

**2585–16**

**Joint ICTP–TWAS School on Coherent State Transforms, Time–  
Frequency and Time–Scale Analysis, Applications**

***2 – 20 June 2014***

**Coherent States and Wavelets: A Unified Approach – III**

S. T. Ali  
*Concordia Univ., Montreal  
Canada*

# Coherent States and Wavelets: A Unified Approach – III

Department of Mathematics and Statistics  
Concordia University  
Montréal, Québec, CANADA H3G 1M8

[twareque.ali@concordia.ca](mailto:twareque.ali@concordia.ca)

Joint ICTP - TWAS School  
on Coherent State Transforms, Time-Frequency  
and Time-Scale Analysis, Applications

The Abdus Salam International Centre for Theoretical Physics,  
Miramare, Trieste, ITALY

June 2 – 21, 2014

# Abstract

*By analogy with the one- and two-dimensional wavelet groups, we introduce the quaternionic affine group, look at some of its properties, its representations on complex and quaternionic Hilbert spaces, the associated wavelet transforms and coherent states.*

# Contents

- 1 Preliminaries
  - Contents
- 2 Some quaternionic facts
  - Action of  $\mathbb{H}^*$  on  $\mathbb{H}$
- 3 The quaternionic affine group
- 4 UIR of  $G_{\text{aff}}^{\mathbb{H}}$  in a quaternionic Hilbert space
- 5 Wavelets and reproducing kernels

## Useful facts

We list some useful facts about quaternions and the matrix representation that we shall use.

Let  $\mathbb{H}$  denote the field of all quaternions and  $\mathbb{H}^*$  the group (under quaternionic multiplication) of all invertible quaternions. A general quaternion can be written as

$$\mathbf{q} = q_0 + q_1\mathbf{i} + q_2\mathbf{j} + q_3\mathbf{k}, \quad q_0, q_1, q_2, q_3 \in \mathbb{R},$$

where  $\mathbf{i}, \mathbf{j}, \mathbf{k}$  are the three quaternionic imaginary units, satisfying  $\mathbf{i}^2 = \mathbf{j}^2 = \mathbf{k}^2 = -1$  and  $\mathbf{ij} = \mathbf{k} = -\mathbf{ji}$ ,  $\mathbf{jk} = \mathbf{i} = -\mathbf{kj}$ ,  $\mathbf{ki} = \mathbf{j} = -\mathbf{ik}$ . The quaternionic conjugate of  $\mathbf{q}$  is

$$\bar{\mathbf{q}} = q_0 - \mathbf{i}q_1 - \mathbf{j}q_2 - \mathbf{k}q_3.$$

We shall use the  $2 \times 2$  matrix representation of the quaternions, in which

$$\mathbf{i} = \sqrt{-1}\sigma_1, \quad \mathbf{j} = -\sqrt{-1}\sigma_2, \quad \mathbf{k} = \sqrt{-1}\sigma_3,$$

and the  $\sigma$ 's are the three Pauli matrices,

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix},$$

## Useful facts

to which we add

$$\sigma_0 = \mathbb{I}_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}.$$

We shall also use the matrix valued vector  $\boldsymbol{\sigma} = (\sigma_1, -\sigma_2, \sigma_3)$ . Thus, in this representation,

$$\mathbf{q} = q_0 \sigma_0 + i\mathbf{q} \cdot \boldsymbol{\sigma} = \begin{pmatrix} q_0 + iq_3 & -q_2 + iq_1 \\ q_2 + iq_1 & q_0 - iq_3 \end{pmatrix}, \quad \mathbf{q} = (q_1, q_2, q_3).$$

In this representation, the quaternionic conjugate of  $\mathbf{q}$  is given by  $\mathbf{q}^\dagger$ .

Introducing two complex variables, which we write as

$$z_1 = q_0 + iq_3, \quad z_2 = q_2 + iq_1,$$

we may also write

$$\mathbf{q} = \begin{pmatrix} z_1 & -\bar{z}_2 \\ z_2 & \bar{z}_1 \end{pmatrix}. \tag{2.1}$$

## Useful facts

From this it is clear that the group  $\mathbb{H}^*$  is isomorphic to the *affine  $SU(2)$  group*, i.e.,  $\mathbb{R}^{>0} \times SU(2)$ , which is the group  $SU(2)$  together with all (non-zero) dilations.

As a set  $\mathbb{H}^* \simeq \mathbb{R}^{>0} \times S(4)$ , where  $S(4)$  is the surface of the sphere in  $\mathbb{R}^4$ , or more simply,  $\mathbb{H}^* \simeq \mathbb{R}^4 \setminus \{\mathbf{0}\}$ .

## Action of $\mathbb{H}^*$ on $\mathbb{H}$

Consider the action of  $\mathbb{H}^*$  on  $\mathbb{H}$  by right (or left) quaternionic (in our representation matrix) multiplication. It is clear that there are only two orbits under this action,  $\{\mathbf{o}\}$  (the zero quaternion) and  $\mathbb{H}^*$ . Furthermore, this latter orbit is *open and free*. Let

$$\mathbf{a} = \begin{pmatrix} w_1 & -\bar{w}_2 \\ w_2 & \bar{w}_1 \end{pmatrix} \in \mathbb{H}^* \quad \text{and} \quad \mathbf{r} = \begin{pmatrix} z_1 & -\bar{z}_2 \\ z_2 & \bar{z}_1 \end{pmatrix} \in \mathbb{H}.$$

Then under left action

$$\mathbf{r} \mapsto \mathbf{r}' = \mathbf{a}\mathbf{r} = \begin{pmatrix} w_1 z_1 - \bar{w}_2 z_2 & -\bar{w}_2 \bar{z}_1 - w_1 \bar{z}_2 \\ w_2 z_1 + \bar{w}_1 z_2 & \bar{w}_1 \bar{z}_1 - w_2 \bar{z}_2 \end{pmatrix}. \quad (2.2)$$

We take  $w_1 = a_0 + ia_3$ ,  $w_2 = a_2 + ia_1$  and  $z_1 = x_0 + ix_3$ ,  $z_2 = x_2 + ix_1$  and consider  $\mathbf{r}$  as the vector

$$\mathbf{x} = \begin{pmatrix} x_0 \\ x_3 \\ x_2 \\ x_1 \end{pmatrix} := \begin{pmatrix} \mathbf{x}_1 \\ \mathbf{x}_2 \end{pmatrix} \in \mathbb{R}^4. \quad (2.3)$$



## Action of $\mathbb{H}^*$ on $\mathbb{H}$

On this vector, the left action (2.2) is easily seen to lead to the matrix left action

$$\mathbf{x} \mapsto \mathbf{x}' = A\mathbf{x} = \begin{pmatrix} a_0 & -a_3 & -a_2 & -a_1 \\ a_3 & a_0 & a_1 & -a_2 \\ a_2 & -a_1 & a_0 & a_3 \\ a_1 & a_2 & -a_3 & a_0 \end{pmatrix} \begin{pmatrix} x_0 \\ x_3 \\ x_2 \\ x_1 \end{pmatrix} = \begin{pmatrix} A_1 & -A_2^T \\ A_2 & A_1^T \end{pmatrix} \begin{pmatrix} \mathbf{x}_1 \\ \mathbf{x}_2 \end{pmatrix}, \quad (2.4)$$

on  $\mathbb{R}^4$ .

The matrices  $A_1$  and  $A_2$  are rotation-dilation matrices, and may be written in the form

$$A_1 = \lambda_1 \begin{pmatrix} \cos \theta_1 & -\sin \theta_1 \\ \sin \theta_1 & \cos \theta_1 \end{pmatrix} = \lambda_1 R(\theta_1), \quad A_2 = \lambda_2 \begin{pmatrix} \cos \theta_2 & -\sin \theta_2 \\ \sin \theta_2 & \cos \theta_2 \end{pmatrix} = \lambda_2 R(\theta_2) \quad (2.5)$$

where

$$\theta_1 = \tan^{-1} \left( \frac{a_3}{a_0} \right), \quad \theta_2 = \tan^{-1} \left( \frac{a_1}{a_2} \right), \quad \lambda_1 = \sqrt{a_0^2 + a_3^2}, \quad \lambda_2 = \sqrt{a_1^2 + a_2^2} \quad \text{and} \quad \lambda_1^2 + \lambda_2^2 \neq 0 \quad (2.6)$$

and  $R(\theta)$  is the  $2 \times 2$  rotation matrix

$$R(\theta) = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}. \quad (2.7)$$

## Action of $\mathbb{H}^*$ on $\mathbb{H}$

From the above it is clear that when  $\mathbb{H}$  is identified with  $\mathbb{R}^4$ , the action of  $\mathbb{H}^*$  on  $\mathbb{H}$  is that of two two-dimensional rotation-dilation groups (rotations of the two-dimensional plane together with radial dilations, where at least one of the dilations is non-zero) acting on  $\mathbb{R}^4$ .

Consequently, we shall consider elements in  $\mathbb{H}^*$  interchangeably as  $2 \times 2$  complex matrices of the type

$$\mathbf{a} = \begin{pmatrix} w_1 & -\bar{w}_2 \\ w_2 & \bar{w}_1 \end{pmatrix}, \quad \det[\mathbf{a}] = |\mathbf{a}|^2 \neq 0$$

or  $4 \times 4$  real matrices of the type  $A$  in (2.4),

$$A = \begin{pmatrix} \lambda_1 R(\theta_1) & -\lambda_2 R(-\theta_2) \\ \lambda_2 R(\theta_2) & \lambda_1 R(-\theta_1) \end{pmatrix}, \quad \det[A] = |\mathbf{a}|^4 = [\lambda_1^2 + \lambda_2^2]^2 \neq 0. \quad (2.8)$$

## Quaternionic affine group

Let us look at the three affine groups,  $G_{\text{aff}}^{\mathbb{R}}$ ,  $G_{\text{aff}}^{\mathbb{C}}$  and  $G_{\text{aff}}^{\mathbb{H}}$ , of the real line, the complex plane and the quaternions, respectively. These groups are defined as the semi-direct products

$$G_{\text{aff}}^{\mathbb{R}} = \mathbb{R} \rtimes \mathbb{R}^*, \quad G_{\text{aff}}^{\mathbb{C}} = \mathbb{C} \rtimes \mathbb{C}^*, \quad G_{\text{aff}}^{\mathbb{H}} = \mathbb{H} \rtimes \mathbb{H}^*.$$

Let  $\mathbb{K}$  denote any one of the three fields  $\mathbb{R}$ ,  $\mathbb{C}$  or  $\mathbb{H}$  and write  $G_{\text{aff}}^{\mathbb{K}} = \mathbb{K} \rtimes \mathbb{K}^*$ . A generic element in  $G_{\text{aff}}^{\mathbb{K}}$  can be written as

$$g = (b, a) = \begin{pmatrix} a & b \\ 0 & 1 \end{pmatrix}, \quad a \in \mathbb{K}^*, \quad b \in \mathbb{K}.$$

Of these,  $G_{\text{aff}}^{\mathbb{R}}$  is the *one-dimensional wavelet group* and  $G_{\text{aff}}^{\mathbb{C}}$ , which is isomorphic to the similitude group of the plane (translations, rotations and dilations of the 2-dimensional plane), is the *two-dimensional wavelet group*.

## Quaternionic affine group

By analogy we shall call the quaternionic affine group  $G_{\text{aff}}^{\mathbb{H}}$  the *quaternionic wavelet group*, which we now analyse in some detail. In the  $2 \times 2$  matrix representation of the quaternions introduced earlier, we shall represent an element of  $G_{\text{aff}}^{\mathbb{H}}$  as the  $3 \times 3$  complex matrix

$$g := (\mathbf{b}, \mathbf{a}) = \begin{pmatrix} \mathbf{a} & \mathbf{b} \\ \mathbf{0}^T & 1 \end{pmatrix}, \quad \mathbf{a} \in \mathbb{H}^*, \quad \mathbf{b} \in \mathbb{H}, \quad \mathbf{0}^T = (0, 0). \quad (3.1)$$

Alternatively, if  $A$  is the  $4 \times 4$  real matrix corresponding to  $\mathbf{a}$ , through (2.2), and  $\mathbf{b} \in \mathbb{R}^4$  the vector made out of the components  $b_0, b_1, b_2, b_3$  of  $\mathbf{b}$  (see (2.3)),

$$\mathbf{b} = \begin{pmatrix} b_0 \\ b_3 \\ b_2 \\ b_1 \end{pmatrix},$$

then  $g$  may also be written as the  $5 \times 5$  real matrix,

$$g := (\mathbf{b}, A) = \begin{pmatrix} A & \mathbf{b} \\ \mathbf{0}^T & 1 \end{pmatrix}, \quad \mathbf{0}^T = (0, 0, 0, 0). \quad (3.2)$$

## Quaternionic affine group

In this real form  $G_{\text{aff}}^{\mathbb{H}}$  may be called the *group of dihedral similitude transformations of  $\mathbb{R}^4$* . We shall use both representations of  $G_{\text{aff}}^{\mathbb{H}}$  interchangeably.

For each one of these groups  $G_{\text{aff}}^{\mathbb{K}} = \mathbb{K} \rtimes \mathbb{K}^*$  there is exactly one non-trivial orbit of  $\mathbb{K}^*$  in the dual of  $\mathbb{K}$  and this orbit is open and free. Hence on a complex Hilbert space, each one of these groups has exactly one irreducible representation.

The irreducible representations of  $G_{\text{aff}}^{\mathbb{R}}$  and  $G_{\text{aff}}^{\mathbb{C}}$  are well known.

We shall compute below the one irreducible representation of  $G_{\text{aff}}^{\mathbb{H}}$ , both in a complex and in a quaternionic Hilbert space.

## Invariant measures of $\mathbb{H}^*$ and $G_{\text{aff}}^{\mathbb{H}}$

The group  $\mathbb{H}^*$  is unimodular. The Haar measure is,

$$d\mu_{\mathbb{H}^*} = \frac{d\mathbf{x}}{\|\mathbf{x}\|^4} = \frac{d\mathbf{x}}{(\det[\mathbf{x}])^2} = \frac{dx}{\|x\|^4}, \quad \text{where } d\mathbf{x} = dx = dx_0 dx_3 dx_2 dx_1. \quad (3.3)$$

The group  $G_{\text{aff}}^{\mathbb{H}}$  is non-unimodular. The left invariant measure is

$$d\mu_{\ell}(\mathbf{b}, A) = \frac{d\mathbf{b}}{\|\mathbf{a}\|^4} d\mu_{\mathbb{H}^*}(A) = \frac{d\mathbf{b} da}{\|\mathbf{a}\|^8} := \frac{d\mathbf{b} dA}{(\det[A])^2}, \quad (3.4)$$

which we shall also write as

$$d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = \frac{d\mathbf{b} da}{(\det[\mathbf{a}])^4}. \quad (3.5)$$

Similarly, the right Haar measure is

$$d\mu_r(\mathbf{b}, A) = d\mathbf{b} d\mu_{\mathbb{H}^*}(A) = \frac{d\mathbf{b} da}{\|\mathbf{a}\|^4} := \frac{d\mathbf{b} dA}{\det[A]}, \quad (3.6)$$

or alternatively written,

$$d\mu_r(\mathbf{b}, \mathbf{a}) = \frac{d\mathbf{b} da}{(\det[\mathbf{a}])^2}. \quad (3.7)$$

The *modular function*  $\Delta$ , such that  $d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = \Delta(\mathbf{b}, \mathbf{a}) d\mu_r(\mathbf{b}, \mathbf{a})$ , is

$$\Delta(\mathbf{b}, \mathbf{a}) = \frac{1}{(\det[\mathbf{a}])^2} = \frac{1}{|\mathbf{a}|^4} = \frac{1}{\|\mathbf{a}\|^4} = \frac{1}{\det[A]} := \Delta(\mathbf{b}, A). \quad (3.8)$$

## UIR of $G_{\text{aff}}^{\mathbb{H}}$ in a complex Hilbert space

From the general theory of semi-direct products of the type  $\mathbb{R}^n \rtimes H$ , where  $H$  is a subgroup of  $GL(n, \mathbb{R})$ , and which has open free orbits in the dual of  $\mathbb{R}^n$ , we know that  $G_{\text{aff}}^{\mathbb{H}}$  has exactly one unitary irreducible representation on a complex Hilbert space and moreover, this representation is square-integrable. Consider the Hilbert space  $\mathfrak{H}_{\mathbb{C}} = L^2_{\mathbb{C}}(\mathbb{R}^4, dx)$  and on it define the representation  $G_{\text{aff}}^{\mathbb{H}} \ni (\mathbf{b}, A) \mapsto U_{\mathbb{C}}(\mathbf{b}, A)$ ,

$$(U_{\mathbb{C}}(\mathbf{b}, A)f)(\mathbf{x}) = \frac{1}{(\det[A])^{\frac{1}{2}}} f(A^{-1}(\mathbf{x} - \mathbf{b})), \quad f \in \mathfrak{H}_{\mathbb{C}}. \quad (3.9)$$

This representation is unitary and irreducible.

The *Duflo-Moore operator*  $C$  is given in the Fourier domain as the multiplication operator

$$(\widehat{Cf})(\mathbf{k}) = C(\mathbf{k})\widehat{f}(\mathbf{k}), \quad \text{where } C(\mathbf{k}) = \left[ \frac{2\pi}{\|\mathbf{k}\|} \right]^2. \quad (3.10)$$

A vector  $f \in \mathfrak{H}_{\mathbb{C}}$  is *admissible* if it is in the domain of  $C$  i.e., if its Fourier transform  $\widehat{f}$  satisfies

$$(2\pi)^4 \int_{\mathbb{R}^4} \frac{|\widehat{f}(\mathbf{k})|^2}{\|\mathbf{k}\|^4} d\mathbf{k} < \infty.$$

## UIR of $G_{\text{aff}}^{\mathbb{H}}$ in a complex Hilbert space

Thus, for any two vectors  $\eta_1, \eta_2$  in the domain of  $C$  and for arbitrary  $f_1, f_2 \in \mathfrak{H}_{\mathbb{C}}$ , we have the *orthogonality relation*,

$$\int_{G_{\text{aff}}^{\mathbb{H}}} \langle f_1 | U_{\mathbb{C}}(\mathbf{b}, A)\eta_1 \rangle \langle \eta_2 | U_{\mathbb{C}}(\mathbf{b}, A)^* f_2 \rangle d\mu_{\ell}(\mathbf{b}, A) = \langle C\eta_2 | C\eta_1 \rangle \langle f_1 | f_2 \rangle, \quad (3.11)$$

which is the same as the operator equation

$$\int_{G_{\text{aff}}^{\mathbb{H}}} U_{\mathbb{C}}(\mathbf{b}, A)|\eta_1\rangle \langle \eta_2| U_{\mathbb{C}}(\mathbf{b}, A)^* d\mu_{\ell}(\mathbf{b}, A) = \langle C\eta_2 | C\eta_1 \rangle I_{\mathfrak{H}_{\mathbb{C}}}. \quad (3.12)$$

If  $\langle C\eta_2 | C\eta_1 \rangle \neq 0$ , we have the *resolution of the identity*

$$\frac{1}{\langle C\eta_2 | C\eta_1 \rangle} \int_{G_{\text{aff}}^{\mathbb{H}}} U_{\mathbb{C}}(\mathbf{b}, A)|\eta_1\rangle \langle \eta_2| U_{\mathbb{C}}(\mathbf{b}, A)^* d\mu_{\ell}(\mathbf{b}, A) = I_{\mathfrak{H}_{\mathbb{C}}}. \quad (3.13)$$

Given an admissible vector  $\eta$ , such that  $\|C\eta\|^2 = 1$ , we define the family of *coherent states or wavelets* as

$$\mathfrak{S}_{\mathbb{C}} = \{ \eta_{\mathbf{b}, A} = U_{\mathbb{C}}(\mathbf{b}, A)\eta \mid (\mathbf{b}, A) \in G_{\text{aff}}^{\mathbb{H}} \}, \quad (3.14)$$



## UIR of $G_{\text{aff}}^{\mathbb{H}}$ in a complex Hilbert space

which then satisfies the resolution of the identity,

$$\int_{G_{\text{aff}}^{\mathbb{H}}} |\eta_{\mathbf{b},A}\rangle \langle \eta_{\mathbf{b},A}| d\mu_{\ell}(\mathbf{b}, A) = I_{\mathfrak{H}_{\mathbb{C}}} . \quad (3.15)$$

The above representation could also be realized on the Hilbert space  $\mathfrak{K}_{\mathbb{C}} = L^2_{\mathbb{C}}(\mathbb{H}, d\mathbf{x})$  over the quaternions. We simply transcribe Eqs. (3.9) – (3.15) into this framework.

Thus, we define the representation  $G_{\text{aff}}^{\mathbb{H}} \ni (\mathbf{b}, \mathbf{a}) \mapsto U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})$ ,

$$(U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})f)(\mathbf{x}) = \frac{1}{\det[\mathbf{a}]} f(\mathbf{a}^{-1}(\mathbf{x} - \mathbf{b})), \quad f \in \mathfrak{K}_{\mathbb{C}}. \quad (3.16)$$

The *Duflo-Moore operator*  $C$  is given in the Fourier domain as the multiplication operator

$$(\widehat{Cf})(\mathfrak{k}) = C(\mathfrak{k})\widehat{f}(\mathfrak{k}), \quad \text{where } C(\mathfrak{k}) = \left[ \frac{2\pi}{|\mathfrak{k}|} \right]^2. \quad (3.17)$$

The admissibility condition is now

$$(2\pi)^4 \int_{\mathbb{R}^4} \frac{|\widehat{f}(\mathfrak{k})|^2}{|\mathfrak{k}|^4} d\mathfrak{k} < \infty,$$

## UIR of $G_{\text{aff}}^{\mathbb{H}}$ in a complex Hilbert space

and for any two vectors  $\eta_1, \eta_2$  in the domain of  $C$  and arbitrary  $f_1, f_2 \in \mathfrak{K}_{\mathbb{C}}$ , the orthogonality relation becomes

$$\int_{G_{\text{aff}}^{\mathbb{H}}} \langle f_1 | U_{\mathbb{C}}(\mathbf{b}, \mathbf{a}) \eta_1 \rangle \langle \eta_2 | U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})^* f_2 \rangle d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = \langle C\eta_2 | C\eta_1 \rangle \langle f_1 | f_2 \rangle, \quad (3.18)$$

with its operator version

$$\int_{G_{\text{aff}}^{\mathbb{H}}} U_{\mathbb{C}}(\mathbf{b}, \mathbf{a}) |\eta_1\rangle \langle \eta_2| U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})^* d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = \langle C\eta_2 | C\eta_1 \rangle I_{\mathfrak{K}_{\mathbb{C}}}. \quad (3.19)$$

Similarly, for  $\langle C\eta_2 | C\eta_1 \rangle \neq 0$ ,

$$\frac{1}{\langle C\eta_2 | C\eta_1 \rangle} \int_{G_{\text{aff}}^{\mathbb{H}}} U_{\mathbb{C}}(\mathbf{b}, \mathbf{a}) |\eta_1\rangle \langle \eta_2| U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})^* d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = I_{\mathfrak{K}_{\mathbb{C}}}. \quad (3.20)$$

The family of *coherent states or wavelets* are

$$\mathfrak{S}_{\mathbb{C}} = \{ \eta_{\mathbf{b}, \mathbf{a}} = U_{\mathbb{C}}(\mathbf{b}, \mathbf{a}) \eta | (\mathbf{b}, \mathbf{a}) \in G_{\text{aff}}^{\mathbb{H}} \}, \quad (3.21)$$

with the resolution of the identity,

$$\int_{G_{\text{aff}}^{\mathbb{H}}} |\eta_{\mathbf{b}, \mathbf{a}}\rangle \langle \eta_{\mathbf{b}, \mathbf{a}}| d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = I_{\mathfrak{K}_{\mathbb{C}}}. \quad (3.22)$$

## A quaternionic Hilbert space

We now proceed to construct a unitary irreducible representation of the quaternionic affine group  $G_{\text{aff}}^{\mathbb{H}}$  on a quaternionic Hilbert space. It will turn out that this representation has an intimate connection with the representation  $U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})$  in (3.16) on  $\mathfrak{K}_{\mathbb{C}}$ .

We consider the Hilbert space  $\mathfrak{H}_{\mathbb{H}}$ , of quaternionic valued functions over the quaternions. An element  $\mathbf{f} \in \mathfrak{H}_{\mathbb{H}}$  has the form

$$\mathbf{f}(\mathbf{x}) = \begin{pmatrix} f_1(\mathbf{x}) & -\overline{f_2(\mathbf{x})} \\ f_2(\mathbf{x}) & \overline{f_1(\mathbf{x})} \end{pmatrix}, \quad \mathbf{x} \in \mathbb{H}, \quad (4.1)$$

The norm is given by

$$\|\mathbf{f}\|_{\mathfrak{H}_{\mathbb{H}}}^2 = \int_{\mathbb{H}} \mathbf{f}(\mathbf{x})^\dagger \mathbf{f}(\mathbf{x}) \, d\mathbf{x} = \int_{\mathbb{H}} |\mathbf{f}(\mathbf{x})|^2 \, d\mathbf{x} = \left[ \int_{\mathbb{H}} (|f_1(\mathbf{x})|^2 + |f_2(\mathbf{x})|^2) \, d\mathbf{x} \right] \sigma_0, \quad (4.2)$$

the finiteness of which implies that both  $f_1$  and  $f_2$  have to be elements of  $\mathfrak{K}_{\mathbb{C}} = L_{\mathbb{C}}^2(\mathbb{H}, d\mathbf{x})$ , so that we may write

$$\|\mathbf{f}\|_{\mathfrak{H}_{\mathbb{H}}}^2 = (\|f_1\|_{\mathfrak{K}_{\mathbb{C}}}^2 + \|f_2\|_{\mathfrak{K}_{\mathbb{C}}}^2) \sigma_0.$$

In view of this, we may also write  $\mathfrak{H}_{\mathbb{H}} = L_{\mathbb{H}}^2(\mathbb{H}, d\mathbf{x})$ .

## UIR of $G_{\text{aff}}^{\mathbb{H}}$ in a quaternionic Hilbert space

In using the “bra-ket” notation we shall use the notation and convention:

$$(\mathbf{f} | = \begin{pmatrix} \langle f_1 | & \langle f_2 | \\ -\langle \bar{f}_2 | & \langle \bar{f}_1 | \end{pmatrix}, \quad \text{and} \quad | \mathbf{f} \rangle = \begin{pmatrix} |f_1\rangle & -|\bar{f}_2\rangle \\ |f_2\rangle & |\bar{f}_1\rangle \end{pmatrix}, \quad (4.3)$$

The scalar product of two vectors  $\mathbf{f}, \mathbf{f}' \in \mathfrak{H}_{\mathbb{H}}$  is

$$\begin{aligned} (\mathbf{f} | \mathbf{f}') &= \int_{\mathbb{H}} \mathbf{f}(\mathbf{x})^\dagger \mathbf{f}'(\mathbf{x}) d\mathbf{x} \\ &= \begin{pmatrix} \langle f_1 | f'_1 \rangle_{\mathfrak{H}_{\mathbb{C}}} + \langle f_2 | f'_2 \rangle_{\mathfrak{H}_{\mathbb{C}}} & -\langle f_2' | \bar{f}_1 \rangle_{\mathfrak{H}_{\mathbb{C}}} + \langle f_1' | \bar{f}_2 \rangle_{\mathfrak{H}_{\mathbb{C}}} \\ \langle \bar{f}'_2 | f_1 \rangle_{\mathfrak{H}_{\mathbb{C}}} - \langle \bar{f}'_1 | f_2 \rangle_{\mathfrak{H}_{\mathbb{C}}} & \langle f_1' | f_1 \rangle_{\mathfrak{H}_{\mathbb{C}}} + \langle f_2' | f_2 \rangle_{\mathfrak{H}_{\mathbb{C}}} \end{pmatrix} \end{aligned} \quad (4.4)$$

Note that

$$(\mathbf{f} | \mathbf{f}')^\dagger = (\mathbf{f}' | \mathbf{f}).$$

Multiplication by quaternions on  $\mathfrak{H}_{\mathbb{H}}$  is defined from the right:

$$(\mathfrak{H}_{\mathbb{H}} \times \mathbb{H}) \ni (\mathbf{f}, \mathbf{q}) \longmapsto \mathbf{f}\mathbf{q}, \quad \text{such that} \quad (\mathbf{f}\mathbf{q})(\mathbf{x}) = \mathbf{f}(\mathbf{x})\mathbf{q},$$

i.e., we take  $\mathfrak{H}_{\mathbb{H}}$  to be a right quaternionic Hilbert space.

## A quaternionic Hilbert space

This convention is consistent with the scalar product (4.4) in the sense that

$$(\mathbf{f} | \mathbf{f}') = (\mathbf{f} | \mathbf{f}')\mathbf{q} \quad \text{and} \quad (\mathbf{f}\mathbf{q} | \mathbf{f}') = \mathbf{q}^\dagger (\mathbf{f} | \mathbf{f}').$$

On the other hand, the action of operators  $\mathbf{A}$  on vectors  $\mathbf{f} \in \mathfrak{H}_{\mathbb{H}}$  will be from the left  $(\mathbf{A}, \mathbf{q}) \mapsto \mathbf{A}\mathbf{f}$ . In particular, an operator  $A$  on  $\mathfrak{R}_{\mathbb{C}}$  defines an operator  $\mathbf{A}$  on  $\mathfrak{H}_{\mathbb{H}}$  as,

$$(\mathbf{A}\mathbf{f})(\mathbf{x}) = \begin{pmatrix} (Af_1)(\mathbf{x}) & -\overline{(Af_2)(\mathbf{x})} \\ (Af_2)(\mathbf{x}) & \overline{(Af_1)(\mathbf{x})} \end{pmatrix}.$$

Multiplication of operators by quaternions will also be from the left. Thus,  $\mathbf{q}\mathbf{A}$  acts on the vector  $\mathbf{f}$  in the manner

$$(\mathbf{q}\mathbf{A}\mathbf{f})(\mathbf{x}) = \mathbf{q}(\mathbf{A}\mathbf{f})(\mathbf{x}).$$

We shall also need the “rank-one operator”

$$\begin{aligned} |\mathbf{f}\rangle\langle\mathbf{f}'| &= \begin{pmatrix} |f_1\rangle & -|\bar{f}_2\rangle \\ |f_2\rangle & |\bar{f}_1\rangle \end{pmatrix} \begin{pmatrix} \langle f'_1| & \langle f'_2| \\ -\langle \bar{f}'_2| & \langle \bar{f}'_1| \end{pmatrix} \\ &= \begin{pmatrix} |f_1\rangle\langle f'_1| + |\bar{f}_2\rangle\langle \bar{f}'_2| & |f_1\rangle\langle f'_2| - |\bar{f}_2\rangle\langle \bar{f}'_1| \\ -|\bar{f}_1\rangle\langle \bar{f}'_2| + |f_2\rangle\langle f'_1| & |\bar{f}_1\rangle\langle \bar{f}'_1| + |f_2\rangle\langle f'_2| \end{pmatrix} \end{aligned} \quad (4.5)$$

## A quaternionic Hilbert space

An orthonormal basis in  $\mathfrak{H}_{\mathbb{H}}$  can be built using an orthonormal basis in  $\mathfrak{K}_{\mathbb{C}}$ . Indeed, let  $\{\phi_n\}_{n=0}^{\infty}$  be an orthonormal basis of  $\mathfrak{K}_{\mathbb{C}} = L_{\mathbb{C}}^2(\mathbb{H}, d\mathbf{x})$ . Define the vectors

$$|\Phi_n\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} |\phi_n\rangle & |\phi_n\rangle \\ -|\bar{\phi}_n\rangle & |\bar{\phi}_n\rangle \end{pmatrix}, \quad n = 0, 1, 2, \dots, \quad (4.6)$$

in  $\mathfrak{H}_{\mathbb{H}}$ . It is easy to check that these vectors are orthonormal in  $\mathfrak{H}_{\mathbb{H}}$ . The fact that they form a basis follows from the fact that the vectors  $\{\phi_n\}_{n=0}^{\infty}$  are a basis of  $L_{\mathbb{C}}^2(\mathbb{H}, d\mathbf{x})$ . Indeed, writing

$$|\mathbf{f}\rangle = \begin{pmatrix} |f_1\rangle & -|\bar{f}_2\rangle \\ |f_2\rangle & |\bar{f}_1\rangle \end{pmatrix} \in L_{\mathbb{H}}^2(\mathbb{H}, d\mathbf{x}),$$

we easily verify that

$$|\mathbf{f}\rangle = \sum_{n=0}^{\infty} |\Phi_n\rangle q_n,$$

where

$$q_n = (\Phi_n | \mathbf{f}) = \frac{1}{\sqrt{2}} \begin{pmatrix} \langle \phi_n | f_1 \rangle_{\mathfrak{K}_{\mathbb{C}}} - \langle \bar{\phi}_n | f_2 \rangle_{\mathfrak{K}_{\mathbb{C}}} & -\langle f_2 | \bar{\phi}_n \rangle_{\mathfrak{K}_{\mathbb{C}}} - \langle f_1 | \phi_n \rangle_{\mathfrak{K}_{\mathbb{C}}} \\ \langle \bar{f}_2 | \phi_n \rangle_{\mathfrak{K}_{\mathbb{C}}} + \langle \bar{f}_1 | \bar{\phi}_n \rangle_{\mathfrak{K}_{\mathbb{C}}} & \langle f_1 | \phi_n \rangle_{\mathfrak{K}_{\mathbb{C}}} - \langle f_2 | \bar{\phi}_n \rangle_{\mathfrak{K}_{\mathbb{C}}} \end{pmatrix}.$$

## Representation of $G_{\text{aff}}^{\mathbb{H}}$ on $\mathfrak{H}_{\mathbb{H}}$

A representation of  $G_{\text{aff}}^{\mathbb{H}}$  on  $\mathfrak{H}_{\mathbb{H}}$  can be obtained by simply transcribing (3.16) into the present context. We define the operators  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$  on  $\mathfrak{H}_{\mathbb{H}}$ :

$$(\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f})(\mathbf{x}) = \frac{1}{\det[\mathbf{a}]} \mathbf{f}(\mathbf{a}^{-1}(\mathbf{x} - \mathbf{b})), \quad \mathbf{f} \in \mathfrak{H}_{\mathbb{H}}, \quad (4.7)$$

which by (3.16) and (4.3) can also be written as

$$|\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f}\rangle = \begin{pmatrix} |U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})f_1\rangle & -|\overline{U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})f_2}\rangle \\ U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})|f_2\rangle & |U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})f_1\rangle \end{pmatrix}. \quad (4.8)$$

The unitarity of this representation is easy to verify. Indeed,

$$\|\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f}\|^2 = \int_{\mathbb{H}} (|(U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})f_1)(\mathbf{x})|^2 + |(U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})f_2)(\mathbf{x})|^2) d\mathbf{x} \sigma_0,$$

which, by the unitarity of the representation  $U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})$  on  $\mathfrak{K}_{\mathbb{C}}$  gives

$$\|\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f}\|_{\mathfrak{H}_{\mathbb{H}}}^2 = (\|f_1\|_{\mathfrak{K}_{\mathbb{C}}}^2 + \|f_2\|_{\mathfrak{K}_{\mathbb{C}}}^2) \sigma_0 = \|\mathbf{f}\|_{\mathfrak{H}_{\mathbb{H}}}^2.$$

Similarly, the irreducibility of  $U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})$  on  $\mathfrak{K}_{\mathbb{C}}$  leads to the irreducibility of  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$ .

## Square-integrability of $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$

Using the Duflo-Moore operator  $C$  in (3.17) for the representation  $U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})$  (see (3.16)), we define the Duflo-Moore operator  $\mathbf{C}$  for the representation  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$ :

$$(\mathbf{C}\mathbf{f})(\mathbf{x}) = \begin{pmatrix} (Cf_1)(\mathbf{x}) & -\overline{(Cf_2)(\mathbf{x})} \\ (Cf_2)(\mathbf{x}) & \overline{(Cf_1)(\mathbf{x})} \end{pmatrix}.$$

We say that the vector  $\mathbf{f}$  is *admissible for the representation*  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$  if it is in the domain of  $\mathbf{C}$ , i.e., if both  $f_1$  and  $f_2$  are admissible for the representation  $U_{\mathbb{C}}(\mathbf{b}, \mathbf{a})$ . It is then easy to see that the set of admissible vectors is dense in  $\mathfrak{H}_{\mathbb{H}}$ .

Let  $\mathbf{f}$  and  $\mathbf{f}'$  be two admissible vectors. Then from (4.8), (4.5) and (3.19) we get

$$\int_{G_{\text{aff}}^{\mathbb{H}}} |\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f})(\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f}')| d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = \mathbf{q} I_{\mathfrak{H}_{\mathbb{H}}}, \quad (4.9)$$

where  $\mathbf{q}$  denotes the operator of multiplication from the left, on the Hilbert space  $\mathfrak{H}_{\mathbb{H}}$ , by the quaternion

$$\mathbf{q} = \begin{pmatrix} \langle Cf_1' | Cf_1 \rangle_{\mathfrak{H}_{\mathbb{C}}} + \langle \overline{Cf_2'} | \overline{Cf_2} \rangle_{\mathfrak{H}_{\mathbb{C}}} & \langle Cf_2' | Cf_1 \rangle_{\mathfrak{H}_{\mathbb{C}}} - \langle \overline{Cf_1'} | \overline{Cf_2} \rangle_{\mathfrak{H}_{\mathbb{C}}} \\ \langle Cf_1' | Cf_2 \rangle_{\mathfrak{H}_{\mathbb{C}}} - \langle \overline{Cf_2'} | \overline{Cf_1} \rangle_{\mathfrak{H}_{\mathbb{C}}} & \langle \overline{Cf_1'} | \overline{Cf_1} \rangle_{\mathfrak{H}_{\mathbb{C}}} + \langle Cf_2' | Cf_2 \rangle_{\mathfrak{H}_{\mathbb{C}}} \end{pmatrix}. \quad (4.10)$$



## Square-integrability of $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$

Equation (4.9) expresses the square-integrability condition for the representation  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$ .

In particular, with  $\mathbf{f} = \mathbf{f}'$ , we get the resolution of the identity,

$$\left[ \|\mathbf{C}\mathbf{f}\|_{\mathfrak{H}_{\mathbb{H}}}^2 \right]^{-1} \int_{G_{\mathbf{aff}}^{\mathbb{H}}} |\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f}|(\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\mathbf{f})| d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = I_{\mathfrak{H}_{\mathbb{H}}}. \quad (4.11)$$

## Wavelets and reproducing kernels

Let  $\eta \in \mathfrak{S}_{\mathbb{H}}$  be an admissible vector for the representation  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$ , normalized so that

$$\|\mathbf{C}\eta\|^2 = 1.$$

We define the *quaternionic wavelets or coherent states* to be the vectors

$$\mathfrak{S}_{\mathbb{H}} = \{\eta_{\mathbf{b}, \mathbf{a}} = \mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})\eta \mid (\mathbf{b}, \mathbf{a}) \in G_{\text{aff}}^{\mathbb{H}}\}, \quad (5.1)$$

By virtue of (4.9) they satisfy the resolution of the identity

$$\int_{G_{\text{aff}}^{\mathbb{H}}} |\eta_{\mathbf{b}, \mathbf{a}}\rangle \langle \eta_{\mathbf{b}, \mathbf{a}}| d\mu_{\ell}(\mathbf{b}, \mathbf{a}) = I_{\mathfrak{S}_{\mathbb{H}}}. \quad (5.2)$$

There is the associated reproducing kernel  $\mathbf{K} : G_{\text{aff}}^{\mathbb{H}} \times G_{\text{aff}}^{\mathbb{H}} \rightarrow \mathbb{H}$ ,

$$\mathbf{K}(\bar{\mathbf{b}}, \bar{\mathbf{a}}; \mathbf{b}', \mathbf{a}') = (\eta_{\mathbf{b}, \mathbf{a}} \mid \eta_{\mathbf{b}', \mathbf{a}'})_{\mathfrak{S}_{\mathbb{H}}}, \quad (5.3)$$

with the usual properties,

$$\begin{aligned} \mathbf{K}(\bar{\mathbf{b}}, \bar{\mathbf{a}}; \mathbf{b}', \mathbf{a}') &= \overline{\mathbf{K}(\bar{\mathbf{b}'}, \bar{\mathbf{a}'}; \mathbf{b}, \mathbf{a})}, & \mathbf{K}(\bar{\mathbf{b}}, \bar{\mathbf{a}}; \mathbf{b}, \mathbf{a}) &> 0, \\ \int_{G_{\text{aff}}^{\mathbb{H}}} \mathbf{K}(\bar{\mathbf{b}}, \bar{\mathbf{a}}; \mathbf{b}'', \mathbf{a}'') \mathbf{K}(\bar{\mathbf{b}'}, \bar{\mathbf{a}'}; \mathbf{b}', \mathbf{a}') d\mu_{\ell}(\mathbf{b}'', \mathbf{a}'') &= \mathbf{K}(\bar{\mathbf{b}}, \bar{\mathbf{a}}; \mathbf{b}', \mathbf{a}'). \end{aligned} \quad (5.4)$$

## Other questions to consider:

1. Unitary embedding of  $\mathfrak{S}_{\mathbb{H}}$  into  $L^2_{\mathbb{H}}(G_{\text{aff}}^{\mathbb{H}}, d\mu_{\ell})$ .
2. Extensions of the representation  $\mathbf{U}_{\mathbb{H}}(\mathbf{b}, \mathbf{a})$ , e.g. by multiplying from the right by the  $SU(2)$  part of  $\mathbf{a}$ .
3. Discretization