

2583-10

Workshop on Coherent Phenomena in Disordered Optical Systems

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**Superfluid Stiffness of a Driven Dissipative Condensate with
Disorder**

Alexander JANOT
*Institute for Theoretical Physics
University Leipzig
Germany*

Superfluid Stiffness of a Driven Dissipative Condensate with Disorder

Alexander Janot

Timo Hyart, Paul Eastham and Bernd Rosenow

A. Janot, T. Hyart, P. R. Eastham, B. Rosenow, PRL **111**, 230403 (2013)

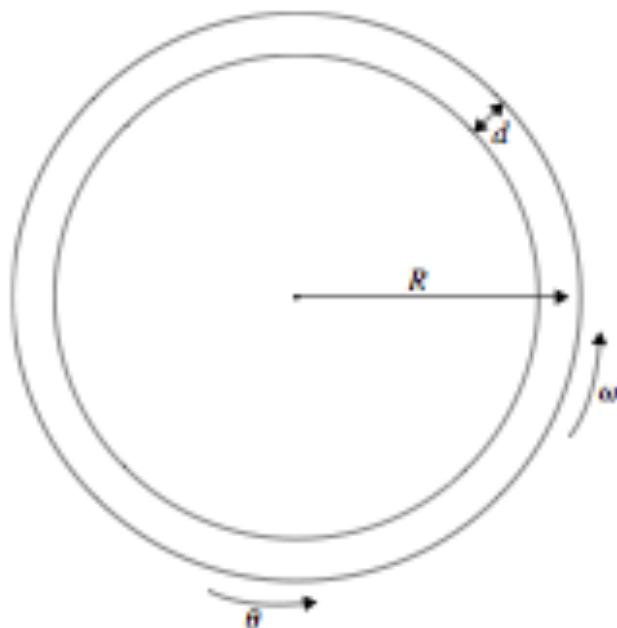
Institute for Theoretical Physics, University Leipzig, Germany

Superfluidity

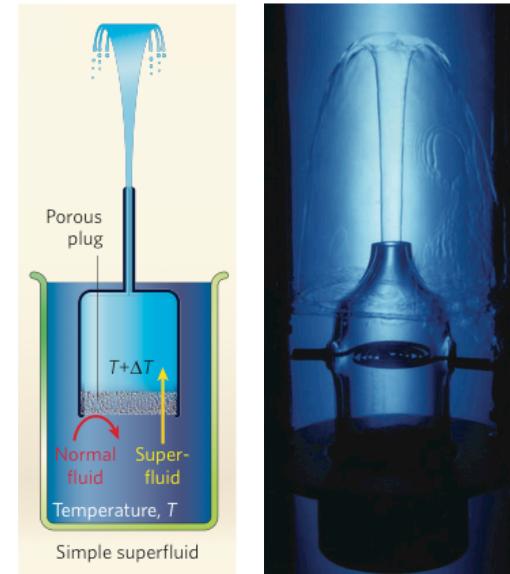
Rotating bucket experiment

- metastable superflow
- Hess-Fairbank effect

→ “frictionless flow”



taken from [A. J. Leggett, Quantum Liquids, Oxford University Press, 2006]



fountain effect

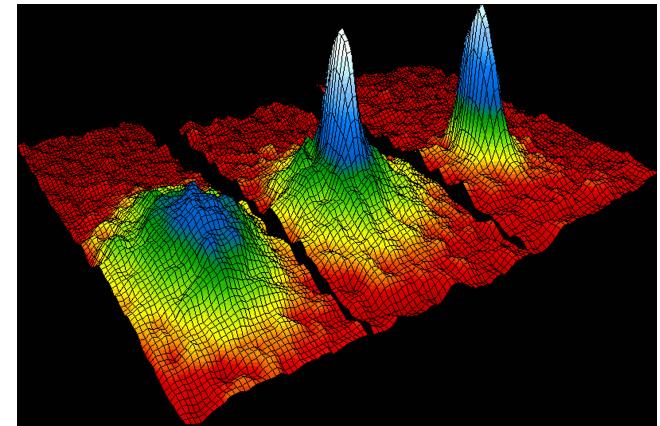
taken from [S. Fisher and G. Pickett, Nature 444, 832 (2006), S. Balibar, Nature 464, 176 (2010)]

Bose Einstein Condensate

- macroscopic occupation of the ground state
- off-diagonal long-range order

Condensate wavefunction

$$\Psi = \sqrt{n} e^{i\phi - i\epsilon t}$$



BEC of ultra-cold Rb-atoms

taken from [E. A. Cornell, C. E. Wieman and W. Ketterle, "Bose-Einstein Condensates in a Dilute Gas", Nobel lecture (2001)]

Superfluid velocity: $v_s \equiv \frac{\hbar}{m} \nabla \phi$

Criteria for Superfluidity

Landau criterion

Creation of elementary excitations?
critical velocity:

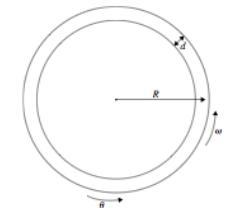
$$v_c = \min \left(\frac{\epsilon(p)}{p} \right)$$



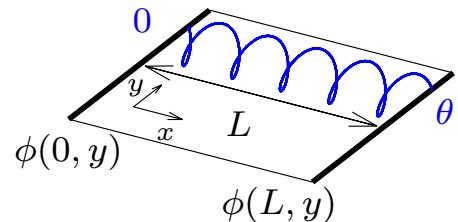
Transverse Current response

$$\vec{J} = \chi \vec{A} \sim ((n_s + n_n) P_{\parallel} + n_n P_{\perp}) \vec{A}$$

\vec{A} – external vector field
 χ – response function



Superfluid Stiffness – Phase twist



$$\text{tbc} \Leftrightarrow \text{pbc} \text{ “+“ } \vec{A}_\theta = \frac{\theta}{L} \vec{e}_\theta \text{ – twist current}$$

$$\text{tbc: } \phi_\theta(\vec{x} + L\vec{e}_\theta) - \phi_\theta(\vec{x}) = \theta$$

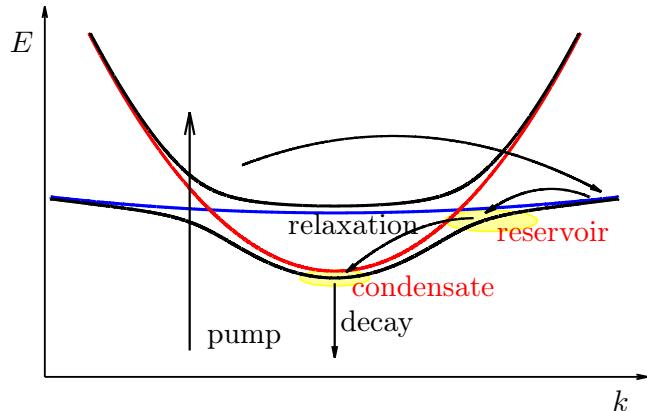
$$f_s = \frac{2mL^2}{\hbar^2} \lim_{\theta \rightarrow 0} \frac{\epsilon(\theta) - \epsilon(0)}{\theta^2}$$

[Leggett, RRL **25**, 1543 (1970); Fisher et al., RRA **8**, 1111 (1973); Roth and Burnett, PRA **67**, 031602(R) (2003)]

Polaritons – Driven Dissipative Condensate

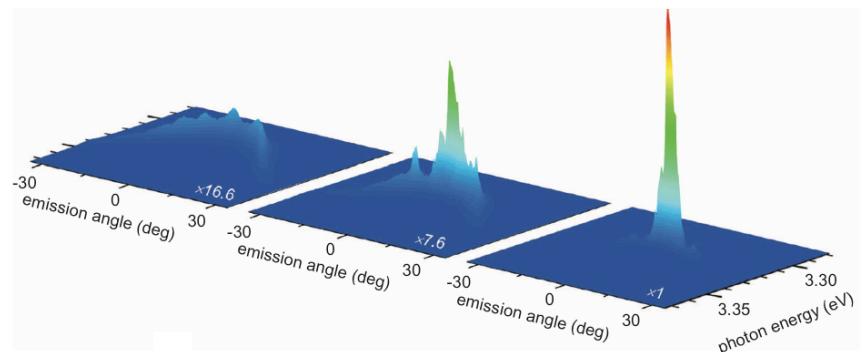
Polariton:

- mixed matter-light particle
- finite life-time → gain & loss physics “=” driven dissipative
- boson



Condensate:

- macroscopic occupation of the ground state



[Franke et al., NJP **14**, 013037 (2012)]

- build up of long-range spatial correlations

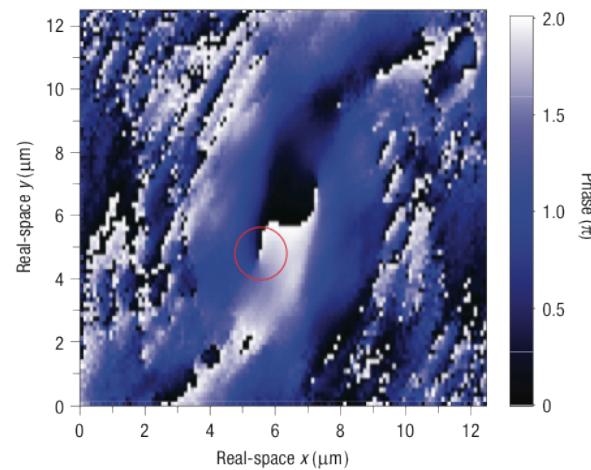
[Kasprzak et al., Nature **443**, 409 (2006), ...]

See also: [Deng et al., Rev. Mod. Phys. **82**, 1489 (2010); Carusotto and Ciuti, Rev. Mod. Phys. **85**, 299 (2013)]

Polariton Condensate – Superfluidity?

Experimental Signatures

- suppression of scattering from defects
[Amo et al., Nat. Phys. **5**, 805 (2009)]
- quantized vortices
[Lagoudakis et al., Nat. Phys. **4**, 706 (2008); Sanvitto et al., Nat. Phys. **6**, 527 (2010)]



Vortex visible in a real space
condensate phase profile
[Lagoudakis et al., Nat. Phys. **4**, 706 (2008)]

Driven vs. Equilibrium condensate

Effects of quantum and classical fluctuations

Similarities

- perturbatively in condensate fluctuations: correlation functions agree with equilibrium results ($d=2,3$)
→ long-range order (?)
[Chiocchetta and Carusotto, *Europhys. Lett.* **102**, 67007 (2013); Roumpos et al., *Proc. Natl. Acad. Sci.* **109**, 6467 (2012)]
- low-frequency and long-wavelength effective **thermalization**, ($d=3$)
[Sieberer et al., *PRL* **110**, 195301 (2013)]

Differences

- **universal low-frequency decoherence, new dynamical critical exponent,** ($d=3$)
[Sieberer et al., *PRL* **110**, 195301 (2013); Sieberer et al., *PRB* **89**, 134310 (2014); Täuber and Diehl, *PRX* **4**, 021010 (2014)]
- **no algebraic long-range order, exponential long-distance correlation function ($d=2$)**
[Altman et al., arXiv 1311.0876 (2013)]

- Experiments suggest **superfluid behavior** of polariton condensate
[Lagoudakis et al., Nat. Phys. 4, 706 (2008); Amo et al., Nat. Phys. 5, 805 (2009); Sanvitto et al., Nat. Phys. 6, 527 (2010)]
- Experiments involve a significant amount of **disorder**
[Franke et al., NJP 14, 013037 (2012)]
- Equilibrium: **Superfluid—Bose-glass** phase transition
[Fisher et al., PRB 40, 546 (1989)]
- A **driven dissipative condensate** is **different**
[Sieberer et al., PRL 110, 195301 (2013), Altman et al., arXiv 1311.0876 (2013)]

Is a Driven Quantum Condensate in a Disordered Environment a Superfluid?

Superfluidity in a Driven Condensate

Steady state condensate wavefunction: $\Psi = \sqrt{n} e^{i\phi - i\omega t}$, $\vec{v}_s = \frac{\hbar}{m} \nabla \phi$

Landau criteria

excitations: $\delta\Psi_k \sim e^{\pm i\omega_k t/\hbar - \gamma t}$

$$\omega_k = \sqrt{\frac{\hbar^2}{2m} k^2 \left(\frac{\hbar^2}{2m} k^2 + 2n_0 U \right) - (\hbar\gamma)^2}$$

→ Diffusive mode: $v_c = 0$???

- generalized Landau criteria:
→ Superfluid behavior
- onset of drag force at

$$v_c \approx \sqrt{\frac{n_0 U}{m}}$$

[Wouters and Carusotto, PRL 105, 020602 (2010); Amo et al., Nat. Phys. 5, 805 (2009)]

Transverse Current response

Greens function:

$$G_k \sim \frac{1}{k^2} \quad \text{for } k \ll 1$$

→ expect superfluidity in principle

[Keeling, PRL 107, 080402 (2011)]

Superfluid Stiffness

$$f_s = \frac{2mL^2}{\hbar^2} \lim_{\theta \rightarrow 0} \frac{\omega(\theta) - \omega(0)}{\theta^2}$$

“our aim”

Synchronization & Desynchronization

Driving & Disorder

Multi mode condensate

- Desynchronization
[Wouters, PRB **77**, 121302 (2008);
Eastham, PRB **78**, 035319 (2008)]

Single mode condensate

- Steady state (synchronized)
- Neglecting gain & loss:
Superfluid—Bose-glass phase
transition predicted for polaritons
[Malpuech et al., PRL **98**, 206402 (2007)]

Phenomenological Description of a Driven Dissipative Condensate

Extended Gross Pitaevskii Equation (eGPE):

$$i\hbar\partial_t\Psi = \left(-\frac{\hbar^2}{2m}\nabla^2 + V(\vec{x}) + U|\Psi|^2 \right)\Psi + i(\hbar\gamma - \Gamma|\Psi|^2)\Psi$$

GPE – equilibrium condensate Driving mechanism

The diagram illustrates the components of the Extended Gross-Pitaevskii Equation. The equation itself is centered, with curly braces above it grouping the 'GPE – equilibrium condensate' terms and the 'Driving mechanism' terms. Below the equation, arrows point from labels to specific terms:

- condensate wavefunction** points to $\Psi(\vec{x}, t)$.
- kinetic mass** points to $-\frac{\hbar^2}{2m}\nabla^2$.
- external potential** points to $V(\vec{x})$.
- interaction outside potential constant** points to $U|\Psi|^2$.
- gain inscattering rate** points to $i(\hbar\gamma - \Gamma|\Psi|^2)$.
- loss gain depletion** points to $i\hbar\gamma$.

Disorder Potential

Disorder potential:

- δ -correlated

$$\langle\langle V(\vec{x}) \rangle\rangle = 0$$

- Gaussian distributed

$$\langle\langle V(\vec{x})V(\vec{y}) \rangle\rangle = V_0^2 \delta\left(\frac{\vec{x} - \vec{y}}{\xi_V}\right)$$

$$i\hbar\partial_t\Psi = \left(-\frac{\hbar^2}{2m}\nabla^2 + V(\vec{x}) + U|\Psi|^2\right)\Psi + i(\hbar\gamma - \Gamma|\Psi|^2)\Psi$$

density

$$n_0 \equiv \frac{\hbar\gamma}{\Gamma}$$

blue shift

$$n_0U$$

healing length

$$\xi \equiv \sqrt{\frac{\hbar^2}{2m n_0 U}}$$

Dimensionless eGPE

$$\mathrm{i}\partial_t\psi = (-\nabla^2 + \vartheta(\vec{x}) + |\psi|^2)\psi + \mathrm{i}\alpha(1 - |\psi|^2)\psi$$

- energy in units of blue shift
- length in units of healing length
- density in units of n_0

$$\psi(\vec{x}, t) \equiv \frac{\Psi(\vec{x}, t)}{\sqrt{n_0}}, \quad \vartheta(\vec{x}) \equiv \frac{V(\vec{x})}{n_0 U}$$

$$\langle\langle \vartheta(\vec{x})\vartheta(\vec{y}) \rangle\rangle = \kappa^2 \delta(\vec{x} - \vec{y})$$

Effectively two parameters

- disorder strength: $\kappa \equiv \frac{\xi_V V_0}{\xi n_0 U}$
- non-equilibrium: $\alpha \equiv \frac{\Gamma}{U}$

Steady state ansatz: $\psi(\vec{x}, t) = \sqrt{n(\vec{x})} e^{\mathrm{i}\phi(\vec{x}) - \mathrm{i}\omega t}$

Random sources and sinks

Hydrodynamic Equations

$$\omega = (\nabla \phi)^2 + \frac{1}{4} \frac{(\nabla n)^2}{n^2} - \frac{1}{2} \frac{\nabla^2 n}{n} + n + \vartheta$$

$$0 = \nabla(n\nabla\phi) + \alpha n(n-1)$$

$n(\vec{x}) < 1$: local source

$n(\vec{x}) > 1$: local sink

condensate emission
energy

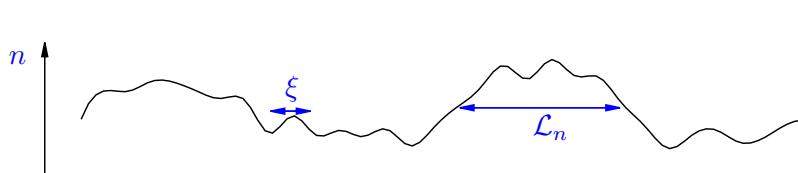
continuity equation with
sources and sinks

*density
fluctuations
induce*

random currents

Density and Phase Fluctuations

- Density fluctuations

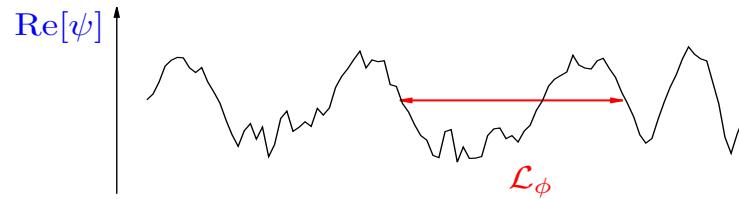


$$\alpha = 0$$

Larkin length: $\mathcal{L}_n \sim \frac{1}{\kappa}$

[Imry and Ma, PRL 35, 1399 (1975); Nattermann and Pokrovsky, PRL 100, 060402 (2008)]

- Phase fluctuations



$$\alpha > 0$$

Phase correlation length: $\mathcal{L}_\phi \sim \frac{1}{\alpha \kappa}$

Random currents: $\nabla \phi \sim \frac{1}{\mathcal{L}_\phi}$

Long-Range Order Destroyed by Phase Fluctuations

Condensate correlation function

$$\langle\langle \psi^*(\vec{x})\psi(\vec{0}) \rangle\rangle \approx e^{-\frac{1}{2}\langle\langle (\phi(\vec{x}) - \phi(\vec{0}))^2 \rangle\rangle}$$

$$\sim e^{-\vec{x}^2/\mathcal{L}_\phi^2}$$

- + random sources and sinks
- + random currents
- + phase fluctuations

→ NO Long-Range Order

Stiffness – Perturbative results

Superfluid Stiffness

- probed by phase twist θ

- solve

$$\omega = (\nabla \phi)^2 + \frac{1}{4} \frac{(\nabla n)^2}{n^2} - \frac{1}{2} \frac{\nabla^2 n}{n} + n + \vartheta$$

$$0 = \nabla(n\nabla\phi) + \alpha n(n-1)$$

perturbatively

- evaluate stiffness

$$f_s = L^2 \lim_{\theta \rightarrow 0} \frac{\omega(\theta) - \omega(0)}{\theta^2}$$

Consider weak disorder limit

- $\kappa \ll 1$
- expansion of density, current and frequency in powers of κ

Stiffness – Perturbative results

Result:

$$f_s \approx 1 - \{c_1 + g_1(L) \alpha^2 + (g_2(L) + c_2 L^2) \alpha^4\} \kappa^2 + \mathcal{O}(\kappa^4)$$

$$c_1, c_2 = \text{const}$$

$$g_1(L), g_2(L) \sim -\log L$$

$$f_s \approx 0 \quad \text{at} \quad \mathcal{L}_s \sim \frac{1}{\alpha^2 \kappa}$$

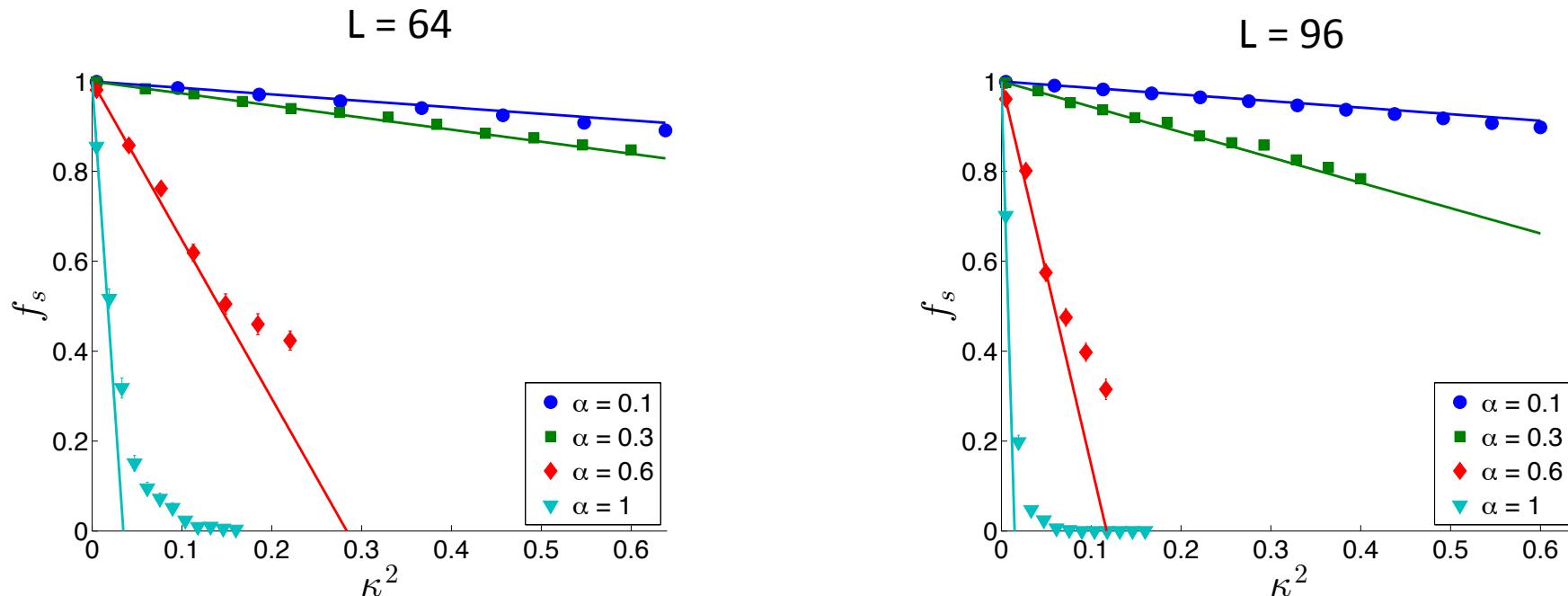
- reproduces equilibrium results for $\alpha = 0$ [Meng, PRB 49, 1205 (1994)]
- for finite $\alpha > 0$, divergence in thermodynamic limit for any $\kappa > 0$

- fastest divergence sets
superfluid depletion length

Expect:

- **superfluid behavior** for $L < \mathcal{L}_s$
- **no superfluid** for $L \rightarrow \infty$

Stiffness – Perturbative vs Numeric results



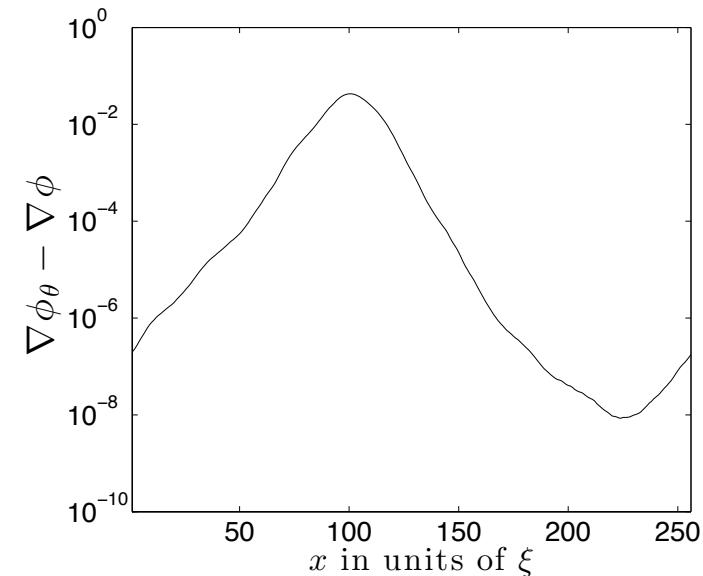
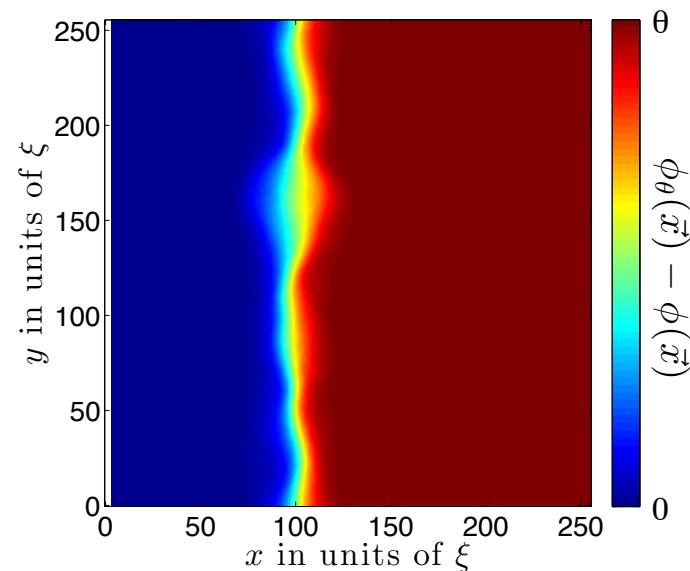
- for $f_s \lesssim 1$: perturbative result agrees with numerics for all α
- for $f_s \approx 0$: discrepancy for $\alpha \approx O(1)$ even if $\kappa \ll 1$

→ Condensate response for $\mathcal{L}_s \gg L$?

→ Mechanism of superfluid depletion ?

Domain Wall Formation

- twist θ along x-direction
- phase response: $\phi \xrightarrow{\vec{A}_\theta} \phi_\theta$, $\nabla \phi \xrightarrow{\vec{A}_\theta} \nabla \phi_\theta$
- one disorder realization with $\mathcal{L}_s \ll L$



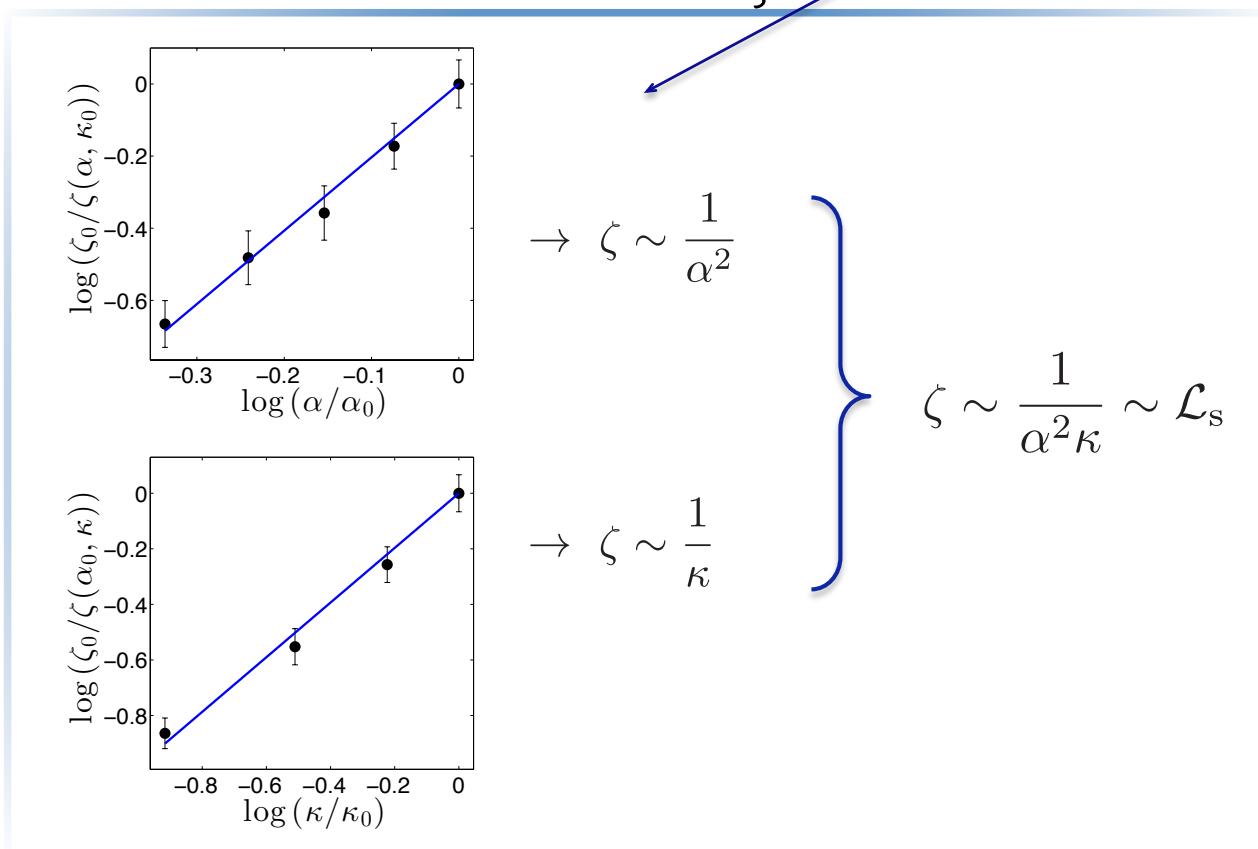
→ response exponentially localized

→ no superfluidity

Domain Wall Thickness

Ansatz: $\nabla \phi_\theta - \nabla \phi = \frac{\theta}{2\zeta(1 - e^{-\frac{L}{2\zeta}})} e^{-\frac{|x-x_0|}{\zeta}} \vec{e}_\theta = \begin{cases} \frac{\theta}{L} \vec{e}_\theta & \text{for } \zeta \rightarrow \infty \\ \theta \delta(x - x_0) \vec{e}_\theta & \text{for } \zeta \rightarrow 0 \end{cases}$

↑
decay length \sim thickness; x_0 – random position



Generic properties of the Domain wall:

- randomly pinned
- thickness is set by superfluid depletion length

Stiffness – Scaling

- vanishing stiffness explained by domain wall formation
 - \mathcal{L}_s is only relevant length scale
- $f_s \sim f_s(L/\mathcal{L}_s)$ for $L \gg \mathcal{L}_s$?

Scaling ansatz:

$$f_s = e^{-c_2 \alpha^4 \kappa^2 L^2} (1 - g) \sim e^{-L^2 / \mathcal{L}_s^2}$$

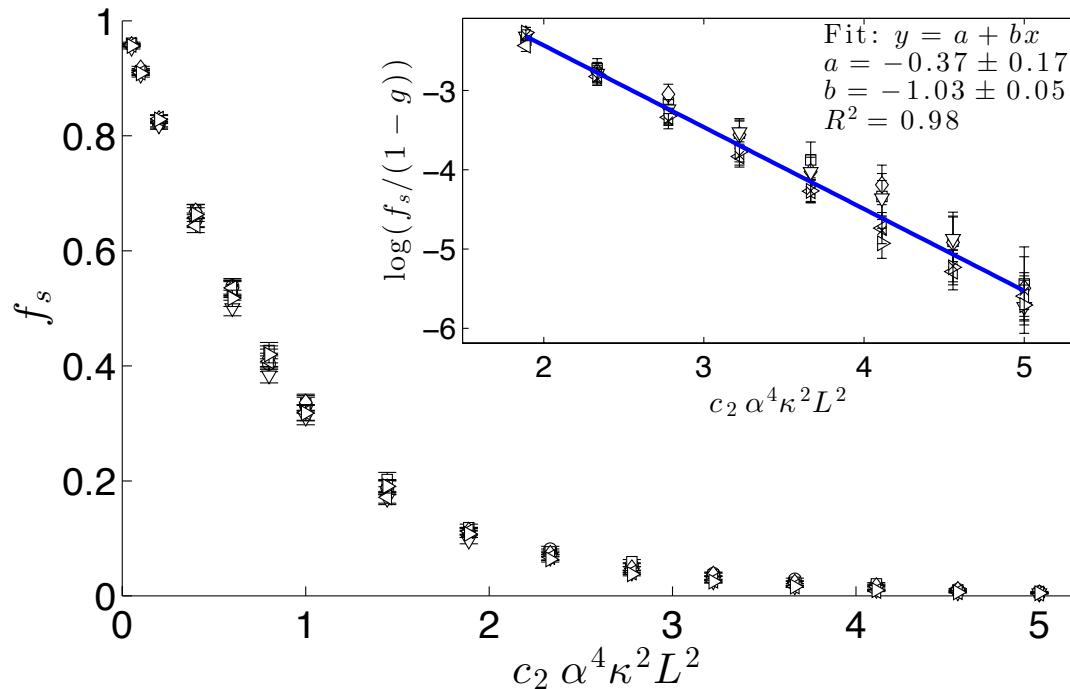
c_2 from perturbation
theory

logarithmic correction
 $g(\alpha, \kappa, \log L)$

- reproduces perturbative results if expanded in κ

Stiffness – Scaling

Check ansatz by numerical simulations:



$$f_s = e^{-c_2 \alpha^4 \kappa^2 L^2} (1 - g) \sim e^{-L^2 / \mathcal{L}_s^2}$$



Generalization to arbitrary dimensions

Perturbatively calculated stiffness:

$$f_s \approx 1 - \dots - c \alpha^4 \kappa^2 L^{4-d}$$

$c = \text{const.}$

leading divergence

Superfluid depletion length:

$$\mathcal{L}_s \sim \left(\frac{1}{\alpha^2 \kappa} \right)^{\frac{2}{4-d}}$$

- upper critical dimension $d_c = 4$
- expect depletion of superfluidity for all $d < d_c$

Driven Disordered Condensate

- no quasi long-range order

$$\langle\langle \psi^*(\vec{x})\psi(\vec{0}) \rangle\rangle \sim e^{-\vec{x}^2/\mathcal{L}_\phi^2}$$

$$\mathcal{L}_\phi \sim \frac{1}{\alpha\kappa}$$

α – driving

κ – disorder

ψ – condensate

- no superfluid in the thermodynamic limit

- superfluid behavior for length scales

$$L \ll \mathcal{L}_s$$

$$\mathcal{L}_s \sim \frac{1}{\alpha^2\kappa}$$

- exponential depletion of stiffness

$$f_s \sim e^{-L^2/\mathcal{L}_s^2}$$

- domain wall formation as response to phase twist

Acknowledgement

I like to thank ...

... my co-workers

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Institut für Theoretische Physik



... my graduate school



... and you!

Phase Fluctuations – Phase Correlation Length

Generalized Imry-Ma argument

- integration of non-equilibrium continuity equation over volume $\Omega_\phi \sim \mathcal{L}_\phi^d$

$$\int_{\Omega_\phi} \nabla(n \nabla \phi) = \alpha \int_{\Omega_\phi} n(1 - n)$$

$$\int_{\partial \Omega_\phi} \cdot (n \nabla \phi) =$$

$$\int_{\partial \Omega_\phi} \cdot \nabla \phi \approx$$

$$\mathcal{L}_\phi^{d-1} \frac{1}{\mathcal{L}_\phi} \sim$$

fluctuation at scale ξ

$n \approx 1 + \mathcal{O}(\kappa)$

$\xi \ll \mathcal{L}_\phi$

$\sim \alpha \kappa \sqrt{\left(\frac{\mathcal{L}_\phi}{\xi}\right)^d}$

boundary

current

$$\mathcal{L}_\phi \sim \left(\frac{1}{\alpha \kappa}\right)^{\frac{2}{4-d}}$$

- gives phase fluctuation length
- result agrees with perturbation theory for $\kappa \ll 1$

Supplemental Material:

Simulation parameters

Domain wall, single snap shot, page 16: $L = 256 a_L$; $\alpha = 0.5$; $\kappa = 0.5$; $\Theta = 1$; $a_L = \xi$

Domain wall scaling, page 17: $L = 256 a_L$; $\alpha_0 = 0.7$; $\kappa_0 = 0.25$; $\Theta = 1$; $a_L = \xi$; 72 disorder realization

Stiffness scaling, page 19: $L = 64,96 a_L$; $\alpha = 0.9, 1, 1.2$; $a_L = \xi$; up to 1320 disorder realization