



# Superconductivity and quantum phase transition at oxides interface

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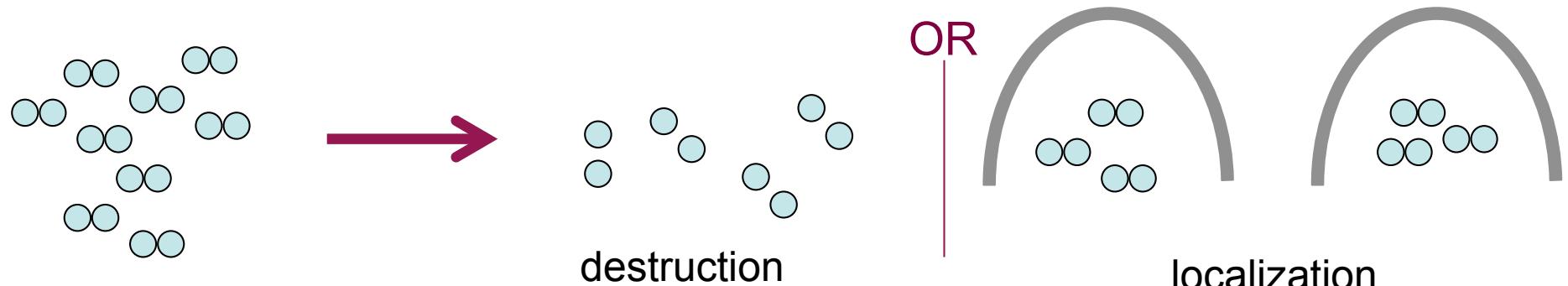
# Quantum Phase Transition in a 2D superconductor

## ■ Superconductor-Insulator Quantum Phase Transition (QPT) in 2D

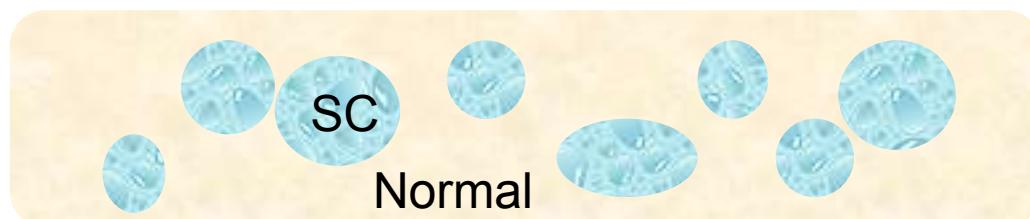
→ Many experimental results

→ Open questions

- Nature of the "Insulating State" : fermionic vs bosonic scenario



- Intrinsic mesoscopic disorder at the transition : homogeneous media ?



Mason et al, PRB 2001

Kapitulnik et al PRB 2001

Feigel'man et al, PRL 2001

Spivak et al, PRB 2008

Multiple Quantum Criticalities ?

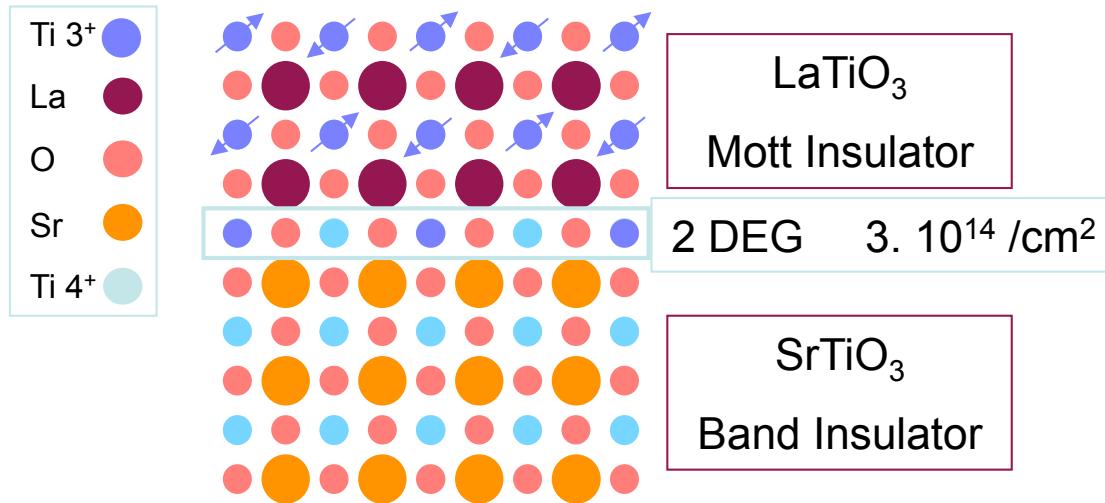
- Towards a 1D filamentary structure at the transition ?

Ioffe-Mezard PRL 2010, Goetz-Benfatto-Castellani PRL 2012

← cf Castellani's & Grilli's talks

# 2 DEG at oxides interfaces

## $\text{LaTiO}_3 / \text{SrTiO}_3$

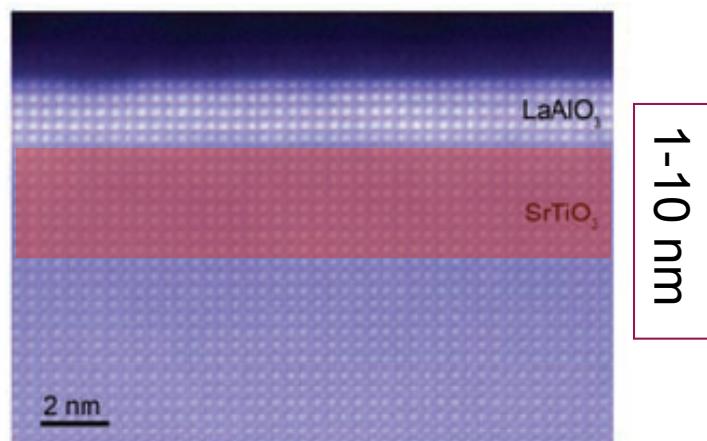


Ohtomo et al Nature 2002

Hwang et al Physica E 2004

2 DEG at the interface

## $\text{LaAlO}_3 / \text{SrTiO}_3$



- High mobility ( $10^4 \text{ cm}^2/\text{Vs}$ )
- Tunable electron density
- Superconductivity

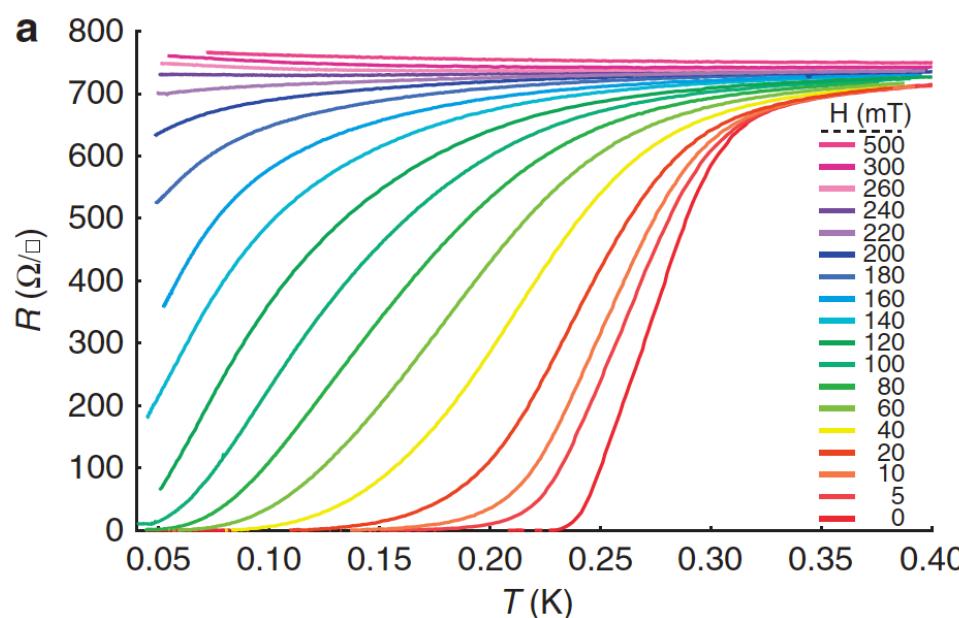


2D SC Physics

Mannhart et al, Science 2010

# An extreme two-dimensional superconductor

## ■ Size smaller than the coherence length



$$H_{\perp}(T) = \frac{\phi_0}{2\pi\xi_{\parallel}(T)}$$

$$H_{\parallel}(T) = \frac{\sqrt{3}\phi_0}{\pi d \xi_{\parallel}(T)}$$

→  $\xi_0 \sim 36 \text{ nm}$

$\xi_0 > d$

2D SC

→  $d < 10 \text{ nm}$

## ■ A true two-dimensional system

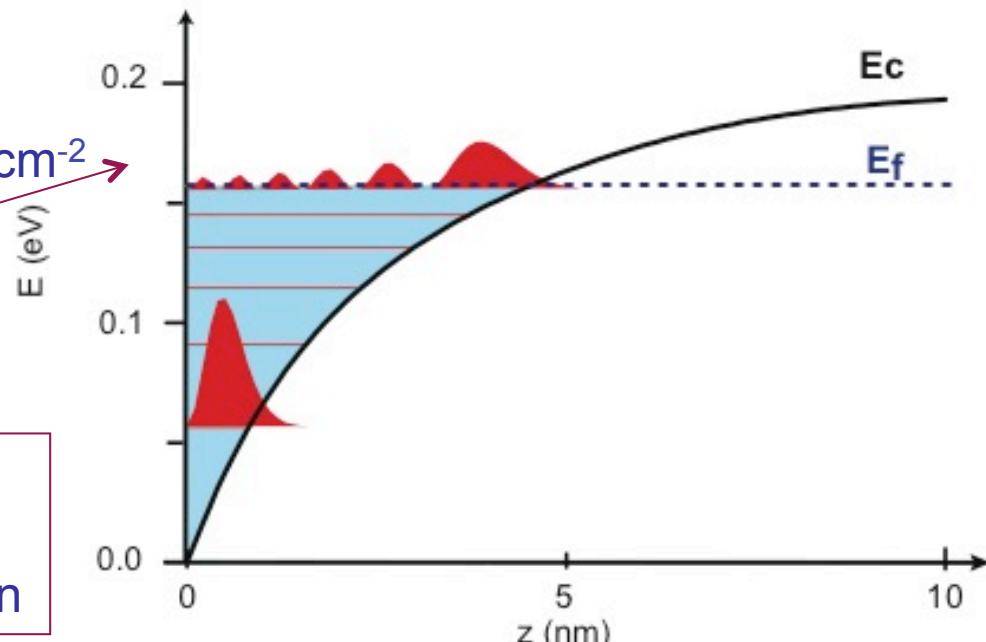
Low electron density  $n = 10^{13}-10^{14} \text{ cm}^{-2}$

$$\lambda_F = \sqrt{\frac{2\pi}{n}}$$

$$E_F = \frac{\hbar^2 \pi n}{m_e}$$

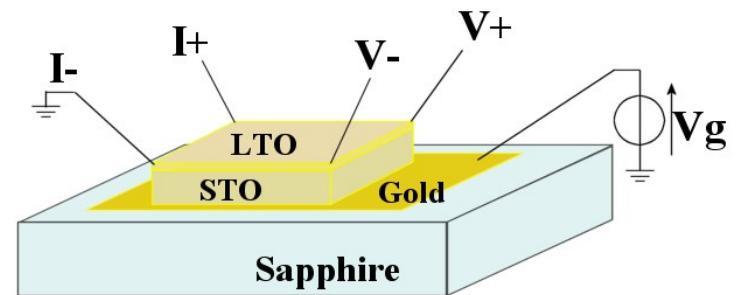
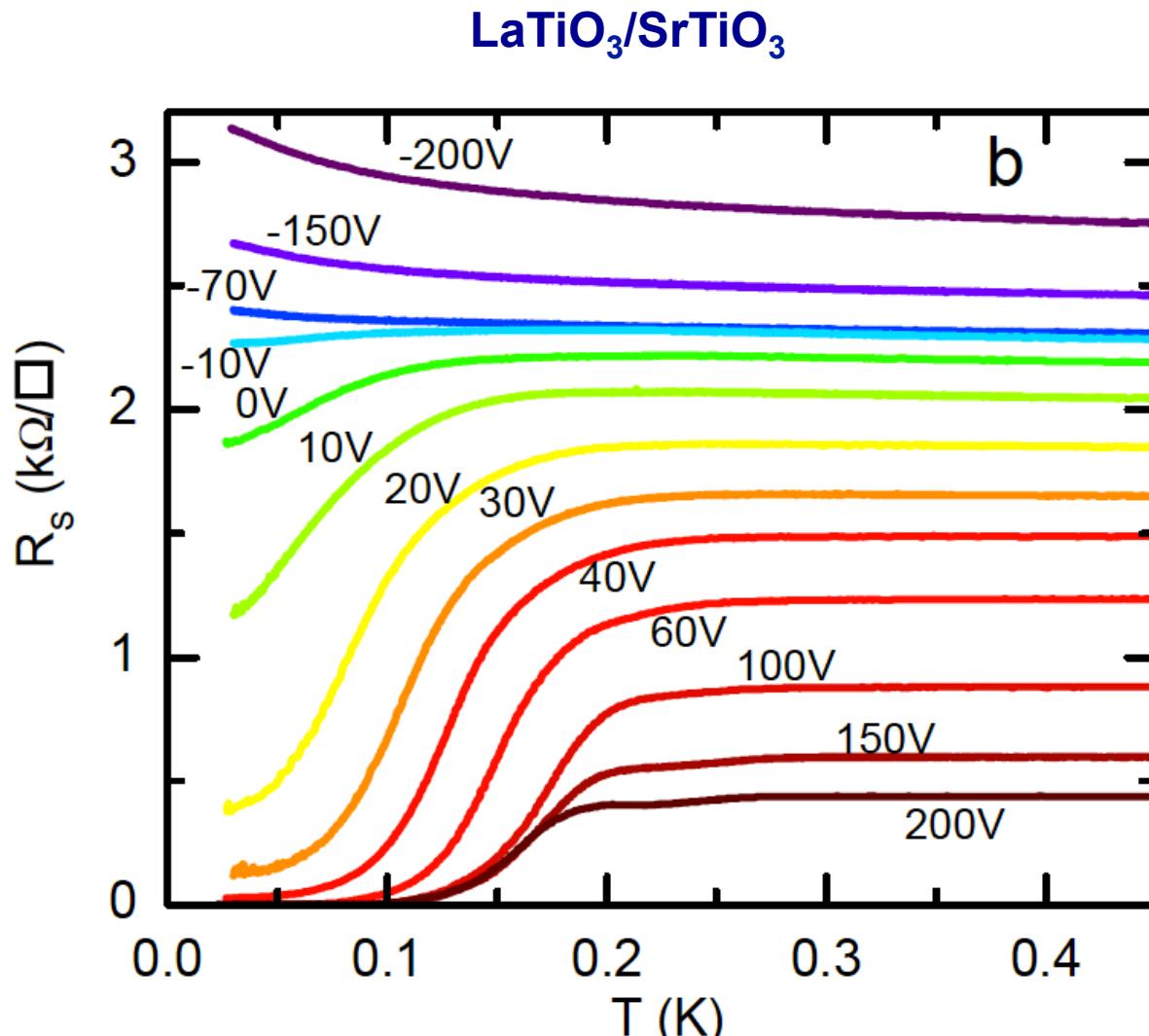
-  $\lambda_F >$  well size  $d$

- quantized levels in the z-direction



# Electric field effect

- Control of the 2-DEG by electrostatic back gate



remove e-

$R_s$  decrease with  $V_g$

$T_c$  goes through  
a maximum

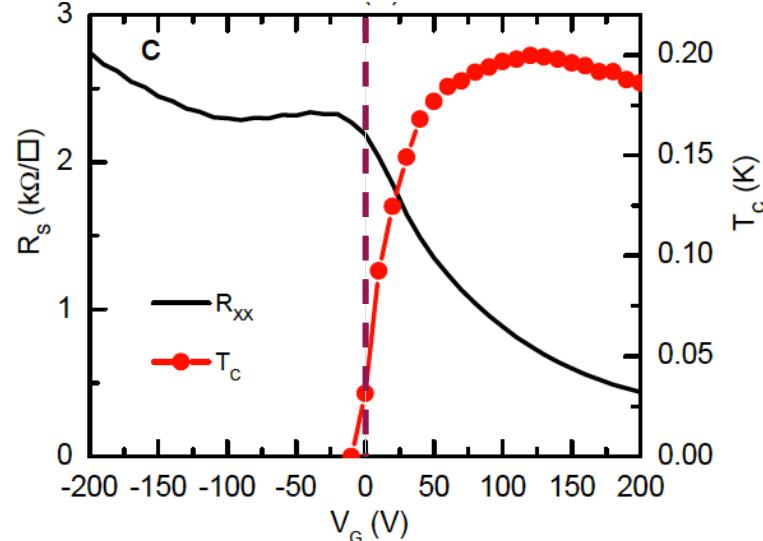
add e-

Biscaras et al, PRL 108, 247004 (2012)

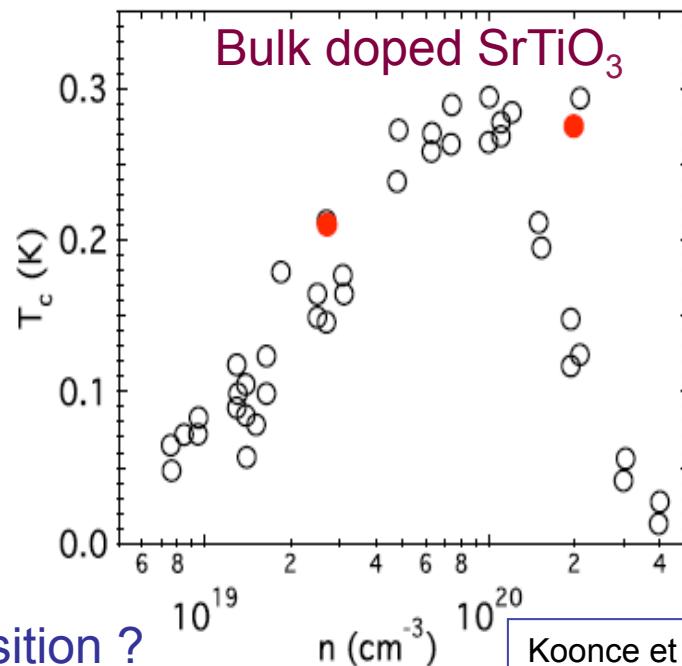
- Superconductor-insulator transition induced by field effect

# Two main questions

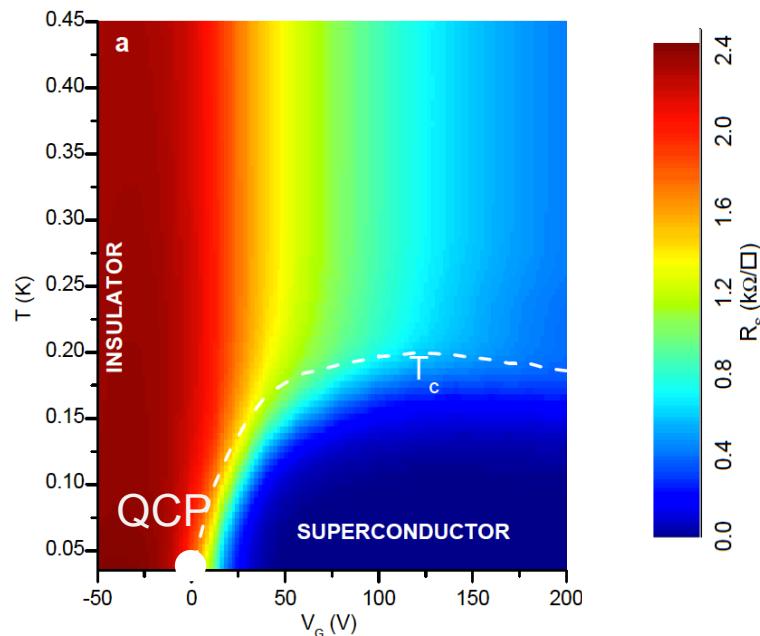
- What is the origin of superconductivity ?



- What type of Superconductor-Insulator transition ?



Koncone et al, PRB (1967)

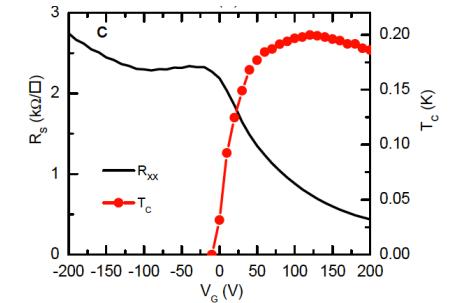


- ➔ Quantum phase transition ?
- ➔ Disorder driven phase transition ?
- ➔ Universality class ?

# Outline

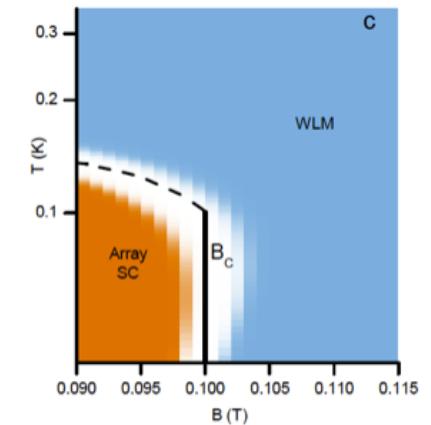
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Tunable superconductivity in oxide 2DEG



Quantum phase transition under magnetic field

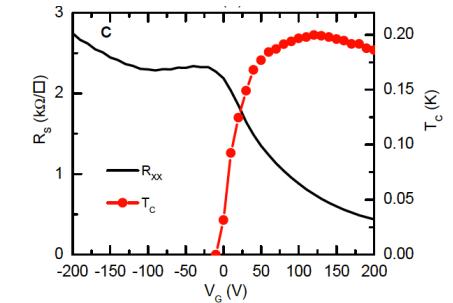
Conclusions



# Outline

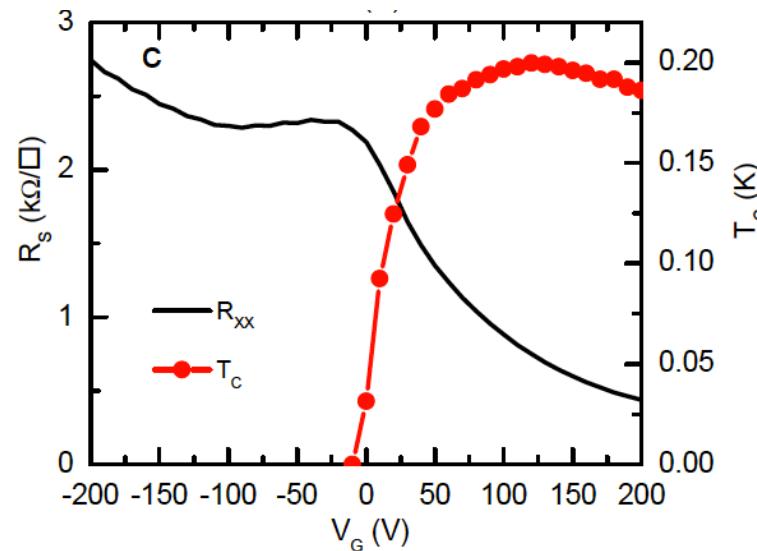
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## Tunable superconductivity in oxide 2DEG

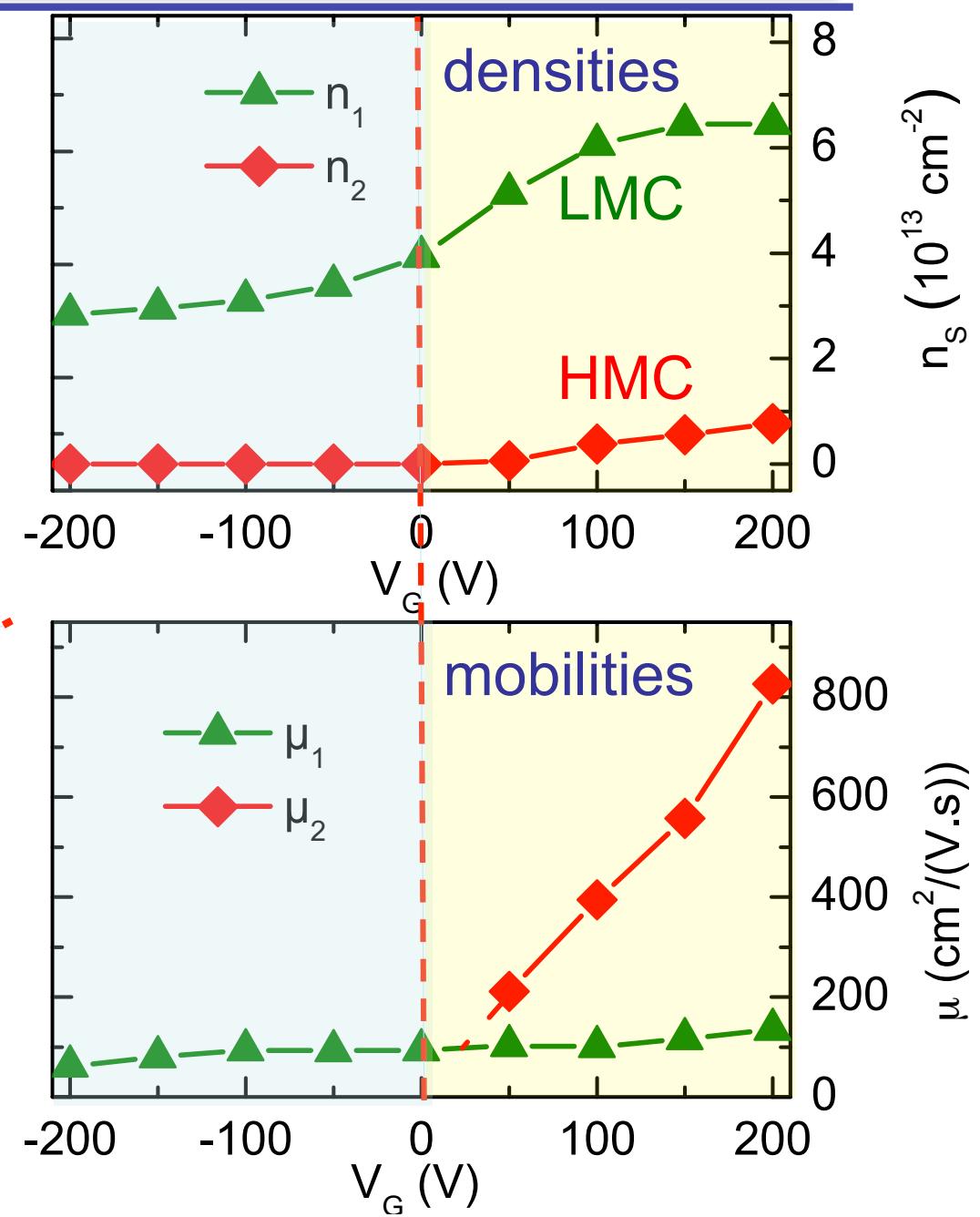
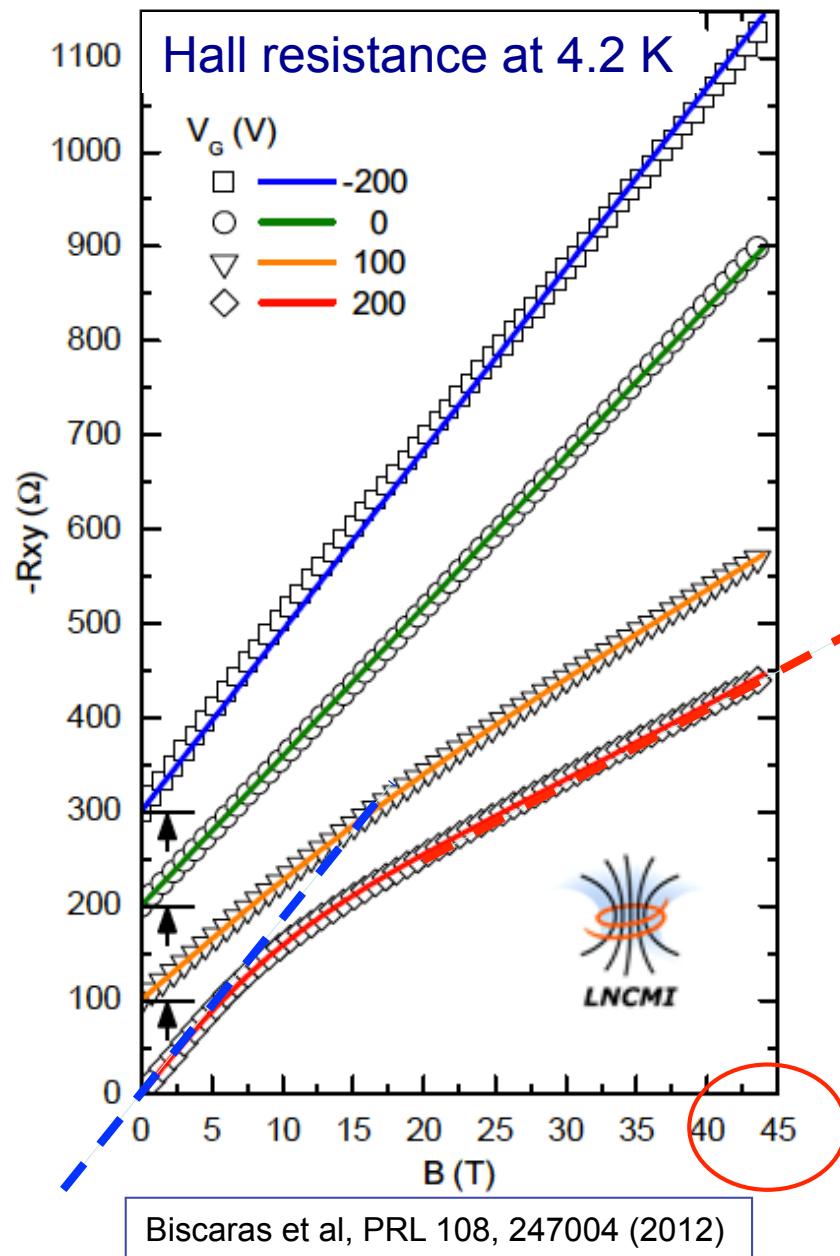


Quantum phase transition under magnetic field

## Conclusions

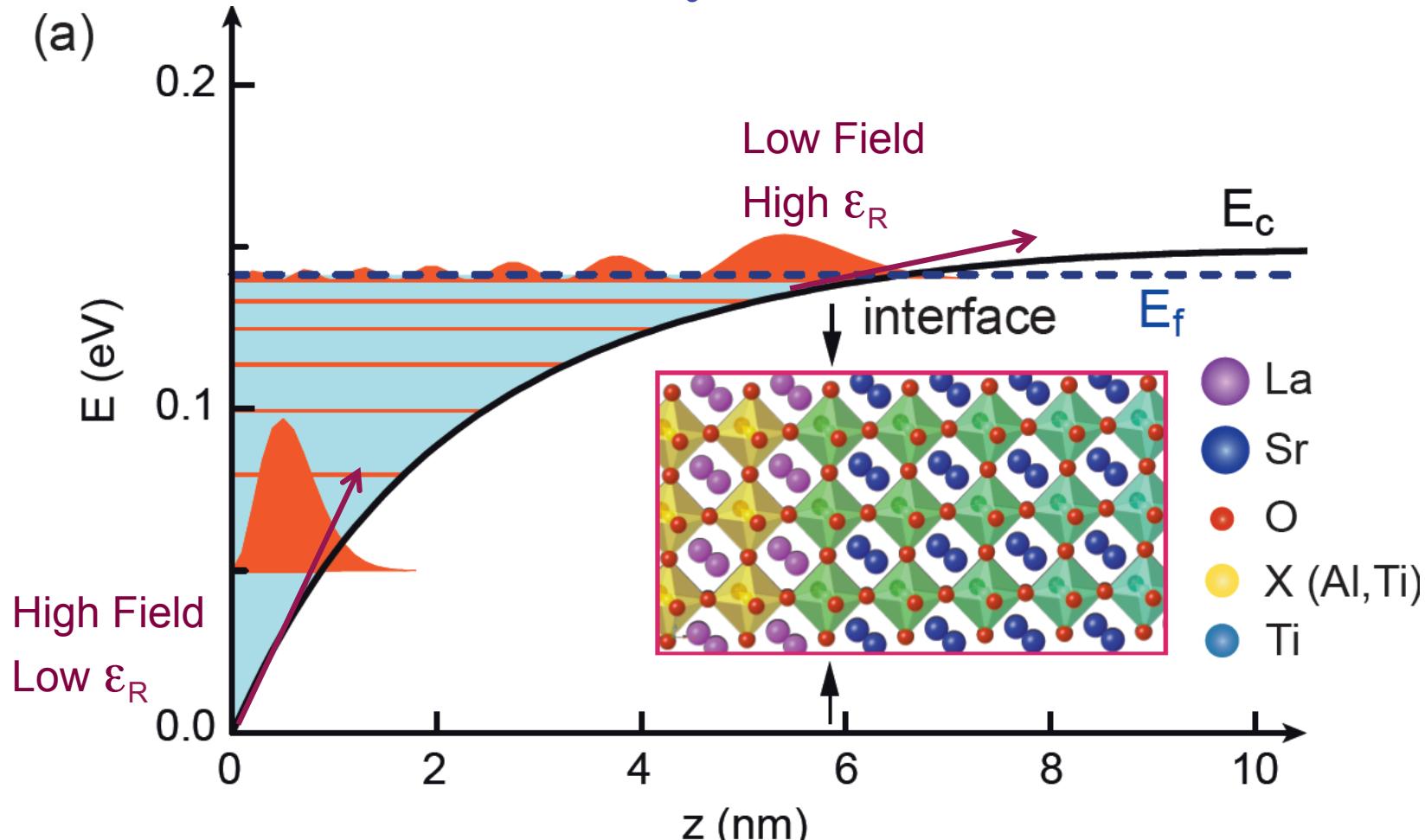


# Two types of carriers for positive gate voltage



# Self consistent calculation : band bending model

- Confinement at the interface : SrTiO<sub>3</sub> conduction band bending



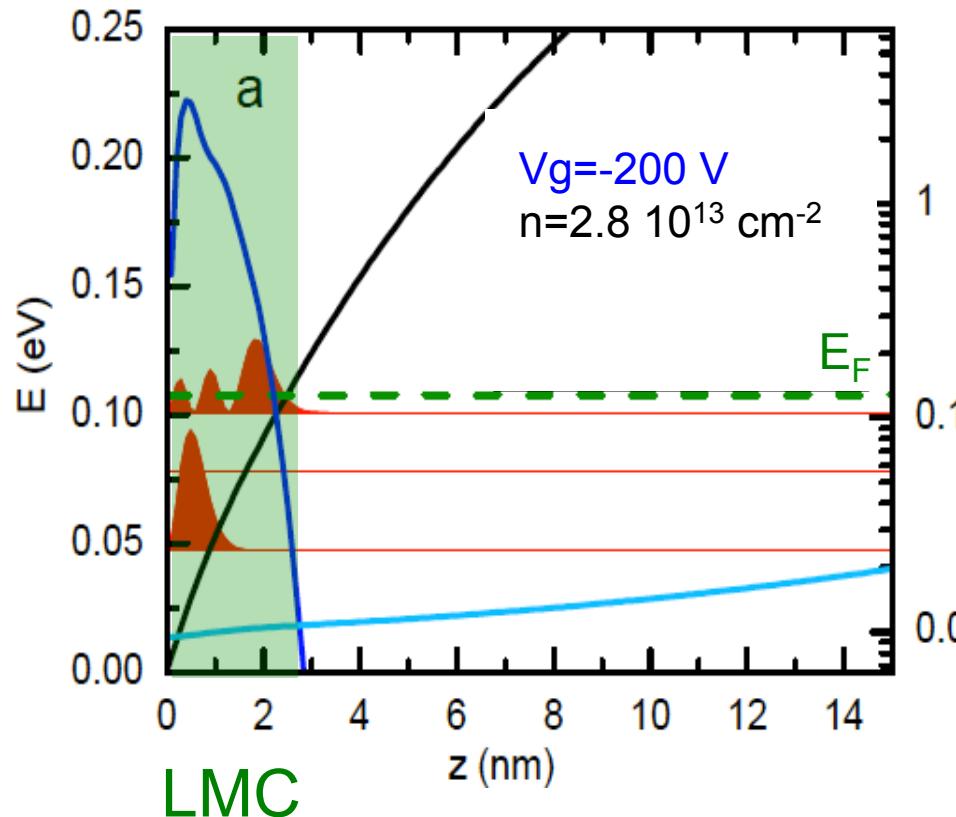
→ Schrödinger Equation

→ Poisson Equation

→ Non linear SrTiO<sub>3</sub> Epsilon

$$\epsilon_R(F) = \epsilon_R(\infty) + \frac{1}{(A+BF)}$$

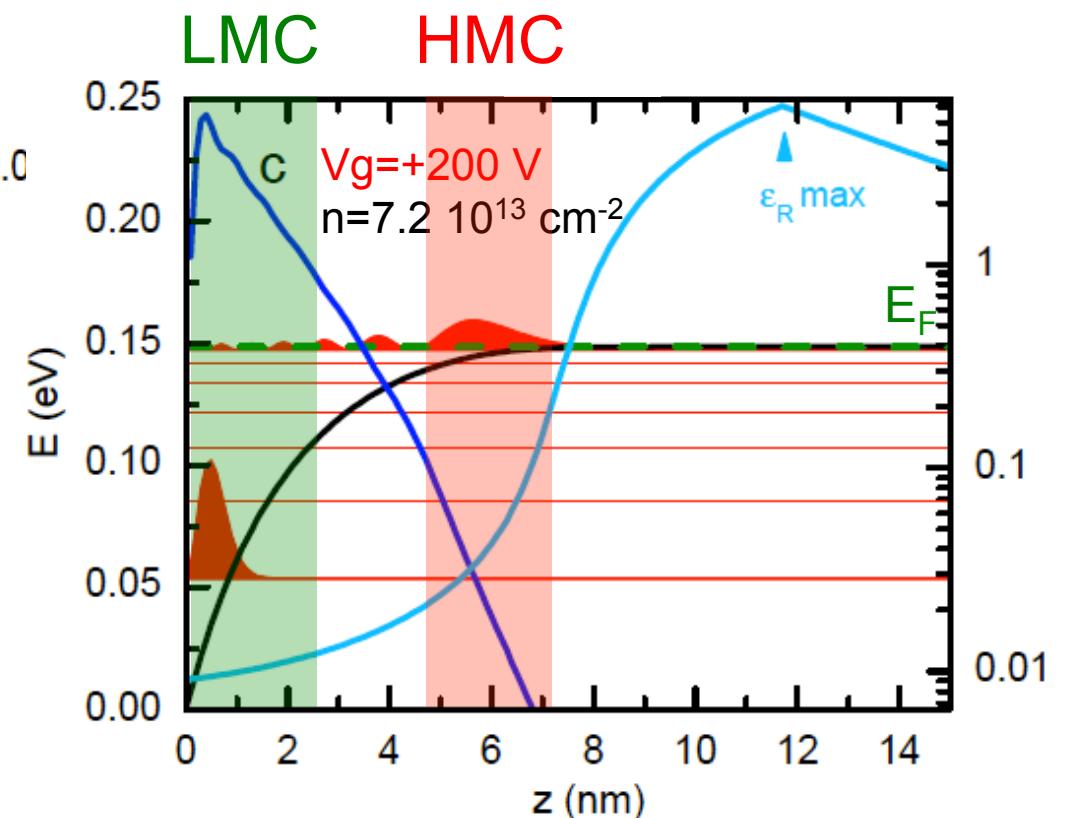
# Solutions of the model ( $V_G \neq 0$ )



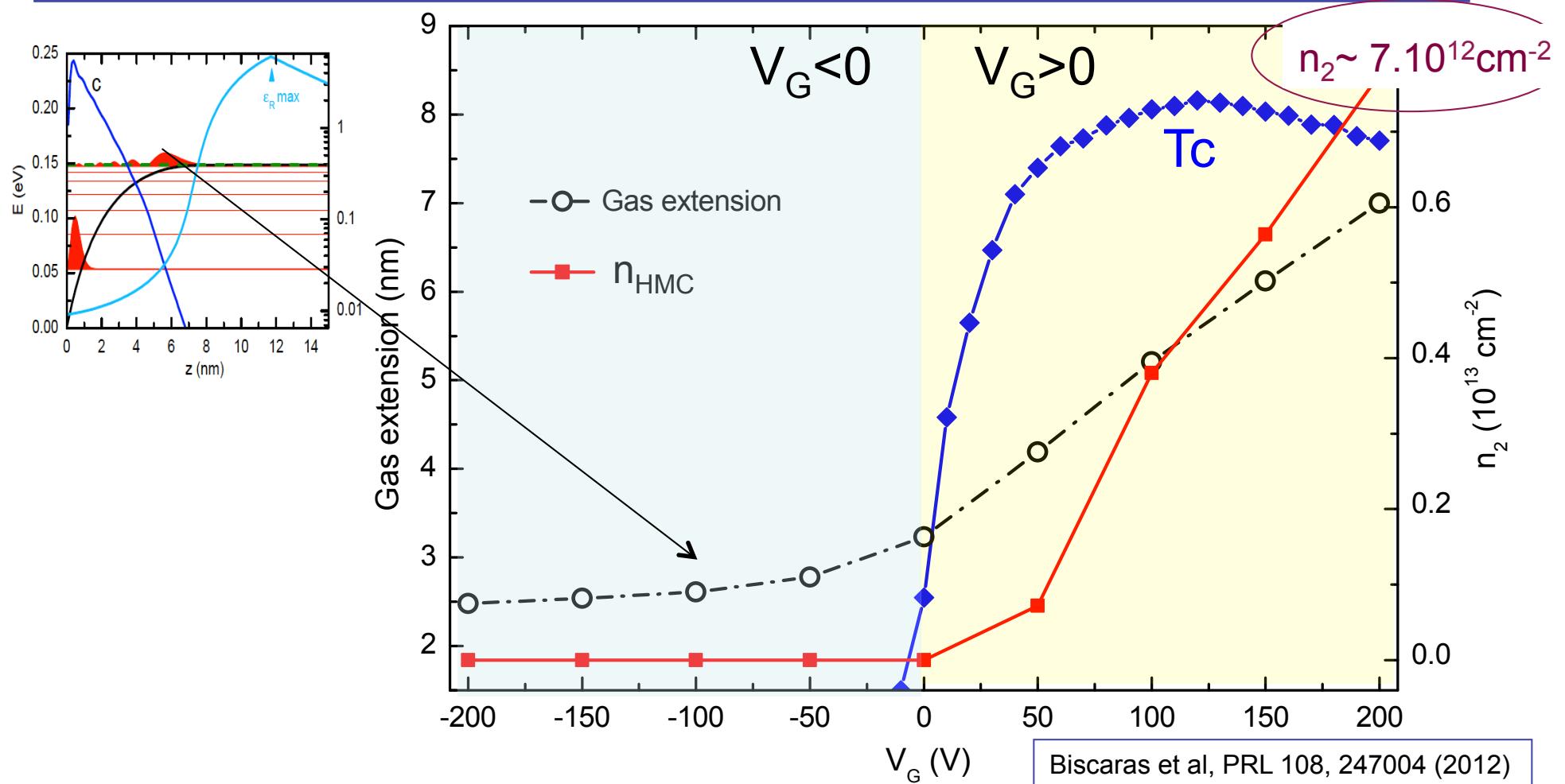
■  $V_g > 0$  carriers extends in SrTiO<sub>3</sub>

► Mobility increases

■  $V_g < 0$  carriers are confined at the interface  
► Low mobility



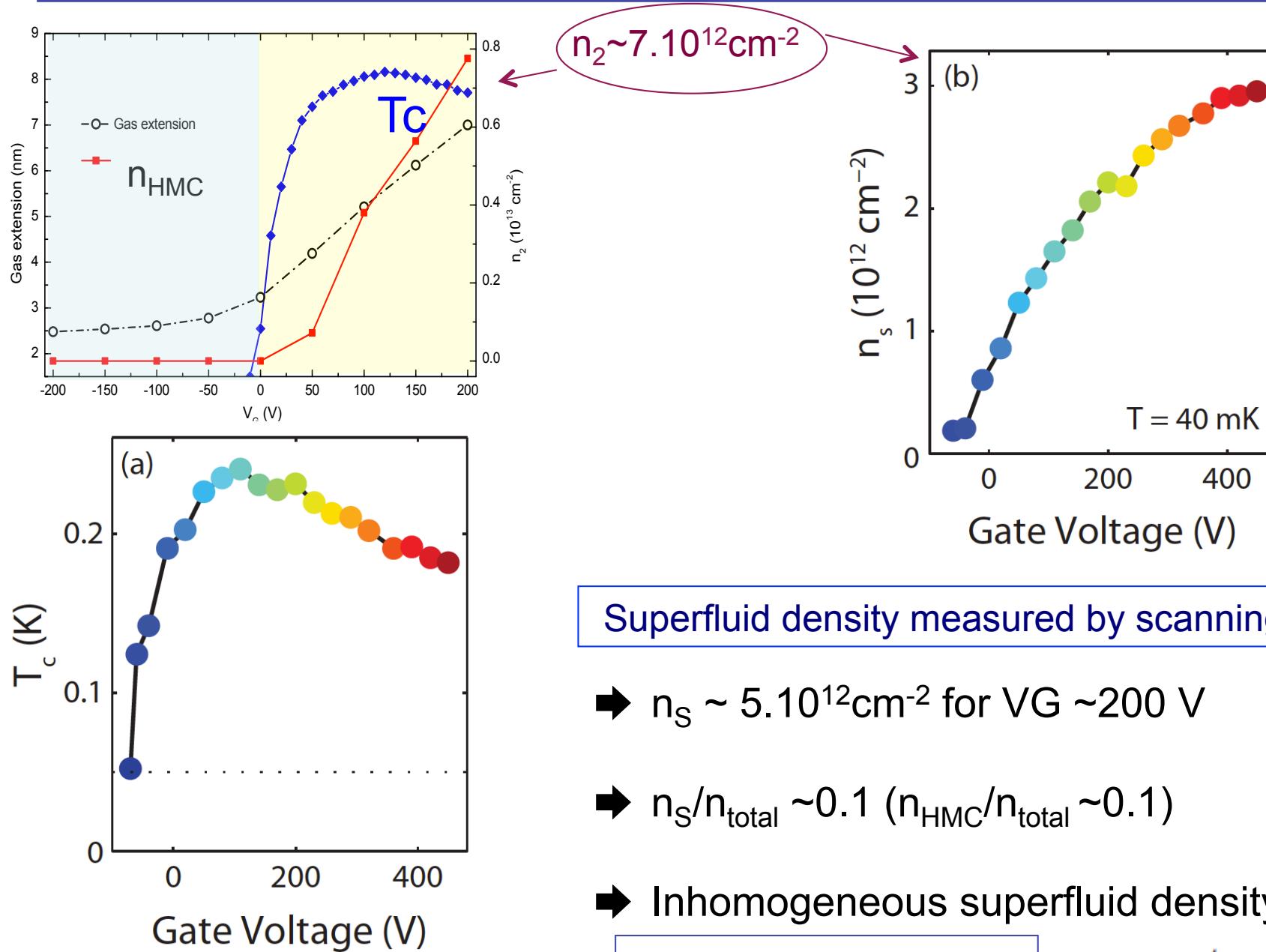
# High Mobility Carriers at the edge : Superconductivity



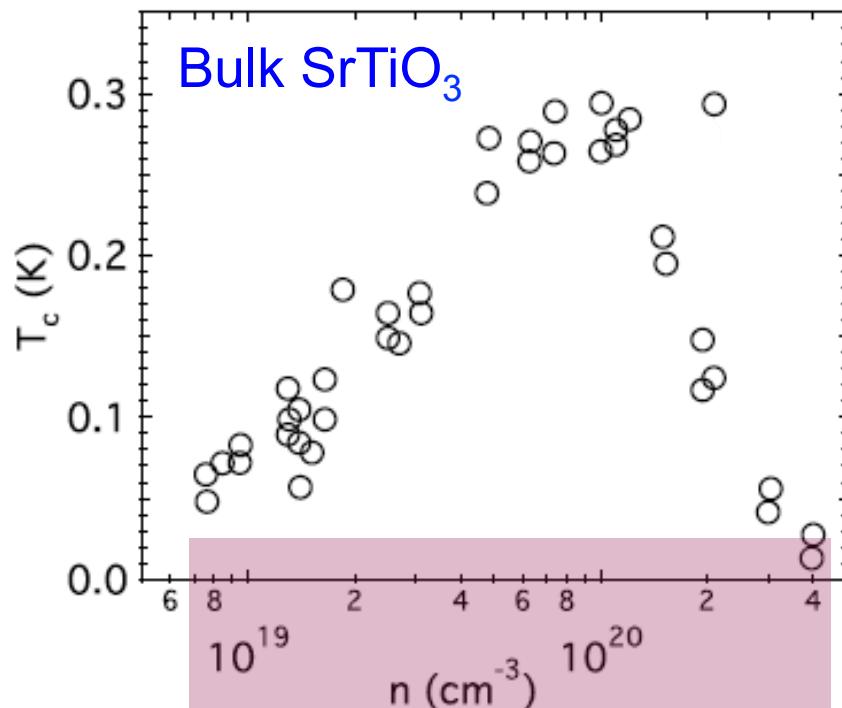
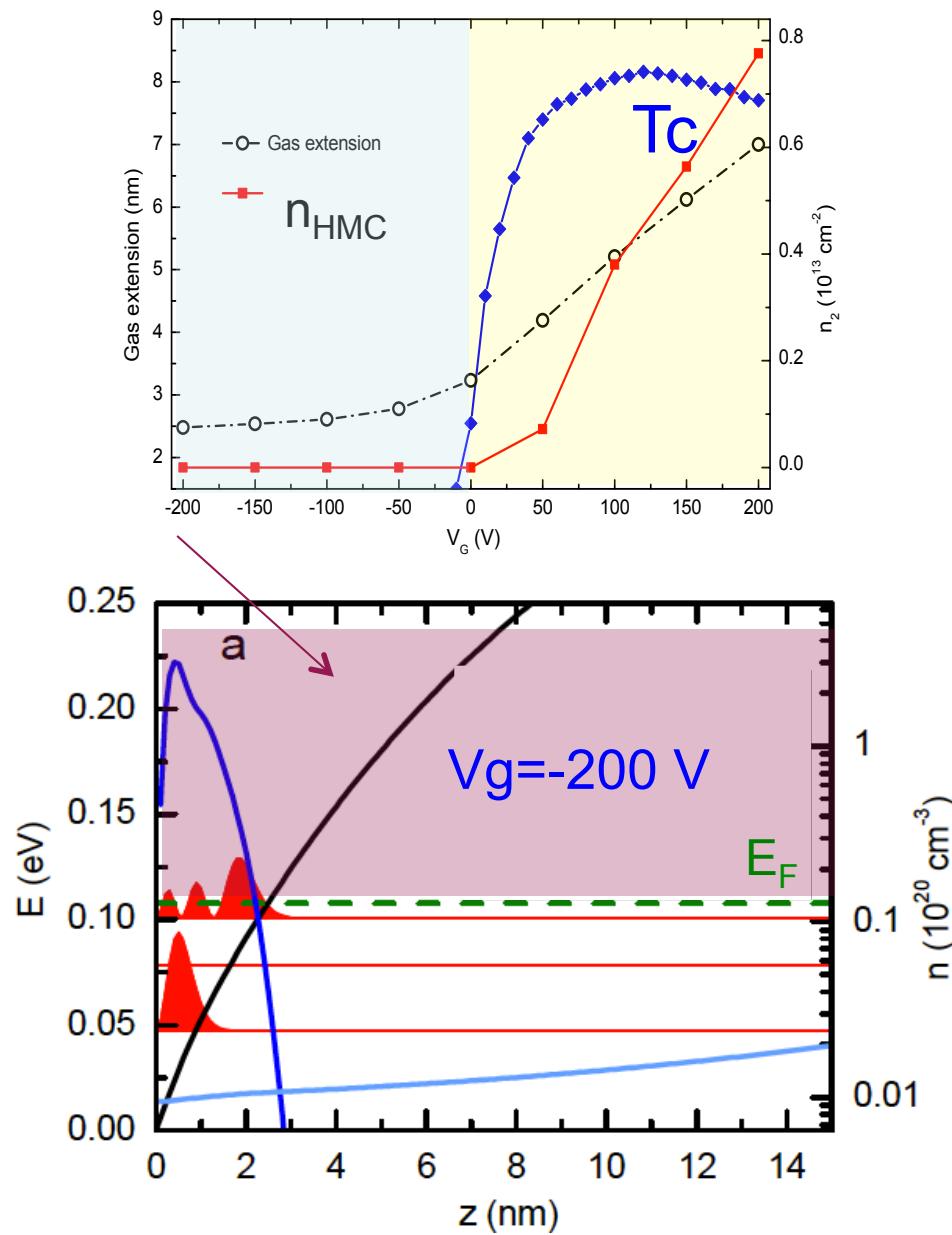
- Superconductivity is triggered by minority mobile carriers that extend in  $\text{SrTiO}_3$
- Superfluid density measured by scanning probe :  $n_s \sim 5 \cdot 10^{12} \text{ cm}^{-2}$

Bert et al, PRB (2012)

# Superconductivity related to mobile carriers ?

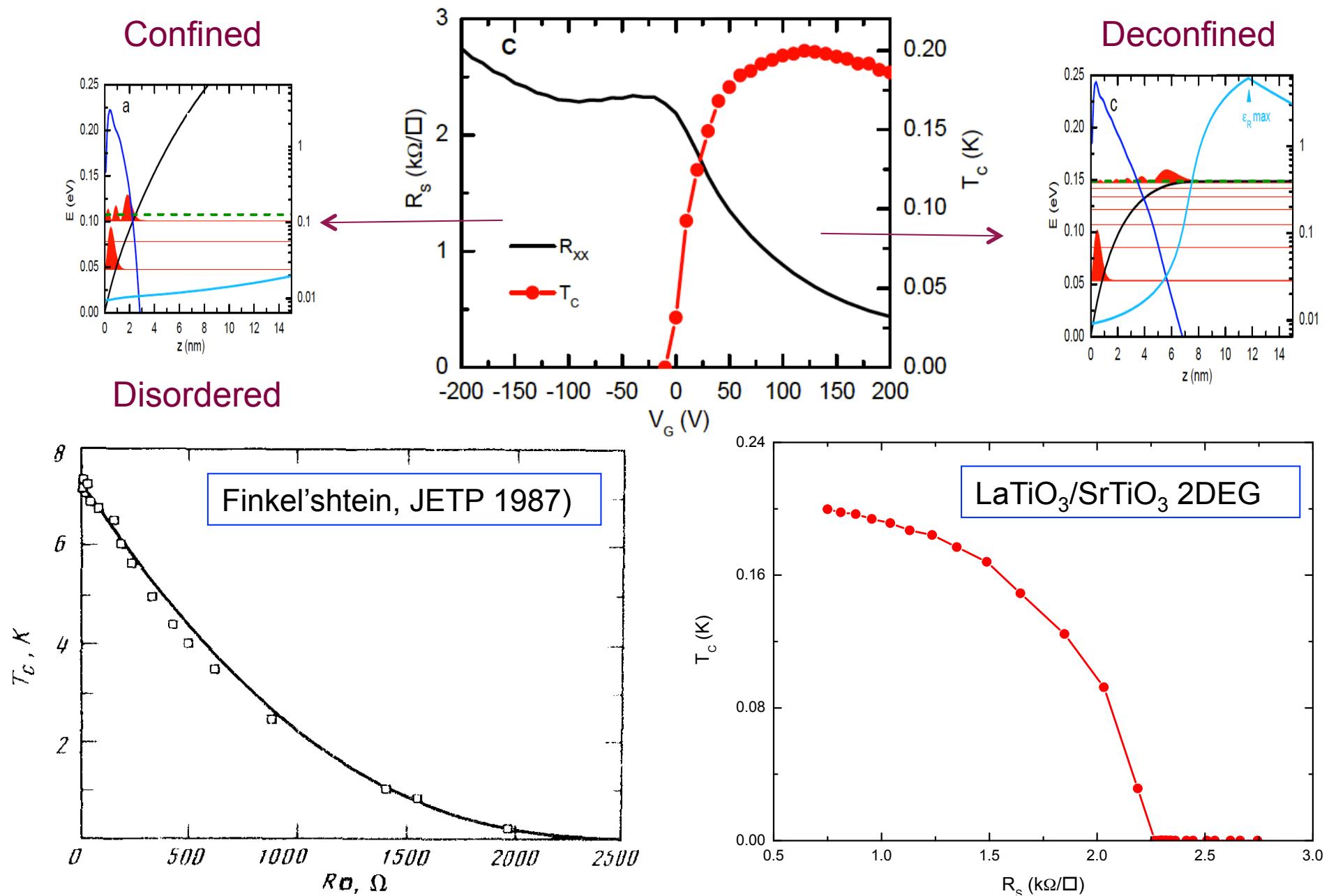


# Superconductivity : not simply the bulk ...

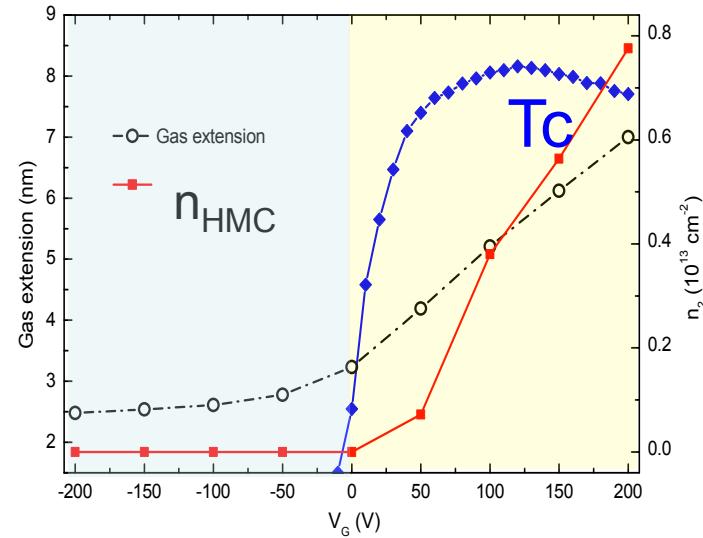


- Density in the well always high enough for bulk STO SC
- Not a direct image of the bulk Superconducting dome

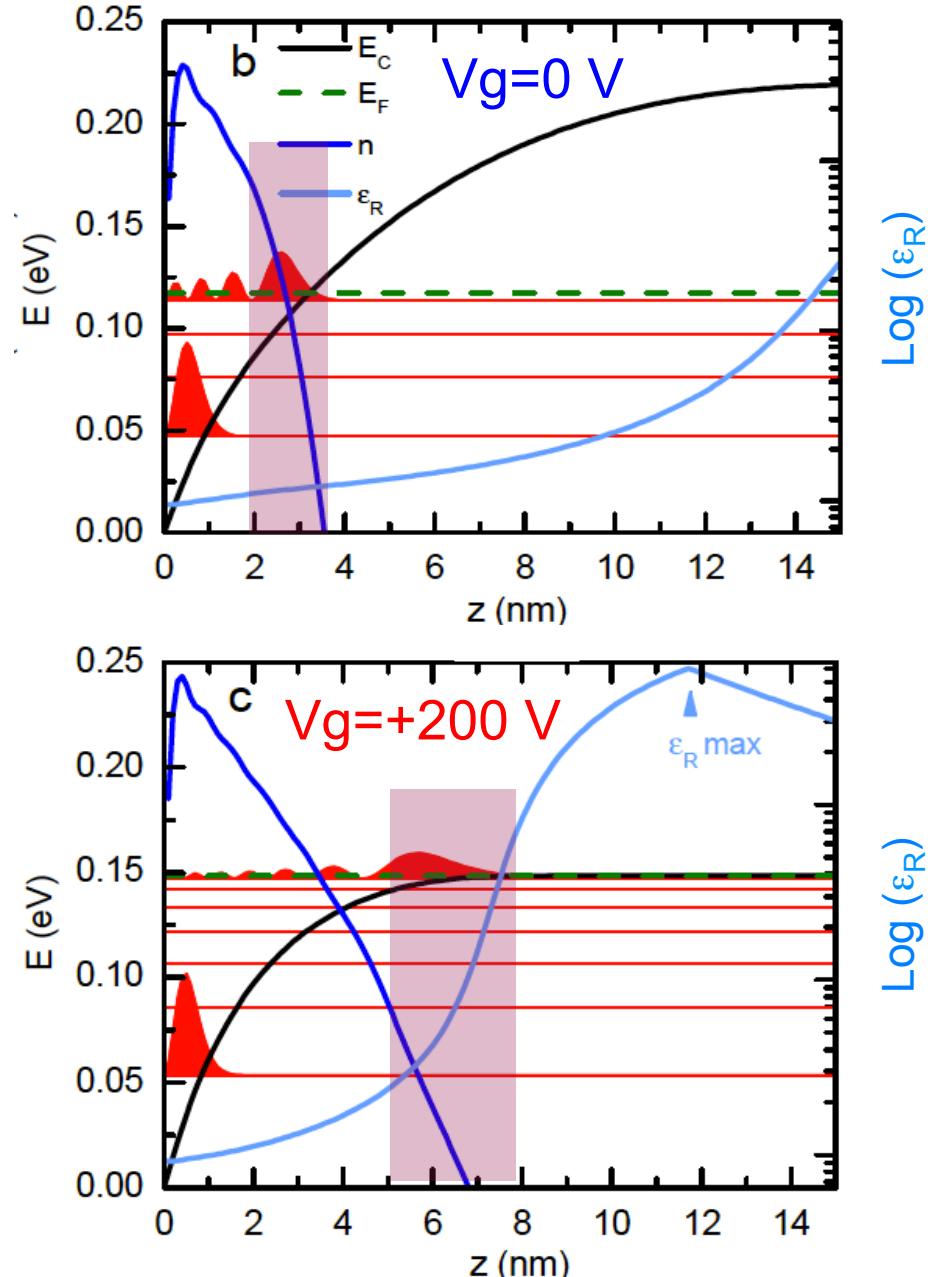
# Superconductivity ... and disorder ?



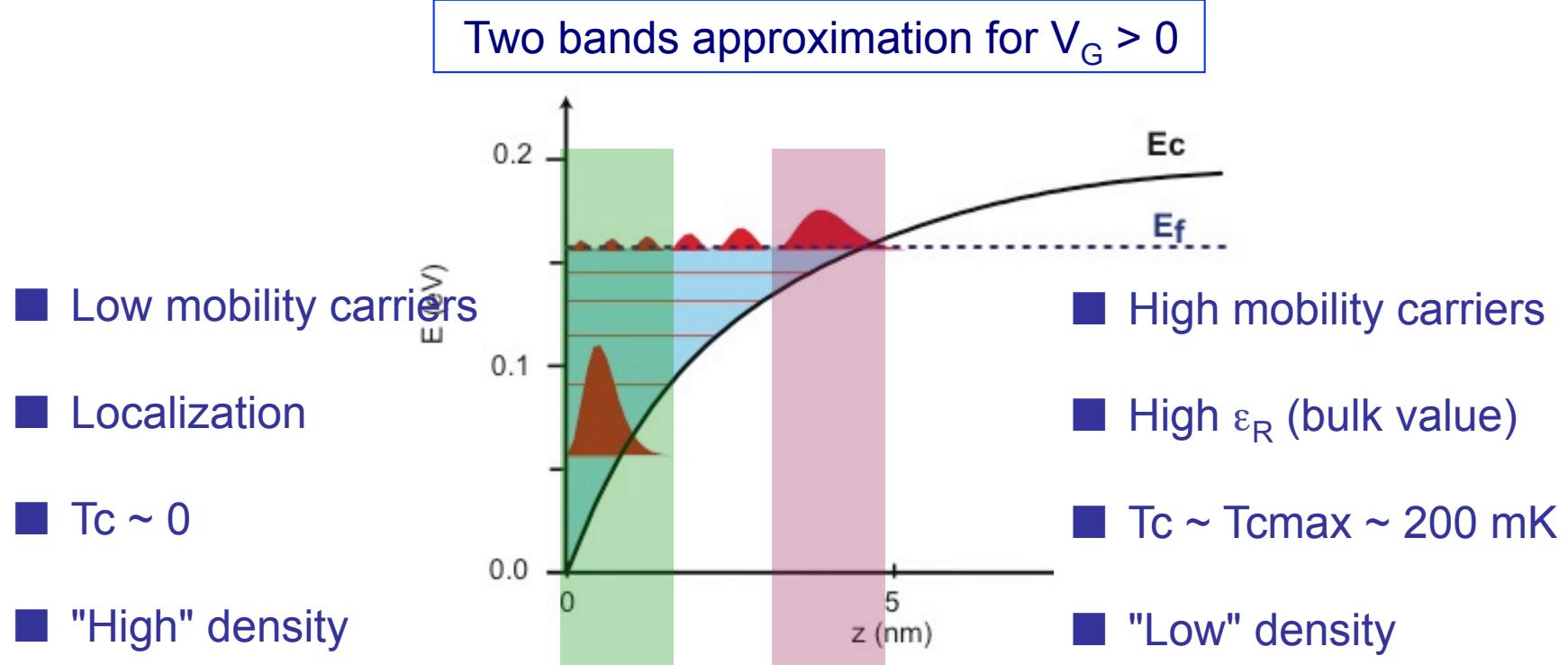
# Superconductivity ... and $\epsilon_R$ ?



- ➡ Bulk STO dielectric constant restored
- ➡ STO Superconductivity ???
- ➡ Polar properties of STO
- ➡ cf Koonce, Appel, Takada ...

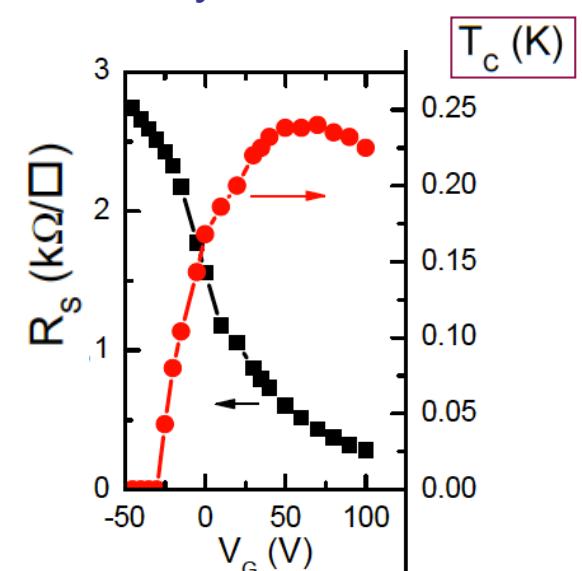


# Superconductivity ... a possible scenario



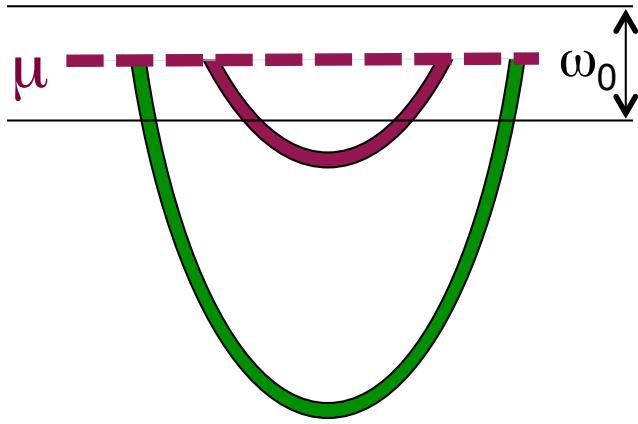
Multiband calculation :  $T_c$  ( chemical potential)  $\sim T_c (V_G)$

S. Caprara et al, Phys Rev B (R) 88, 020504 (2013)



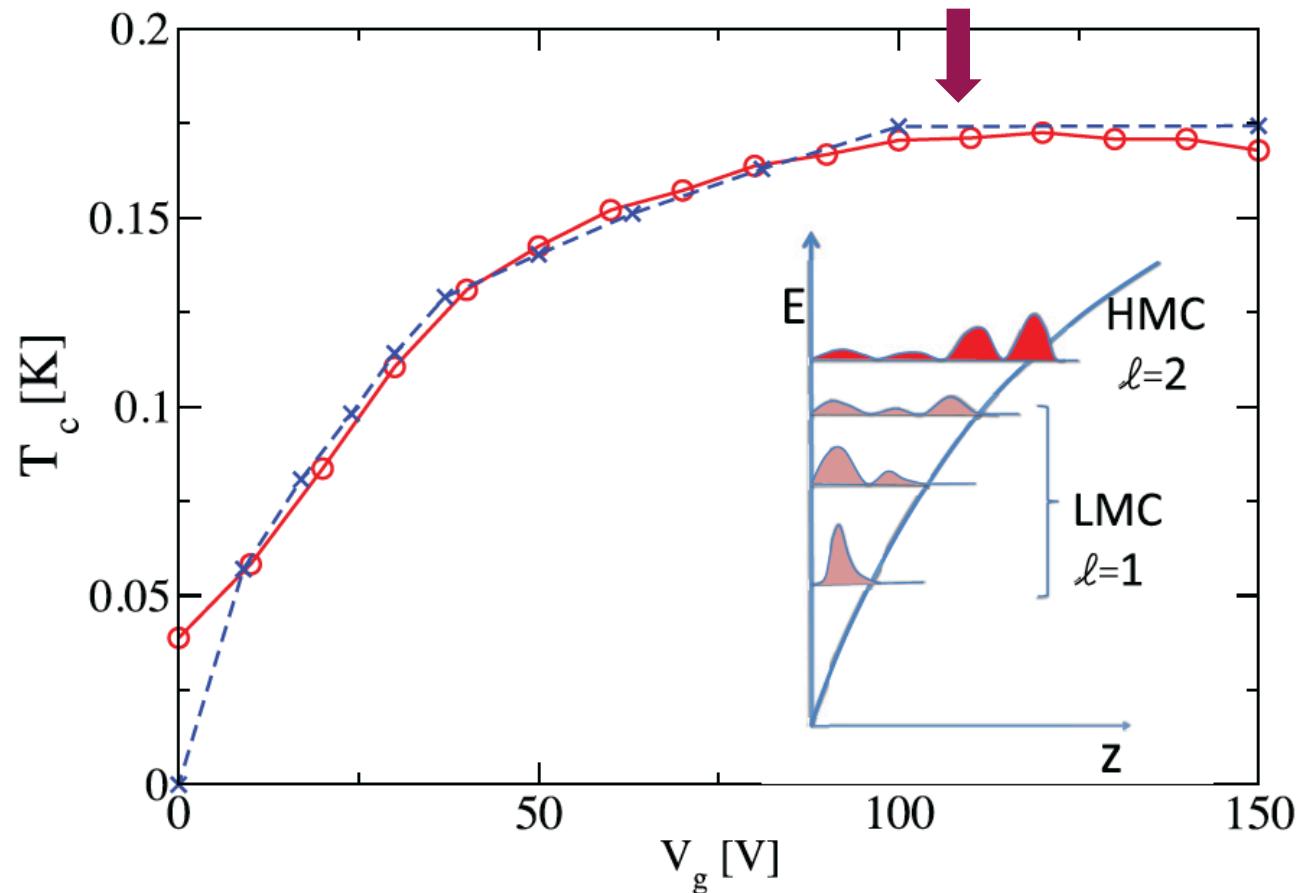
# Superconductivity ... a simple model

Coupling of two bands progressively filled



$$T_c \approx 1.14 \sqrt{\hbar\omega_0\mu} e^{-1/\lambda_2}$$

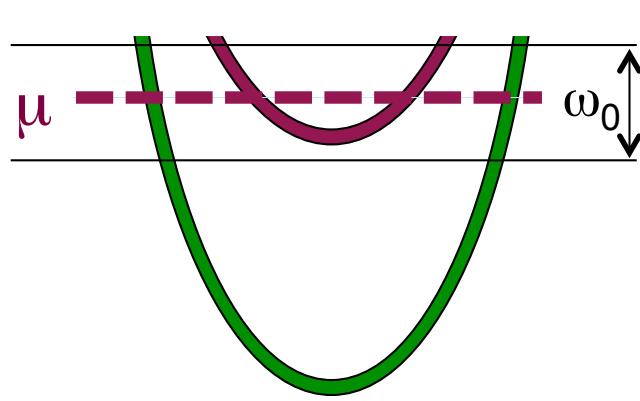
- Good agreement
- $\omega_0 \approx 23$  meV
- $\lambda_2 \approx 0.25$



S. Caprara et al, Phys Rev B (R) 88, 020504 (2013)

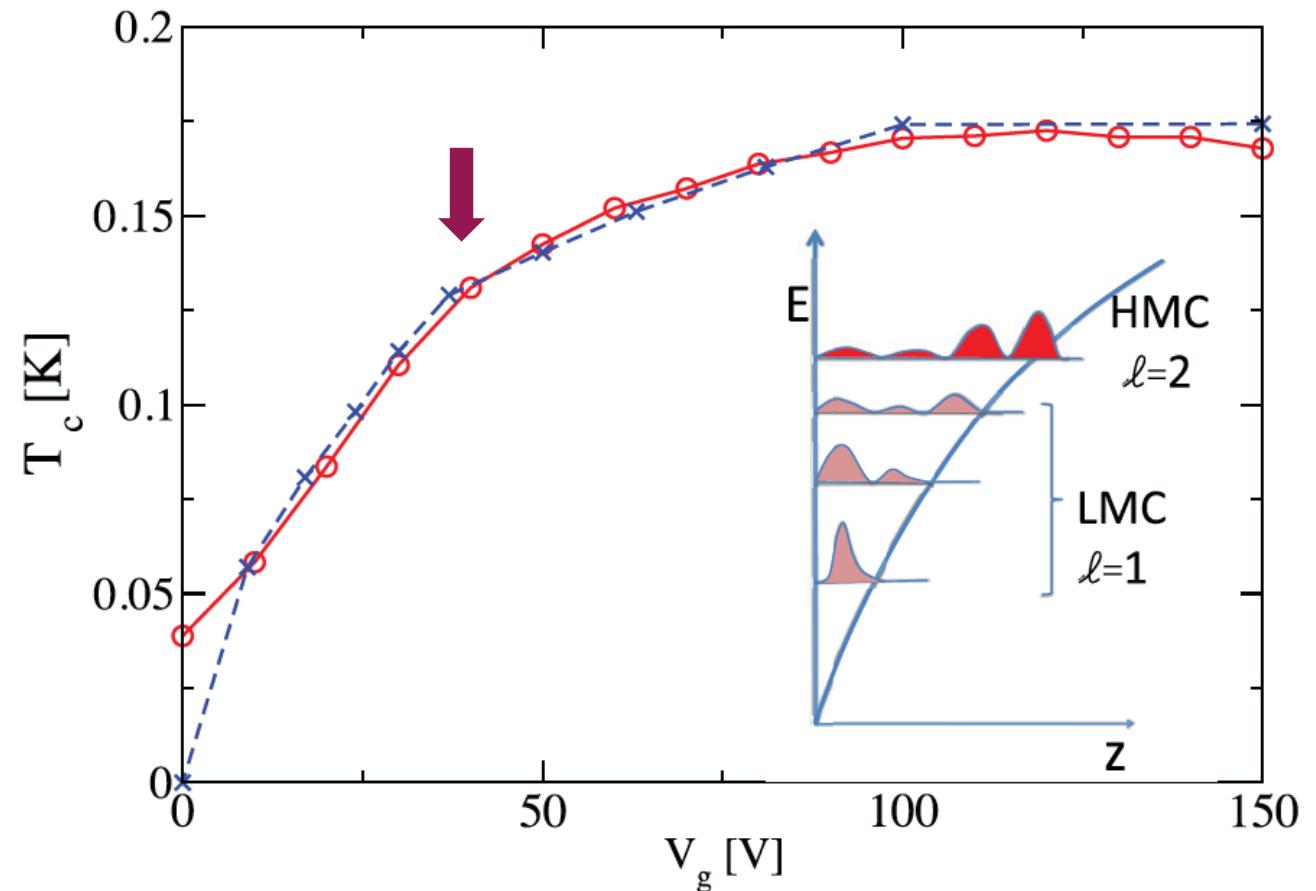
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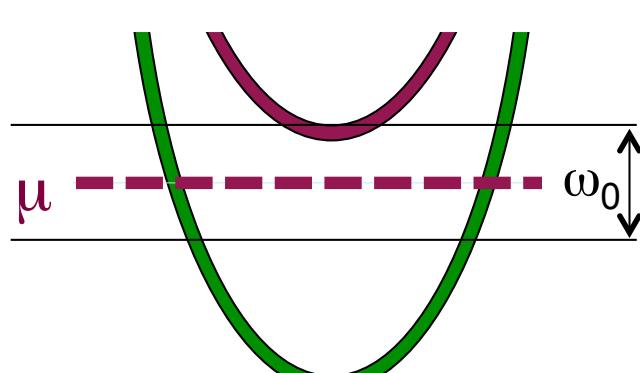
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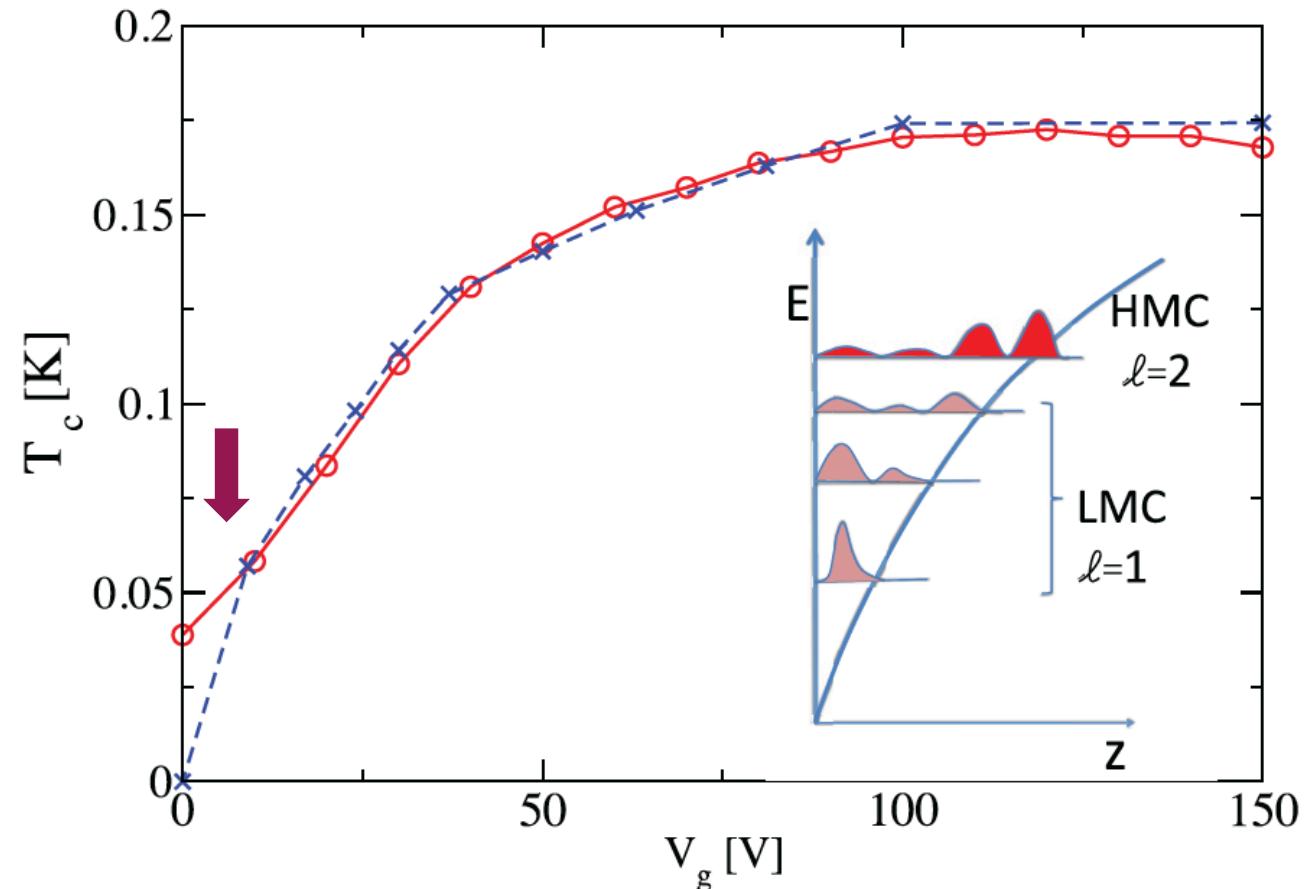


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■ Good agreement

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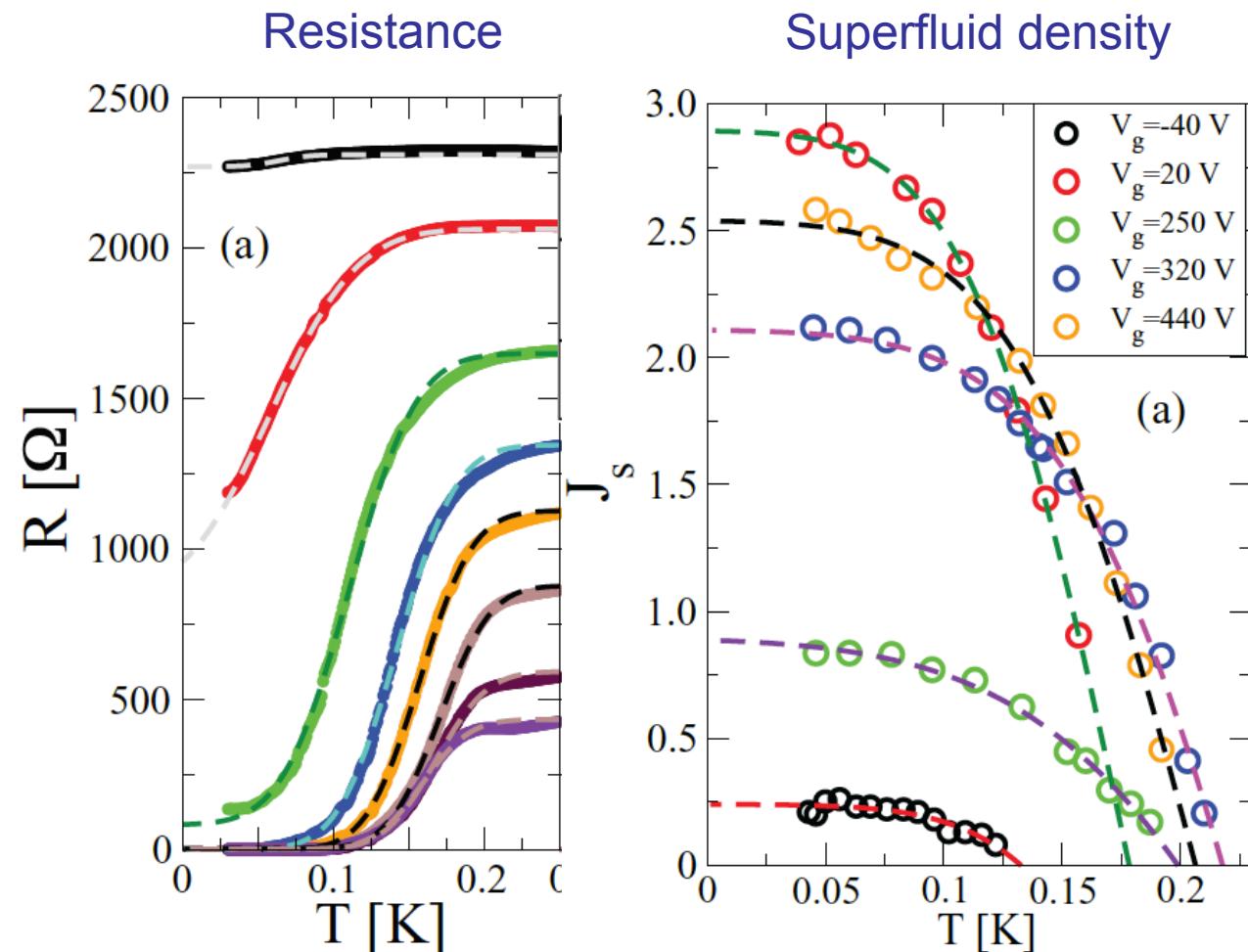
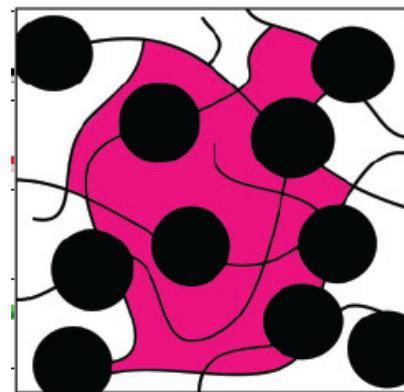
■  $\lambda_2 \approx 0.25$



S. Caprara et al, Phys Rev B (R) 88, 020504 (2013)

# Superconductivity ... inhomogeneous medium

- Effective Medium Theory
- Random Resistance Network
- Filamentary structure



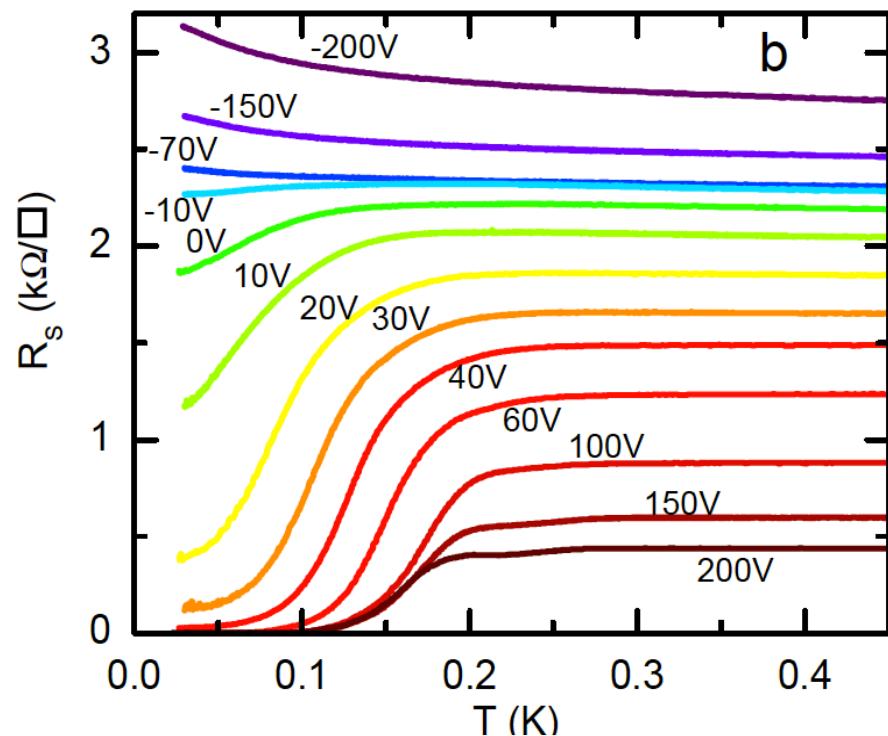
S. Caprara et al, Phys Rev B (R) 88, 020504 (2013)

D. Bucheli et al, New J. of Phys. 15, 023014 (2013)

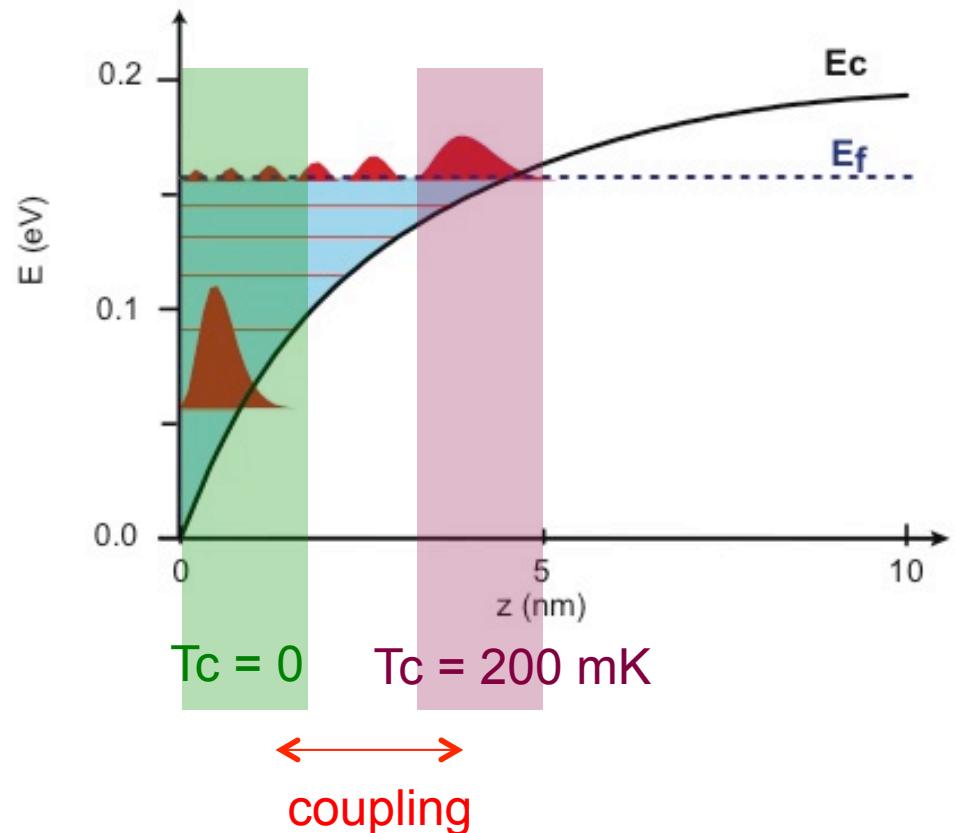
cf M. Grilli's talk

# First conclusions about superconductivity

- Control of the 2-DEG SC by  $V_G$



- Two types of carriers (one with  $T_c \neq 0$ )



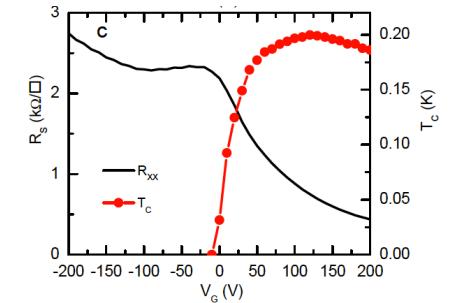
Biscaras et al, PRL 108, 247004 (2012)

- Presence of inhomogeneous superconducting behavior

# Outline

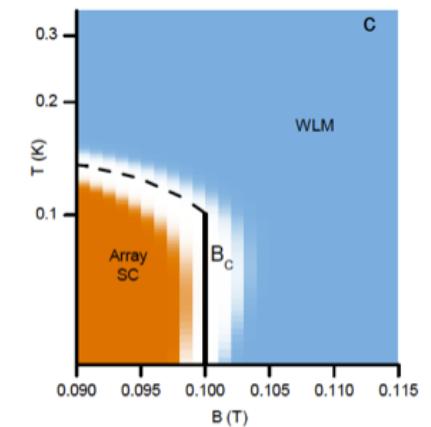
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Tunable superconductivity in oxide 2DEG



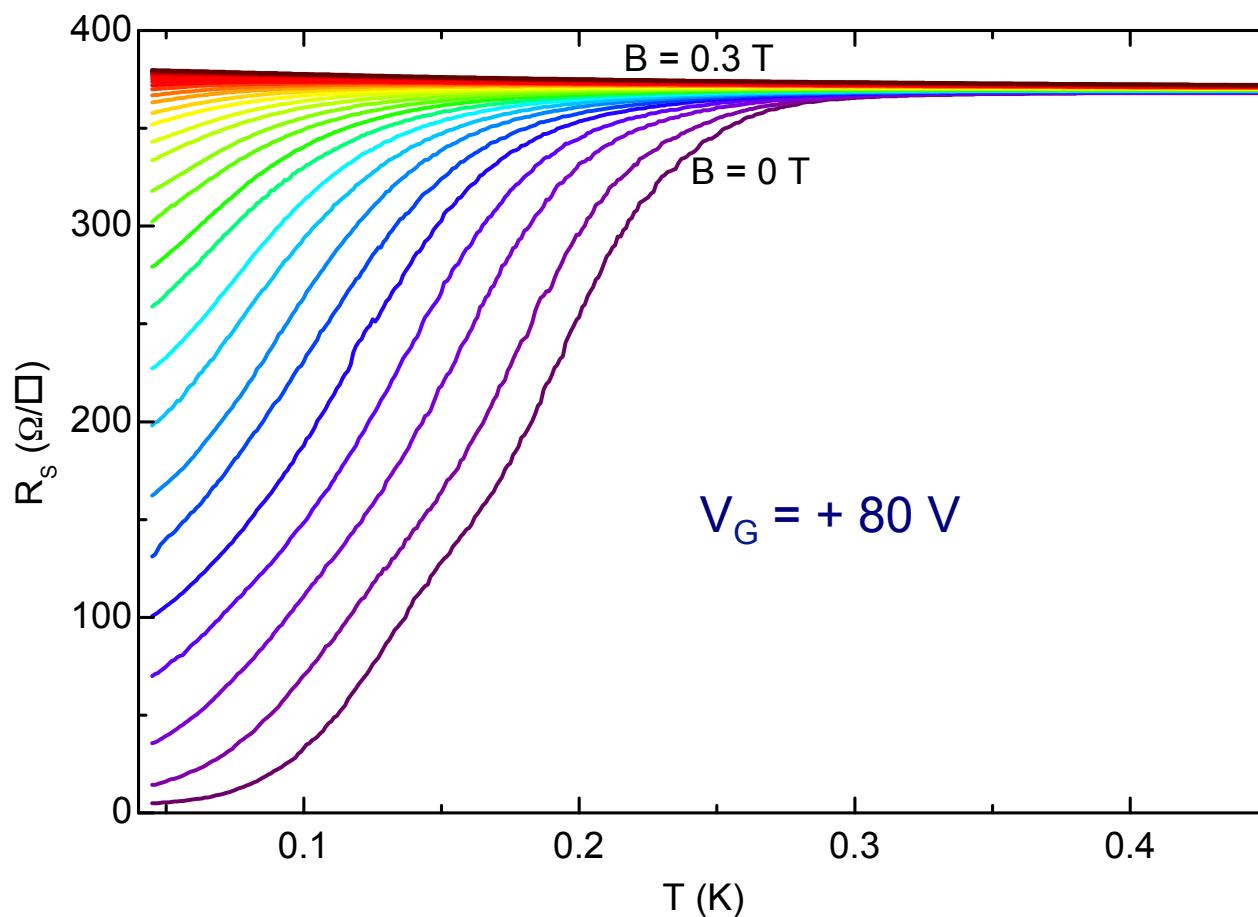
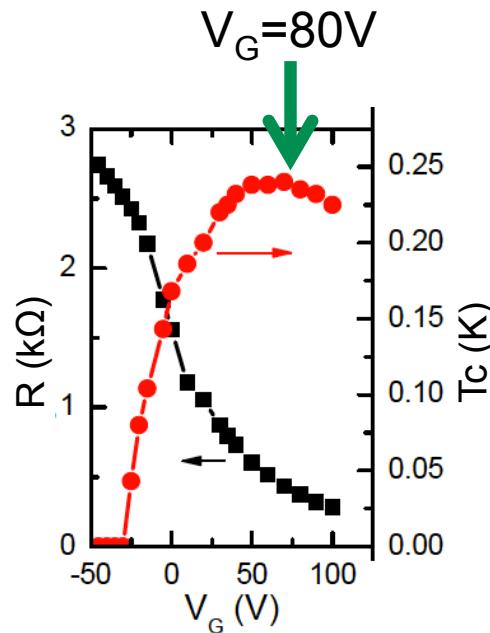
Quantum phase transition under magnetic field

Conclusions



# Magnetic field driven quantum phase transition

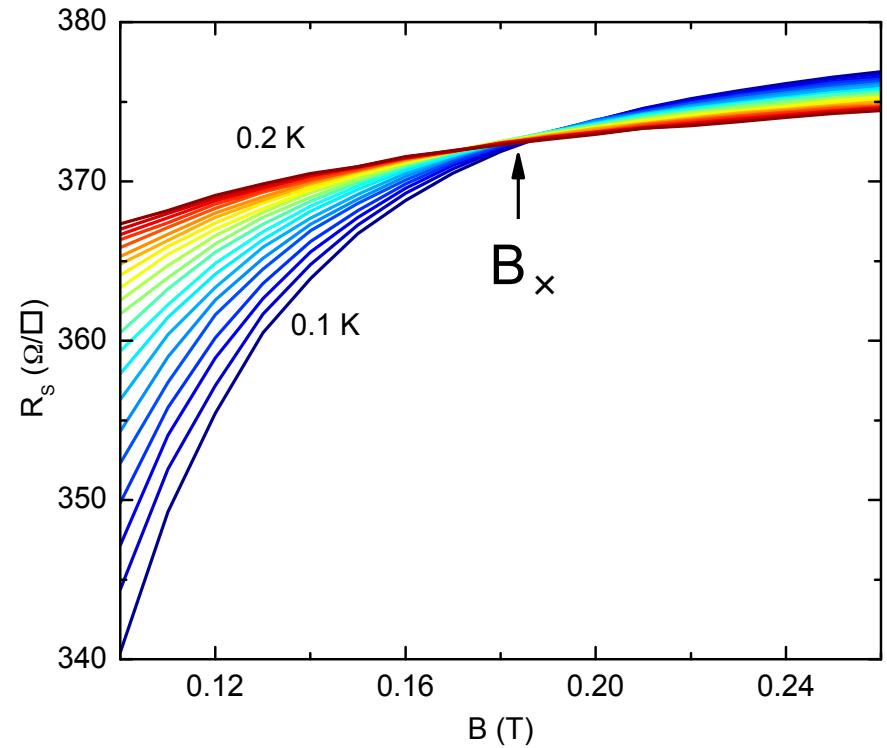
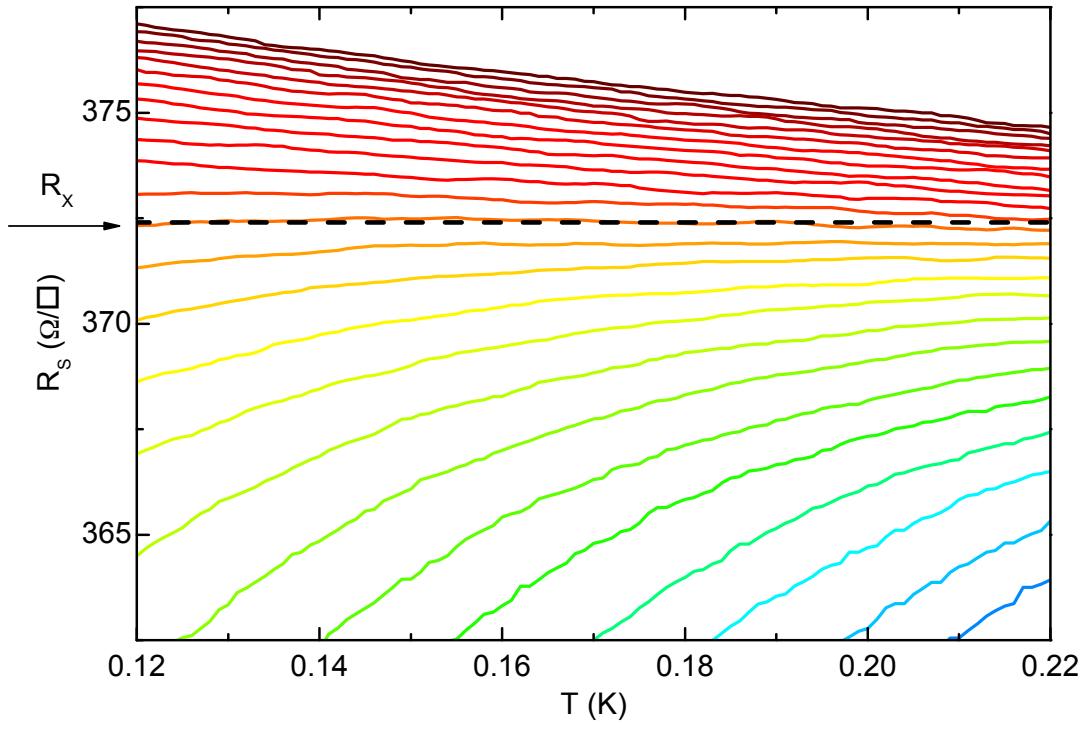
- Suppression of superconductivity by a perpendicular magnetic field at  $V_G=80V$



- ➔ Transition from superconducting to weakly localized metallic state

# Magnetic field driven quantum phase transition

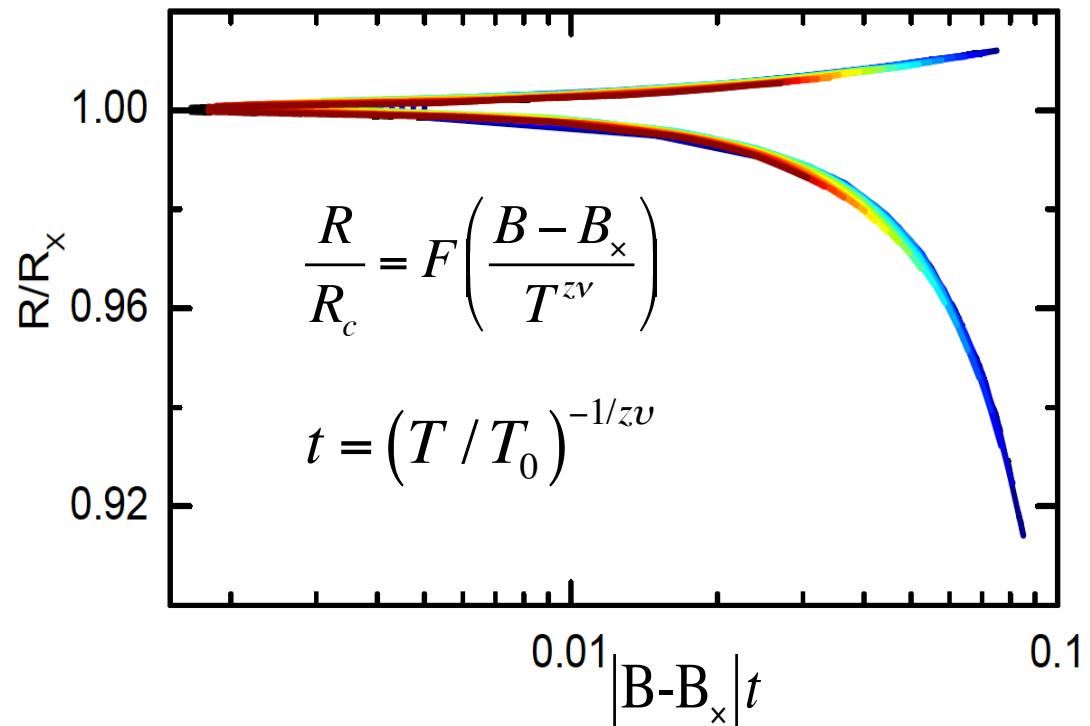
- Suppression of superconductivity by a perpendicular magnetic field at  $V_G=80V$



- ➔ Crossing point at  $B_x$ : a first signature of a quantum phase transition

# Scaling and critical exponents

## ■ Finite size scaling analysis



Correlation length

$$\xi \approx |B - B_x|^{-\nu}$$

Dynamical correlation length

$$\xi_\tau \approx |B - B_x|^{-z\nu}$$

Biscaras et al, Nature Mat. 12, 542 (2013)

Scaling Behaviour with  $z\nu = 2/3$  (as in a-Bi, NbSi, ... )

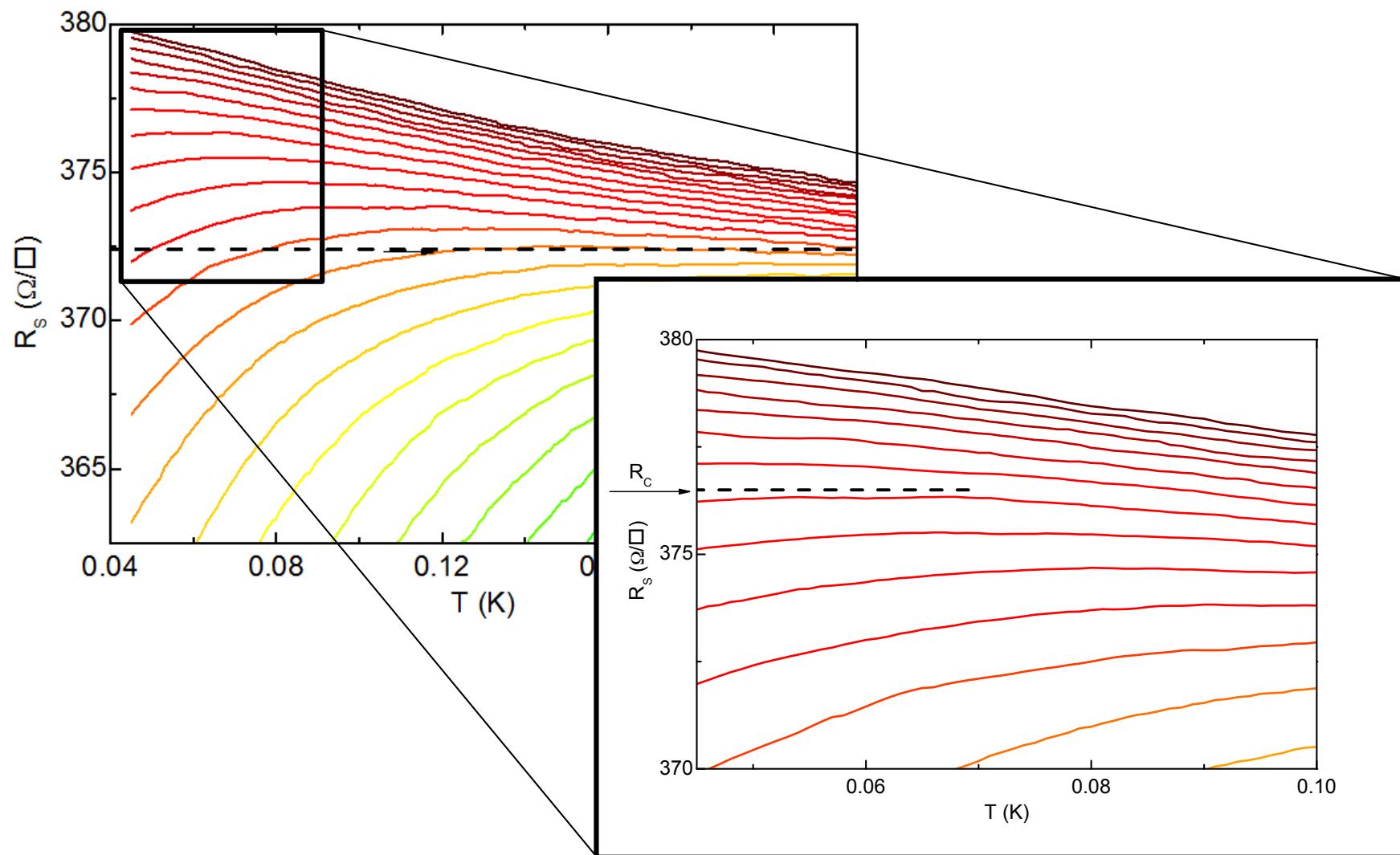
Superfluid transition in charged system :  $z=1$

Universality Class : (2+1)D XY in the **clean limit** :  $\nu = 2/3$  (Quantum Phase Fluctuations)

Disordered systems : **Harris criteria** sets  $\nu > 1$

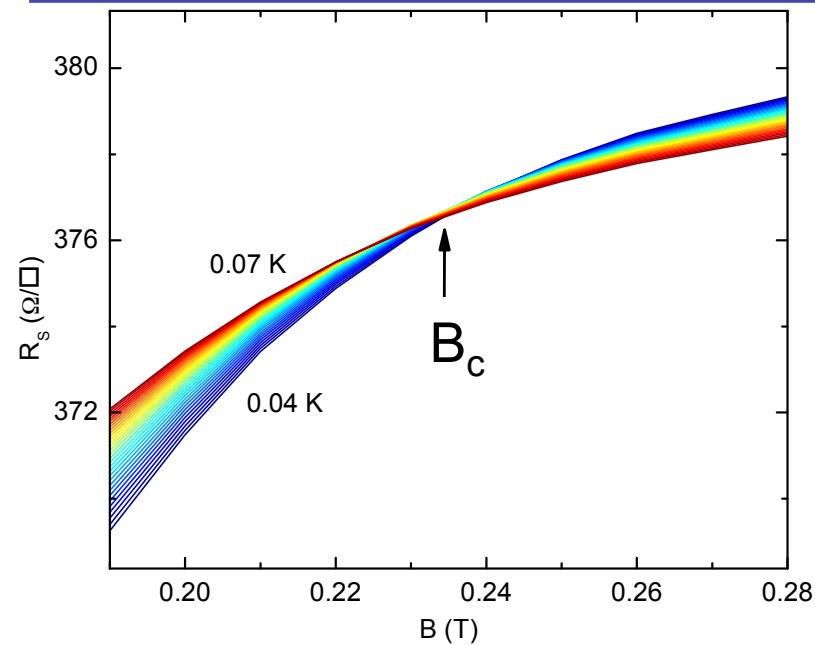
Harris, J Phys. C Solid State, 1974

# A true quantum Phase Transition ?



→ Scaling does not work at low temperature !

# Scaling at lower temperature



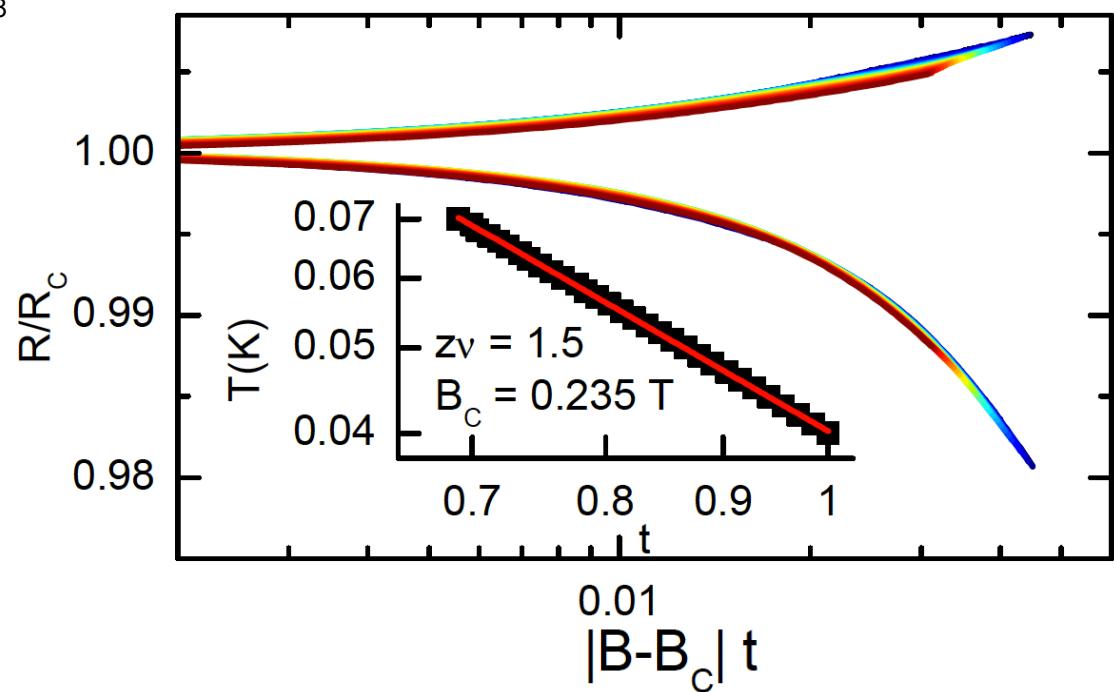
■ Crossing point at  $B_c$

$$B_c > B_x$$

■ Finite size scaling analysis

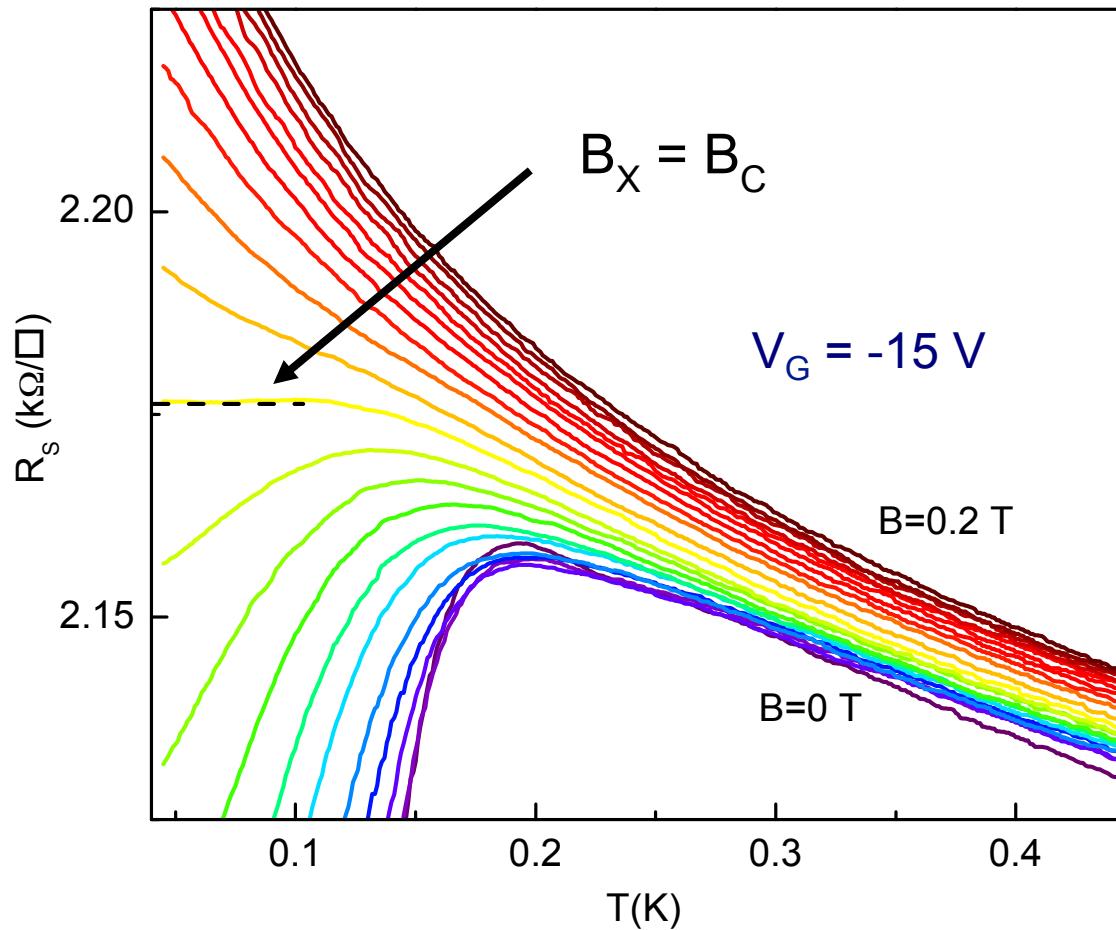
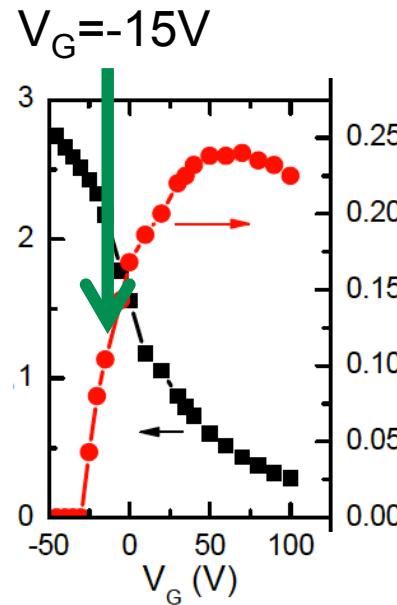
$$t = (T/T_0)^{-1/z\nu}$$

→ Critival exponent  $z\nu \approx 3/2$



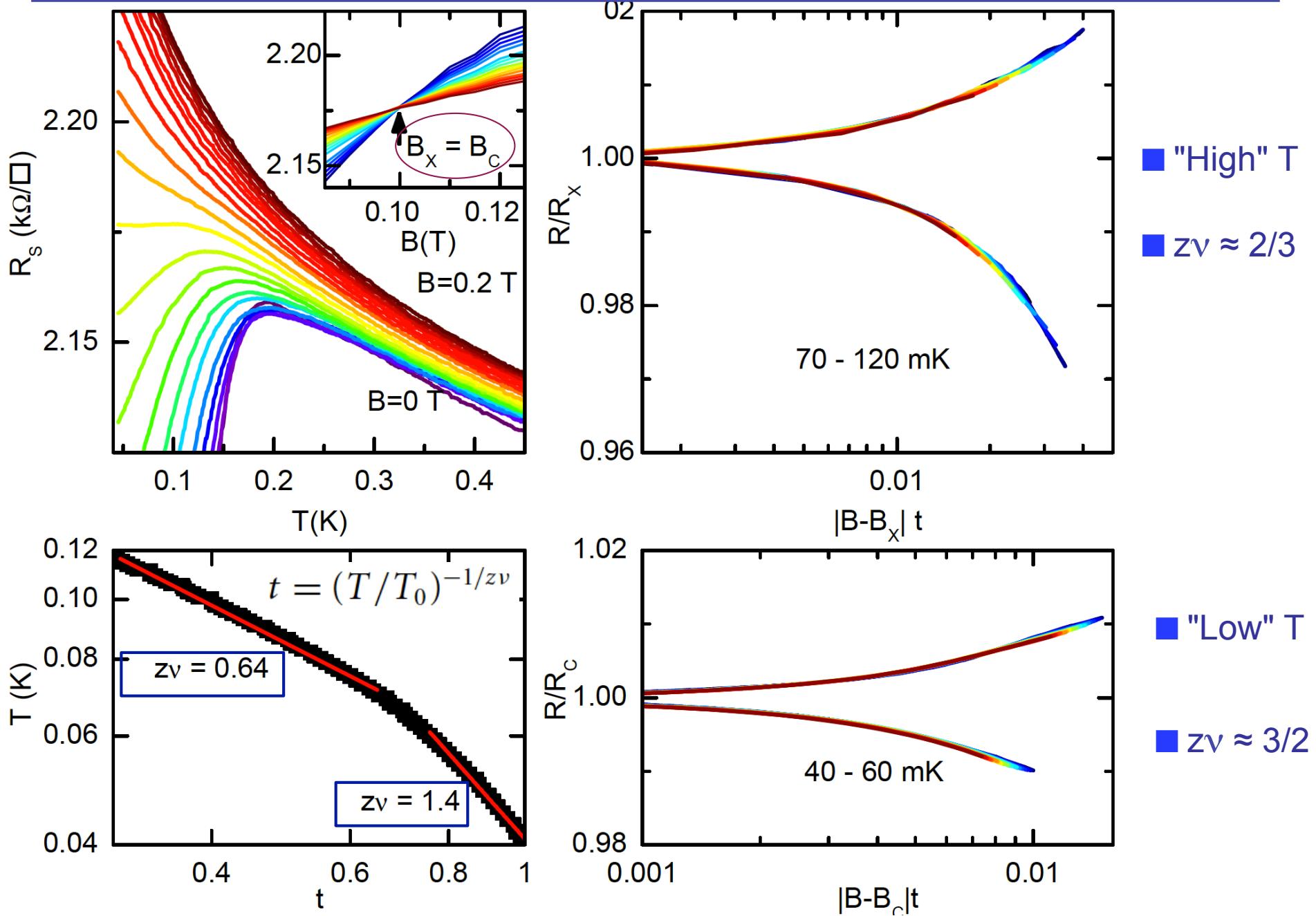
# Magnetic field driven quantum phase transition

- Suppression of superconductivity by a perpendicular magnetic field at  $V_G = -15V$

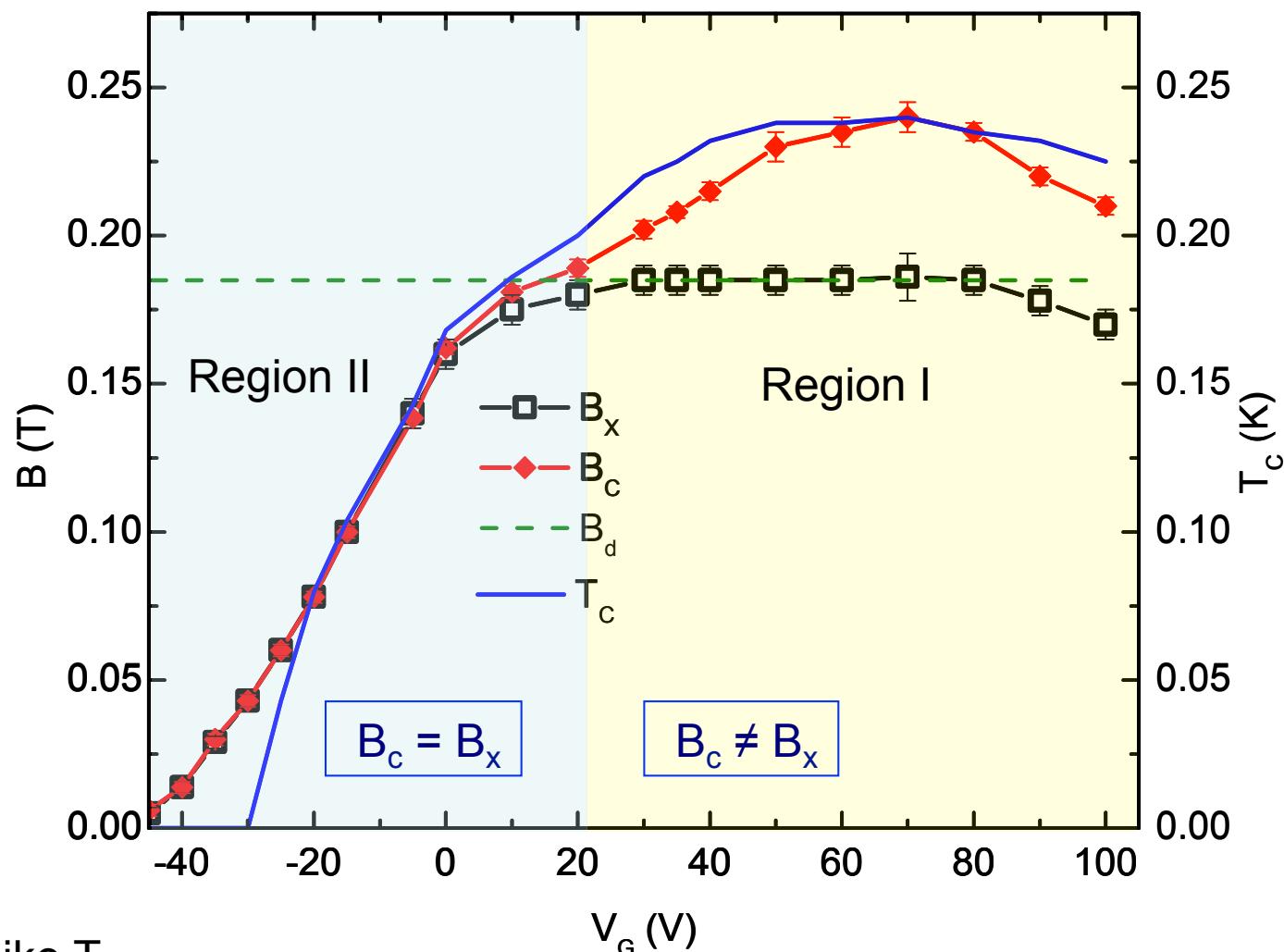


► Plateau in the resistance down to the lowest temperature !

# Scaling and critical exponents

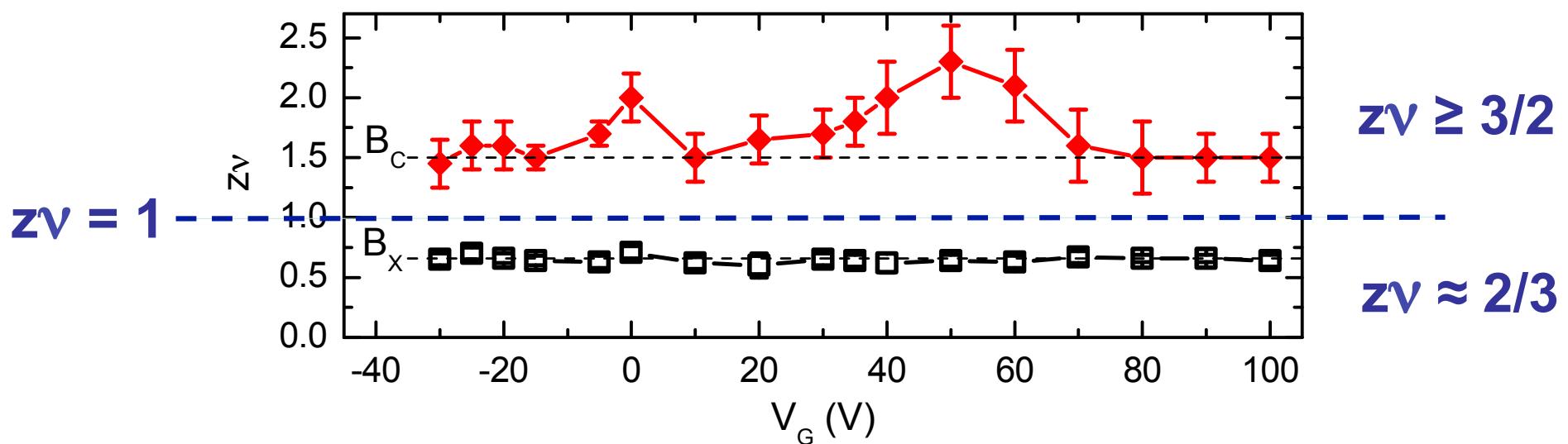
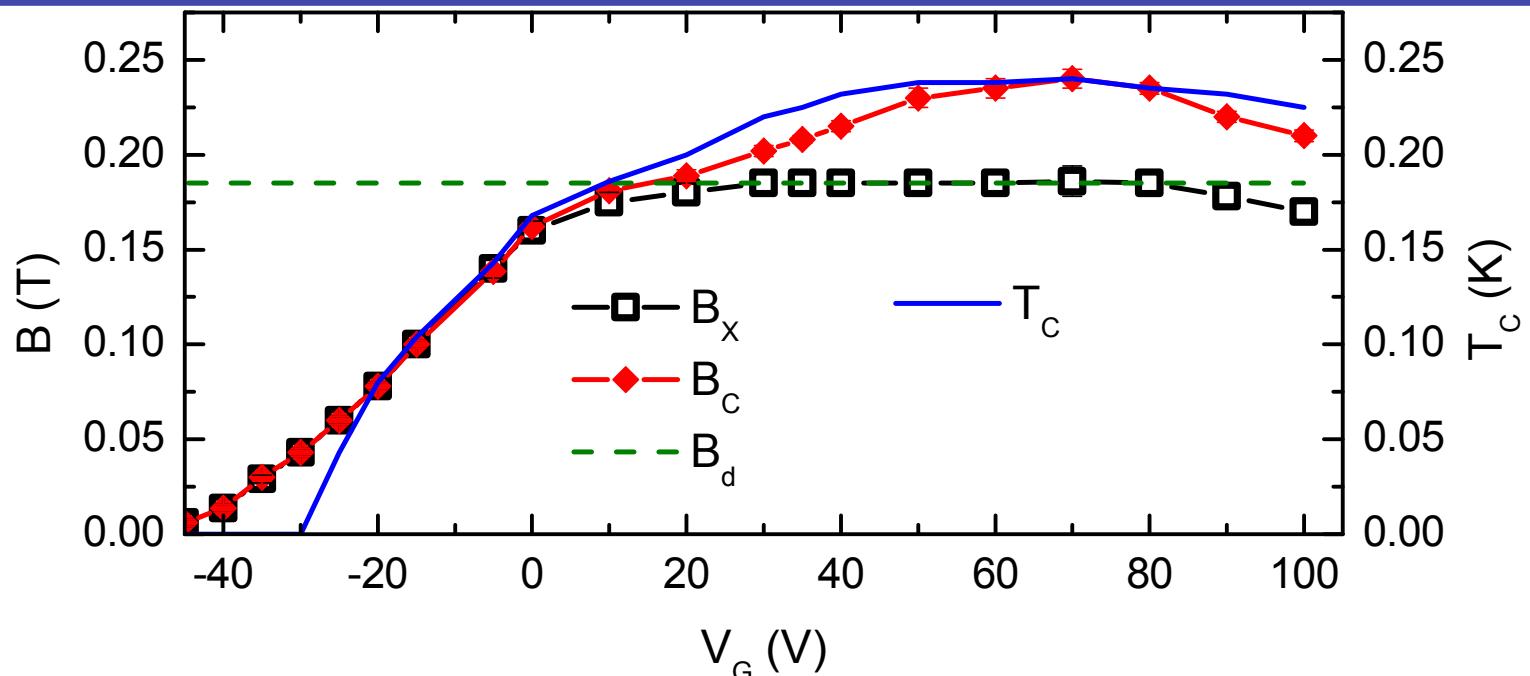


# Critical fields as a function of $V_G$



- $B_c$  scales like  $T_c$
- Two different regions ( $B_x = B_c$  &  $B_x \neq B_c$ )
- $B_x$  is constant in Region I

# Critical exponents as a function $V_G$

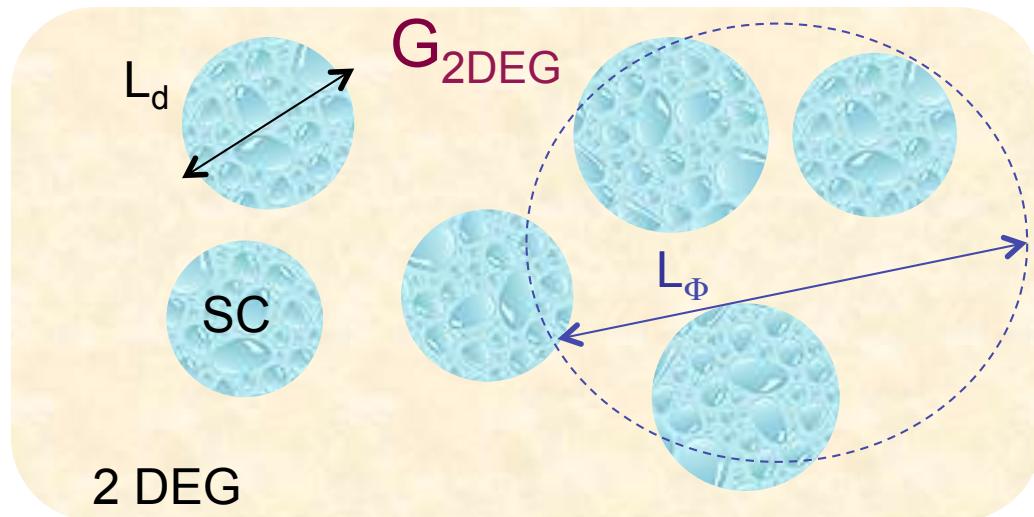


→ Multiple Critical Behavior ( $B_x$  &  $B_c$ ) associated to different critical exponents

# Multiple Quantum Critical Behaviors in 2D SC

Disordered array of SC puddles

Spivak et al, PRB 2008

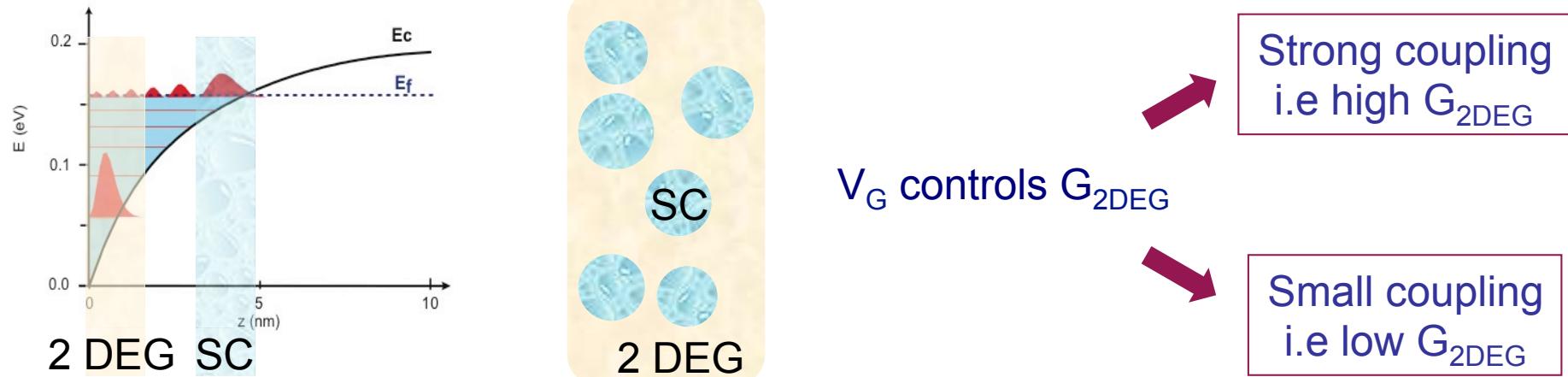


- Characteristic puddle size  $L_d$
- Puddles coupled by Josephson through  $G_{2\text{DEG}}$
- Phase length  $L_\Phi \sim T^{-1/z}$

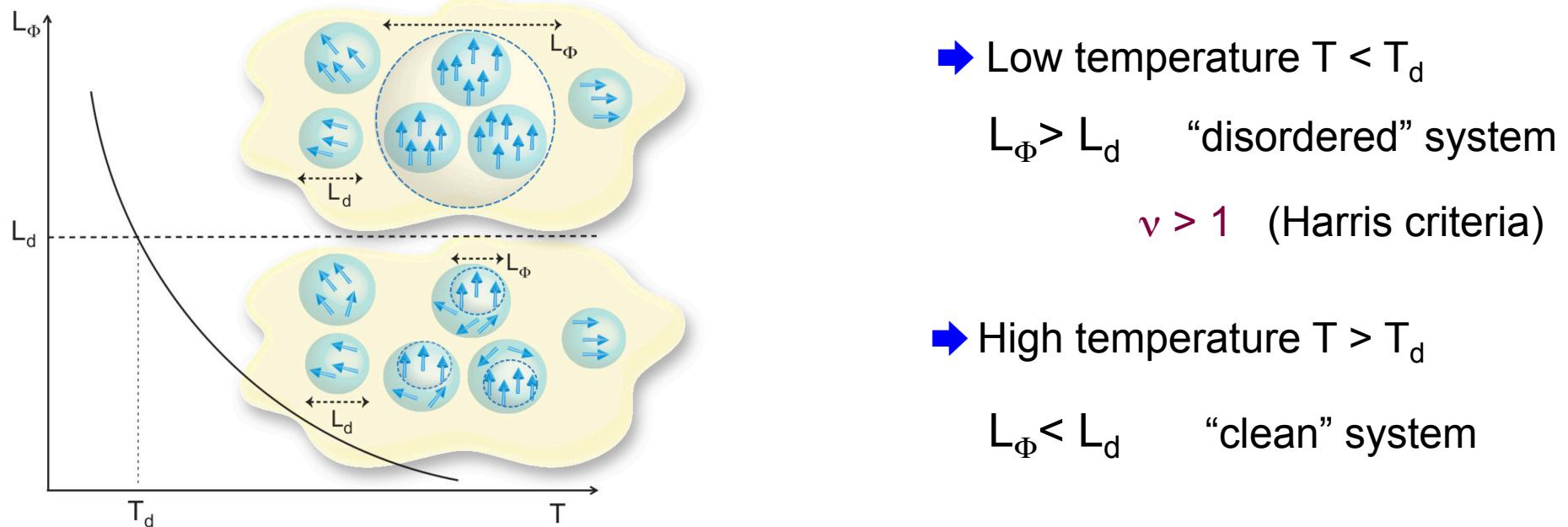
- Phase fluctuations in the puddles and between puddles : two critical fields
- Puddles may be not SC if the coupling is too weak : two regimes ( $G_{2\text{DEG}}$ )
- The critical field  $B_c$  scales like  $T_c$

# Towards a scenario ...

## ■ Superconducting puddles in a 2DEG matrix

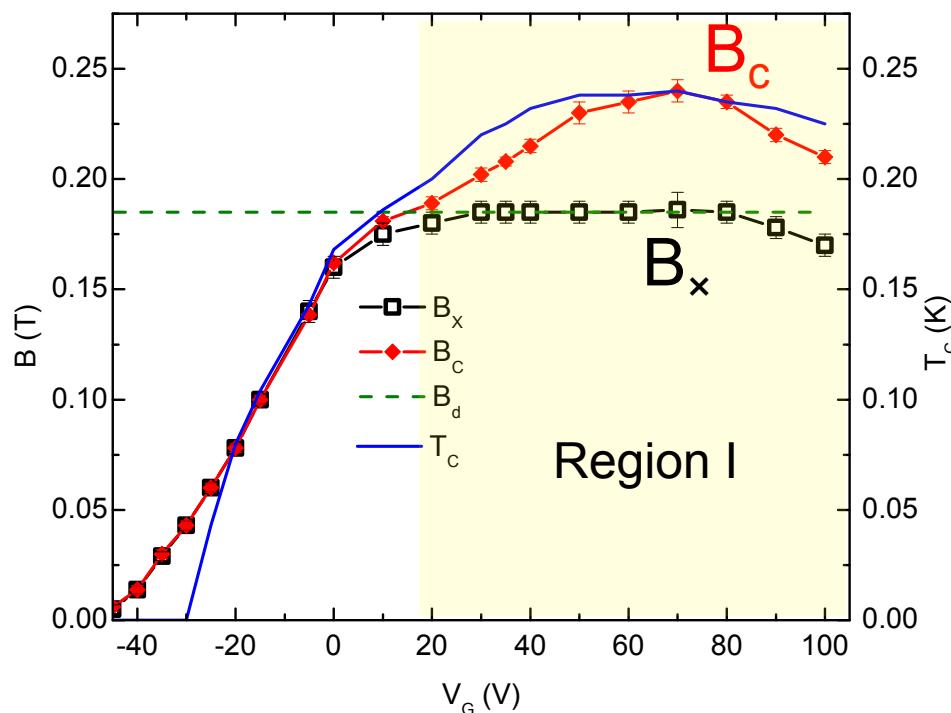


## ■ Divergence of the thermal dephasing length $L_\Phi \sim T^{-1/z}$

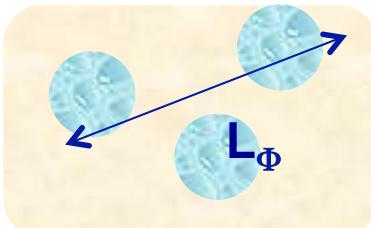


# The puddles scenario

■ High  $G_{2\text{DEG}}$  regime for  $V_G > 20\text{V}$



Low temperature

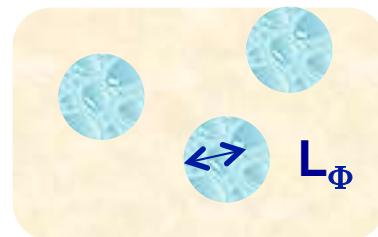


“disordered”

$z\nu \sim 3/2 > 1$

►  $B_c$  destroys superconductivity of the array

High temperature



“clean”  
 $z\nu = 2/3$

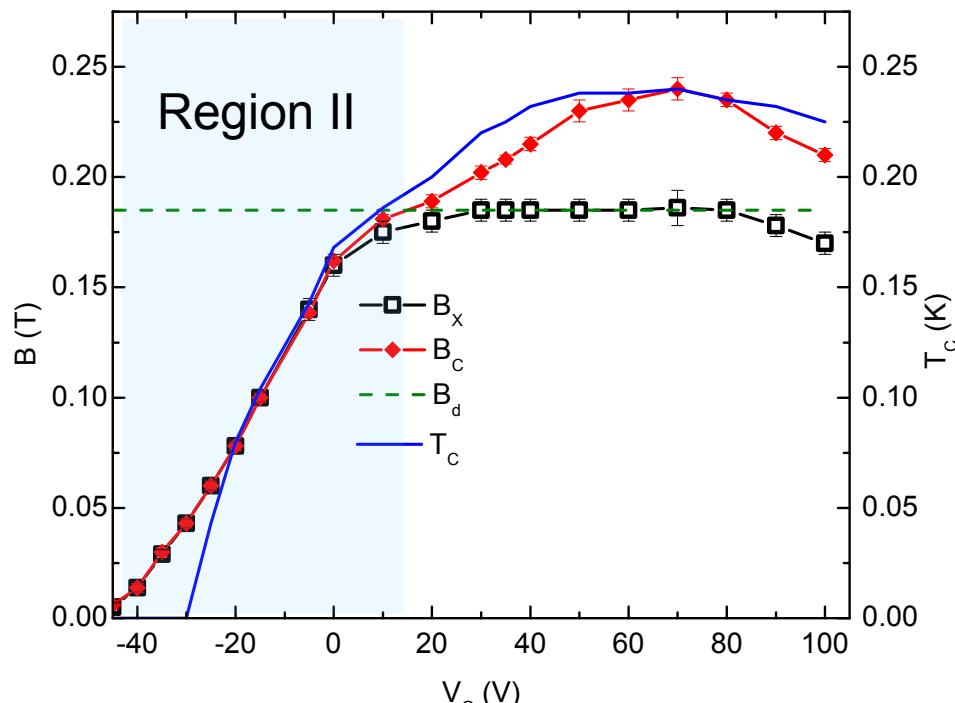
►  $B_x$  destroys superconductivity inside puddles

► Typical size of a puddle

$$L_d = \sqrt{\Phi_0 / B_x} \approx 100\text{nm}$$

# The puddles scenario

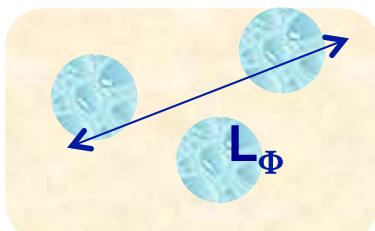
■ Low  $G_{2\text{DEG}}$  regime for  $V_G < 20\text{V}$



→ a single critical field  $B_c = B_x$

→ Puddles are always fluctuating

→  $B_c$  destroys superconductivity of the array



"disordered"

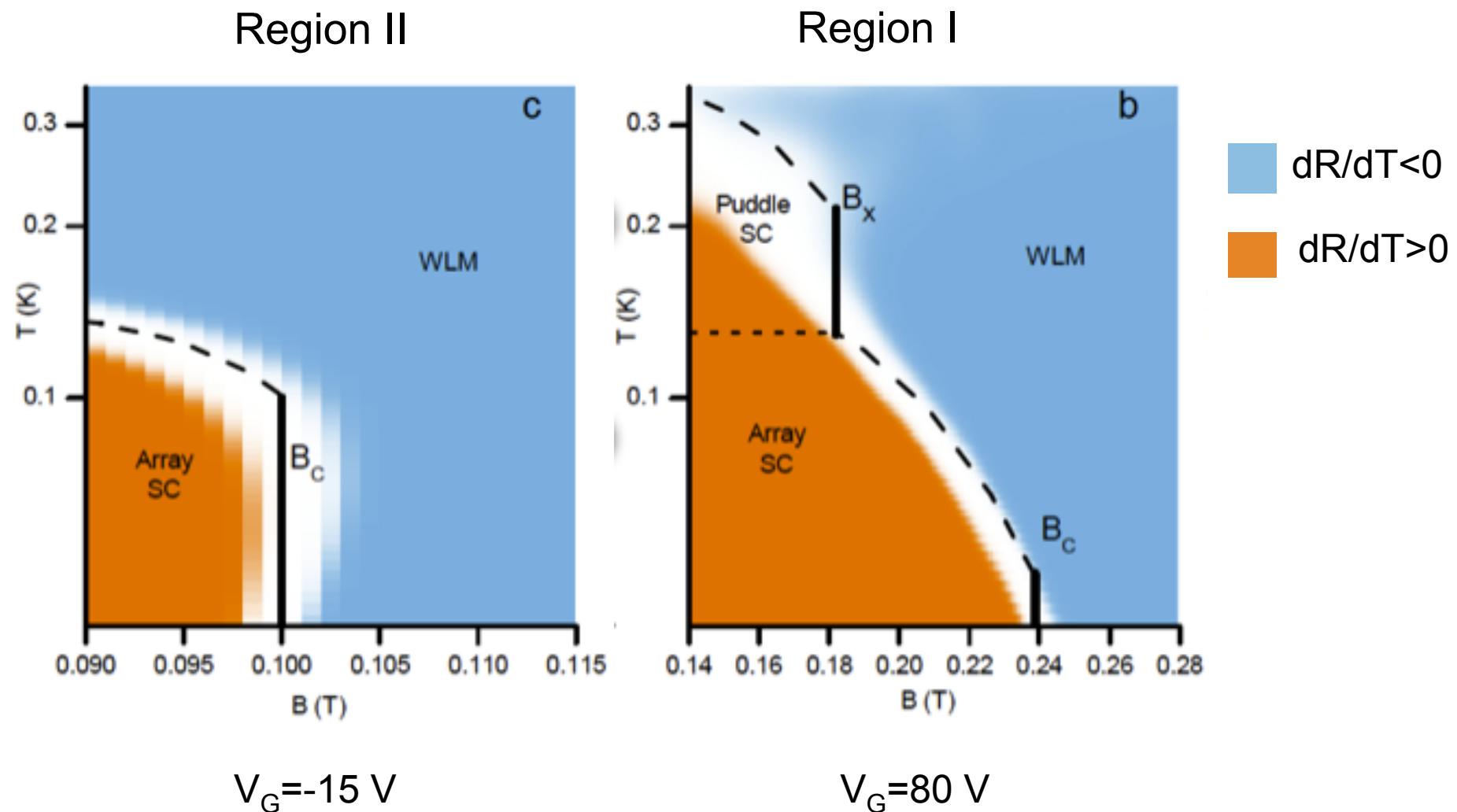
$$z\nu \sim 3/2 > 1$$



"clean"

$$z\nu = 2/3$$

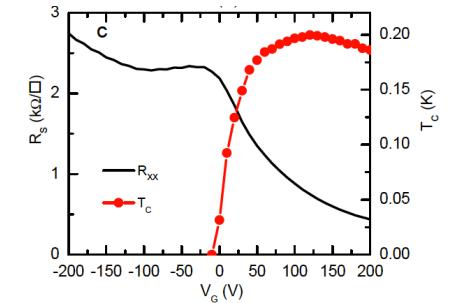
# B-T phase diagrams



# Outline

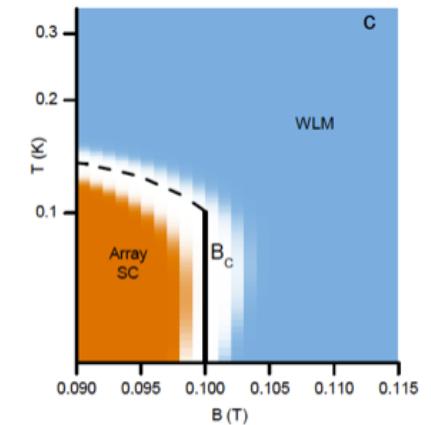
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Tunable superconductivity in oxide 2DEG



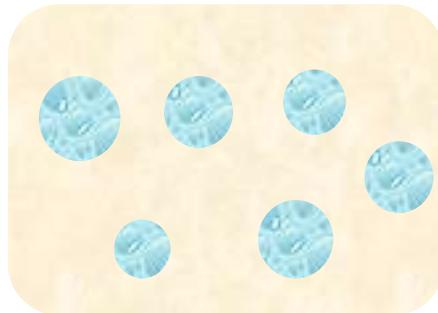
Quantum phase transition under magnetic field

Conclusions

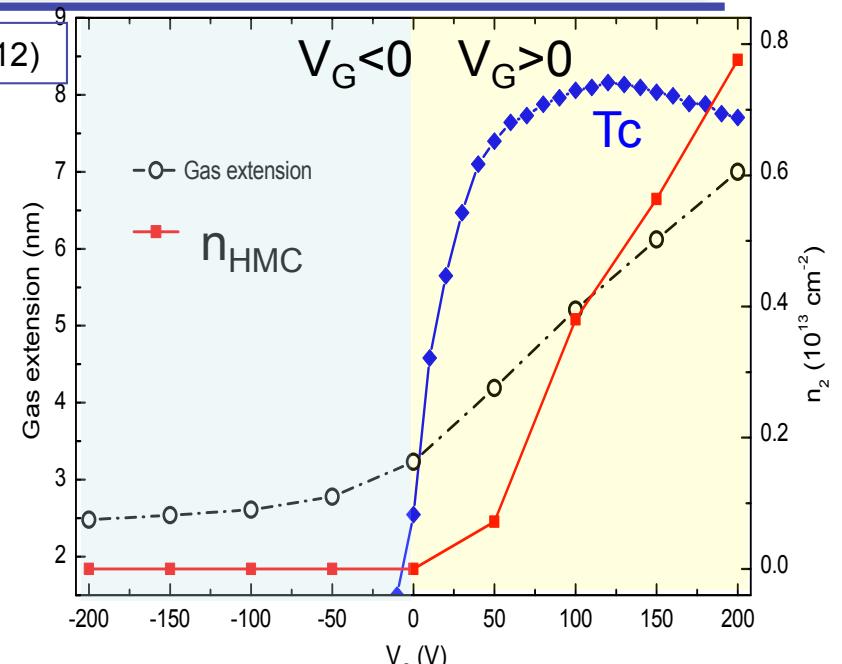


# Conclusions

- Two types of carriers
- Highly Mobile Carriers & Superconductivity
- Inhomogeneous behavior

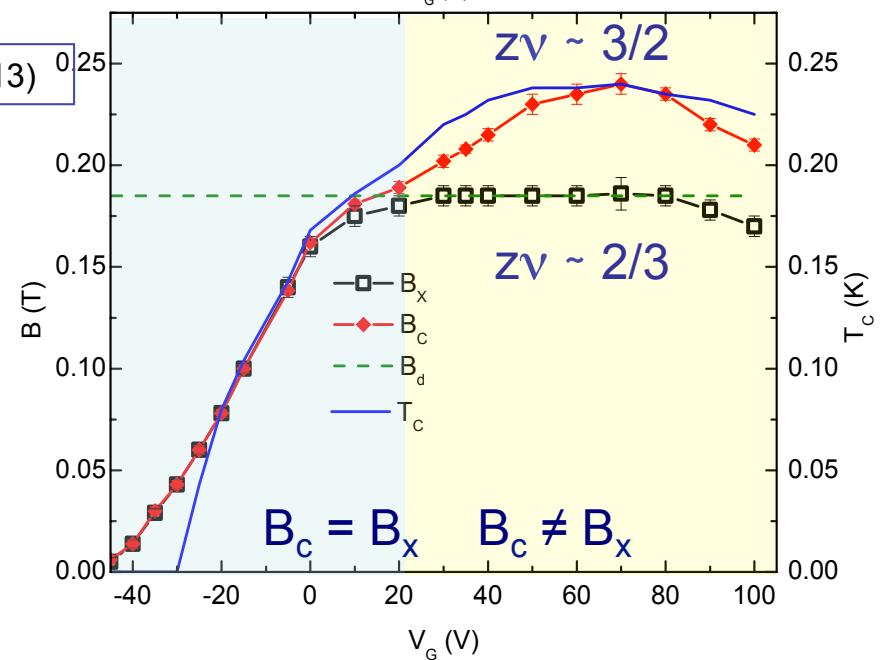


Biscaras et al, PRL 108, 247004 (2012)



Biscaras et al, Nature Mat. 12, 542 (2013)

- SC puddles coupled by a 2DEG
- (2+1)D XY quantum phase transition
- Two Critical Behaviors
- Cross over the "Harris criteria"
- ? Intrinsic filamentary inhomogeneity ?



# People

Phd students : Johan Biscaras – Simon Hurand – Alexis Jouan

Collaborators : N. Bergeal, C. Palma, LPEM (Paris)



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M. Grilli, S. Caprara, L. Benfatto, La Sapienza (Rome)



Biscaras et al, Nature Com. 1, 89 (2010)

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Caprara et al, PRB(R) 88, 020504 (2013)

# Collaborations

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