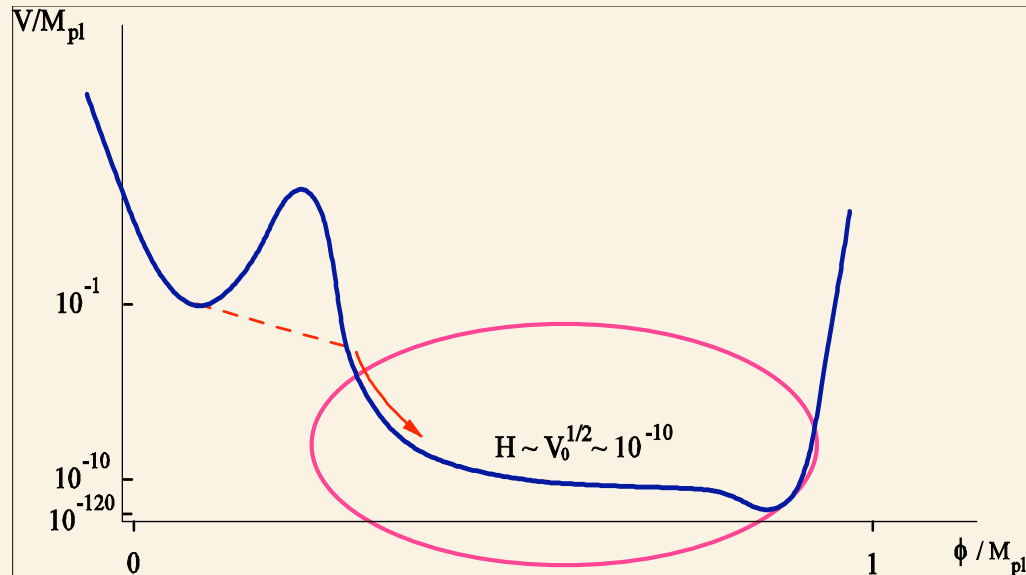
Guth's original idea: **universe trapped in a metastable minimum**
Inflation is eternal and a powerful attractor: some regions of the universe always inflate

How to end inflation (graceful exit)?

First order phase transition, but very difficult to percolate

Other problems: a bubble contains an *open* FRW universe...
and where is the matter?

Slow-roll (new) inflation



A better model: a very flat scalar potential (approximate shift symmetry)

The scalar slowly rolls and does the job: inflates the universe, and at the end it gets converted into radiation

All models constructed are basically EFTs, with some degree of fine-tuning and not simple to embed into a UV-complete theory

Can we do better?

Suppose we want to realize inflation in string theory

Generic ingredients are

- Extra (compact) dimensions
- Extended objects (D-branes), which couple to
- Higher-form fields

In inflation we basically want to slowly decrease vacuum energy.
In d dimensions, a d -form electric flux (d antisymmetric indices) is vacuum energy!

1+1-d: Schwinger model

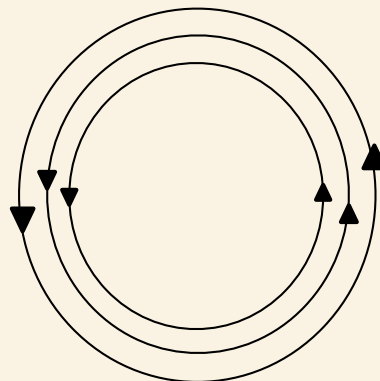
Electric field in 1+1-d is $F_{\mu\nu} = E \epsilon_{\mu\nu}$

Compactify the space to S_1 .

We can have a field flux that “wraps” the circle

Varying the size of the circle does not modify the field or its energy density: vacuum energy $\Lambda \sim E^2$

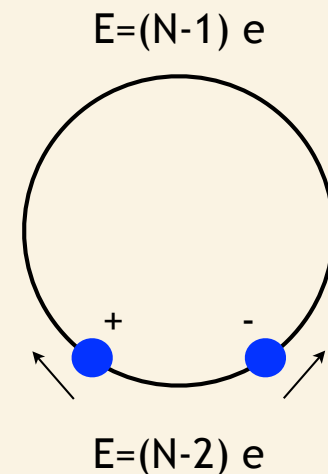
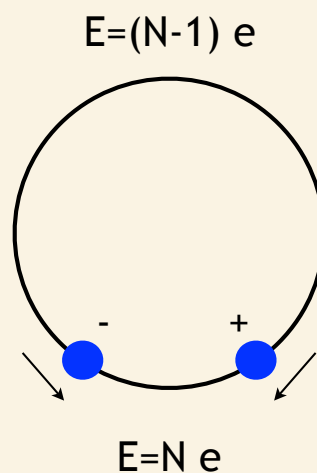
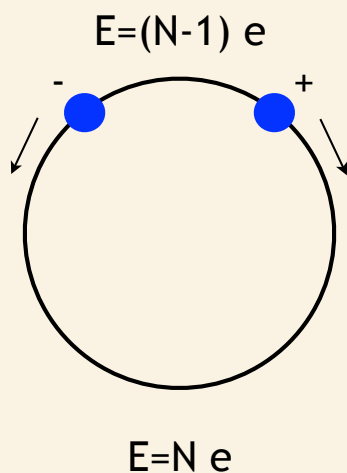
QM can discharge the field by the spontaneous nucleation of charge pairs!



$$E = Ne$$

The story of a pair

- A quantum nucleation event happens: the field is discharged by e between the $e^+ e^-$ pair that appears
- Classically, the field accelerates the charges in opposite directions, until they meet on the opposite side of the circle
- Typically, they pass through each other and continue to accelerate

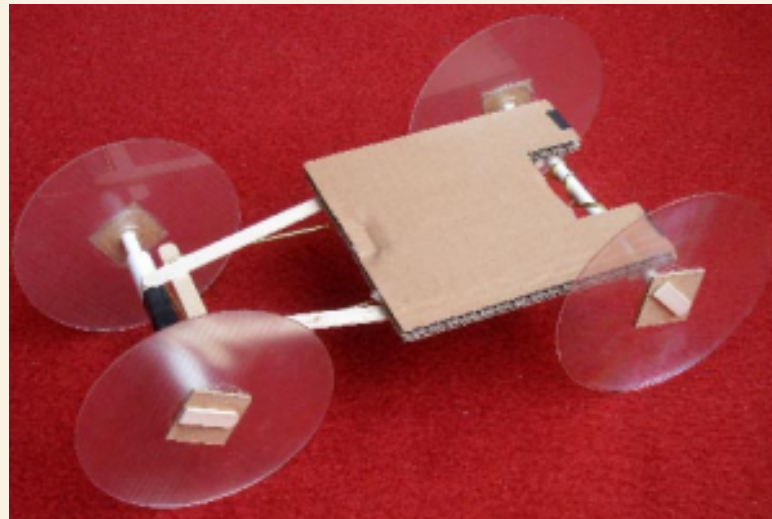


Inflating with a rubber band

Very simple analogy: think of flux lines as a stretched rubber band, wound around the compact dimension(s)

At a certain point, somebody creates a hole in the band: the rubber band begins to unwind!

To some observer not sensitive to the compact dimension(s), this is just a decrease in potential energy



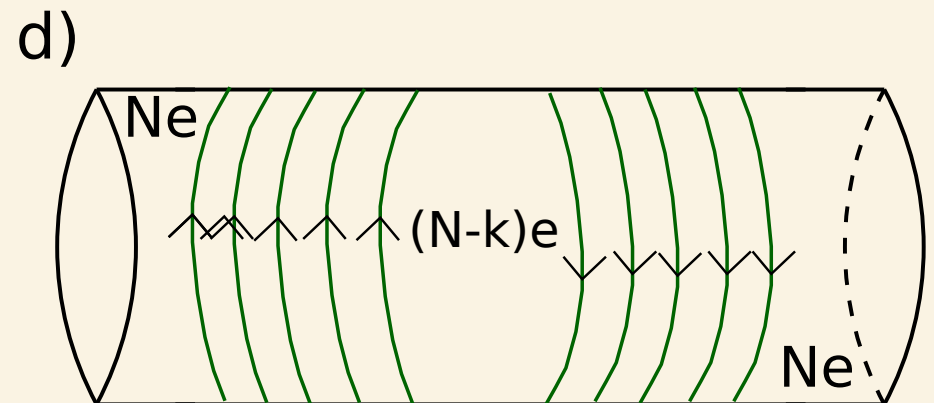
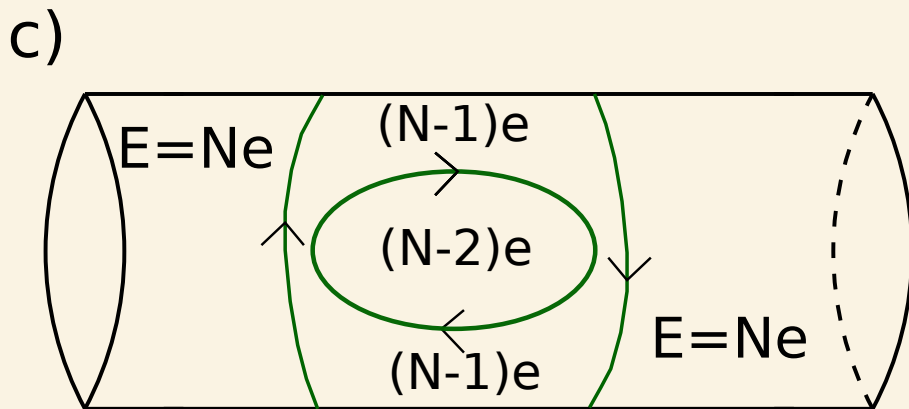
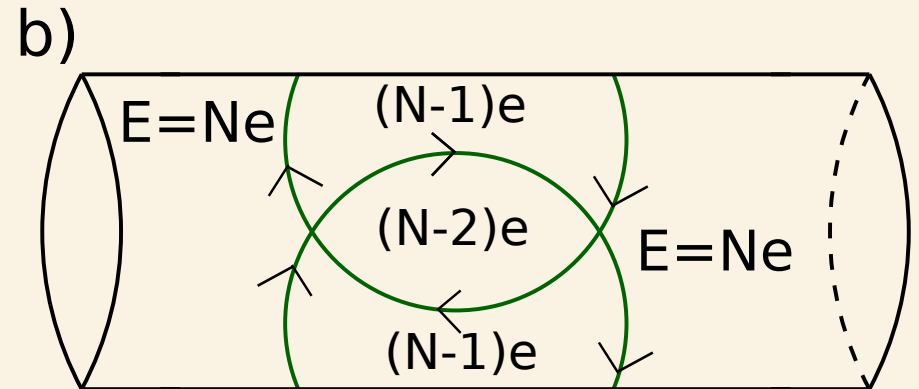
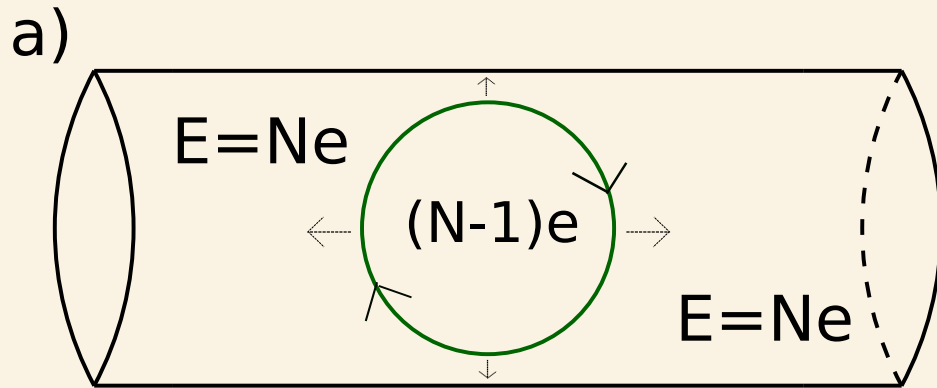
Let's (really) make inflation

- Basic setup: spacetime of the form $dS_4 \times M$.

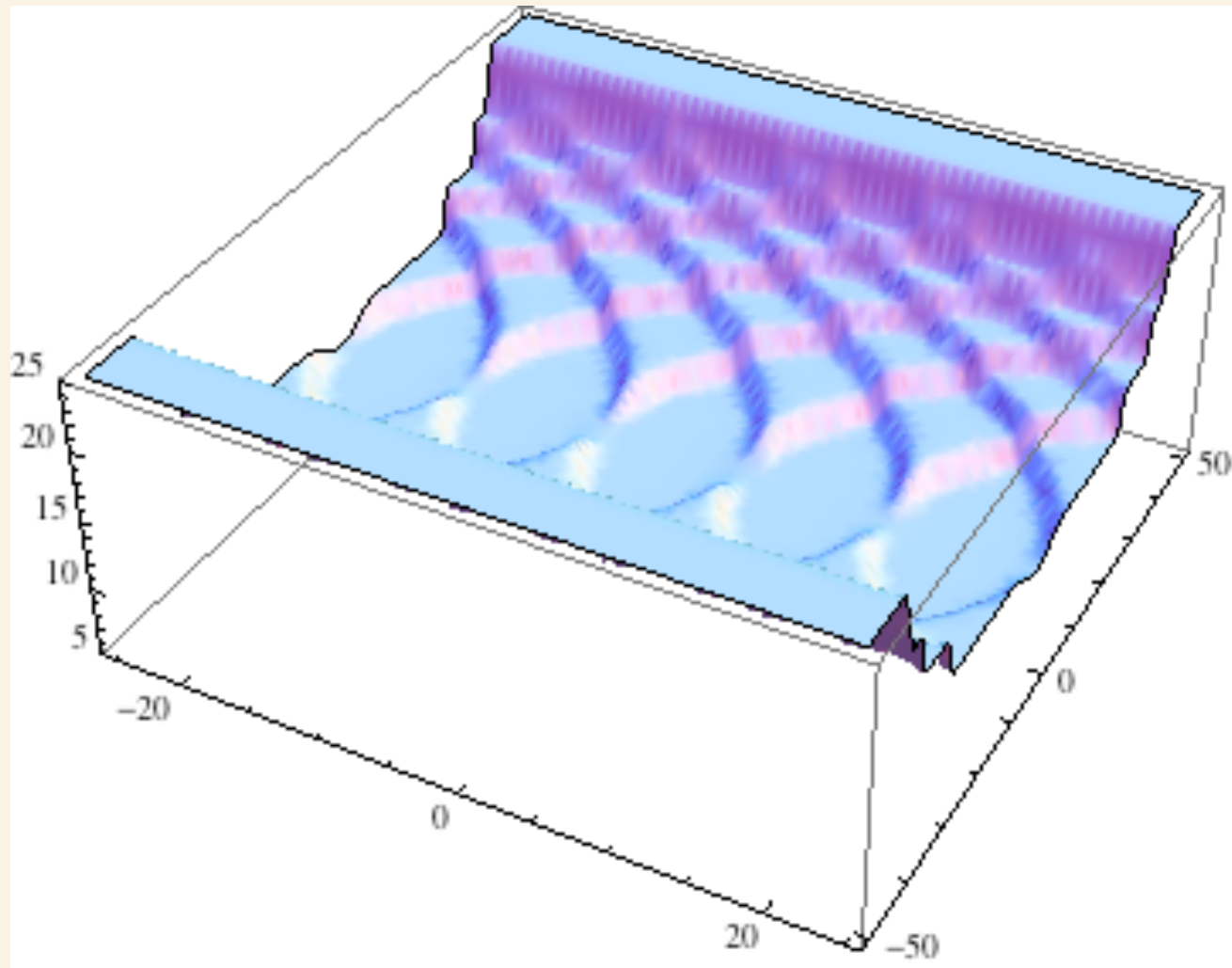
Hubble constant is determined by the amount of flux.

- Configuration unstable to (QM) nucleation of a brane bubble
- The bubble expands in the dS directions and collides with itself in the compact directions, discharging the flux one unit at a time
- The cascade ends when the brane annihilates \rightarrow reheating

Flux cascade in more dim's



Animated evolution



How it works

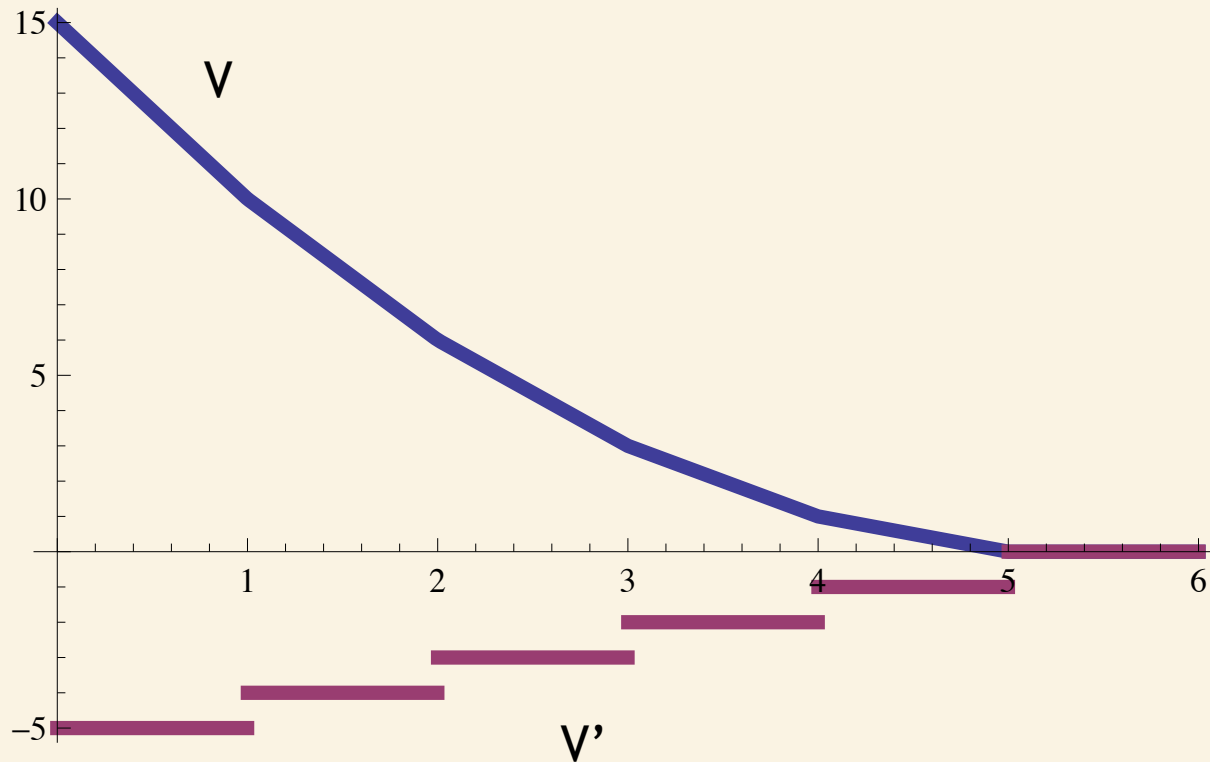
- The bubble contains a *homogeneous and isotropic slightly open FRW universe*. Expansion in dS directions inflates away the curvature!
- A crucial point: *collisions happen at instants of FRW time*, which preserve the full $SO(1,3)$ symmetry of open FRW
- This is an exit from the FVEI, which produces a homogeneous and flat universe
- How rare bubble nucleations are is irrelevant, because *the cascade is classical and doesn't stop once it starts*
- Reheating occurs naturally near zero flux, when the brane slows down and can self-annihilate (by tachyon condensation)

Effective action

4d effective inflaton is the brane separation.

Effective 4d action for simplest model:

$$S_4 = - \int d^3\Omega dt \sinh^3 t \left[2\sigma \sqrt{1 - \dot{z}_0^2} + V(z_0) \right]$$



Fluctuations

The brane separation in the compact dimensions $z(x)$ is a 4d light scalar field, which determines when reheating happens.

There are two sources of perturbations in z :

- de Sitter quantum fluctuations, as in ordinary inflation
- Brane self-collisions produce open strings, and variations in the density of these cause perturbations in z

Power spectrum:

$$\Delta_{\mathcal{R}}^2 = \frac{H^4}{8\pi^2 \sigma \dot{z}_0^2}$$

Tilt:

$$n_s - 1 \sim -\frac{2}{N_e} \sim -0.03$$

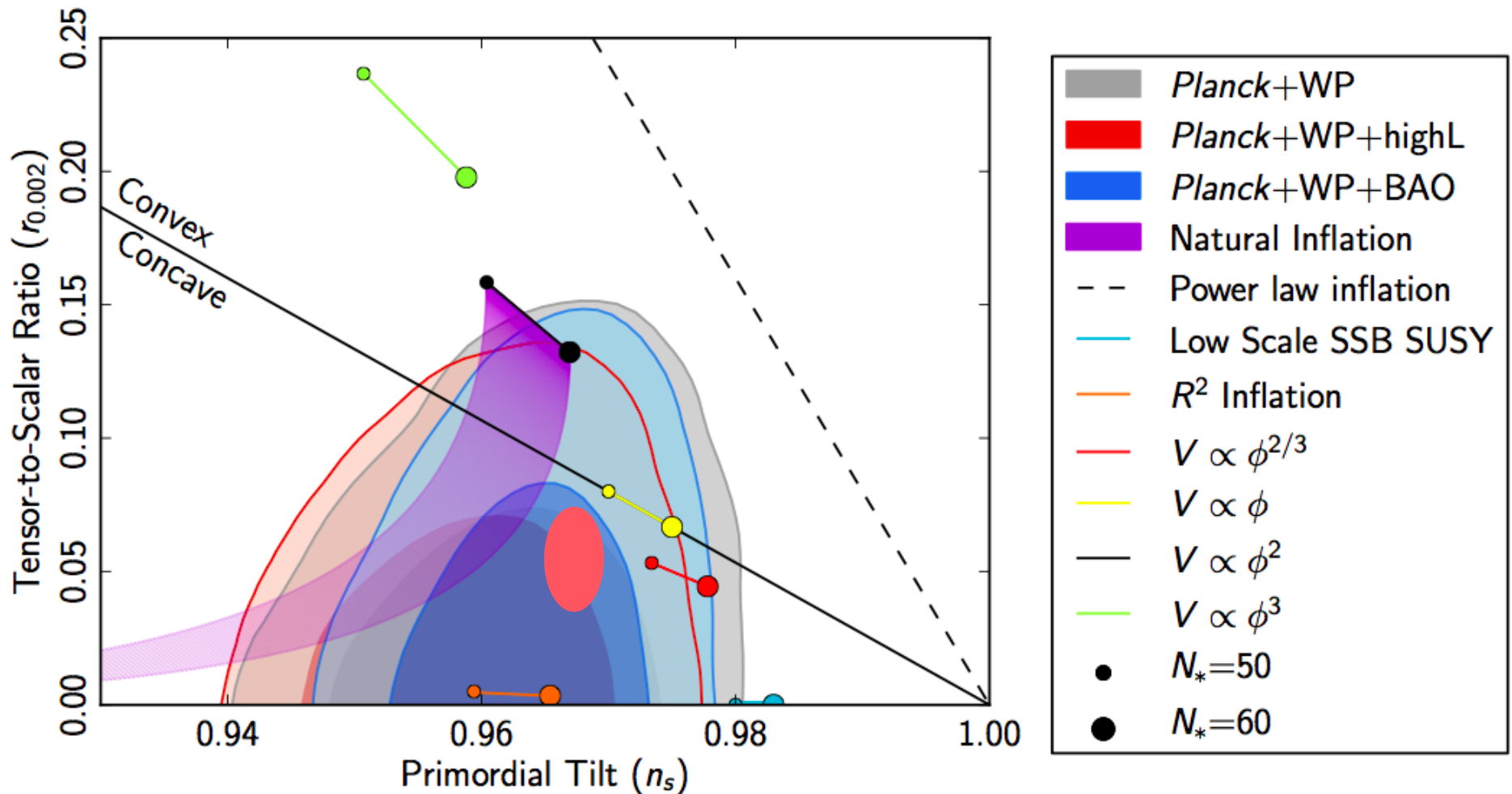
Tensors:

$$\Delta_h^2 = \frac{2H^2}{\pi^2 M_{\text{Pl}}^2}$$

Tensor-to-scalar ratio potentially observable

$$r \sim \frac{1}{Q} \sim 10^{-2}$$

Where are we?



What about non-Gaussianity?

Main source of NG is the DBI kinetic term.

This gives ~equilateral shape and we expect $f_{\text{nl}} \sim 1/c_s^2 \sim \gamma^2$

Given Planck's results, $c_s > 0.07$, we need $\gamma < 14$

Difficult to analyze NG due to fluctuations in string density.

If they are important, **we can expect a folded shape**, as in models with particle production (due to negative frequency modes)

In general, except for the codimension-1 case, we expect **additional light scalars** b_i which describe the position of the branes in transverse dimensions.

They determine the reheating moment, which could translate into local NG.

However, not if the inflationary trajectory has $\langle b_i \rangle = 0$ at reheating

Conclusions and future work

- Inflation is the best model we have for the early universe dynamics, and FVEI a generic prediction of String Theory
- Unwinding inflation is a model which naturally realizes slow-roll type inflation, starting from a FVEI landscape
If you want, a graceful entry scenario!
- At the EFT level, it encompasses different effective models present in the literature (DBI, dissipative, oscillating), but with no fine-tuning of parameters
- Need to build a detailed model in the context of string theory

Thank you!