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Leptonic CP Violation in Neutrino Oscillations

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Based on T. Ohlsson, H. Zhang and S.Z., Phys. Rev. D 87 (2013) 013012; Phys. Rev. D 87 (2013) 053006; Phys. Rev. D 88 (2013) 013001.

From Majorana to LHC: Workshop on the Origin of Neutrino Mass

ICTP, Trieste, Oct. 2 - 5, 2013



Outline

- Introduction
- RG running of Dirac CP-violating phase
- Leptonic CP violation in v oscillations
- NSI Effects @ IceCube (DeepCore & PINGU)
- Summary



Open Questions in v Physics

• Are neutrinos Dirac or Majorana particles?

Lepton number violation, neutrinoless double beta decays

- What is the neutrino mass hierarchy? Normal $(m_1 < m_2 < m_3)$ or inverted $(m_3 < m_1 < m_2)$?
- What is the absolute neutrino mass scale?

Is the lightest v massless? Hierarchical or degenerate?

- What is the origin of neutrino masses and flavor mixing? Seesaw mechanisms, flavor symmetries, ...
- Is there CP violation in the lepton sector?

What is the value of the Dirac CP-violating phase δ ?



Current Status

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Precision measurements of neutrino parameters:

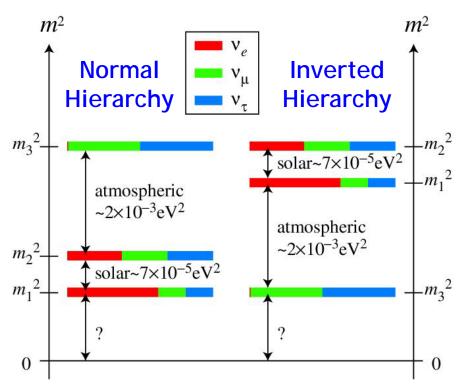
$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
Global-fit analyses:
Forero *et al.*, Phys. Rev. D
86 (2012) 073012
Fogli *et al.*, Phys. Rev. D
86 (2012) 013012

		ctor / erator	Solar / Reactor	Gonzalez-Garc JHEP 1212 (20	,
		bfp $\pm 1\sigma$			3σ range
$\sin^2 \theta_{12}$		0.30 ± 0.013	Pascoli & Schw	etz	$0.27 \rightarrow 0.34$
$ heta_{12}/^{\circ}$		33.3 ± 0.8	Adv. HEP, 2013		$31 \rightarrow 36$
$\sin^2\theta_{23}$		$0.41^{+0.037}_{-0.025} \oplus 0.59$	+0.021 -0.022		$0.34 \rightarrow 0.67$
$\theta_{23}/^{\circ}$		$40.0^{+2.1}_{-1.5} \oplus 50.4$	+1.2 -1.3		$36 \rightarrow 55$
$\sin^2 \theta_{13}$		0.023 ± 0.002	23		$0.016 \ \rightarrow \ 0.030$
$\theta_{13}/^{\circ}$		$8.6^{+0.44}_{-0.46}$			$7.2 \rightarrow 9.5$
$\delta/^{\circ}$	Leptonic CP violation?	300^{+66}_{-138}			$0 \rightarrow 360$
$\Delta m_{21}^2 / 10^{-5} \mathrm{eV}^2$		7.50 ± 0.185	5		$7.00 \rightarrow 8.09$
$\Delta m_{31}^2 / 10^{-3} \mathrm{eV}^2$ (NH)		$2.47^{+0.069}_{-0.067}$	Neutrino mass	hierarchv2	$2.27 \rightarrow 2.69$
$\Delta m_{32}^2 / 10^{-3} \mathrm{eV}^2$ (IH)		$-2.43^{+0.042}_{-0.065}$			$-2.65 \rightarrow -2.24$



Neutrino Mass Hierarchy

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- Medium-baseline reactor experiments: JUNO, RENO-50

- Long-baseline accelerator experiments: T2K, NOvA, LBNO, LBNE
- Huge neutrino telescopes: PINGU, ORCA





Leptonic CP Violation

Gonzalez-Garcia et al.

Branco et al., Rev. Mod.

Forero et al.

Phys. 84 (2012) 515

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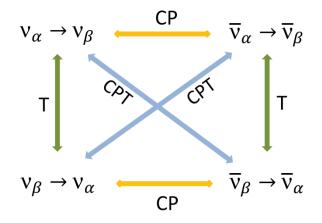
Parameter

Neutrino oscillations in vacuum:

$$A_{\alpha\beta}^{\rm CP} \equiv P(\nu_{\alpha} \rightarrow \nu_{\beta}) - P(\overline{\nu}_{\alpha} \rightarrow \overline{\nu}_{\beta})$$

Fogli et al.

$$=16\underbrace{s_{12}c_{12}s_{23}c_{23}s_{13}c_{13}^{2}\sin\delta}_{I:\text{ Jarlskog Invariant}}\sin\frac{\Delta m_{21}^{2}L}{4E}\sin\frac{\Delta m_{32}^{2}L}{4E}\sin\frac{\Delta m_{31}^{2}L}{4E}$$



- First hint from global fits?

	J		
$\sin^2\theta_{12}$	0.307	0.300	0.320
$\delta m \theta_{12}$	0.291-0.325	0.287-0.313	0.303-0.336
$\sin^2\theta_{13}$	0.0241	0.0230	0.0246
$\sin^2\theta_{13}$	0.0216-0.0266	0.0207-0.0253	0.0218-0.0275
$ain^2 0$	0.386	0.410	0.427
$\sin^2\theta_{23}$	0.365-0.410	0.385-0.447	0.400-0.461
$\Lambda = \frac{2}{10^{-5}} \sqrt{2}$	7.54	7.50	7.62
$\Delta m_{21}^2 / 10^{-5} \text{ eV}^2$	7.32-7.80	7.32–7.69	7.43-7.81
$\Lambda = \frac{2}{10^{-3}} \cdot \frac{10^{-3}}{10^{-3}}$	2.51	2.47	2.55
$\Delta m_{31}^2 / 10^{-3} \text{ eV}^2$	2.41-2.57	2.40-2.54	2.46-2.61
S/	1.08	1.67	0.8
δ/π	0.77-1.36	0.90-2.03	0–2.0

- Predictions from theories
- Compare between theory
 & exp. observation of δ:
 Radiative corrections

- Optimize the exp. setup: CP measures, v-oscillogram

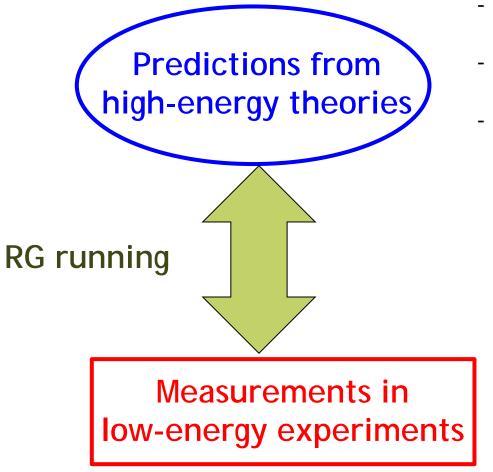
- E.g., NSI effects in the ice



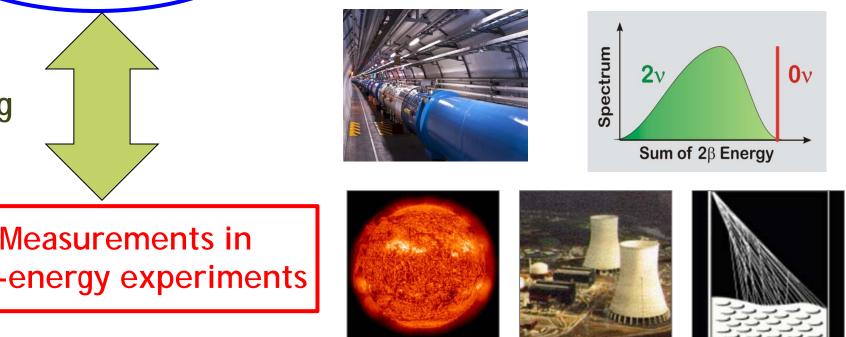
RG Running of Neutrino Parameters

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Why is RG running important?



- Grand Unified Theories: SU(5), SO(10), ...
- Extra-dimensional models: ADD, UED, ...
- Flavor symmetry models: A₄, S₄, ...





RG Running of Neutrino Parameters

The SM as an effective theory: dimension-5 operator Weinber 43 (1970

Weinberg, Phys. Rev. Lett. 43 (1979) 1566

 $\mathcal{L}_{\nu} = \frac{1}{2} (\overline{\ell} H) \cdot \kappa \cdot (H^{T} \ell^{C})$ $[\Lambda]^{-1} : \text{high energy scale } \Lambda$

Renormalization Group Equation:

$$16\pi^2 \frac{\mathrm{d}\kappa}{\mathrm{d}t} = \alpha_{\kappa} + C_{\kappa} [(Y_l Y_l^{\dagger})\kappa + \kappa (Y_l Y_l^{\dagger})^T]$$

$$t = \ln\left(\frac{\mu}{\Lambda_{\rm EW}}\right) \qquad C_{\kappa} = -\frac{3}{2}$$
$$\alpha_{\kappa} = -3g_2^2 + \lambda + 2\mathrm{tr}\left[3\left(Y_{\rm u}Y_{\rm u}^+\right) + 3\left(Y_{\rm d}Y_{\rm d}^+\right) + \left(Y_{l}Y_{l}^+\right)\right]$$

Majorana neutrino mass matrix:

$$M_{\nu} = \kappa v^2$$

Masses, mixing angles, and leptonic CP-violating phases

Babu *et al.*, Phys. Lett. B 319 (1993) 191; Chankowski & Pluciennik, ibid., 316 (1993) 312

Antusch *et al.*, Phys. Lett. B 519 (2001) 238

- Realized in various seesaw models

- Running of masses is dominated by the flavor-diagonal term (gauge, quark Yuk.)

- RG running of mixing angles and CPviolating phases is dominated by charged-lepton Yukawa couplings



MSSM and 5D-UED

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Extensions of the SM: supersymmetry and extra dimensions

$$L_{\nu} = \frac{1}{2} L H_{u} \cdot \kappa \cdot L H_{u}$$

- v mass matrix:
- $M_v = \kappa (v \sin \beta)^2$

RGE coefficients:

 $C_{\kappa}^{\mathrm{MSSM}} = 1$

Tau Yukawa coupling dominated:

$$y_{\tau}^2 = m_{\tau}^2 (1 + \tan^2 \beta) / v^2$$

$$L_{\nu} = \frac{1}{2} LH \cdot \hat{\kappa} \cdot LH$$

v mass matrix: $M_{\nu} = \hat{\kappa} v^2 / (\hat{\kappa})$ radius of the extra spatial dimension RGE coefficients: $C_{\kappa}^{\text{UED}} = C_{\kappa}^{\text{SM}} (1 - s)$ number of excited KK modes Coefficient becomes larger at higher- $s \equiv \left\lfloor \frac{\mu}{\mu_0} \right\rfloor$



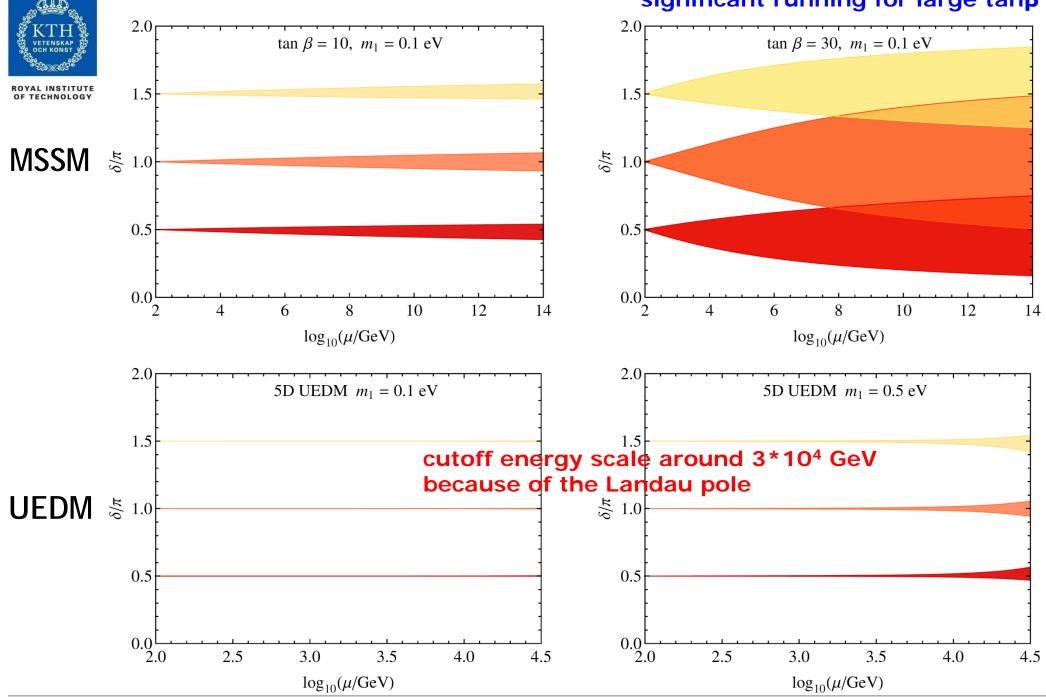
RG Running of Dirac CP-violating Phase

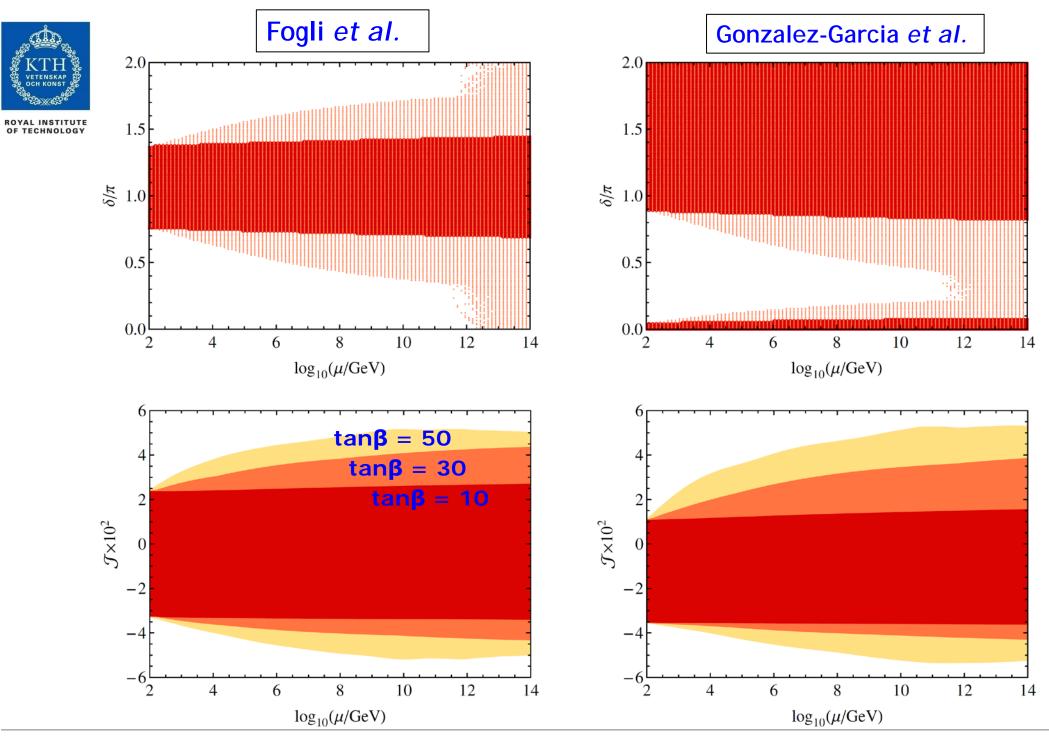
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In the standard parametrization:

- Keep all contributions of the same order.
- RG running depends crucially on the Majorana CP-violating phases; insignificant running for $\rho = \sigma$.
- RG running is in the opposite direction for the MSSM, compared to the SM and 5D-UED.

significant running for large $tan\beta$







Summary for RG Running of $\boldsymbol{\delta}$

The RGE of the leptonic Dirac CP-violating phase is derived, and small contributions of the same order are included.

Very tiny RG running effects in the SM, even for nearly-degenerate neutrino masses

> No significant RG running also in the MSSM and 5-UED, except for a very large tan β in the former case; however, note the dependence on the Majorana phases

> Non-zero δ can be radiatively generated at the low-energy scale even if $\delta = 0$ holds at a high-energy scale.

➢ For Dirac neutrinos, RG running is even smaller because of the absence of Majorana CP-violating phases.

> Constraints on δ from neutrino oscillation experiments can be directly applied to theory at a superhigh-energy scale.



CP Violation and Matter Effects

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Neutrino oscillations in matter:

$$H_{\rm eff} = \frac{1}{2E} \begin{bmatrix} U \begin{pmatrix} 0 & & & \\ & \Delta m_{21}^2 & \\ & & \Delta m_{31}^2 \end{pmatrix} U^+ + 2\sqrt{2}G_{\rm F}N_eE \begin{pmatrix} 1 & & \\ & 0 & \\ & & 0 \end{pmatrix} \end{bmatrix}$$

Oscillation probabilities:

Kimura *et al.*, Phys. Lett. B 537 (2012) 86

 $P_{\mu e}(\delta) \equiv P(v_{\mu} \to v_{e}) = a \cos \delta + b \sin \delta + c$ $\overline{P}_{\mu e}(\delta) \equiv P(\overline{v}_{\mu} \to \overline{v}_{e}) = \overline{a} \cos \delta + \overline{b} \sin \delta + \overline{c} \} \Rightarrow A_{\mu e}^{CP}(\delta) = \Delta a \cos \delta + \Delta b \sin \delta + \Delta c$

- Fake CP violation is induced by matter effects:
 - obscuring the intrinsic CP-violating effects by ${oldsymbol \delta}$
- How to describe leptonic CP violation?
 - working observables based on oscillation probabilities



Working Observables

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Series expansions of the probabilities in terms of $\alpha = \Delta m_{21}^2 / \Delta m_{31}^2$ and s_{13} :

$$\begin{aligned} a &\approx +8\alpha J_r \frac{\sin A\Delta}{A} \frac{\sin (A-1)\Delta}{A-1} \cos \Delta \\ b &\approx -8\alpha J_r \frac{\sin A\Delta}{A} \frac{\sin (A-1)\Delta}{A-1} \sin \Delta \\ c &\approx 4s_{13}^2 s_{23}^2 \frac{\sin^2 (A-1)\Delta}{(A-1)^2} \\ J_r &\equiv s_{12}c_{12}s_{23}c_{23}s_{13}c_{13}^2 \end{aligned} \qquad \begin{aligned} \Delta a &\approx +8\alpha J_r \Theta_- \frac{\sin A\Delta}{A} \cos \Delta \\ \Delta b &\approx -8\alpha J_r \Theta_+ \frac{\sin A\Delta}{A} \sin \Delta \\ \Delta b &\approx -8\alpha J_r \Theta_+ \frac{\sin A\Delta}{A} \sin \Delta \\ \Delta c &\approx 4s_{13}^2 s_{23}^2 \Theta_+ \Theta_- \\ \Theta_{\pm} &\equiv \frac{\sin (A-1)\Delta}{A-1} \pm \frac{\sin (A+1)\Delta}{A+1} \end{aligned}$$

Here $\Delta = \Delta m_{31}^2 L/4E$ and $A = VL/2\Delta$, with V being the matter potential, Δ the oscillation phase, and L the distance between the source and the detector.

$$\Delta P_{\mu e}^{\rm CP}(\delta) \& \Delta P_{\mu e}^{\rm m} \qquad \Delta A_{\mu e}^{\rm CP}(\delta) \& \Delta A_{\mu e}^{\rm m}$$

Definitions:

 $\Delta P_{\mu e}^{\rm CP}(\delta) \equiv P_{\mu e}(\delta) - P_{\mu e}(\delta = 0)$

 $\Delta P_{\mu e}^{\rm m} \equiv \max \left[P_{\mu e}(\delta) \right] - \min \left[P_{\mu e}(\delta) \right]$

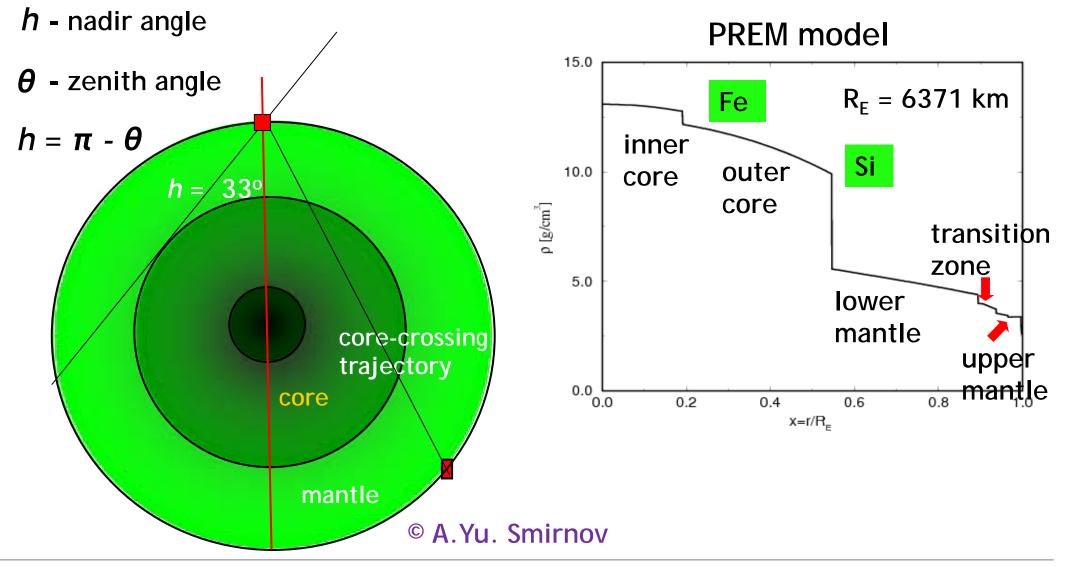
Definitions:

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$$\Delta A_{\mu e}^{\rm m} \equiv \max \left[A_{\mu e}^{\rm CP}(\delta) \right] - \min \left[A_{\mu e}^{\rm CP}(\delta) \right]$$



Neutrino Oscillograms of the Earth

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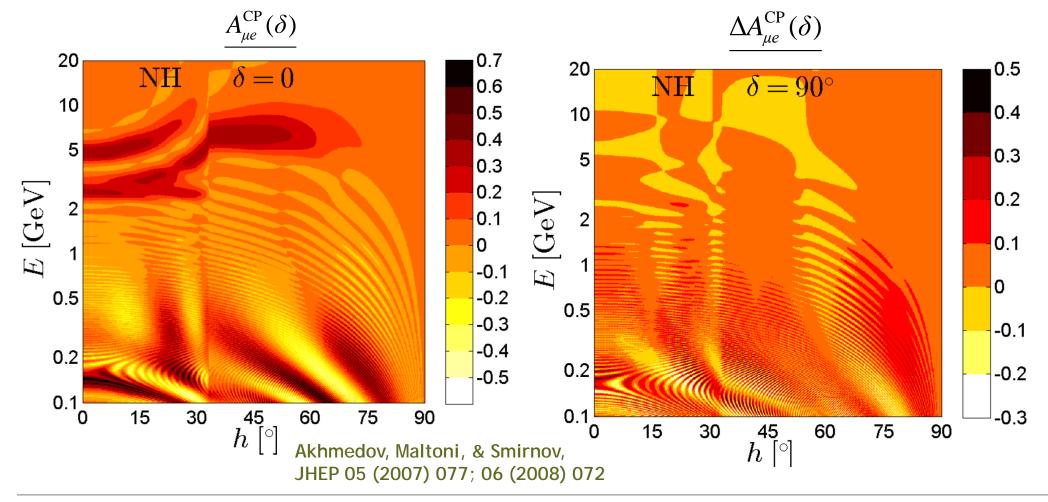


Oscillograms for CP Violation

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Conventional CP asymmetry

Working observable



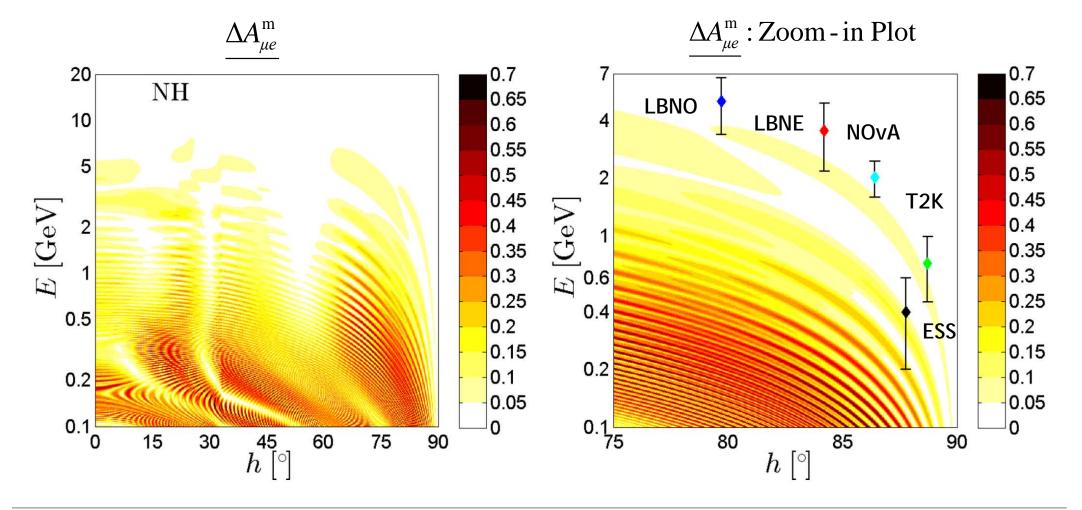


Oscillograms for CP Violation

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Working observable

Locating future experiments





Non-Standard Neutrino Interactions

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Neutrino oscillations in matter with NSIs:

For a review, Ohlsson, Rept. Prog. Phys. 76 (2013) 044201

$$H_{\rm eff} = \frac{1}{2E} \left\{ U \begin{pmatrix} 0 & & \\ & \Delta m_{21}^2 & \\ & & \Delta m_{31}^2 \end{pmatrix} U^+ + 2\sqrt{2}G_{\rm F}N_eE \begin{bmatrix} \begin{pmatrix} 1 & & \\ & 0 & \\ & & 0 \end{pmatrix} + \begin{pmatrix} \mathcal{E}_{ee} & \mathcal{E}_{e\mu} & \mathcal{E}_{e\tau} \\ \mathcal{E}_{e\mu}^* & \mathcal{E}_{\mu\tau} & \mathcal{E}_{\mu\tau} \\ \mathcal{E}_{e\tau}^* & \mathcal{E}_{\mu\tau}^* & \mathcal{E}_{\tau\tau} \end{pmatrix} \right\}$$

Oscillation probability with NSIs:

$$P_{\mu\mu}^{\text{NSI}} \simeq P_{\mu\mu}^{\text{SD}} - |\varepsilon_{\mu\tau}| c_{\phi_{\mu\tau}} \left(s_{2\times23}^3 A \Delta \sin \Delta + 4s_{2\times23} c_{2\times23}^2 A \sin^2 \frac{\Delta}{2} \right)$$
$$+ \left(|\varepsilon_{\mu\mu}| - |\varepsilon_{\tau\tau}| \right) s_{2\times23}^2 c_{2\times23} \left(\frac{A\Delta}{2} \sin \Delta - 2A \sin^2 \frac{\Delta}{2} \right)$$

- Deviation from the std. probability; only mu-tau flavors are involved
- Not suppressed by mixing angle θ_{13} or the mass ratio $\Delta m_{21}^2 / \Delta m_{31}^2$



Neutrino Parameter Mappings

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Leptonic mixing matrix elements in matter:

$$\begin{split} U_{e3}^{\rm m} &= \sin\hat{\theta}_{13} + \frac{\cos\hat{\theta}_{13}}{2\hat{C}} \left\{ \sin 2\hat{\theta}_{13} \left[\alpha s_{12}^2 + A(\varepsilon_{ee} - \tilde{\varepsilon}_{\tau\tau}) \right] + 2A \left[\cos 2\hat{\theta}_{13} \operatorname{Re}(\tilde{\varepsilon}_{e\tau}) + i\operatorname{Im}(\tilde{\varepsilon}_{e\tau}) \right] \right\} \\ U_{e2}^{\rm m} &= -\frac{c_{13}}{2A} \alpha \sin 2\theta_{12} - \tilde{\varepsilon}_{e\mu} + \frac{\tan\theta_{13}}{A} \tilde{\varepsilon}_{\mu\tau}^* , \\ U_{\mu3}^{\rm m} &= s_{23} \cos\hat{\theta}_{13} e^{\mathrm{i}\delta} \left\{ 1 - \frac{A \tan\hat{\theta}_{13}}{\hat{C}} \left[\cos 2\hat{\theta}_{13} \operatorname{Re}(\tilde{\varepsilon}_{e\tau}) + i\operatorname{Im}(\tilde{\varepsilon}_{e\tau}) \right] \right\} + \frac{\alpha c_{23} \sin 2\theta_{12} \sin\hat{\theta}_{13}}{1 + A + \hat{C}} , \end{split}$$

Modified NSI parameters:

 $\tilde{\varepsilon} = \varepsilon c_{\alpha\alpha} - \varepsilon s_{\alpha\alpha}$

Modified mixing angle:

$$\begin{split} \tilde{\varepsilon}_{e\mu} &= \varepsilon_{e\mu} \varepsilon_{23} - \varepsilon_{e\tau} \varepsilon_{23} ,\\ \tilde{\varepsilon}_{e\tau} &= (\varepsilon_{e\mu} s_{23} + \varepsilon_{e\tau} c_{23}) e^{i\delta} ,\\ \tilde{\varepsilon}_{\mu\mu} &= (\varepsilon_{\mu\mu} c_{23}^2 + \varepsilon_{\tau\tau} s_{23}^2) - 2s_{23} c_{23} \text{Re}[\varepsilon_{\mu\tau}] ,\\ \tilde{\varepsilon}_{\tau\tau} &= (\varepsilon_{\mu\mu} s_{23}^2 + \varepsilon_{\tau\tau} c_{23}^2) + 2s_{23} c_{23} \text{Re}[\varepsilon_{\mu\tau}] ,\\ \tilde{\varepsilon}_{\mu\tau} &= \left[(\varepsilon_{\mu\tau} c_{23}^2 - \varepsilon_{\mu\tau}^* s_{23}^2) + (\varepsilon_{\mu\mu} - \varepsilon_{\tau\tau}) s_{23} c_{23} \right] e^{i\delta} \end{split}$$

$$\sin^2 \hat{\theta}_{13} = \frac{\hat{C} - \cos 2\theta_{13} + A}{2\hat{C}}$$

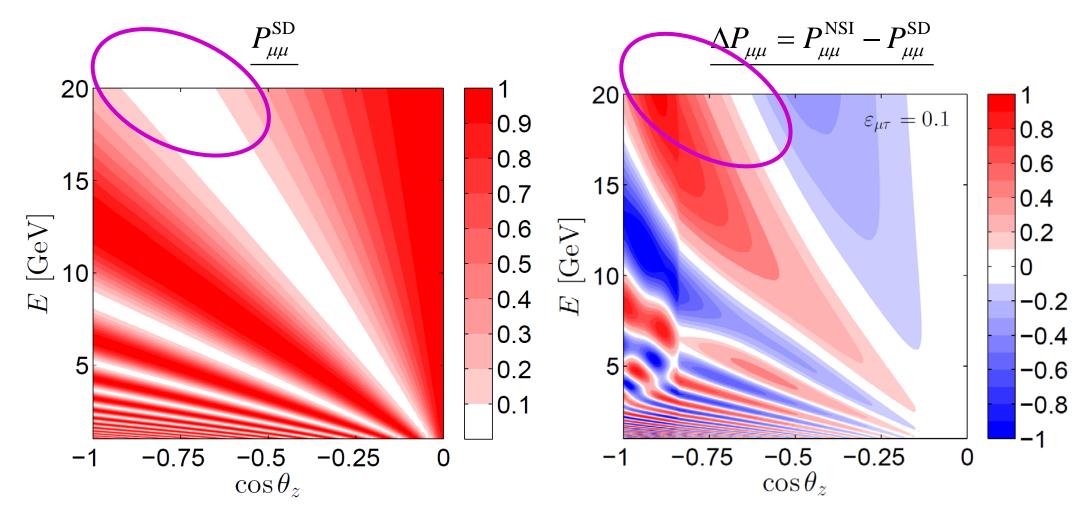
$$\hat{C} = \sqrt{(\cos 2\theta_{13} - A)^2 + \sin^2 2\theta_{13}}$$



Oscillation Probabilities with NSIs

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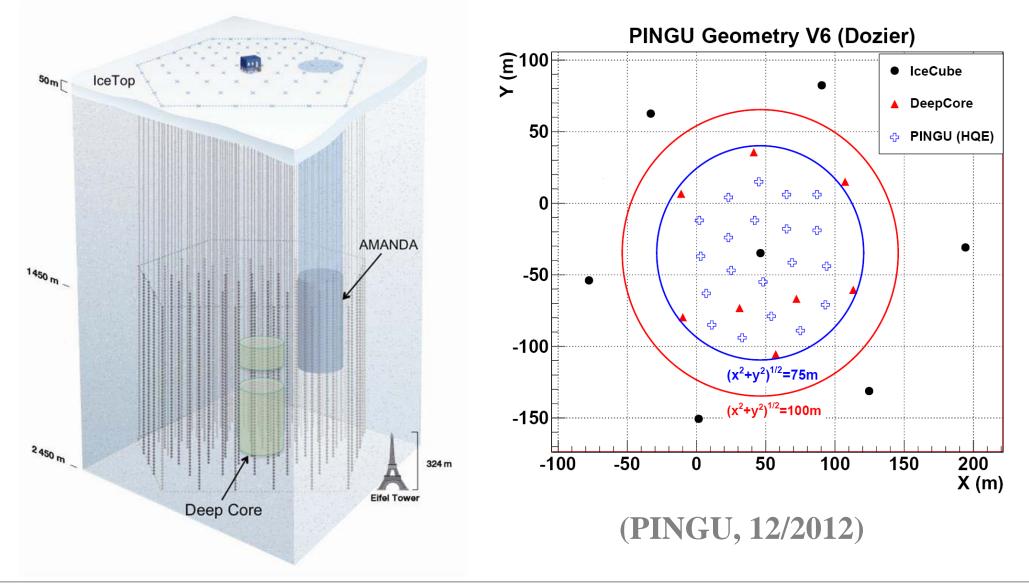
Significant deviation in the high-energy range:





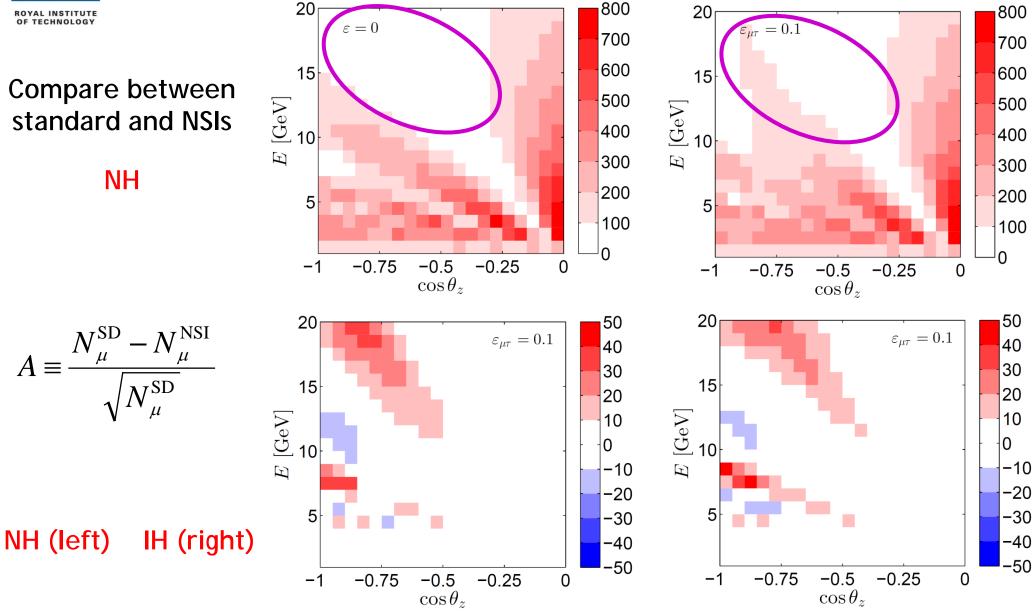
IceCube (DeepCore and PINGU)

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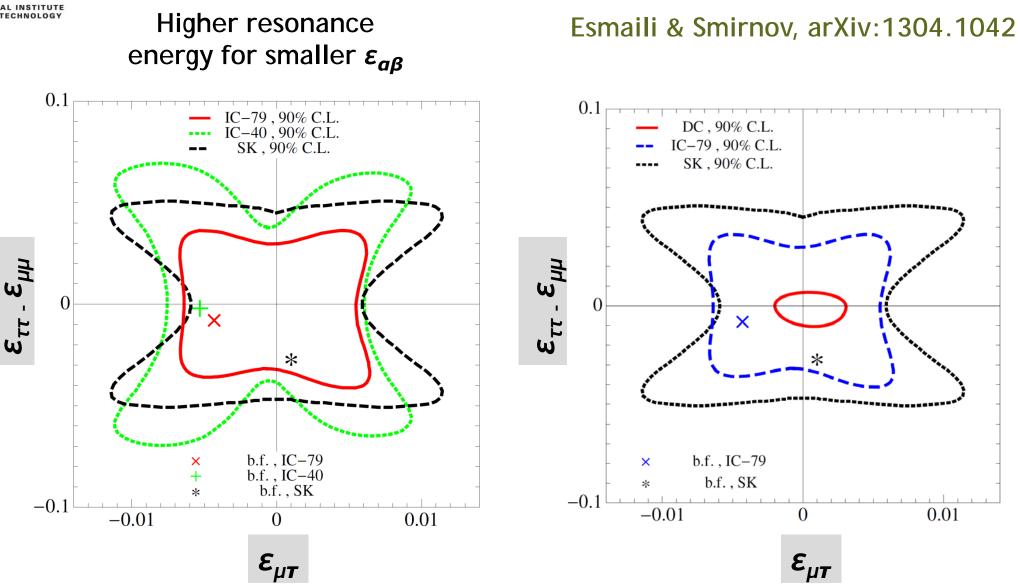
Number of Events @ PINGU

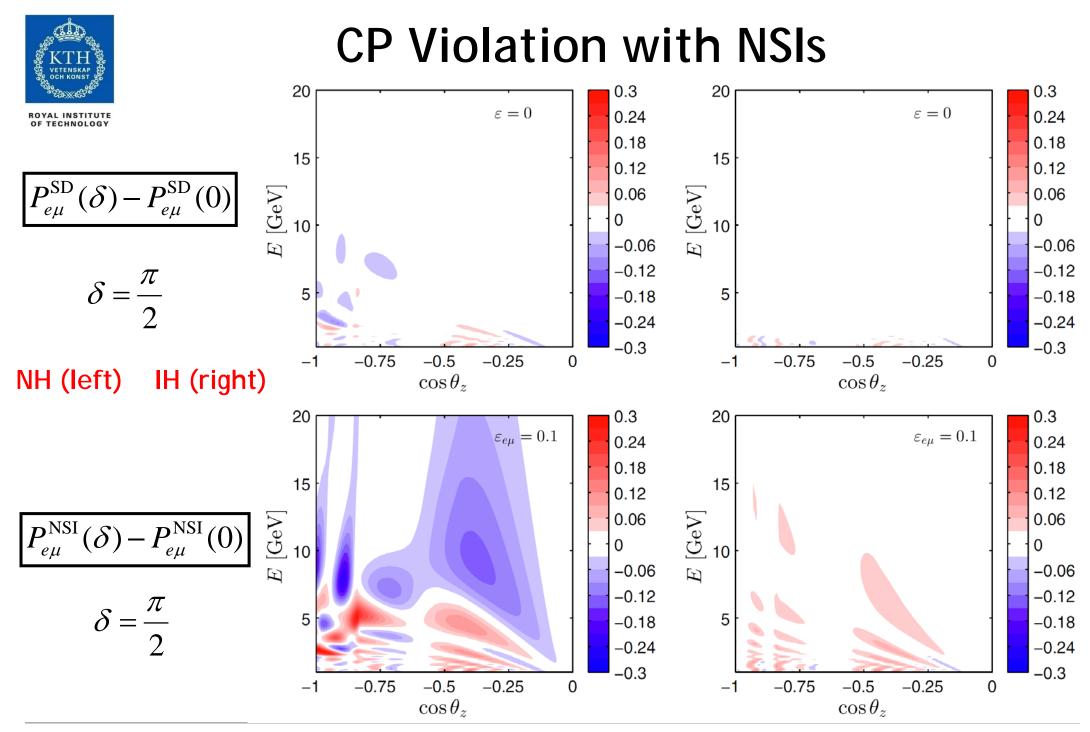




Sensitivity @ IceCube (DeepCore)

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Summary

> Future neutrino oscillation experiments aim to determine the neutrino mass ordering and measure the Dirac CP-violating phase.

➢ Usually theoretical prediction for the CP phase is given at a super-high energy scale, which should be compared with low-energy measurements. So, RG running has to be taken into account.

➤ Two new working observables have been introduced to describe intrinsic leptonic CP violation, disentangle the fake CP-violating effects, and optimize the experimental setup.

The possibility to constrain NSIs has been investigated at IceCube (DeepCore and PINGU).



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Thanks a lot for your attention!