

# Gauge Theories for Baryon and Lepton Numbers

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MAX-PLANCK-GESELLSCHAFT

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ICTP, Trieste, October 2-5, 2013*



MAX-PLANCK-INSTITUT  
FÜR KERNPHYSIK

## Collaborators:

2009 - M. B. Wise (Caltech, USA)

2013 - M. Duerr (MPIK, Germany)

2013 - M. Lindner (MPIK, Germany)

## References:

P. F. P., M. B. Wise,  
Phys. Rev. D 82, 011901 (2010);  
JHEP 08 (2011) 068



M. Duerr, P. F. P., M. B. Wise,  
Physical Review Letters 110,231801 (2013)

M. Duerr, P. F. P., M. Lindner,  
Phys. Rev. D 88, 051701(R) (2013)

M. Duerr, P. F. P., 1309.3970

# The Desert Hypothesis in Particle Physics

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**B and L Violation:**

Seesaw Camel

$$\frac{c}{\Lambda^2} QQQQL \quad (\tau_p > 10^{32-34} \text{ years} \implies \Lambda > 10^{15} \text{ GeV})$$

Standard Model

$$\Lambda_{\text{Weak}} \sim 100 \text{ GeV}$$

GUTs, Strings ?

$$\Lambda \sim 10^{15-19} \text{ GeV}$$

# Can we break B and L at the TeV scale?

LOW  
SCALE

HIGH  
SCALE

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# Aim

Theory for Baryon and Lepton Numbers



Cosmology:

Dark Matter  
Baryogenesis



Neutrino  
Masses



Signatures at  
the LHC

# Outline

- Introduction
- Living without the Great Desert
- Left-Right Symmetry and Type III Seesaw
- Summary

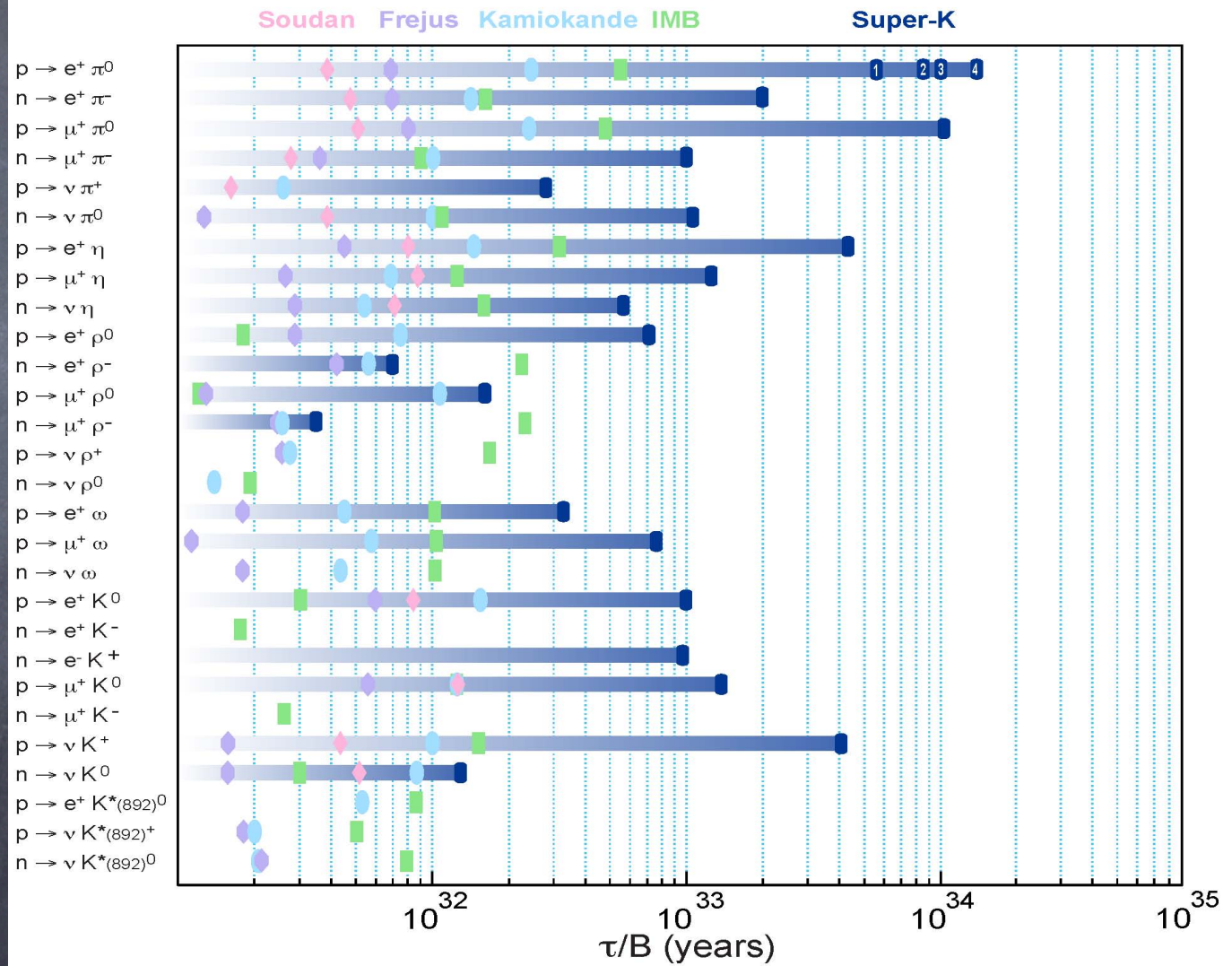
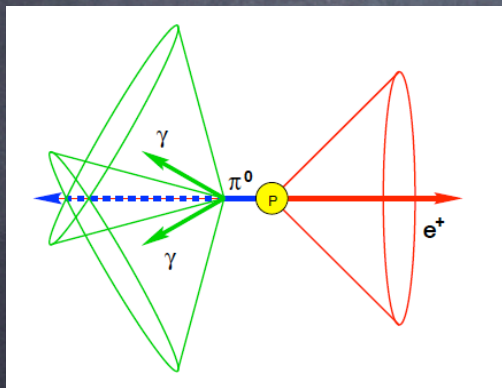
# Introduction



# Experimental Results

- Proton Decay:

$$\Delta B = 1, \Delta L = \text{odd}$$



- Neutrino Oscillations:

$$\Delta L_e \neq 0, \Delta L_\mu \neq 0, \Delta L_\tau \neq 0$$

## Cosmology

Baryon Asymmetry:

$$\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-10}$$



**BARYON NUMBER VIOLATION**

$$\Delta B \neq 0$$

Sakharov's Condition, 1967

# Living without the Great Desert

# Aim: Can we break B and L at the TeV scale?

LOW  
SCALE

HIGH  
SCALE

**B and L Violation:**

Seesaw Camel

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GUTs, Strings ?

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# B and L as Local Symmetries

A. Pais, 1973 (B as a Local Symmetry)

S. Rajpoot, 1987; Foot, Joshi, Lew, 1989

Carone, Murayama, 1995

**Breaking Local Baryon and Lepton Numbers at the TeV Scale (NO Desert !!)**

[P. F. P.](#), M.B. Wise, 2010

[P. F. P.](#), M. B. Wise, JHEP 1108(2011) 068

M. Duerr, [P. F. P.](#), M. B. Wise, arXiv:1304.0576 (Phys. Rev. Lett. 2013)

## Breaking B and L at the TeV scale !



$$G_{SM} \otimes U(1)_B \otimes U(1)_L$$

where:  $U(1)_B$  and  $U(1)_L$  can be broken at the TeV Scale !

$$Q_L \sim (3, 2, 1/6, 1/3, 0), \quad u_R \sim (3, 1, 2/3, 1/3, 0), \quad d_R \sim (3, 1, -1/3, 1/3, 0)$$

$$\ell_L \sim (1, 2, -1/2, 0, 1), \quad e_R \sim (1, 1, -1, 0, 1), \quad \nu_R \sim (1, 1, 0, 0, 1)$$

How to define an anomaly free theory ?

# Anomalies Cancellation

**Baryonic Anomalies:**  $\mathcal{A}_1 (SU(3)^2 \otimes U(1)_B)$ ,  $\mathcal{A}_2 (SU(2)^2 \otimes U(1)_B)$ ,  
 $\mathcal{A}_3 (U(1)_Y^2 \otimes U(1)_B)$ ,  $\mathcal{A}_4 (U(1)_Y \otimes U(1)_B^2)$ ,  
 $\mathcal{A}_5 (U(1)_B)$ ,  $\mathcal{A}_6 (U(1)_B^3)$ ,

**Leptonic Anomalies:**  $\mathcal{A}_7 (SU(3)^2 \otimes U(1)_L)$ ,  $\mathcal{A}_8 (SU(2)^2 \otimes U(1)_L)$ ,  
 $\mathcal{A}_9 (U(1)_Y^2 \otimes U(1)_L)$ ,  $\mathcal{A}_{10} (U(1)_Y \otimes U(1)_L^2)$ ,  
 $\mathcal{A}_{11} (U(1)_L)$ ,  $\mathcal{A}_{12} (U(1)_L^3)$ ,

**Mixed:**  $\mathcal{A}_{13} (U(1)_B^2 \otimes U(1)_L)$ ,  $\mathcal{A}_{14} (U(1)_L^2 \otimes U(1)_B)$ ,  
 $\mathcal{A}_{15} (U(1)_Y \otimes U(1)_L \otimes U(1)_B)$ ,

**In the SM:**  $\mathcal{A}_2 = -\mathcal{A}_3 = 3/2$        $\mathcal{A}_8 = -\mathcal{A}_9 = 3/2$

P. F. P., M. B. Wise, PRD82 (2010)011901; JHEP1108(2011)068

## Possible Solutions

- Sequential Family ( $B = -1, L = -3$ )
- Mirror Family ( $B = 1, L = 3$ )
- Vector-like Family with Seesaw

Now they are in disagreement with LHC Constraints !

What about Fermionic Leptoquarks ?



M. Duerr, **P. F. P.**, M. B. Wise, arXiv:1304.0576 (Phys. Rev. Lett. 2013)

One can define an anomaly free theory using the **Fermionic Lepto-quarks**:

$$\Psi_L \sim (1, 2, -1/2, -3/2, -3/2), \quad \Psi_R \sim (1, 2, -1/2, 3/2, 3/2)$$

$$\eta_L \sim (1, 1, -1, 3/2, 3/2), \quad \eta_R \sim (1, 1, -1, -3/2, -3/2)$$

$$\chi_L \sim (1, 1, 0, 3/2, 3/2), \quad \chi_R \sim (1, 1, 0, -3/2, -3/2)$$

**They can have vector-like masses and cancel all anomalies !**

M. Duerr, **P. F. P.**, M. B. Wise, arXiv:1304.0576 (Phys. Rev. Lett. 2013)

**Interactions:**

$$\begin{aligned} -\mathcal{L} \supset & h_1 \bar{\Psi}_L H \eta_R + h_2 \bar{\Psi}_L \tilde{H} \chi_R + h_3 \bar{\Psi}_R H \eta_L + h_4 \bar{\Psi}_R \tilde{H} \chi_L \\ & + \lambda_1 \bar{\Psi}_L \Psi_R S_{BL} + \lambda_2 \bar{\eta}_R \eta_L S_{BL} + \lambda_3 \bar{\chi}_R \chi_L S_{BL} \\ & + a_1 \chi_L \chi_L S_{BL} + a_2 \chi_R \chi_R S_{BL}^\dagger + \text{h.c.} \end{aligned}$$

$$-\mathcal{L}_\nu = Y_\nu \bar{\ell}_L \tilde{H} \nu_R + \frac{\lambda_R}{2} \nu_R \nu_R S_L + \text{h.c.}$$



Higgses:  $S_{BL} \sim (1, 1, 0, -3, -3), \quad S_L \sim (1, 1, 0, 0, -2)$



M. Duerr, [P. F. P.](#), M. B. Wise, arXiv:1304.0576 (Phys. Rev. Lett. 2013)

## Some Features:

**Symmetry Breaking:**  $S_{BL} \sim (1, 1, 0, -3, -3)$ ,  $S_L \sim (1, 1, 0, 0, -2)$

$$\Delta B = \pm 3, \Delta L = \pm 2, \Delta L = \pm 3$$

**NO Proton Decay !**

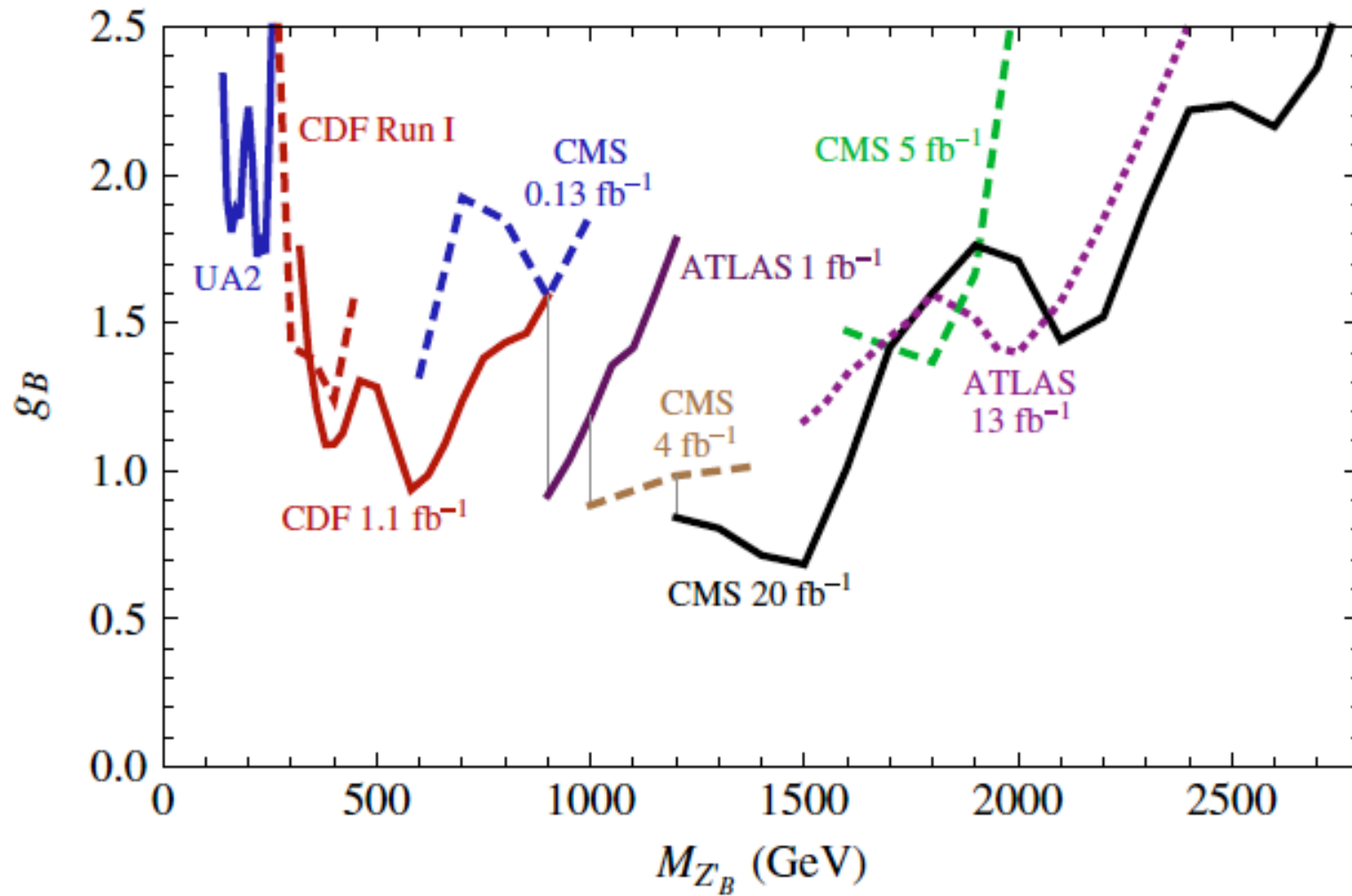
**NO DESERT !**

**Dark Matter:**  $\Psi_{LF}^0$  can be a cold dark matter candidate !

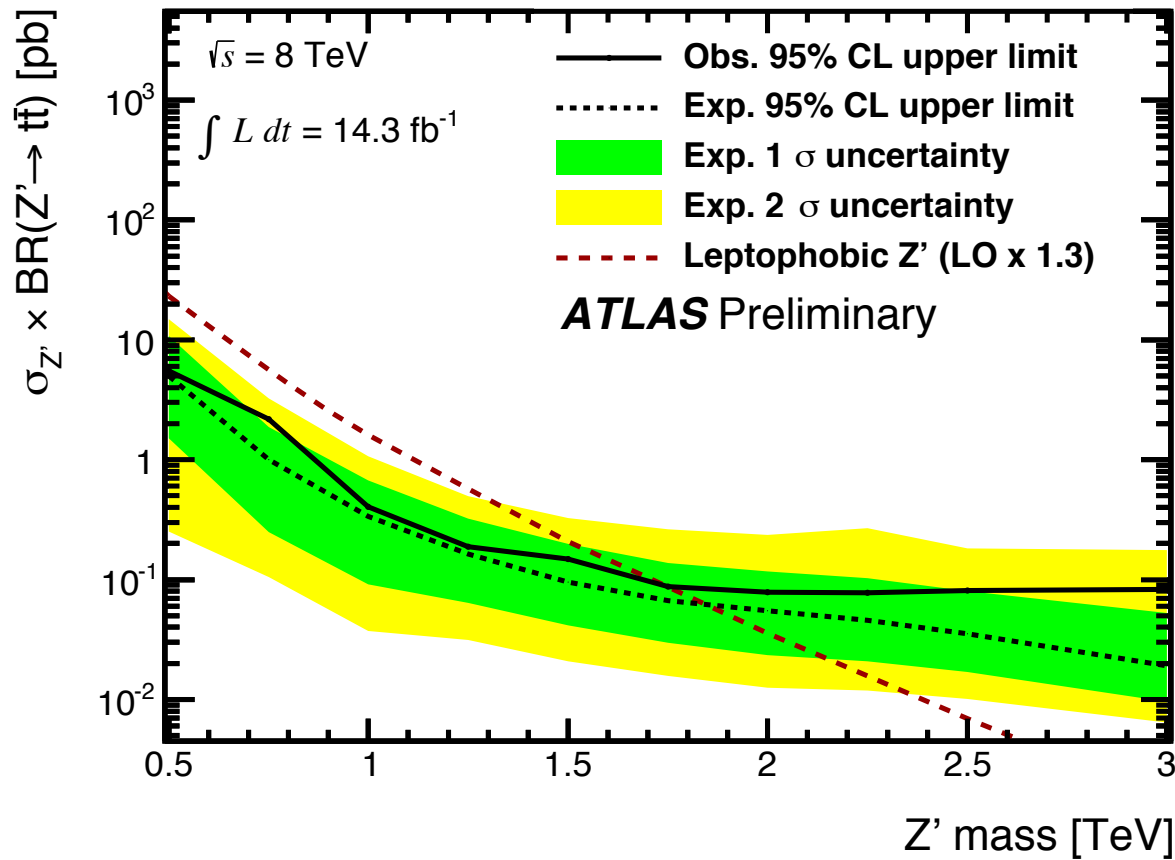
**NO extra Flavour violation !**

**New Gauge Bosons:**  $Z_L, Z_B$

## Bounds on the Baryonic Breaking Scale !

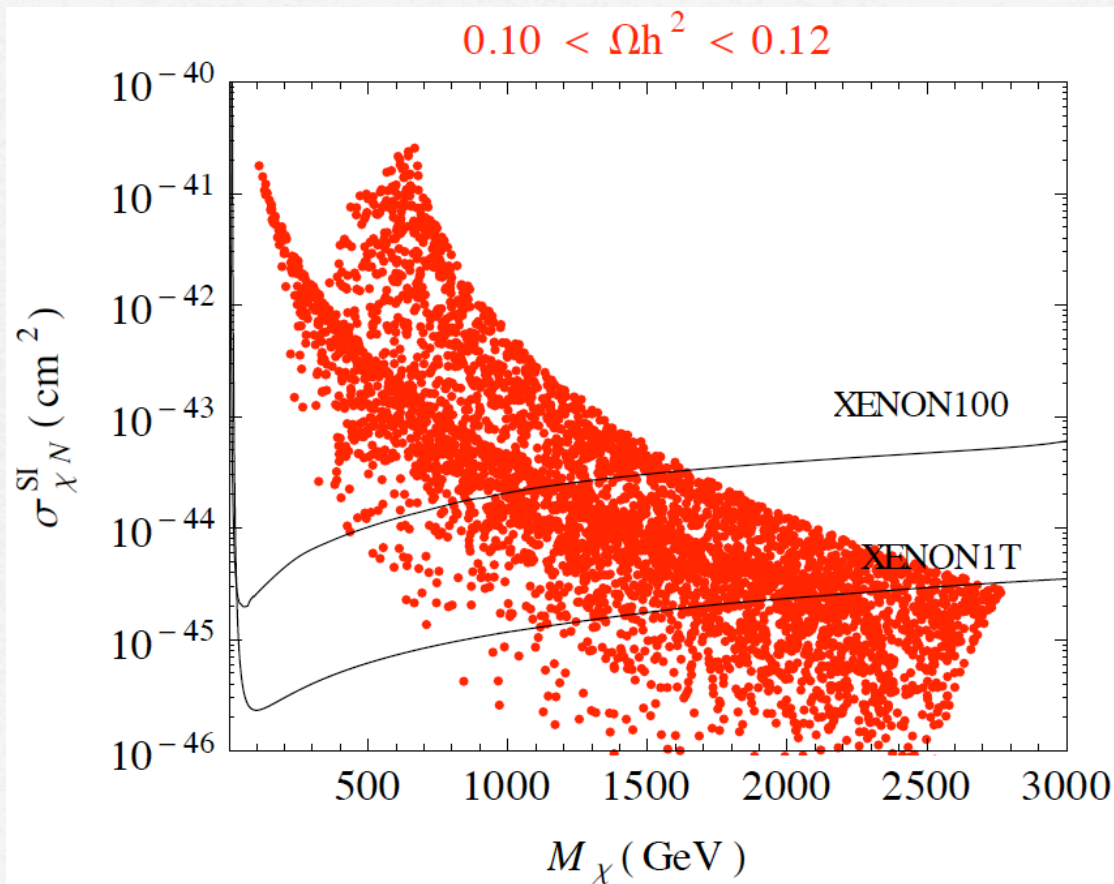


May 13, 2013



(a)  $Z'$  upper cross section limits.

# Baryonic Dark Matter



$g_B, M_{Z_B}, M_\chi$  &  $B$

Annihilation:

$$\bar{\chi}\chi \rightarrow Z_B \rightarrow \bar{q}q$$

Direct Detection:

$$\chi N \rightarrow Z_B \rightarrow \chi N$$

LHC Signatures:

$$pp \rightarrow Z_B \rightarrow \bar{\chi}\chi, \bar{q}q, \bar{\chi}\chi j, \dots$$

# Left-Right Symmetry and Type III Seesaw

# Again: Can we break B and L at the TeV scale?

L  
O  
W  
S  
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E

H  
I  
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H  
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E

**B and L Violation:**

Seesaw Camel

$$\frac{c}{\Lambda^2} QQQQL \quad (\tau_p > 10^{32-34} \text{ years} \implies \Lambda > 10^{15} \text{ GeV})$$

Left-Right Symmetric Theory

GUTs, Strings ?

$$\Lambda \sim 10^{15-19} \text{ GeV}$$



Pati, Salam; Pati, Mohapatra; Senjanovic, Mohapatra

## Left-Right Symmetry

$$SU(3)_C \otimes SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$

- Connection between Neutrino Masses and the Scale of Parity Violation
- Minimal Model has Type I and Type II Seesaw Mechanisms
- Doorway to SO(10) Unification
- If the scale is low one has 'exotic' signals at the LHC

He, Rajpoot, 1990

M. Duerr, [P. F. P.](#), M. Lindner, 1306.0568 (PRD)



$$SU(2)_L \otimes SU(2)_R \otimes U(1)_B \otimes U(1)_L$$

SM Fermions:

$$Q_L \sim (2, 1, 1/3, 0), \quad Q_R \sim (1, 2, 1/3, 0), \quad \ell_L \sim (2, 1, 0, 1), \quad \ell_R \sim (1, 2, 0, 1)$$

Anomalies:

$$\begin{aligned} \mathcal{A}_1 (SU(2)_L^2 \otimes U(1)_B) &= 3/2, \\ \mathcal{A}_2 (SU(2)_L^2 \otimes U(1)_L) &= 3/2, \\ \mathcal{A}_3 (SU(2)_R^2 \otimes U(1)_B) &= -3/2, \\ \mathcal{A}_4 (SU(2)_R^2 \otimes U(1)_L) &= -3/2. \end{aligned}$$

The Simplest Solution is .....



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## Type III Seesaw Fields

$$\rho_L \sim (3, 1, -3/4, -3/4), \quad \& \quad \rho_R \sim (1, 3, -3/4, -3/4),$$

$$\rho_L = \frac{1}{2} \begin{pmatrix} \rho_L^0 & \sqrt{2}\rho_L^+ \\ \sqrt{2}\rho_L^- & -\rho_L^0 \end{pmatrix} \quad \text{and} \quad \rho_R = \frac{1}{2} \begin{pmatrix} \rho_R^0 & \sqrt{2}\rho_R^+ \\ \sqrt{2}\rho_R^- & -\rho_R^0 \end{pmatrix}.$$

The theory is anomaly free !

**Higgs Sector:**  $\Phi \sim (2, 2, 0, 0)$

$$H_L \sim (2, 1, 3/4, -1/4), \quad H_R \sim (1, 2, 3/4, -1/4)$$

$$S_{BL} \sim (1, 1, 3/2, 3/2)$$

**Relevant Interactions for Neutrino Masses:**

$$\begin{aligned} -\mathcal{L} \supset & \bar{\ell}_L \left( Y_3 \Phi + Y_4 \tilde{\Phi} \right) \ell_R \\ & + \lambda_D \left( \ell_L^T C i \sigma_2 \rho_L H_L + \ell_R^T C i \sigma_2 \rho_R H_R \right) \\ & + \lambda_\rho \text{Tr} \left( \rho_L^T C \rho_L + \rho_R^T C \rho_R \right) S_{BL} + \text{h.c.}, \end{aligned}$$

# Type III Seesaw and Left-Right Symmetry

Parity Violation !

$$v_L \ll v_R$$

$$M_{\nu L}^{III} \ll M_{\nu R}^{III}.$$

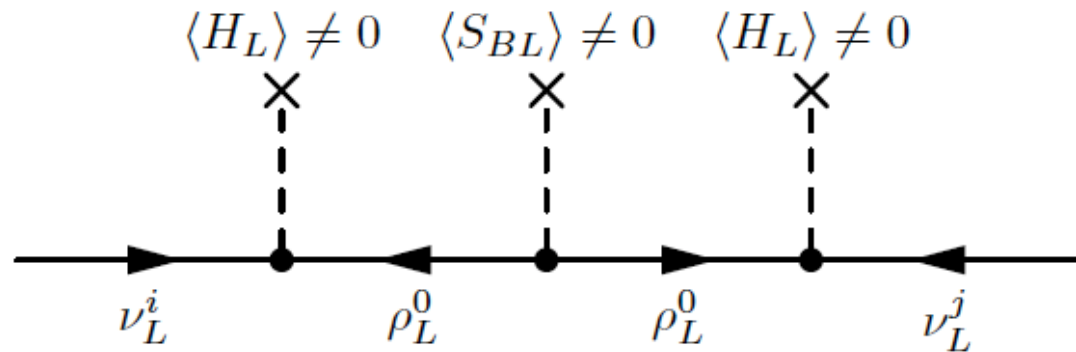


FIG. 1: Type III seesaw for the left-handed neutrinos.

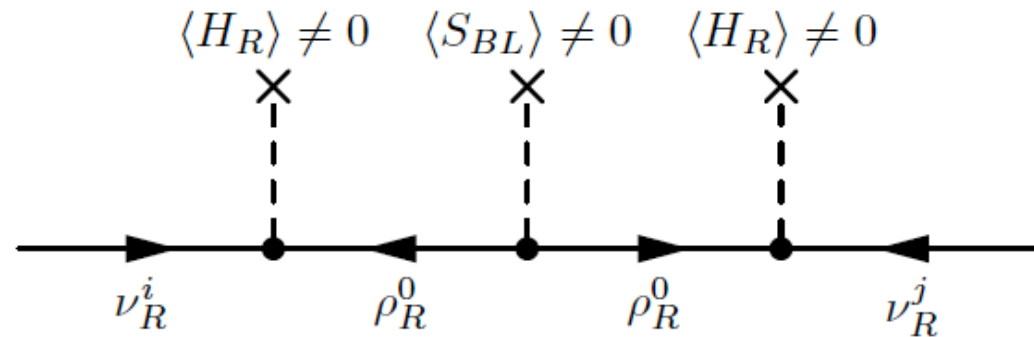


FIG. 2: Type III seesaw for the right-handed neutrinos.

M. Duerr, P. F. P., M. Lindner, 1306.0568 (PRD)

## Neutrino Masses

$$-\mathcal{L}_\nu = \tilde{M}_\nu^D \bar{\nu}_L \nu_R - \frac{1}{2} M_{\nu L}^{III} \nu_L^T C \nu_L - \frac{1}{2} M_R \nu_R^{3T} C \nu_R^3 + \text{h.c.},$$

3 + 2 System

$$-\mathcal{L}_\nu = -\frac{1}{2} M_{\nu L}^{LL} \nu_L^T C \nu_L + \left( \tilde{M}_\nu^D \right)^{i\alpha} \bar{\nu}_L^i \nu_R^\alpha + \text{h.c.},$$

$$\mathcal{M}_\nu^{3+2} = \begin{pmatrix} 0 & 0 & 0 & m_D^1 & m_D^2 \\ 0 & m_1 & 0 & m_D^3 & m_D^4 \\ 0 & 0 & m_2 & m_D^5 & m_D^6 \\ m_D^1 & m_D^3 & m_D^5 & 0 & 0 \\ m_D^2 & m_D^4 & m_D^6 & 0 & 0 \end{pmatrix}.$$

Two light sterile neutrinos at the renormalizable level !!

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M. Duerr, P. F. P., M. Lindner, 1306.0568 (PRD)

## Effect of higher-dimensional operators

$$\mathcal{O}_{\nu_L} = c_L \ell_L \ell_L H_L H_L S_{BL}^\dagger / \Lambda^2$$

$$\mathcal{O}_{\nu_R} = c_R \ell_R \ell_R H_R H_R S_{BL}^\dagger / \Lambda^2.$$

Example:

if  $v_L \sim 1$  GeV, and  $v_{BL} \sim 10$  TeV  $\rightarrow \Lambda \gtrsim 3 \times 10^3$  TeV

Using  $v_R \sim 1$  TeV and  $v_{BL} \sim 10$  TeV  $\rightarrow M_{\nu_R} < 1$  MeV.

# Summary

- The Desert Hypothesis plays a major role in our view of the relation between the physics at the low and high scales. However, **this picture can be WRONG !**
- One can define a consistent theory where B and L are local symmetries broken at the low scale in agreement with the experiments and **there is no need to postulate the Great Desert.** One has a simple theory for dark matter (and baryogenesis) which can be tested at LHC.
- Local B and L Symmetries together with Left-Right Symmetry requires Type III Seesaw. The Minimal Model predicts light sterile neutrinos at the renormalizable level.



# THANK YOU !

This is my contribution to the GoranFest which took place in 2010 !

Sorry, I am a bit late ;)