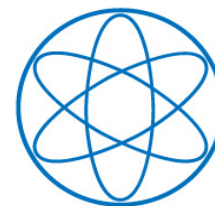


Dark Matter

Alejandro Ibarra

Technische Universität München



Summer School
on Cosmology
ICTP, Trieste
August 2014

Direct detection

DM nucleus \rightarrow DM nucleus



Indirect detection

DM DM $\rightarrow \gamma X, e^+e^- \dots$ (annihilation)

DM $\rightarrow \gamma X, e^+X \dots$ (decay)

Collider searches

pp \rightarrow DM X

Direct detection

DM nucleus \rightarrow DM nucleus

Indirect detection

DM DM $\rightarrow \gamma X, e^+e^- \dots$ (annihilation)

DM $\rightarrow \gamma X, e^+X \dots$ (decay)

Collider searches

pp \rightarrow DM X

Direct detection

DM nucleus \rightarrow DM nucleus

Indirect detection

DM DM $\rightarrow \gamma X, e^+e^- \dots$ (annihilation)

DM $\rightarrow \gamma X, e^+X \dots$ (decay)

Collider searches

pp \rightarrow DM X

Direct detection

DM nucleus \rightarrow DM nucleus

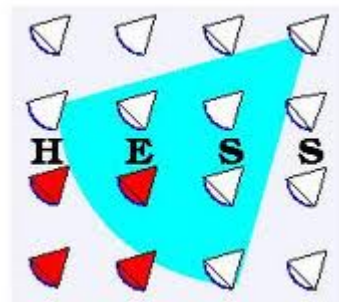
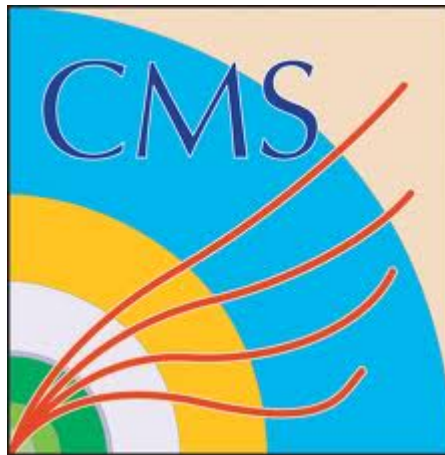
Indirect detection

DM DM $\rightarrow \gamma X, e^+e^- \dots$ (annihilation)

DM $\rightarrow \gamma X, e^+X \dots$ (decay)

Collider searches

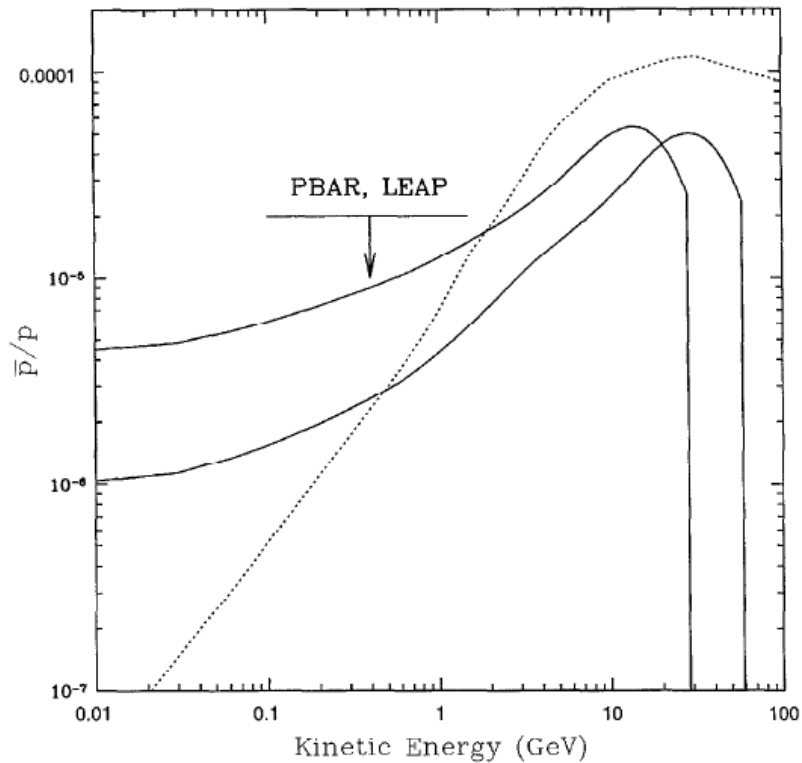
pp \rightarrow DM X



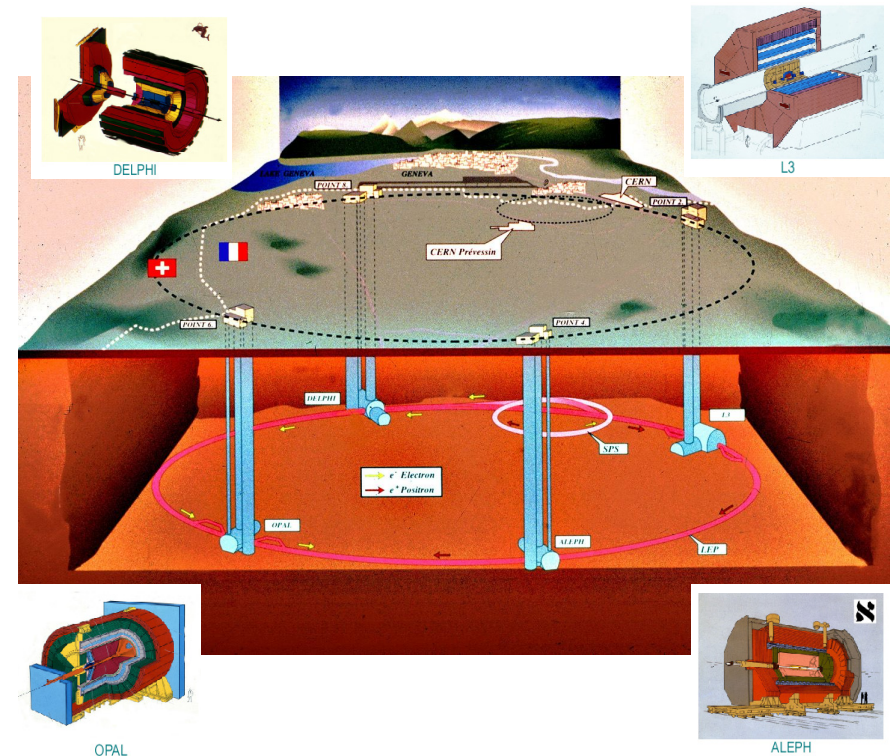
Fifteen years ago...

Cosmic antiprotons

G. Jungman et al. / *Physics Reports* 267 (1996) 195–373



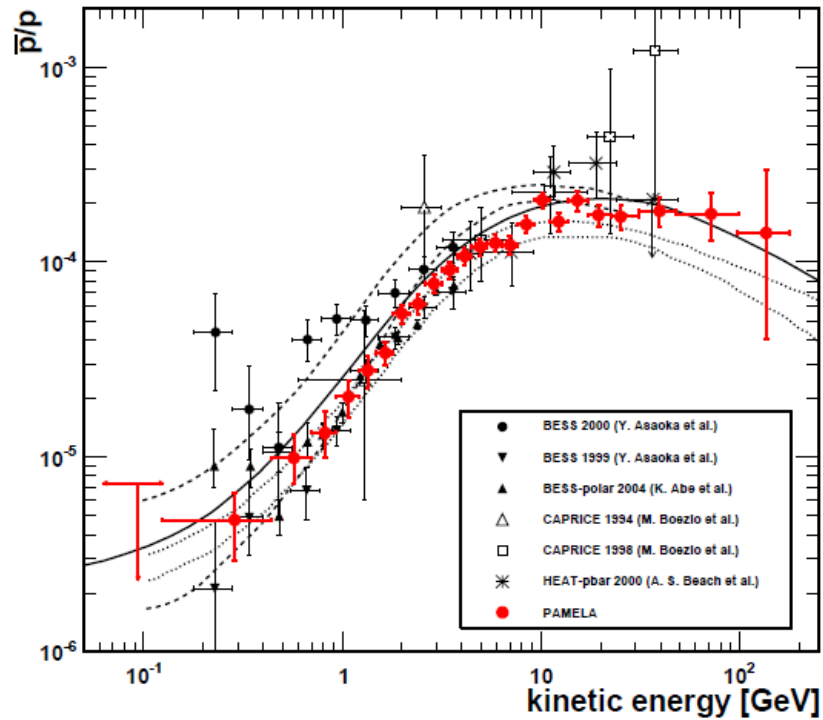
Collider experiments



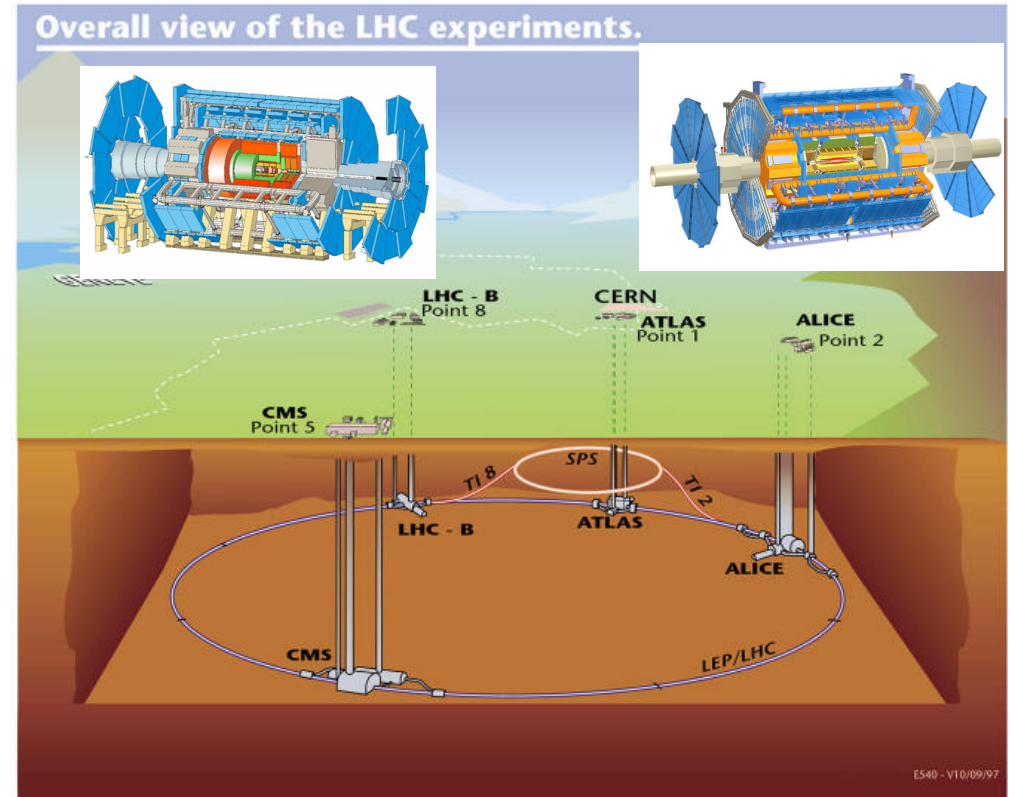
Year	Beam energy [GeV]	Maximum \mathcal{L}_{int}	Total luminosity [pb^{-1}]	Average luminosity rate [$\text{pb}^{-1}/\text{day}$]
1994	45.6	0.045	64	0.31
1995	45.6 – 70.0	0.050	47	0.23
1996	80.5 – 86.0	0.040	25	0.17
1997	91.0 – 92.0	0.055	75	0.66
1998	94.5	0.075	200	1.16
1999	96.0 – 101.0	0.083	254	1.35
2000*	100.0 – 104.3	0.055	71	0.96

Today

Cosmic antiprotons



Collider experiments

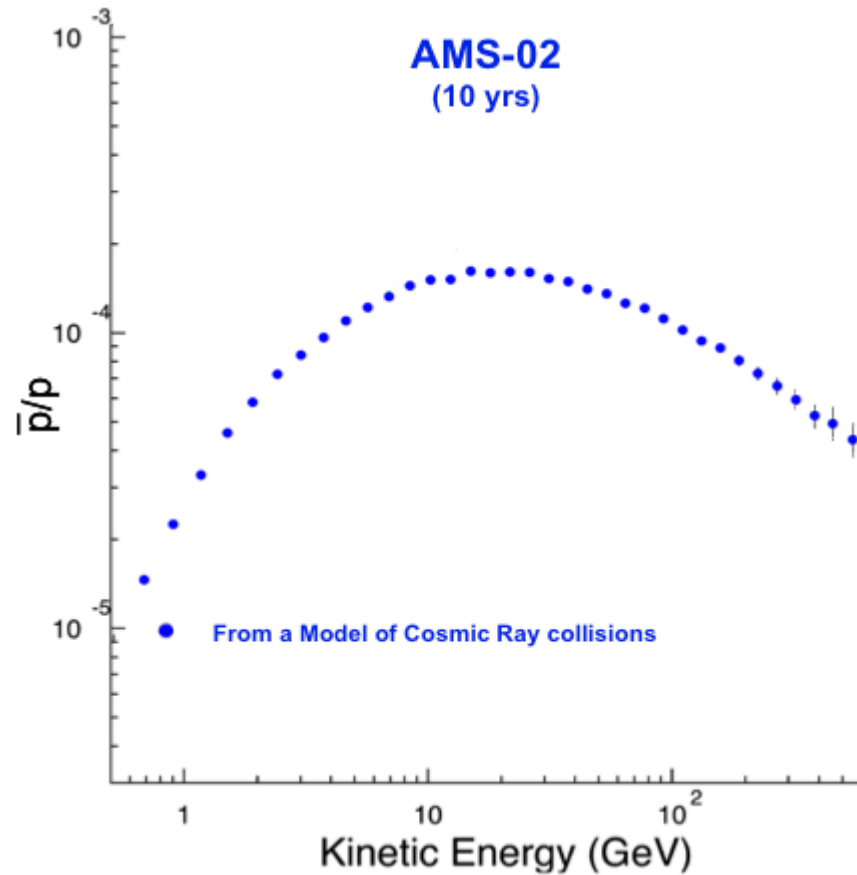


Beam energy: 4000 GeV

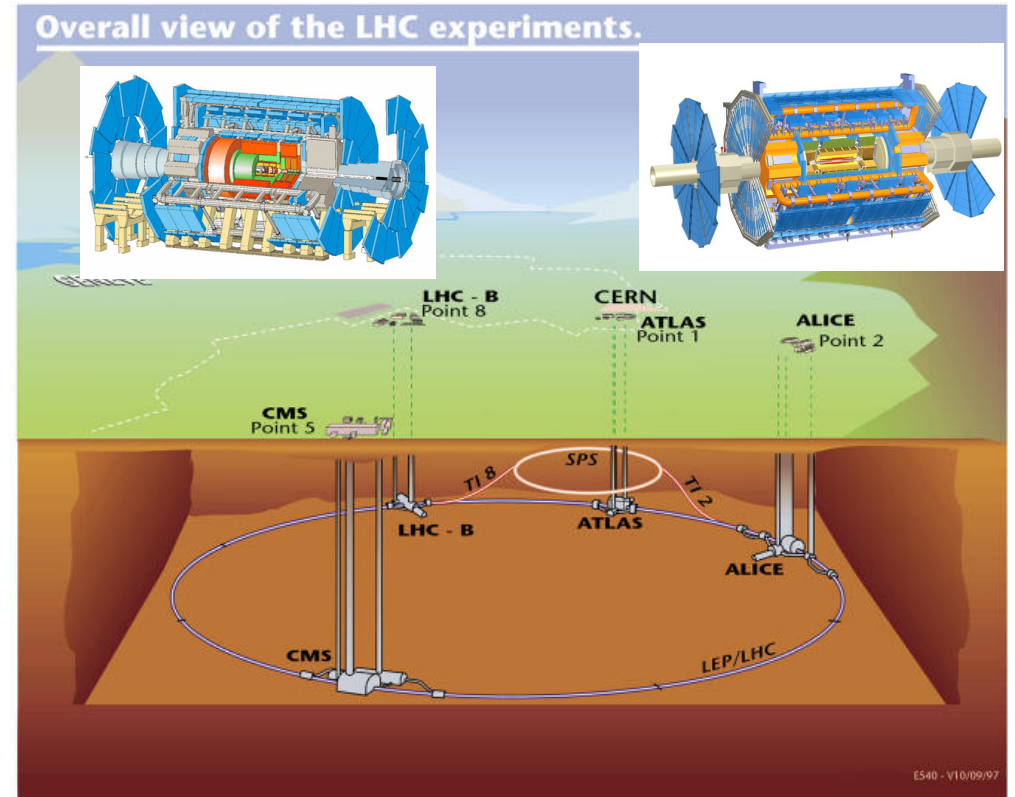
Integrated luminosity: 23.26 fb⁻¹

End of this decade

Cosmic antiprotons



Collider experiments



Beam energy: 6500 GeV

Integrated luminosity: 500 fb⁻¹

Dark matter distribution

The universe $1\ \mu\text{s}$ after the Big Bang

$z = 20.0$

200 million years after the Big Bang

50 Mpc/h



$z = 0.0$

50 Mpc/h



Volker Springel
Max-Planck-Institute
for Astrophysics



z=0.0

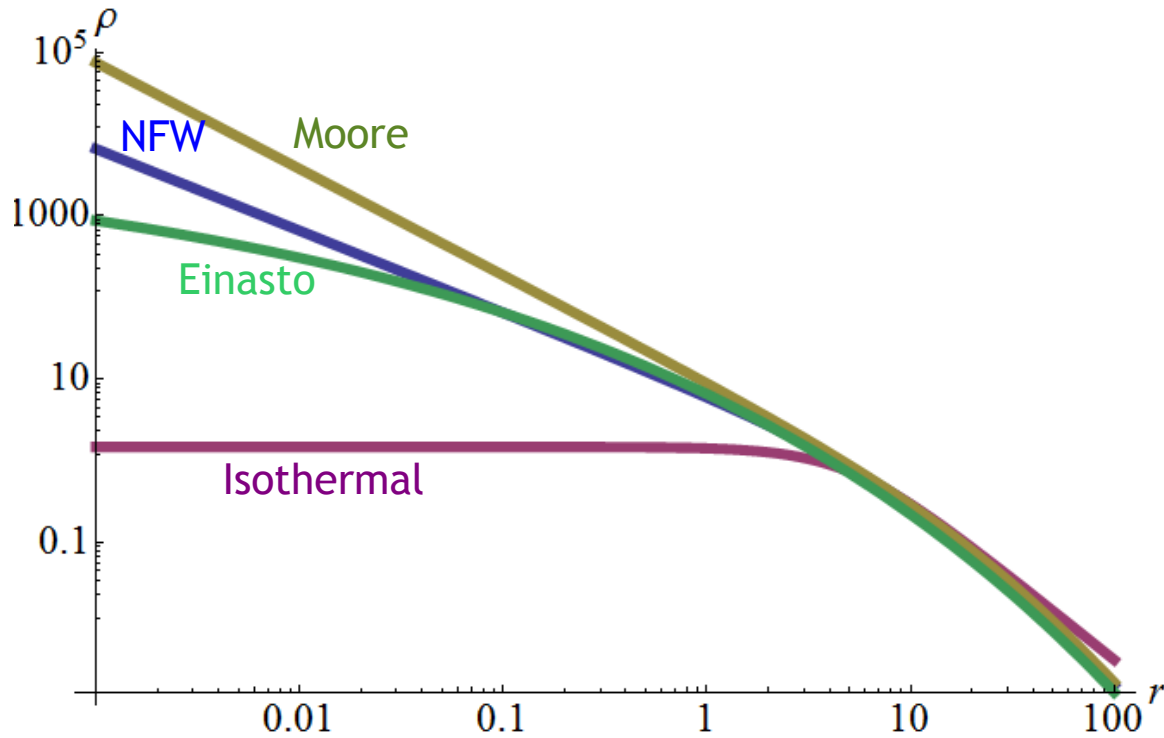
Distance Sun to Milky Way Center ~ 8.5 kpc

8 kpc



Density distribution of dark matter particles:

- Assume spherical symmetry (in a first approximation).
- Radial distribution:



NFW, Isothermal, Moore

$$\rho(r) = \frac{\rho_0}{(r/r_c)^\gamma [1 + (r/r_c)^\alpha]^{(\beta-\gamma)/\alpha}}$$

Halo model	α	β	γ	r_c (kpc)
Navarro, Frenk, White	1	3	1	20
Isothermal	2	2	0	3.5
Moore	1.5	3	1.5	28

Einasto

$$\rho(r) = \rho_0 \exp \left[-\frac{2}{\alpha} \left(\left(\frac{r}{r_s} \right)^\alpha - 1 \right) \right]$$

$$\alpha = 0.17, r_s = 20 \text{ kpc}$$

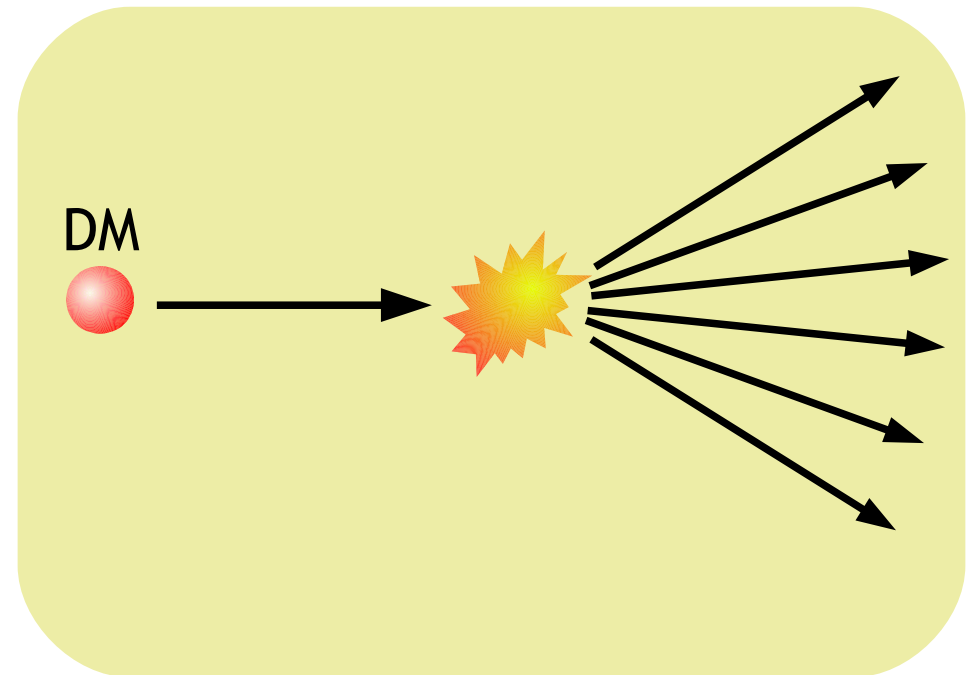
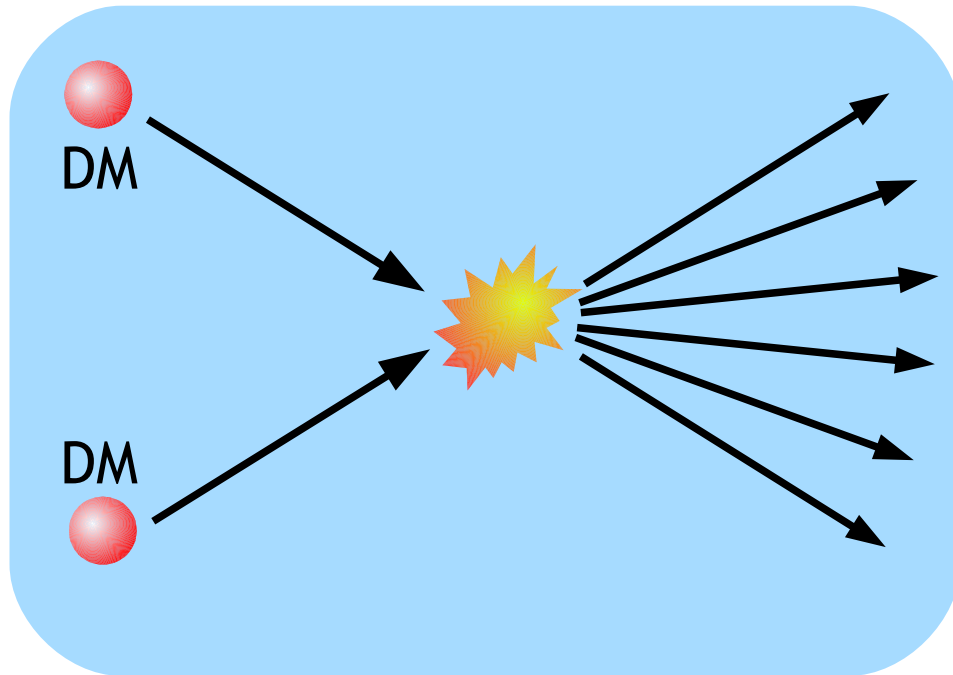
- Normalized such that the local DM density is $\rho(r=8.5 \text{ kpc}) = 0.38 \text{ GeV/cm}^3$

Indirect Dark Matter Searches

Indirect dark matter searches

General idea:

1) Dark matter particles annihilate or decay producing a flux of stable particles: photons, electrons, protons, positrons, antiprotons or (anti-)neutrinos.



Indirect dark matter searches

General idea:

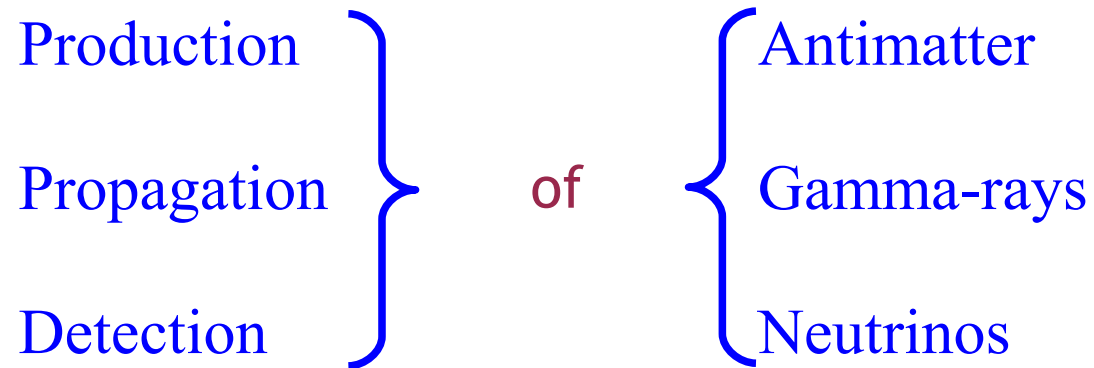
- 1) Dark matter particles annihilate or decay producing a flux of stable particles: photons, electrons, protons, positrons, antiprotons or (anti-)neutrinos.
- 2) These particles propagate through the galaxy and through the Solar System. Some of them will reach the Earth.

Indirect dark matter searches

General idea:

- 1) Dark matter particles annihilate or decay producing a flux of stable particles: photons, electrons, protons, positrons, antiprotons or (anti-)neutrinos.
- 2) These particles propagate through the galaxy and through the Solar System. Some of them will reach the Earth.
- 3) The products of the dark matter annihilations or decays are detected **together with other particles produced in astrophysical processes** (for example, cosmic ray collisions with nuclei in the interstellar medium). The existence of dark matter can then be inferred if there is a significant excess in the fluxes compared to the expected astrophysical backgrounds.

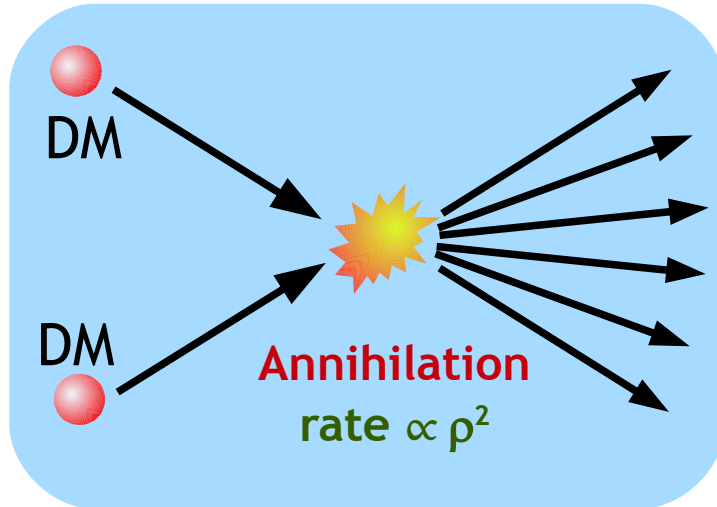
Indirect dark matter searches



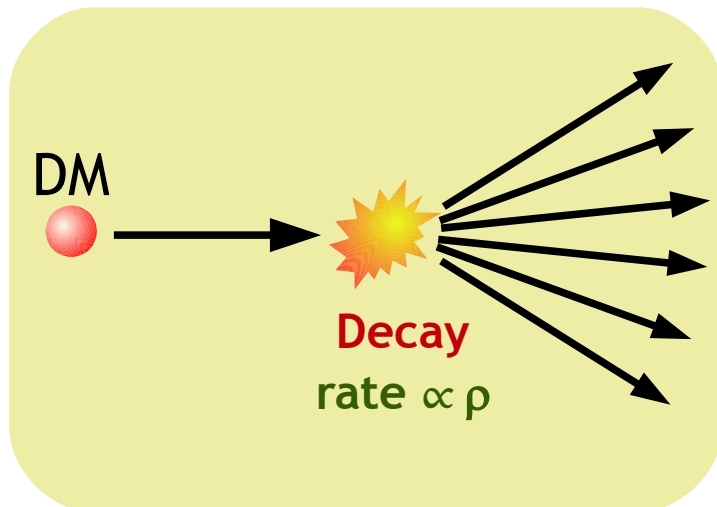
Antimatter

Production

The production is described by the **source function**: number of particles produced at a given position per unit volume, unit time and unit energy.

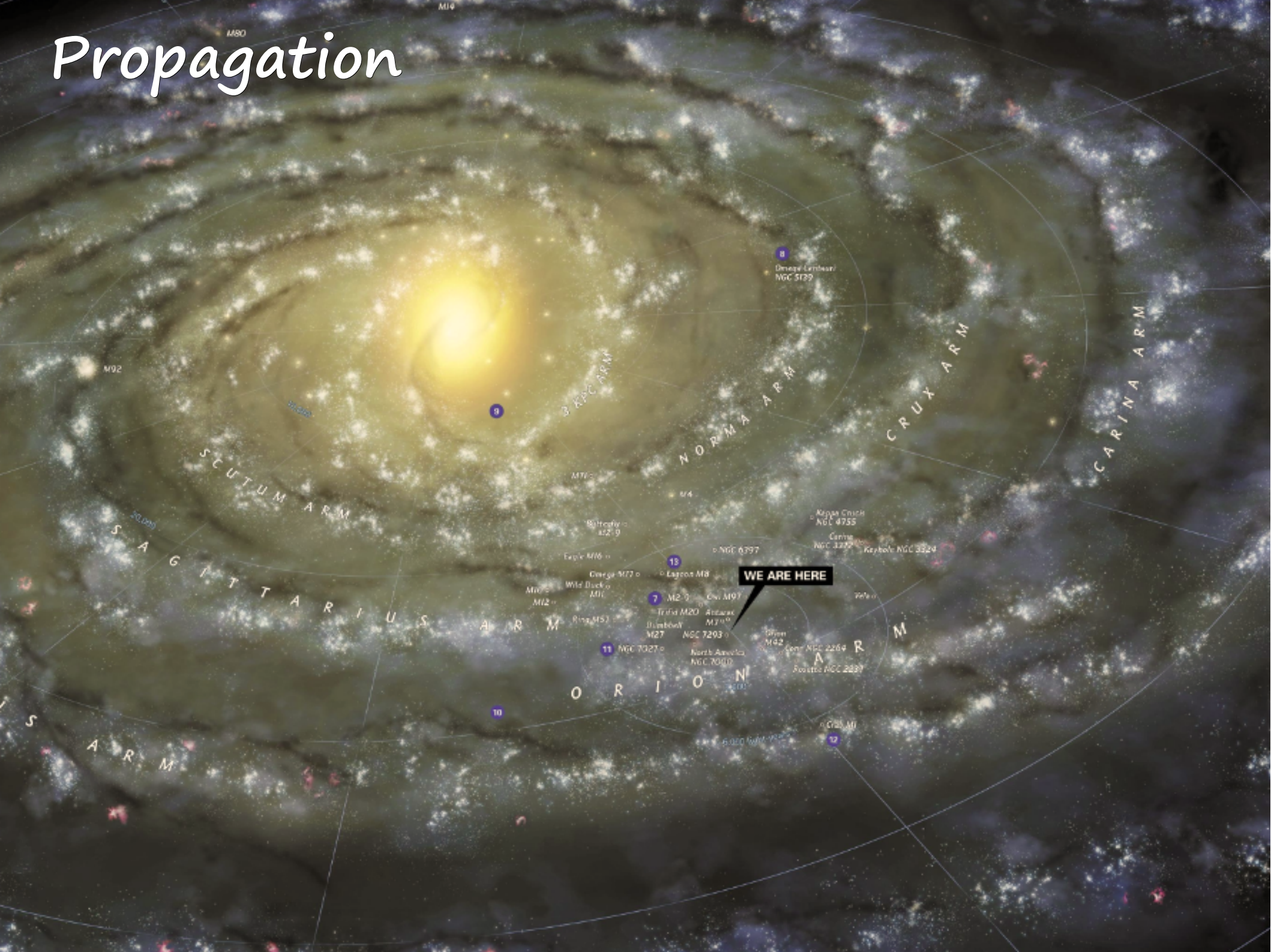


$$Q(E, \vec{r}) = \frac{1}{2} \frac{\rho^2(\vec{r})}{m_{\text{DM}}^2} \langle \sigma v \rangle \frac{dN}{dE}$$

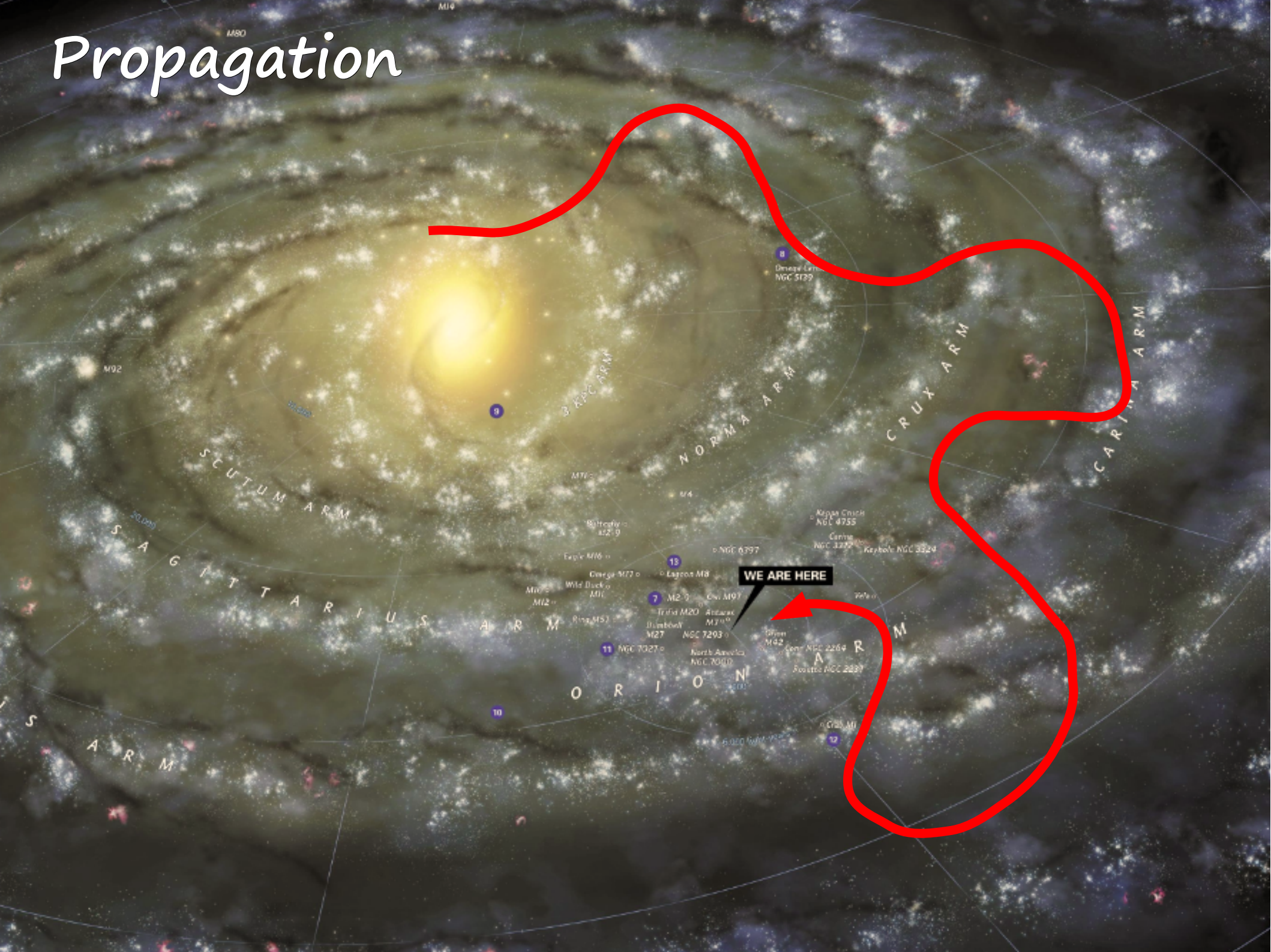


$$Q(E, \vec{r}) = \frac{\rho(\vec{r})}{m_{\text{DM}}} \frac{1}{\tau_{\text{DM}}} \frac{dN}{dE}$$

Propagation



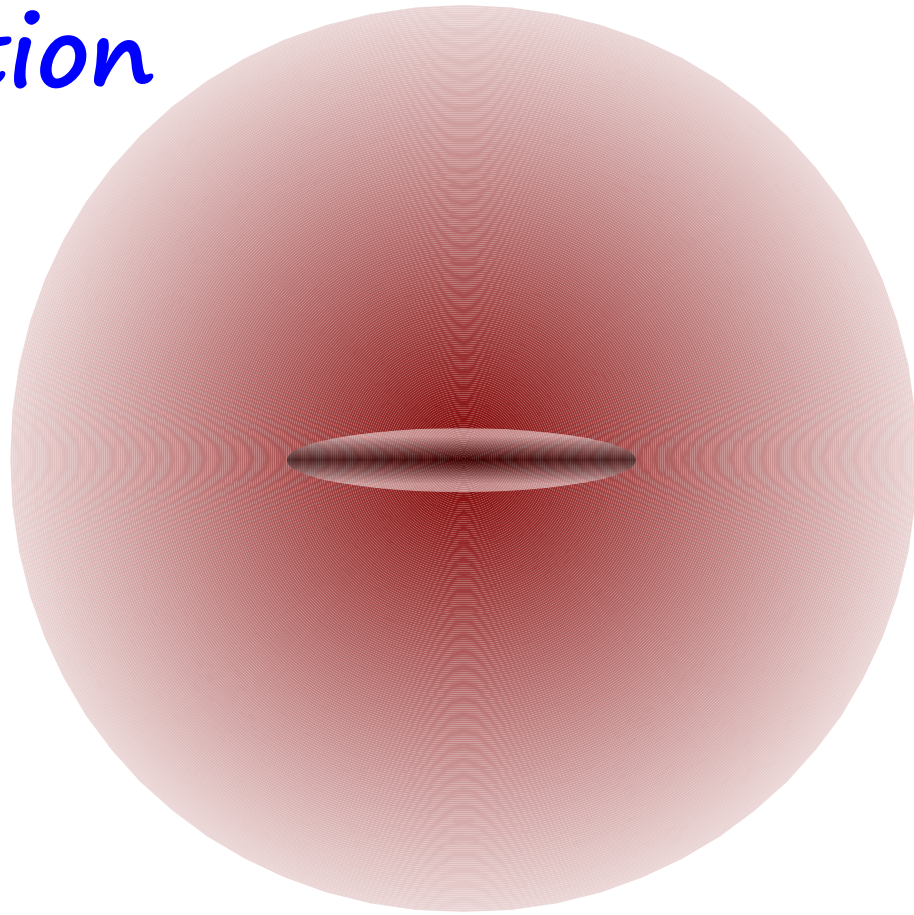
Propagation



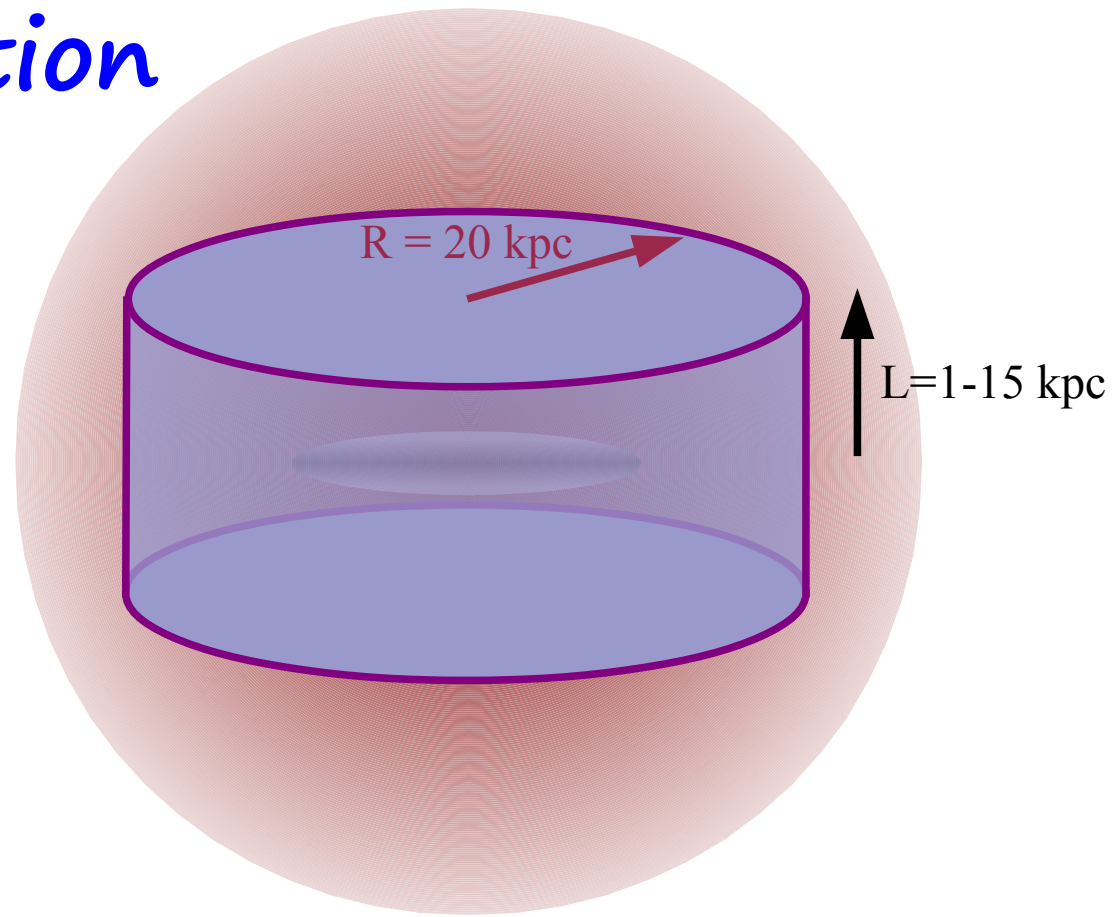
Propagation



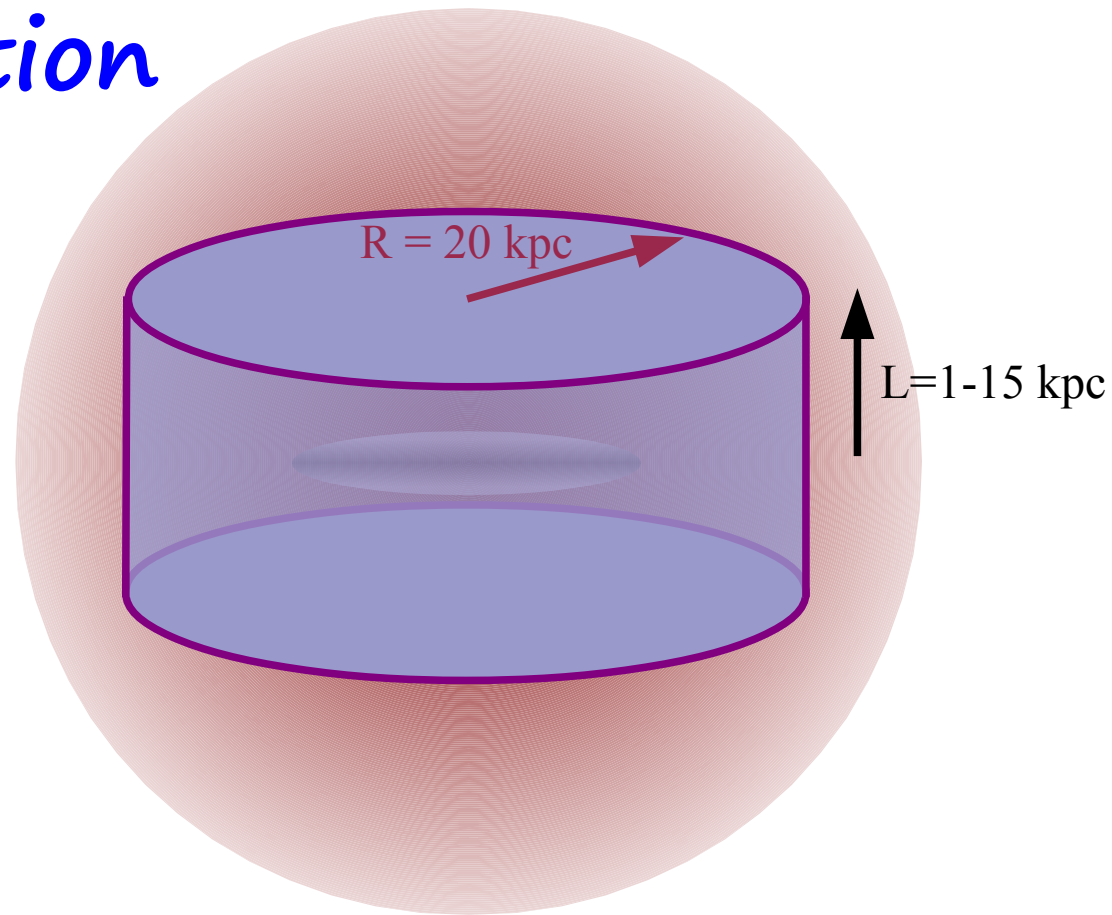
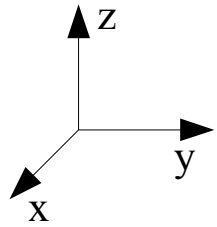
Propagation



Propagation



Propagation



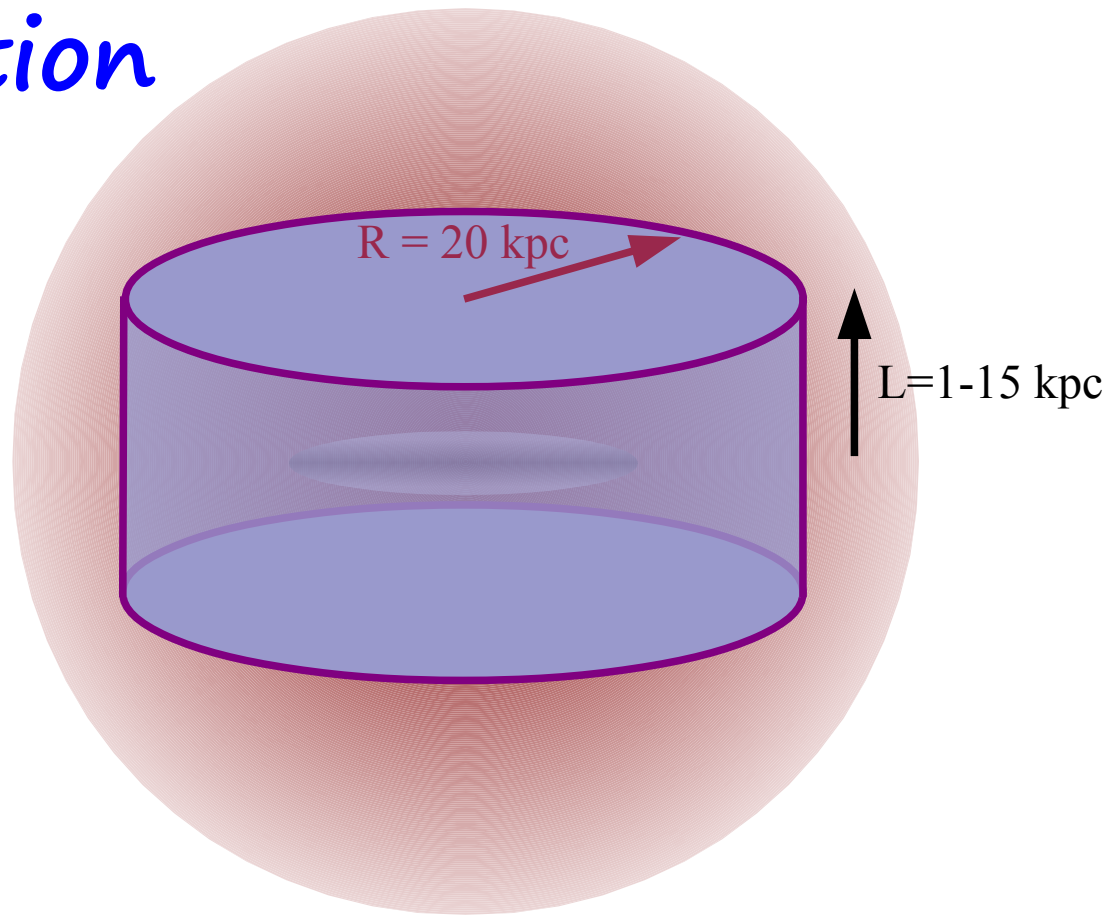
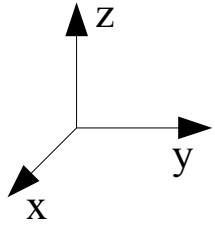
$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

f : number density of antiparticles per unit kinetic energy

interstellar antimatter flux:

$$\Phi^{\text{IS}}(T) = \frac{dN}{dt dS dT d\Omega} = \frac{v}{4\pi} f(T)$$

Propagation

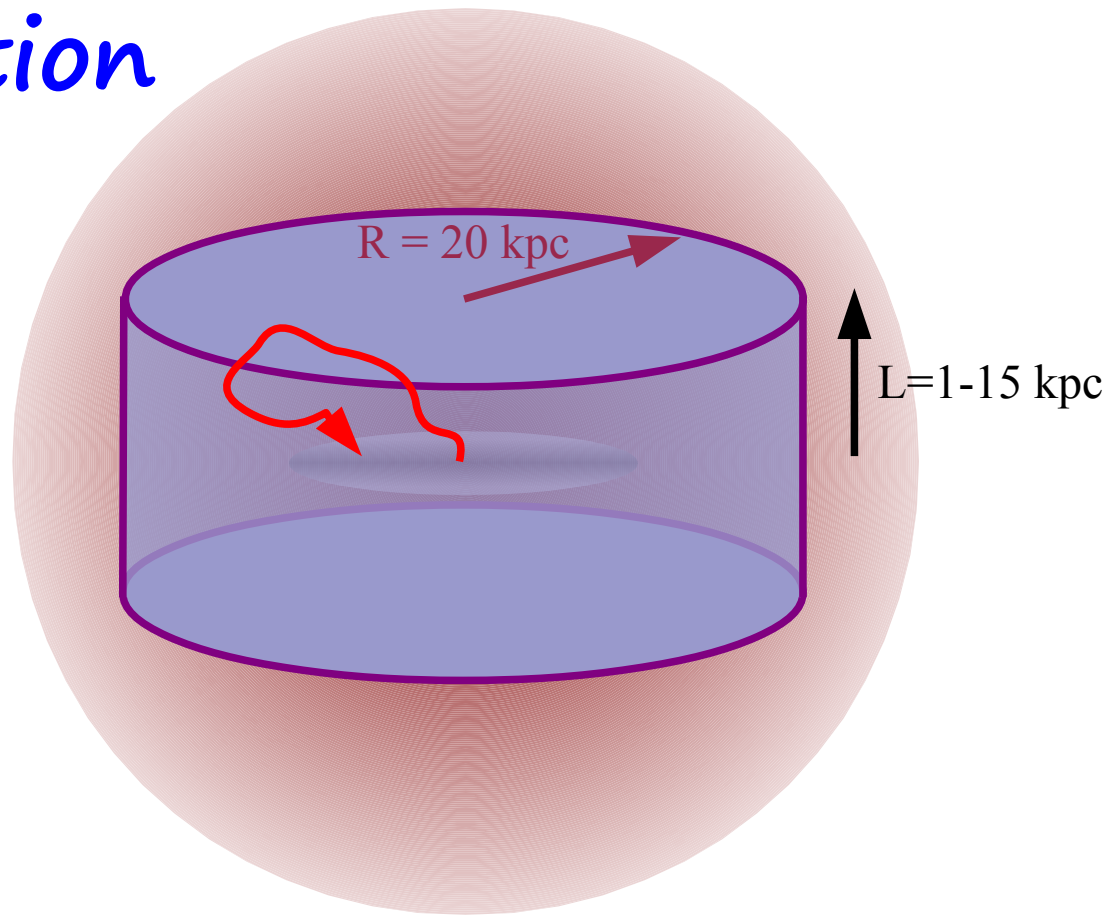
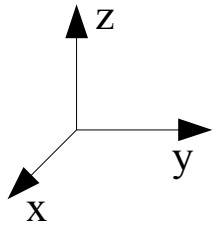


$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

Source term

$$Q(T, \vec{r}) = \begin{cases} \frac{1}{2} \frac{\rho^2(\vec{r})}{m_{\text{DM}}^2} \langle \sigma v \rangle \frac{dN}{dT} & \text{dark matter annihilation} \\ \frac{\rho(\vec{r})}{m_{\text{DM}}} \frac{1}{\tau_{\text{DM}}} \frac{dN}{dE} & \text{dark matter decay} \end{cases}$$

Propagation

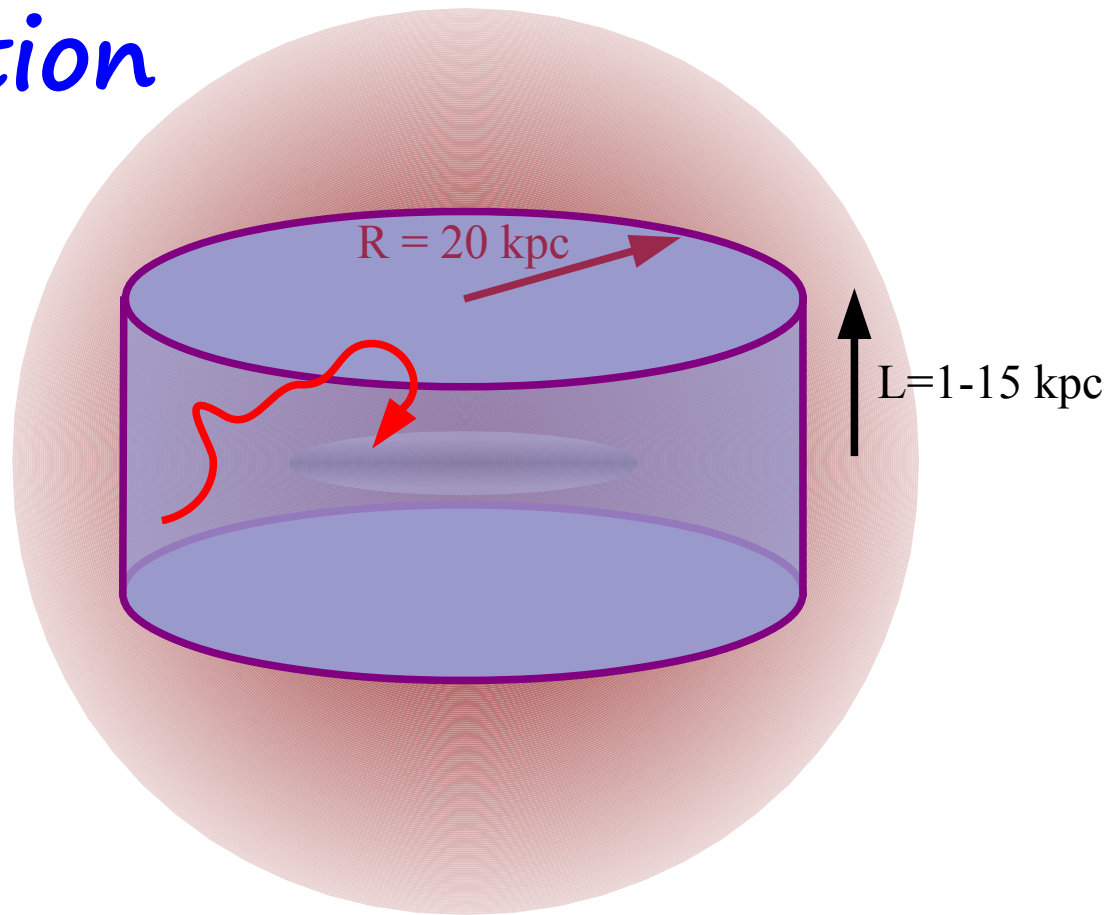
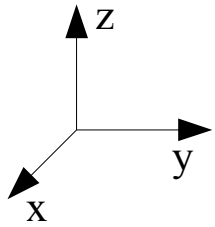


$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}} f + Q(T, \vec{r}) .$$

Source term

$$Q(T, \vec{r}) = \begin{cases} \frac{1}{2} \frac{\rho^2(\vec{r})}{m_{\text{DM}}^2} \langle \sigma v \rangle \frac{dN}{dT} & \text{dark matter annihilation} \\ \frac{\rho(\vec{r})}{m_{\text{DM}}} \frac{1}{\tau_{\text{DM}}} \frac{dN}{dE} & \text{dark matter decay} \end{cases}$$

Propagation

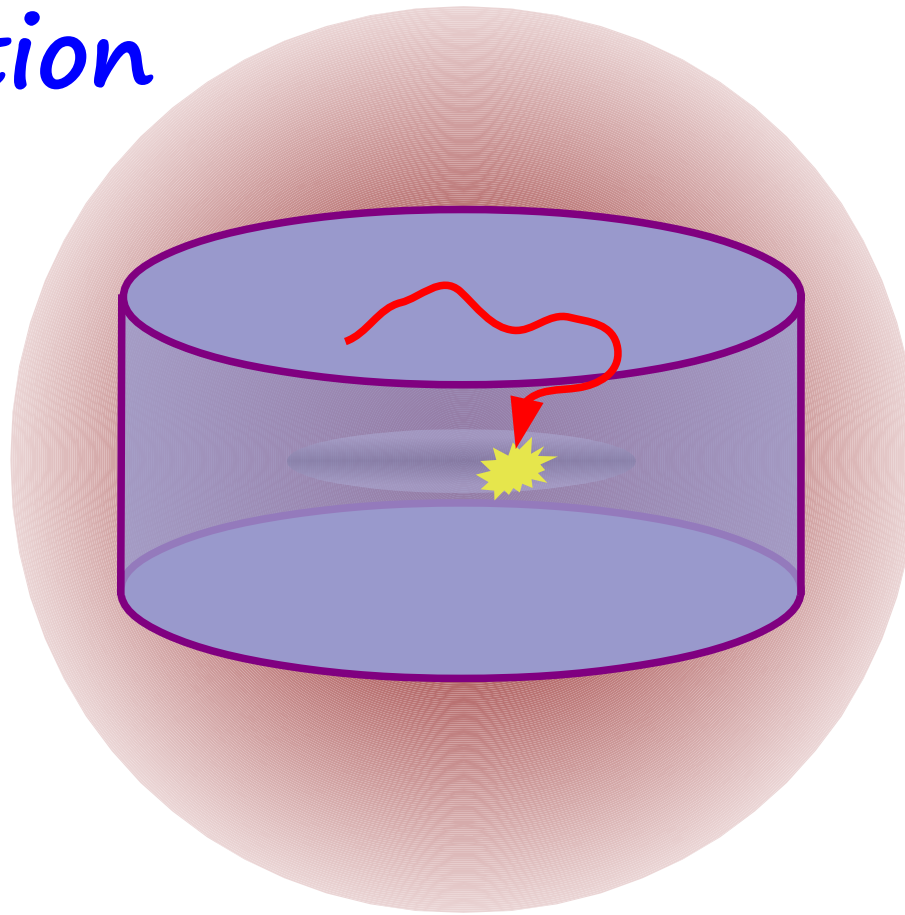
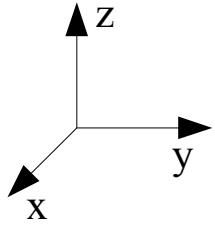


$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

Source term

$$Q(T, \vec{r}) = \begin{cases} \frac{1}{2} \frac{\rho^2(\vec{r})}{m_{\text{DM}}^2} \langle \sigma v \rangle \frac{dN}{dT} & \text{dark matter annihilation} \\ \frac{\rho(\vec{r})}{m_{\text{DM}}} \frac{1}{\tau_{\text{DM}}} \frac{dN}{dE} & \text{dark matter decay} \end{cases}$$

Propagation



$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f - Q(T, \vec{r}) .$$

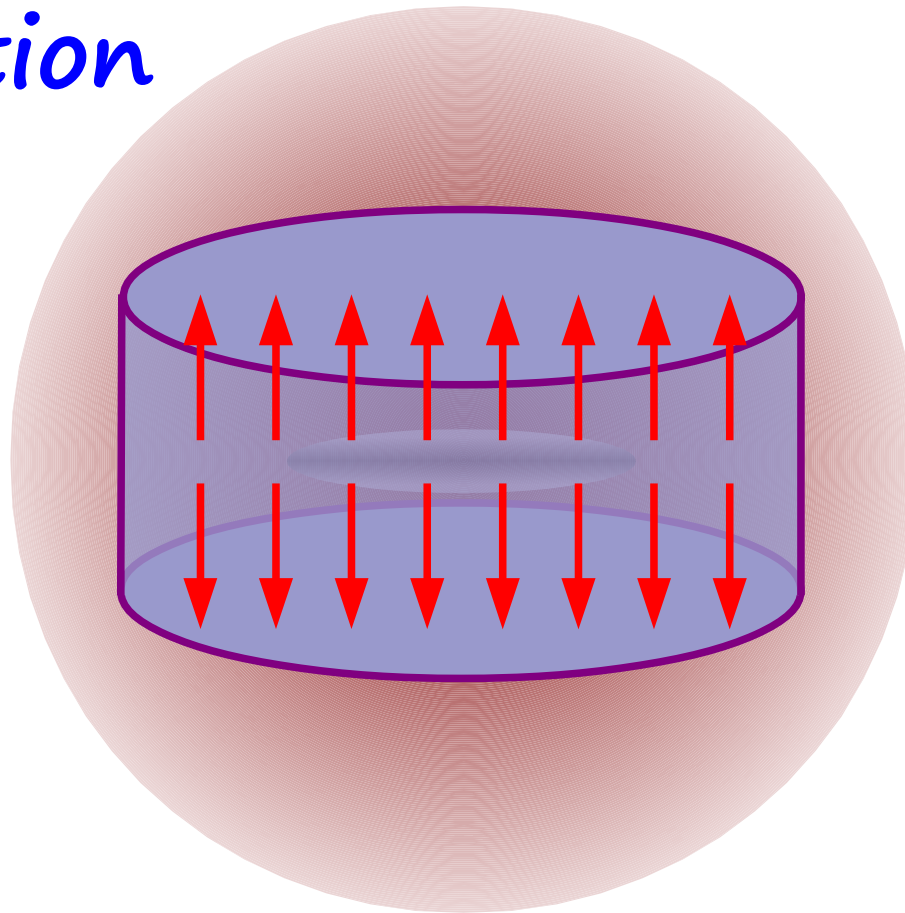
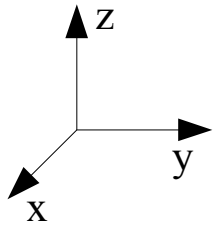
Annihilation term

Negligible for positrons.
For antiprotons,

$$\Gamma_{\text{ann}} = (n_{\text{H}} + 4^{2/3}n_{\text{He}})\sigma_{\bar{p}p}^{\text{ann}}v_{\bar{p}} .$$

$$\sigma_{\bar{p}p}^{\text{ann}}(T) = \begin{cases} 661 (1 + 0.0115 T^{-0.774} - 0.948 T^{0.0151}) \text{ mbarn} , & T < 15.5 \text{ GeV} , \\ 36 T^{-0.5} \text{ mbarn} , & T \geq 15.5 \text{ GeV} , \end{cases}$$

Propagation



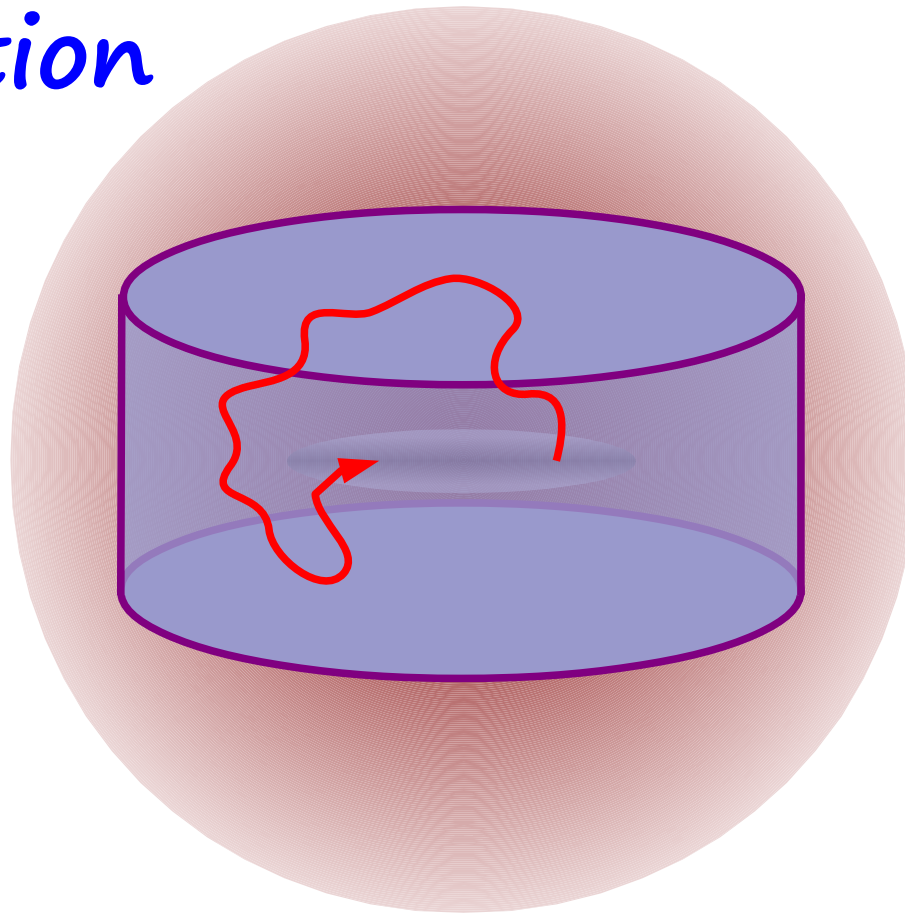
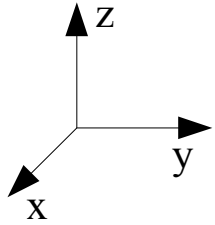
$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

Convection term

- Due to the Milky Way galactic wind.
- It drifts particles away from the Galactic disk.
- **Difficult to model.** Assume:

$$\vec{V}_c(\vec{r}) = V_e \text{sign}(z) \vec{k}$$

Propagation



$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] - \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

Energy loss term

- Due to inverse Compton scattering on the interstellar radiation field (starlight, thermal radiation of dust, CMB) and synchrotron radiation.
- Negligible for antiprotons and antideuterons
- Can be modelled

- Energy loss due to Inverse Compton scattering: $e^+\gamma \rightarrow e^+\gamma$

$$b_{\text{ICS}}(E_e, \vec{r}) = \int_0^\infty d\epsilon \int_\epsilon^{E_\gamma^{\text{max}}} dE_\gamma (E_\gamma - \epsilon) \frac{d\sigma^{\text{IC}}(E_e, \epsilon)}{dE_\gamma} f_{\text{ISRF}}(\epsilon, \vec{r})$$

Number density
of photons in ISRF

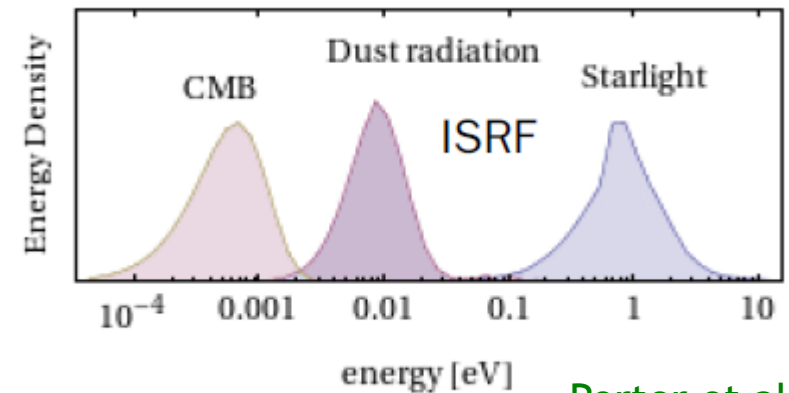
$$\frac{d\sigma^{\text{IC}}(E_e, \epsilon)}{dE_\gamma} = \frac{3}{4} \frac{\sigma_T}{\gamma_e^2 \epsilon} \times \left[2q \ln q + 1 + q - 2q^2 + \frac{1}{2} \frac{(q\Gamma)^2}{1 + q\Gamma} (1 - q) \right]$$

$\gamma_e = E_e/m_e \rightarrow$ Lorentz factor.

$\Gamma_e = 4 \gamma_e \epsilon/m_e$

$q = E_\gamma/\Gamma(E_e - E_\gamma)$

$\sigma_T = 0.67$ barn \rightarrow Compton scattering cross section
in the Thomson limit.



Porter et al.

- Energy loss due to synchrotron radiation:

$$b_{\text{sync}}(E_e, \vec{r}) = \frac{4}{3} \sigma_T \gamma_e^2 \frac{B^2}{2}$$

$$B = 6 \mu\text{G} \exp(-|z|/5\text{kpc} - r/20\text{kpc})$$

Approximately $b(E) = \frac{E^2}{E_0 \tau_E}$, with $E_0 = 1$ GeV and $\tau_E = 10^{16}$ s

- Energy loss due to Inverse Compton scattering: $e^+\gamma \rightarrow e^+\gamma$

$$b_{\text{ICS}}(E_e, \vec{r}) = \int_0^\infty d\epsilon \int_\epsilon^{E_\gamma^{\text{max}}} dE_\gamma (E_\gamma - \epsilon) \frac{d\sigma^{\text{IC}}(E_e, \epsilon)}{dE_\gamma} f_{\text{ISRF}}(\epsilon, \vec{r})$$

Number density of photons in ISRF

$$\frac{d\sigma^{\text{IC}}(E_e, \epsilon)}{dE_\gamma} = \frac{3}{4} \frac{\sigma_T}{\gamma_e^2 \epsilon} \times \left[2q \ln q + 1 + q - 2q^2 + \frac{1}{2} \frac{(q\Gamma)^2}{q\Gamma(1-q)} \right]$$

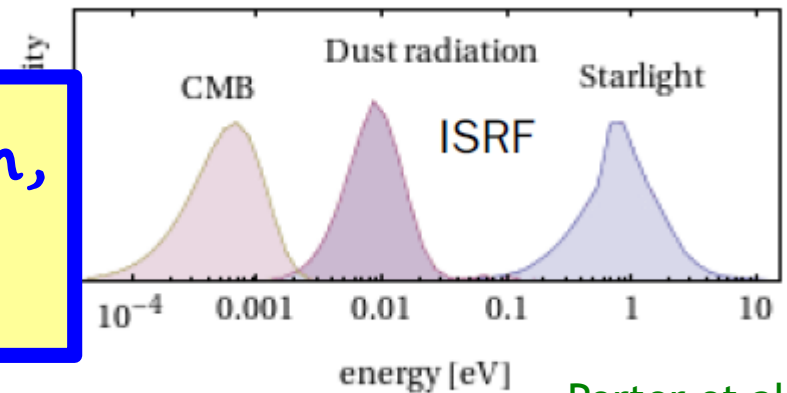
$\gamma_e = E_e/m_e \rightarrow$ Lorentz factor

$\Gamma_e = 4 \gamma_e \epsilon/m_e$

$q = E_\gamma/\Gamma(E_e - E_\gamma)$

$\sigma_T = 0.67$ barn \rightarrow Compton scattering cross section in the Thomson limit.

Not very well known, though...



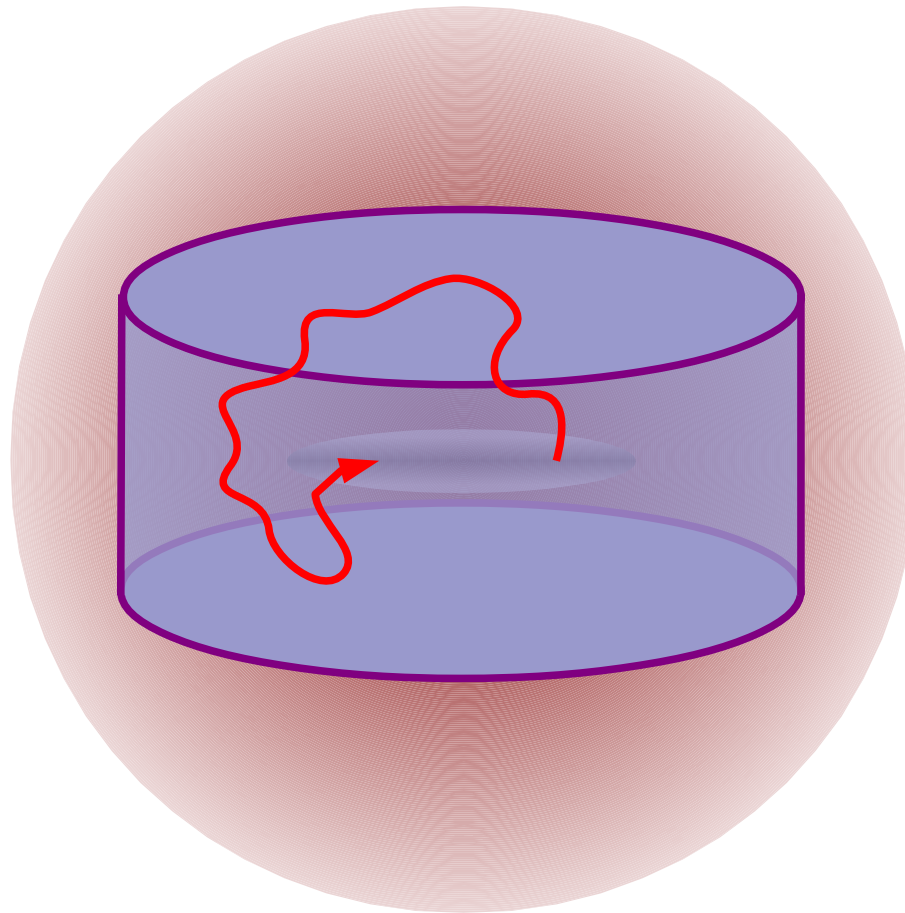
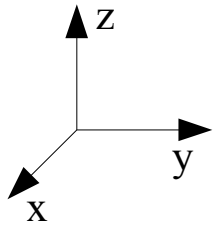
Porter et al.

- Energy loss due to synchrotron radiation:

$$b_{\text{sync}}(E_e, \vec{r}) = \frac{4}{3} \sigma_T \gamma_e^2 \frac{B^2}{2}$$

$$B = 6 \mu\text{G} \exp(-|z|/5\text{kpc} - r/20\text{kpc})$$

Approximately $b(E) = \frac{E^2}{E_0 \tau_E}$, with $E_0 = 1$ GeV and $\tau_E = 10^{16}$ s



$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] - \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}} f + Q(T, \vec{r}) .$$

Diffusion term

- Due to the tangled magnetic field of the Galaxy.
- **Difficult to model.** Assume

$$K(T) = K_0 \beta \mathcal{R}^\delta$$

$$\left(\begin{array}{l} \beta = \text{velocity} \\ \mathcal{R} = \text{rigidity} \end{array} \right)$$

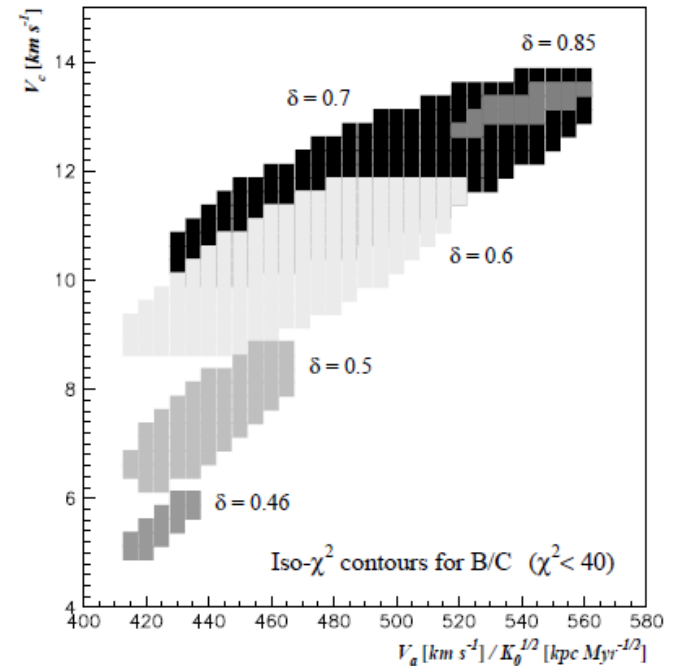
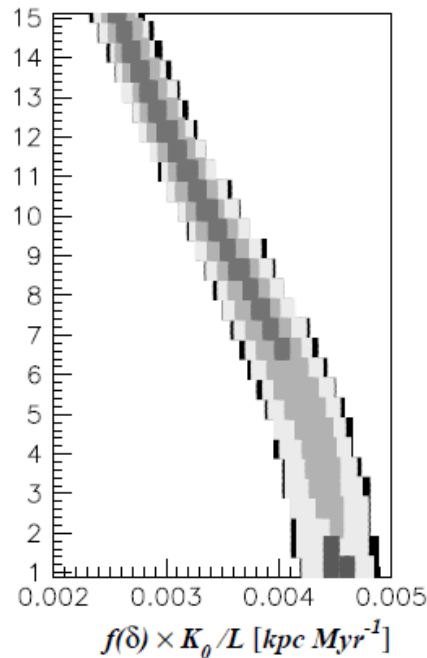
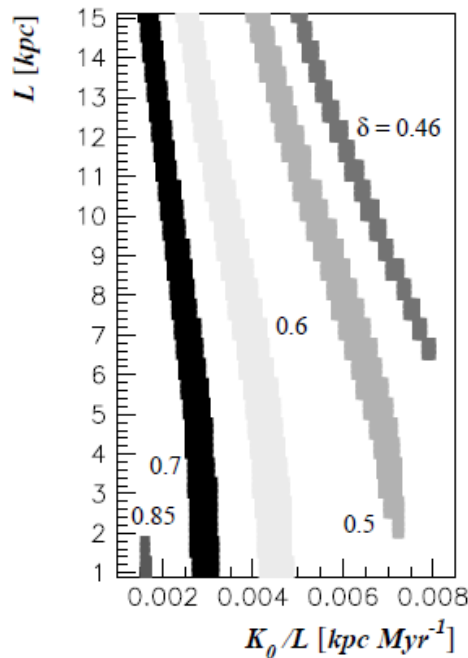
$$0 = \frac{\partial f}{\partial t} - \nabla \cdot [K(T, \vec{r}) \nabla f] - \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}} f + Q(T, \vec{r}) .$$

$$K(T) = K_0 \beta \mathcal{R}^\delta$$

$$\vec{V}_c(\vec{r}) = V_c \text{sign}(z) \vec{k}$$

K_0 , δ , V_c (as well as L) must be determined with measurements of other cosmic ray species (mainly B/C ratio).

Iso- χ^2 contours for B/C ($\chi^2 < 40$)



Model	δ	K_0 (kpc ² /Myr)	L (kpc)	V_c (km/s)
MIN	0.85	0.0016	1	13.5
MED	0.70	0.0112	4	12
MAX	0.46	0.0765	15	5

Maurin, Donato, Taillet, Salati '01



M80

M14

M92

9

8
Draconis
NGC 5129

SCUTUM ARM

ORION ARM

NORMA ARM

CRUX ARM

CARINA ARM

M71

M4

Alphei
NGC 4755

Carina
NGC 3292

Keyhole NGC 3324

WE ARE HERE

Eagle M16

13

George M77

Lagoon M8

M10

Wild Duck M10

M2

Antares M7

M12

Ring M53

Trifid M20

Bumblebee M27

NGC 7293

11

NGC 7027

North America NGC 7000

Orion M42

NGC 2269

Rosette NGC 2237

O R I O N

10

6,000 light years

12

S A G I T T A R I U S

S

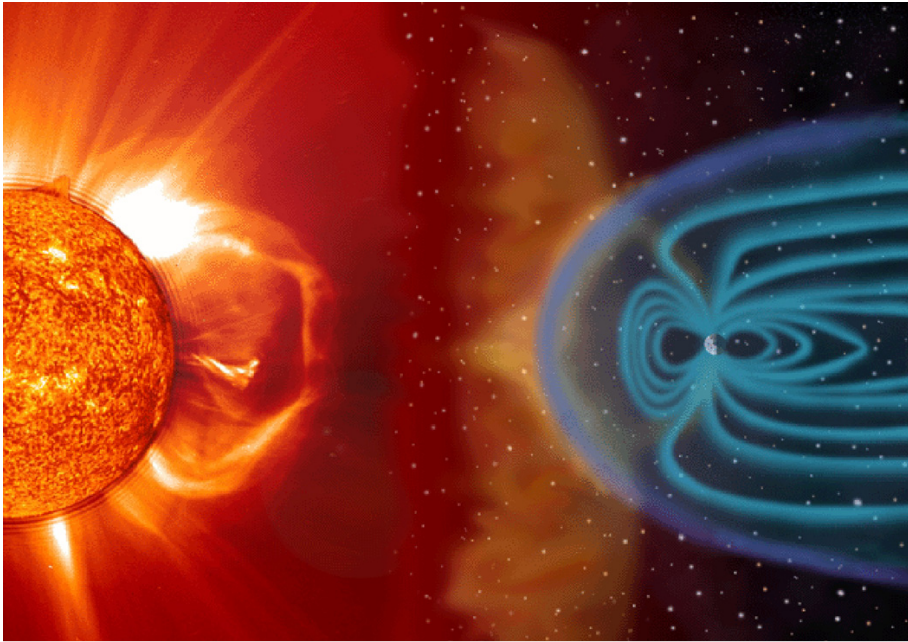
M

A

R

M

Propagation *inside* the Solar System



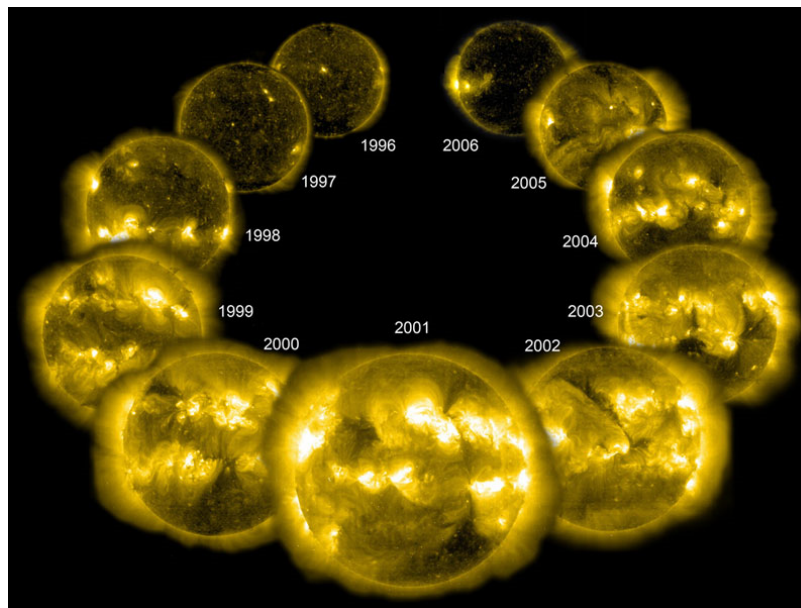
In the “force field approximation”, the flux at the top of the atmosphere (TOA) is related to the interstellar flux (IS) by

$$\Phi_{e^\pm}^{\text{TOA}}(E_{\text{TOA}}) = \frac{E_{\text{TOA}}^2}{E_{\text{IS}}^2} \Phi_{e^\pm}^{\text{IS}}(E_{\text{IS}})$$

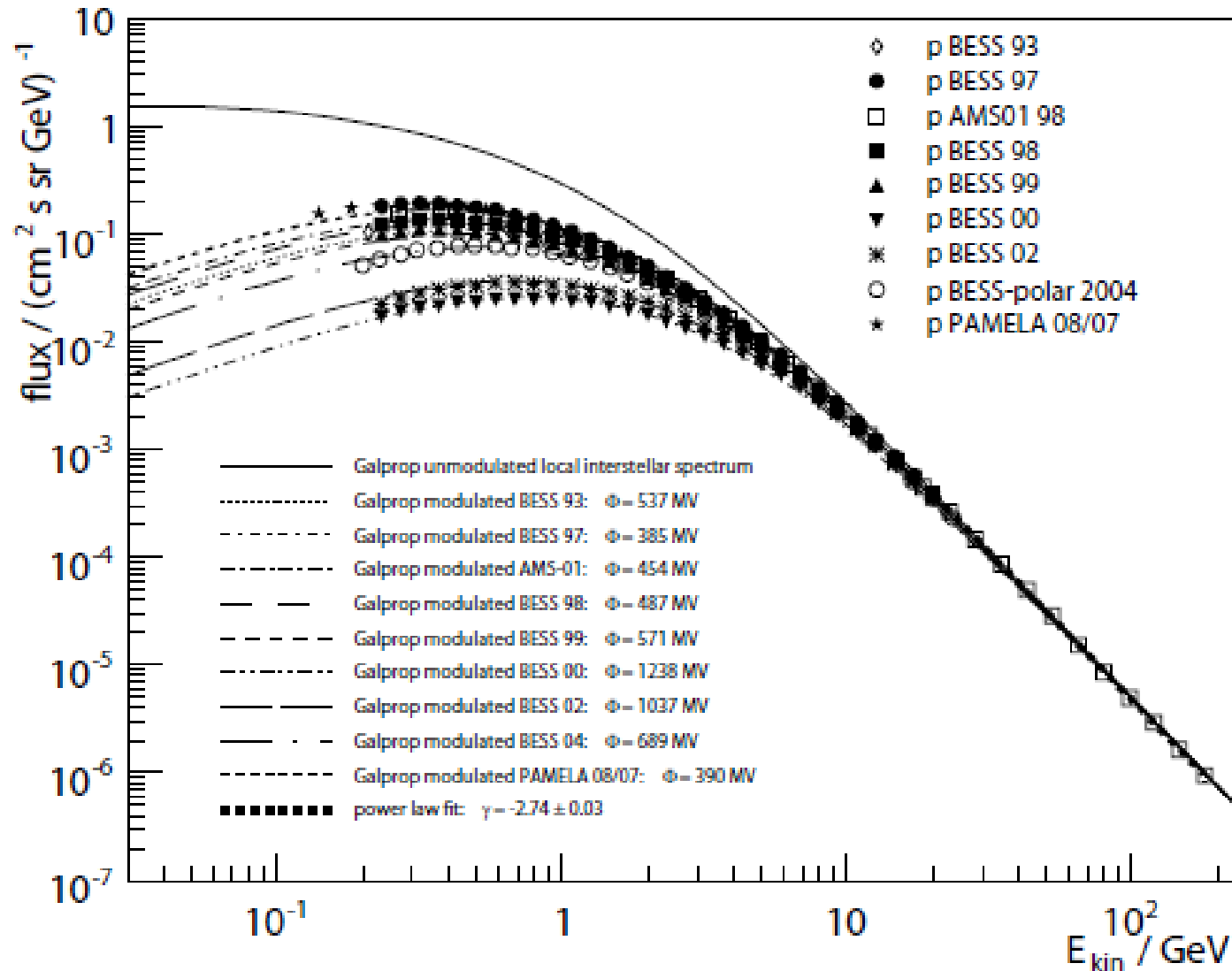
$$E_{\text{IS}} = E_{\text{TOA}} + \phi_F$$

↓
solar modulation parameter

$$\phi_F = 500 \text{ MV} - 1.3 \text{ GV}$$



Cosmic ray **proton** spectrum as measured by BESS, AMS-01 and PAMELA



Gast, Schael '09