

Snapshots of the retarded interaction of charge carriers with ultrafast fluctuations in cuprates

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i-Lamp

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ICTP 27 October 2014



Outline

- Ultrafast spectroscopies of quantum materials

J. Orenstein, *Phys. Today* 65, 44 (2012)

- The “glue” and the retarded-interaction problem in high-T_c superconductors

P.W. Anderson, *Science* 317, 1705 (2007) and reply from D.J. Scalapino

- Snapshots of the retarded interaction with ultrafast fluctuations via 10 fs pulses

S. Dal Conte et al., *Science* 335, 1600 (2012)

S. Dal Conte et al. *submitted to Nature Physics*

- Towards a non-equilibrium phase diagram of cuprates

F. Cilento et al., *Nature Communications* 5:4353 (2014)



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Hunting for fast details

A movie in 4 frames:

<http://moviesinframes.tumblr.com>



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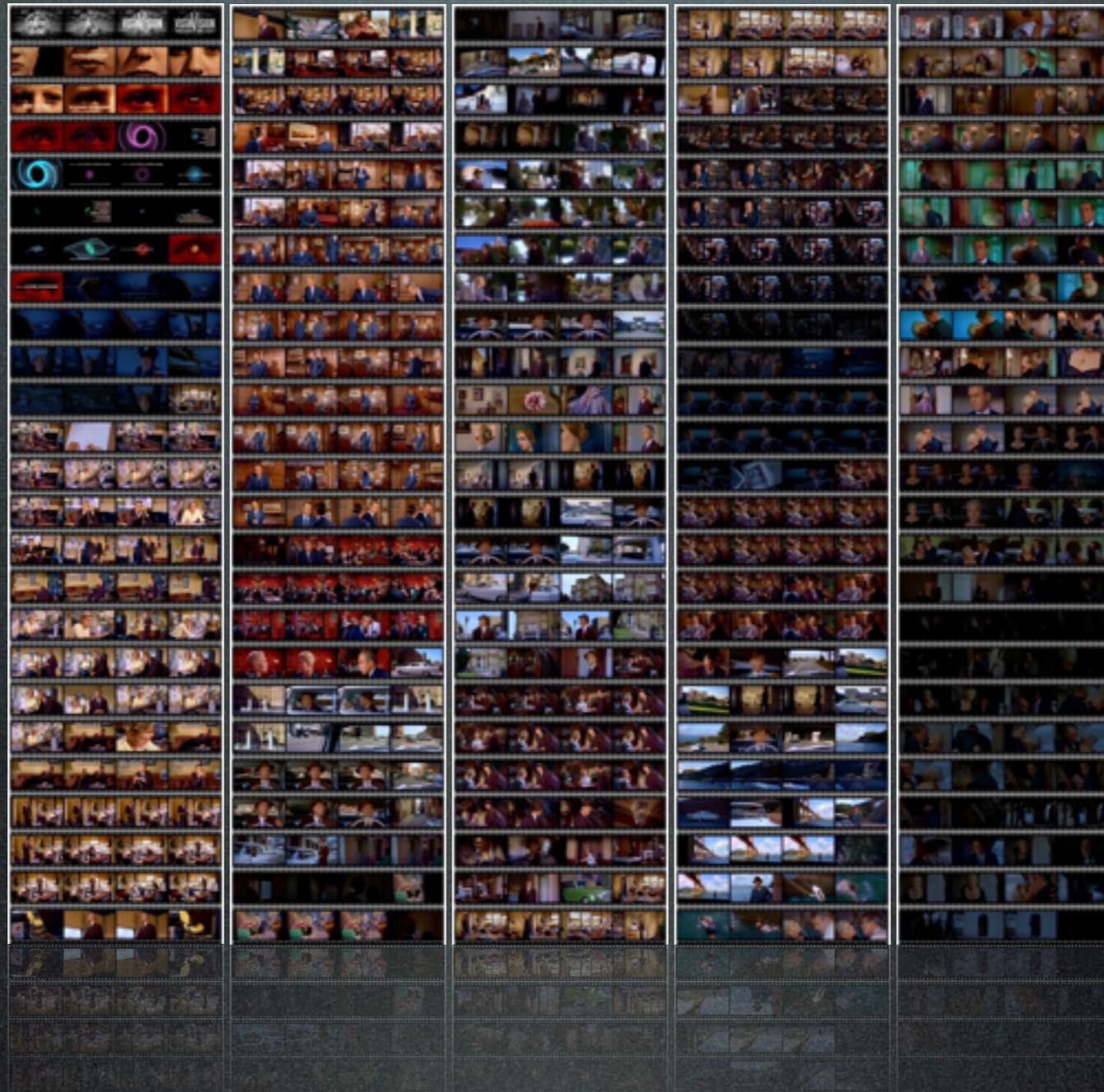


Which sequence is
Inglourious Basterds from Q. Tarantino?



Time-resolved optics

A movie in 1000 frames:

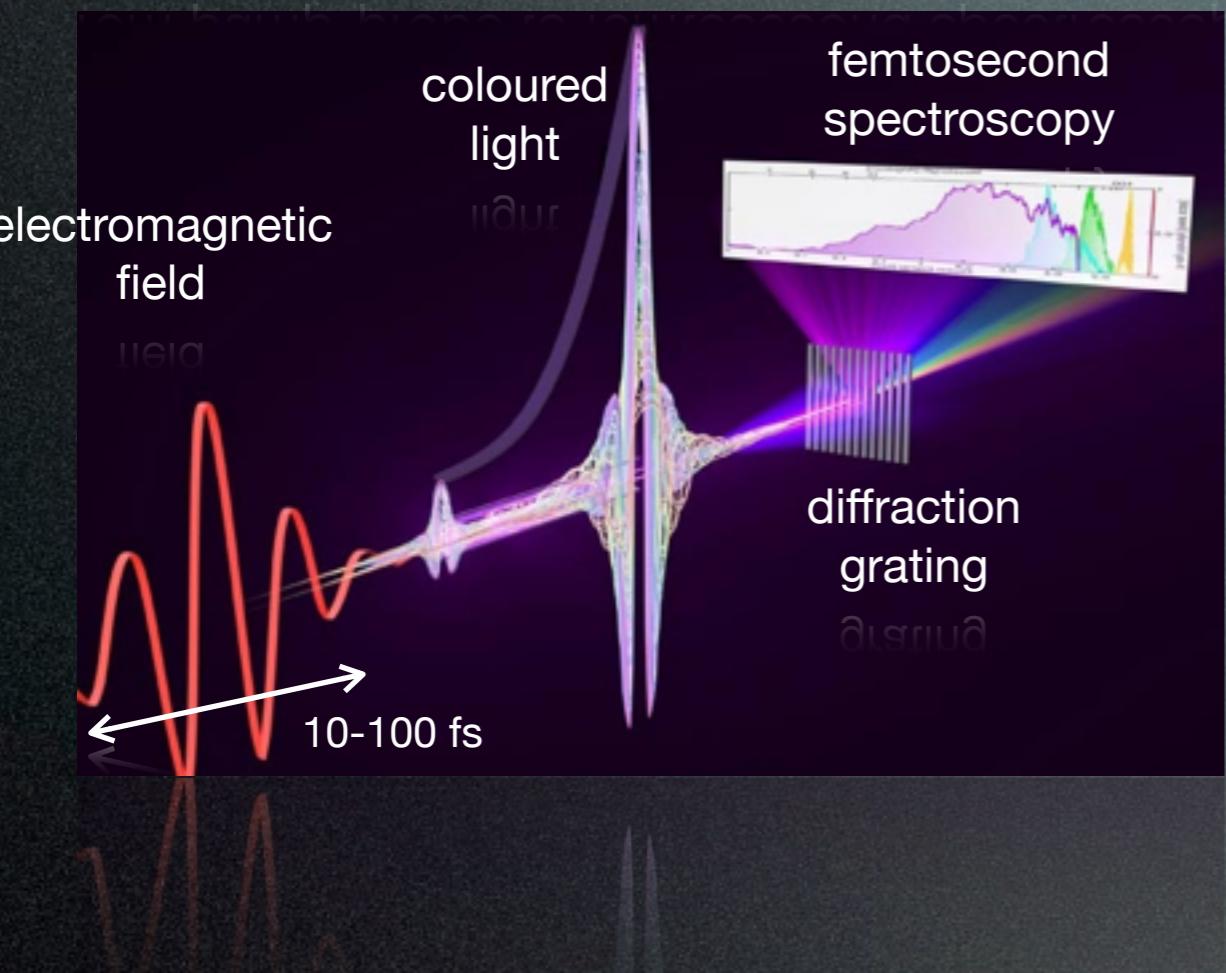


ultrafast shutter for ultrafast details...



Time-resolved optical spectroscopy

from pump-probe to femtosecond spectroscopy

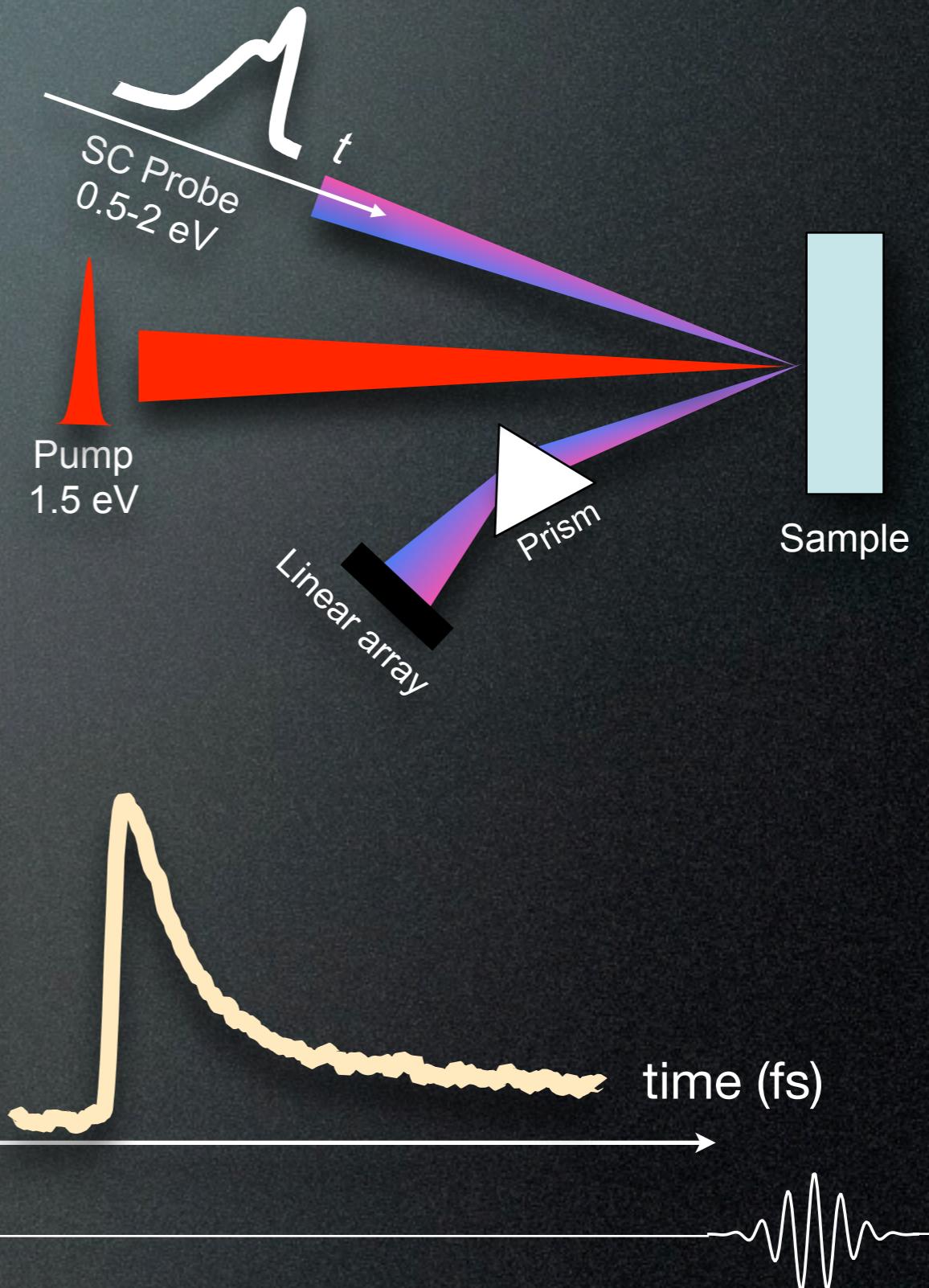
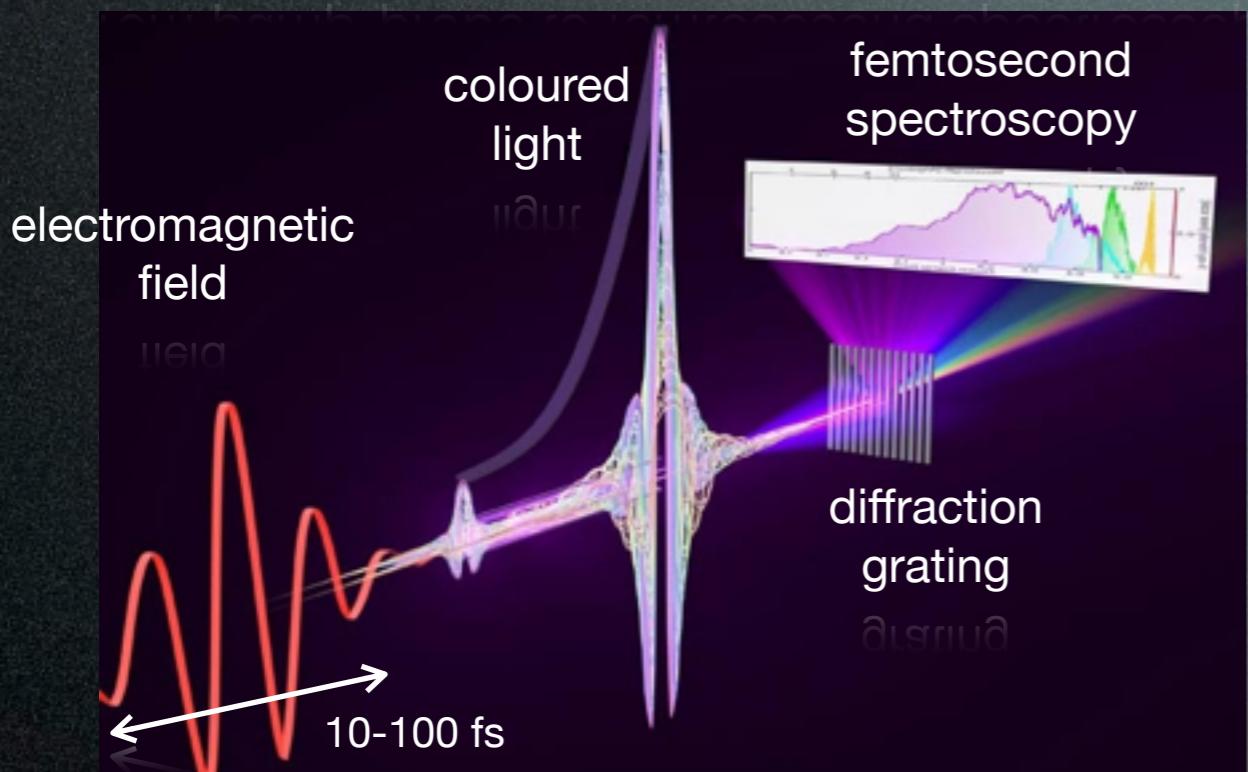


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Time-resolved optical spectroscopy

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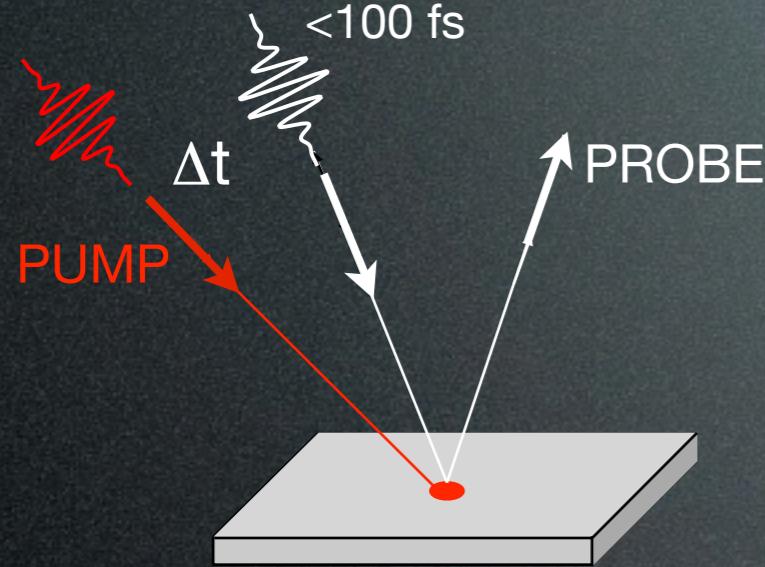


$$\frac{\delta R}{R}(\omega, t) = \frac{R_{exc}(\omega, t) - R_{eq}(\omega)}{R_{eq}(\omega)}$$

time+frequency information

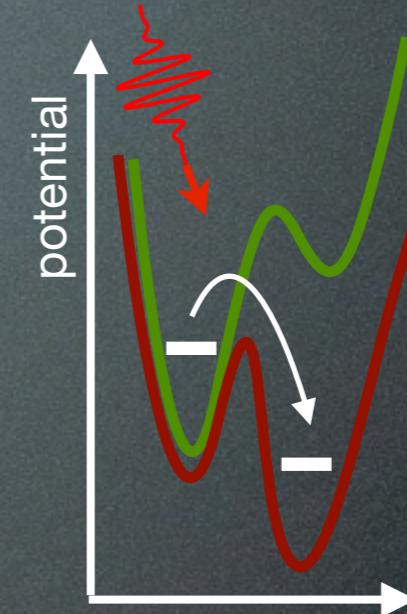
Time-resolved optical spectroscopy

pump probe on quantum materials



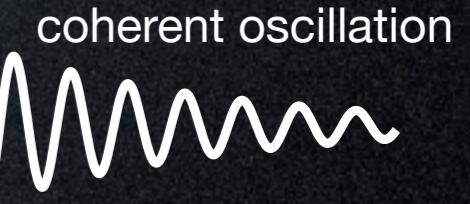
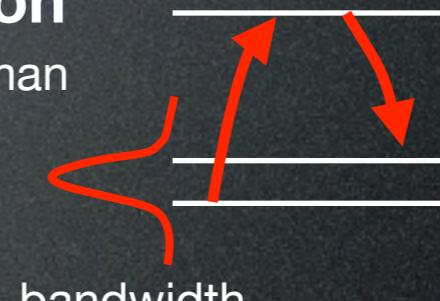
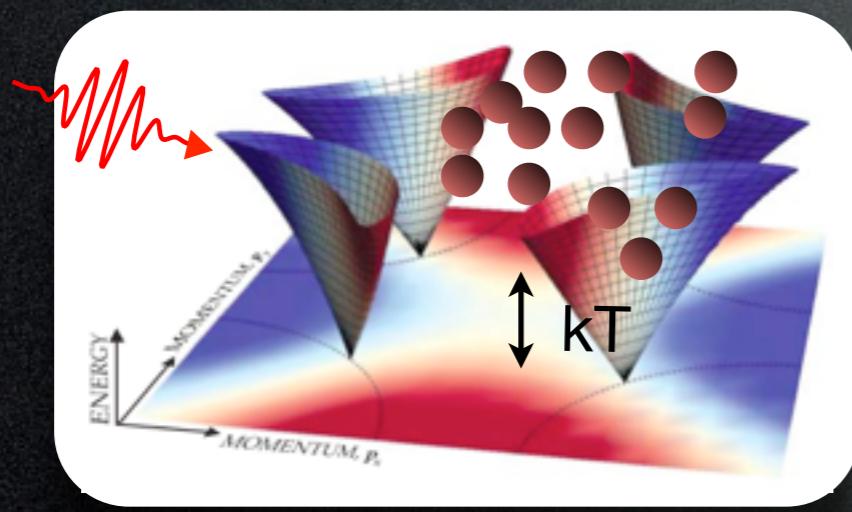
J. Orenstein, *Phys. Today* 65, 44 (2012)

optical control
new transient ground state



coherent excitation

impulsive Raman

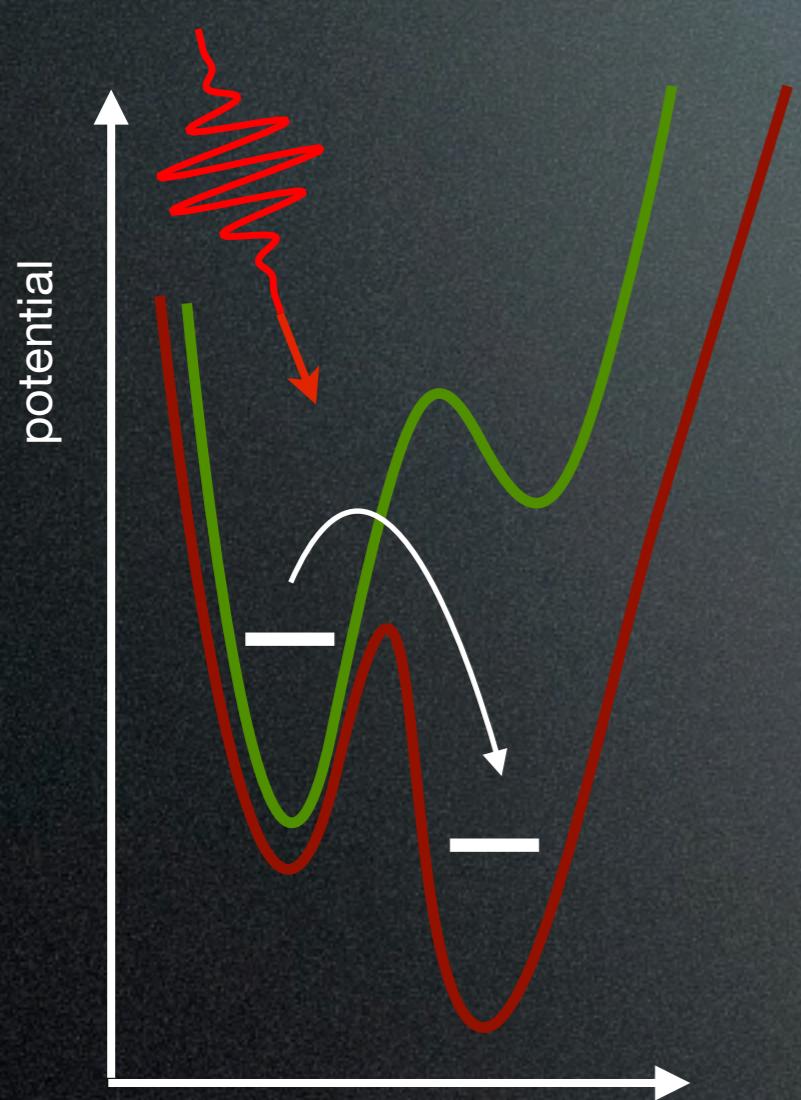


non-equilibrium spectroscopy

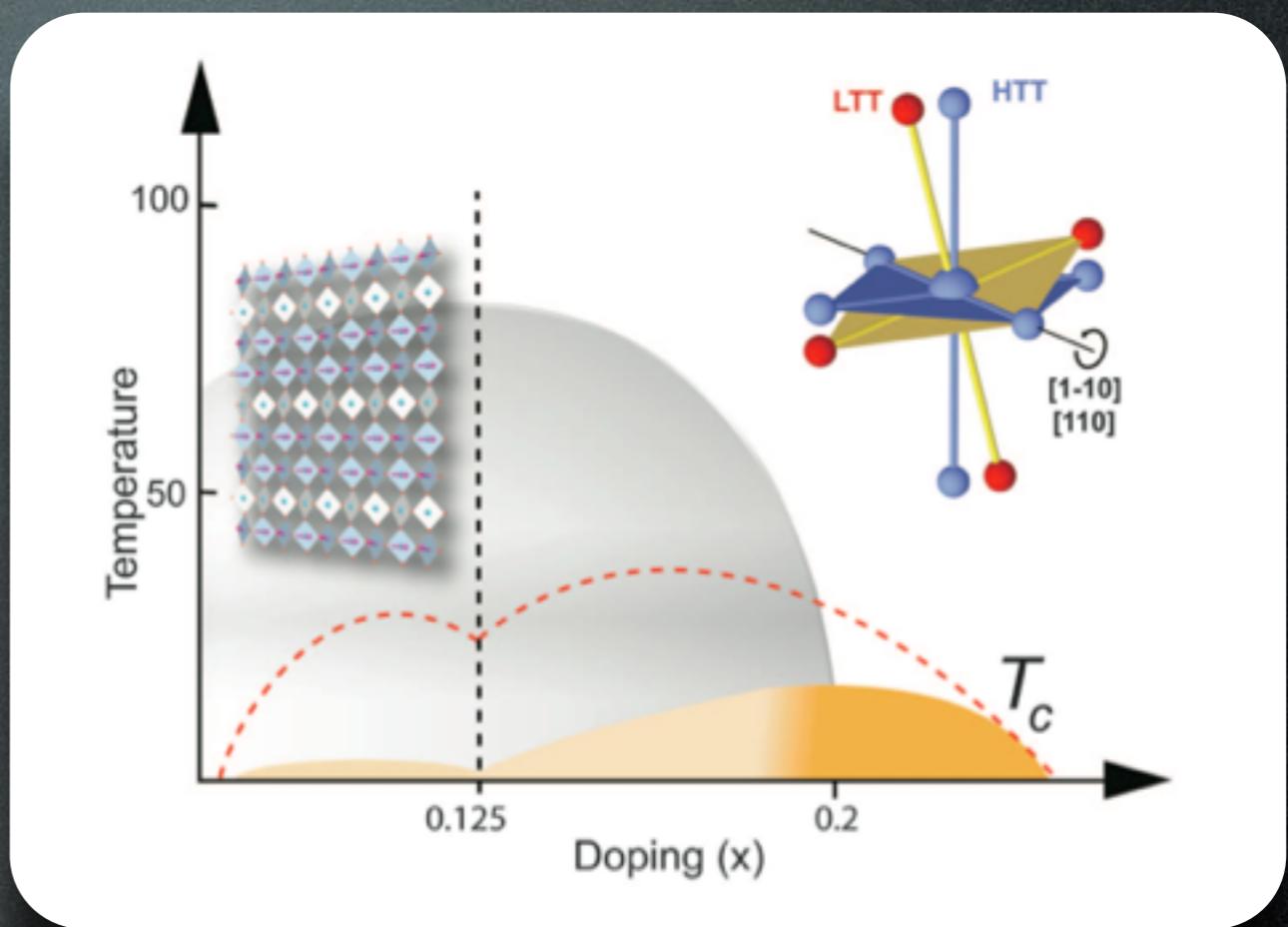
recovery of the ground state



Optical control



Transient superconductivity by removing a competing order



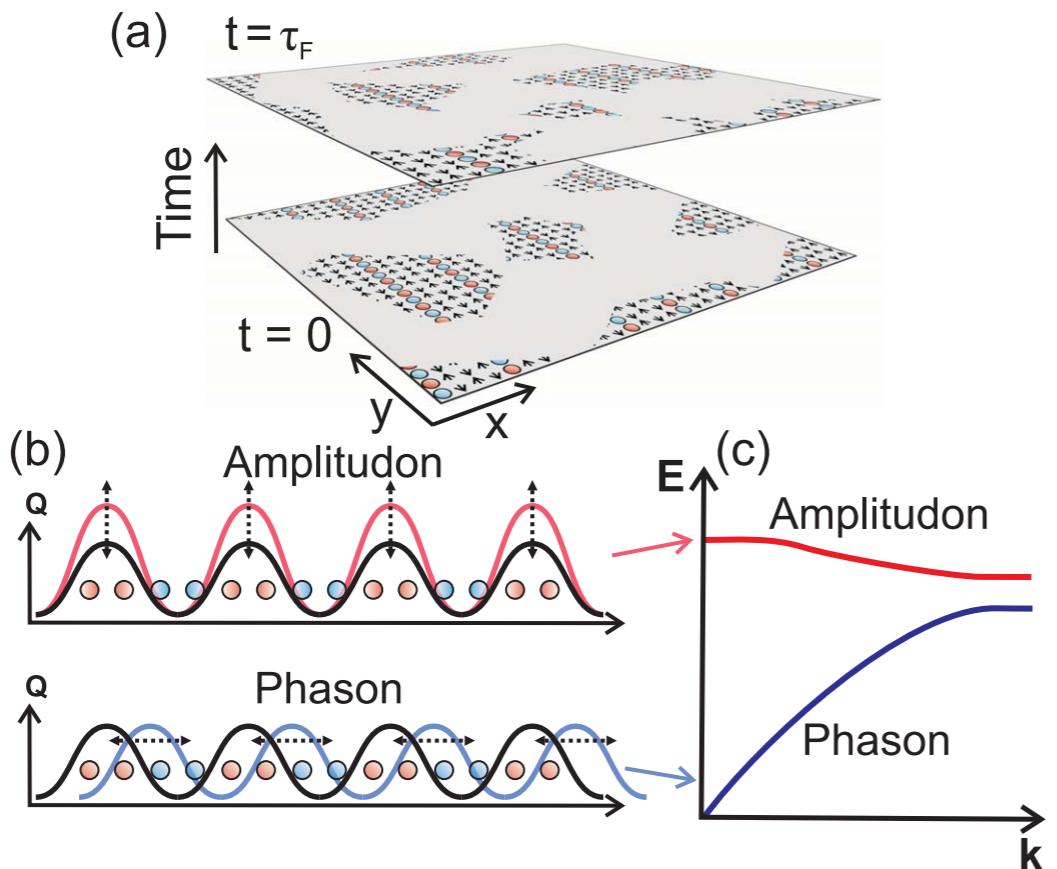
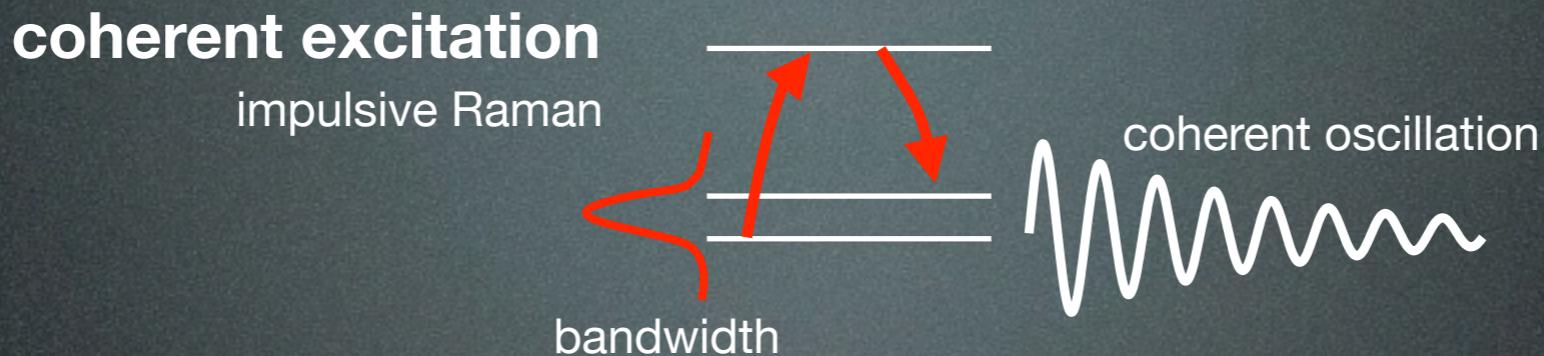
D. Fausti et al., *Science* **331**, 189 (2011)



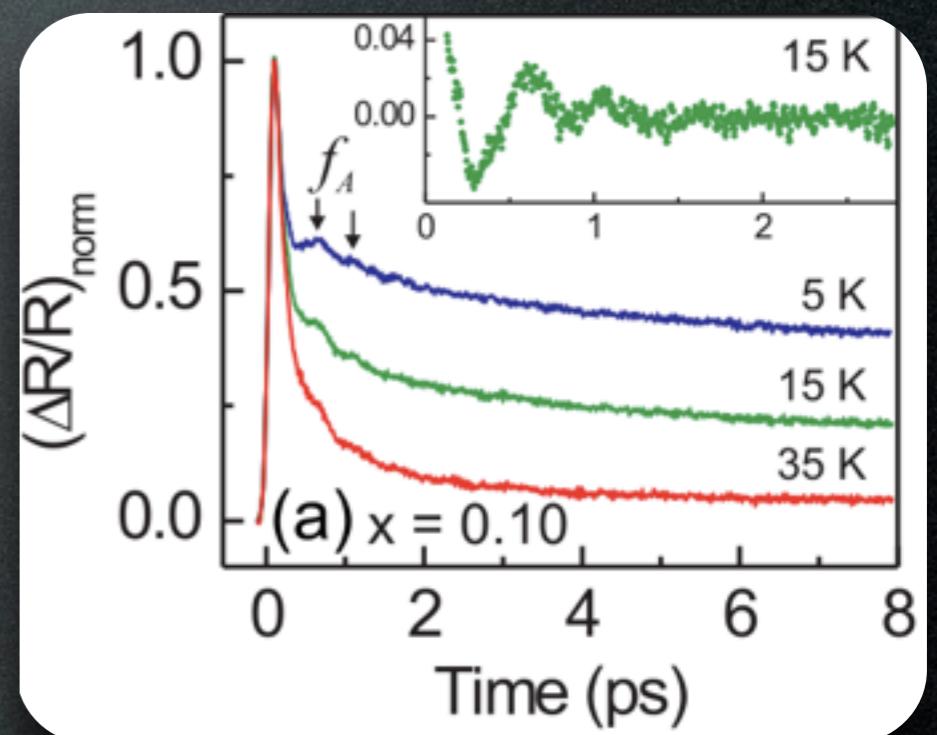
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Coherent excitation



Fluctuating Charge-Density-Wave
in UD $\text{La}_{1.9}\text{Sr}_{0.1}\text{CuO}_4$



D. Torchinsky et al, *Nature Materials* 2013

see talks: N. Gedik and F. Carbone

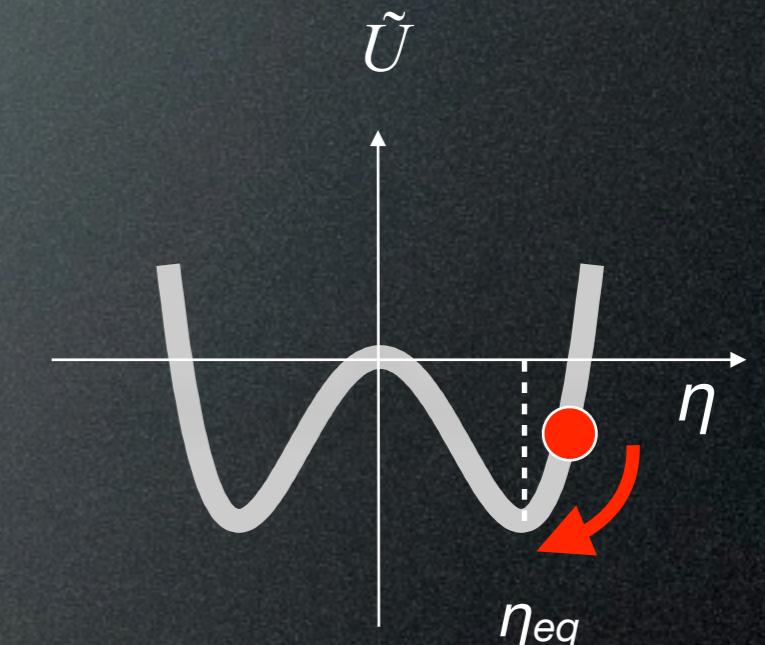
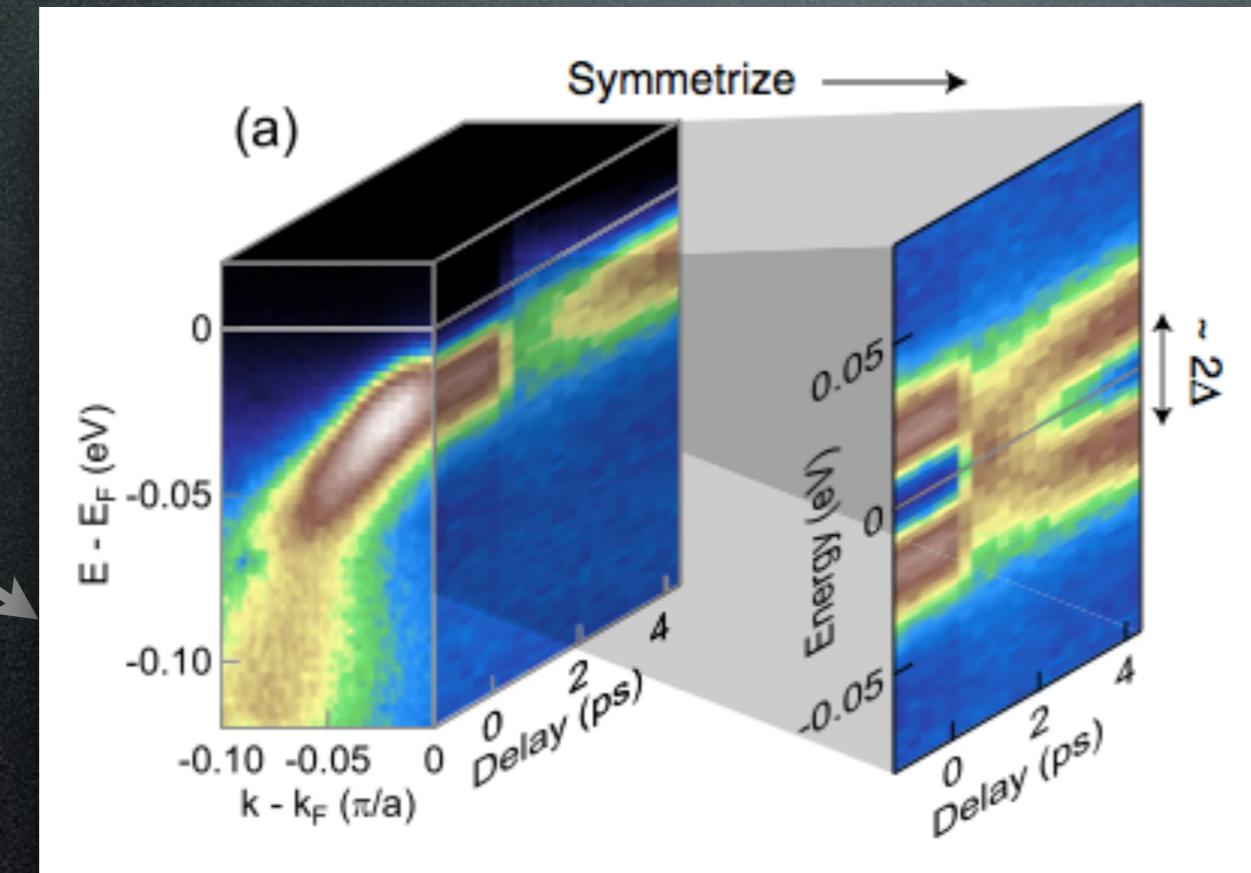


Non-equilibrium spectroscopy

disentangling intertwined degrees of freedom by their dynamics

Gap dynamics in superconducting copper oxides

Time-resolved
photoemission
spectroscopy



C. Smallwood et al. et al, *Phys. Rev. B* **89**, 115126 (2014)

see A. Lanzara's talk



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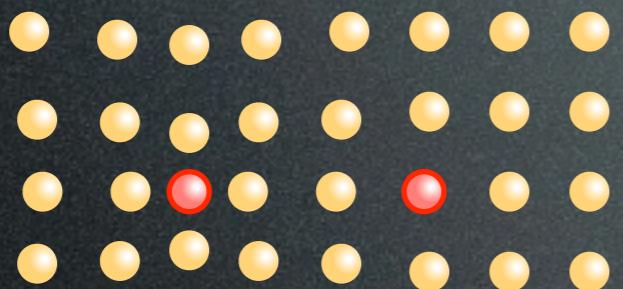


retarded vs instantaneous interaction

BCS

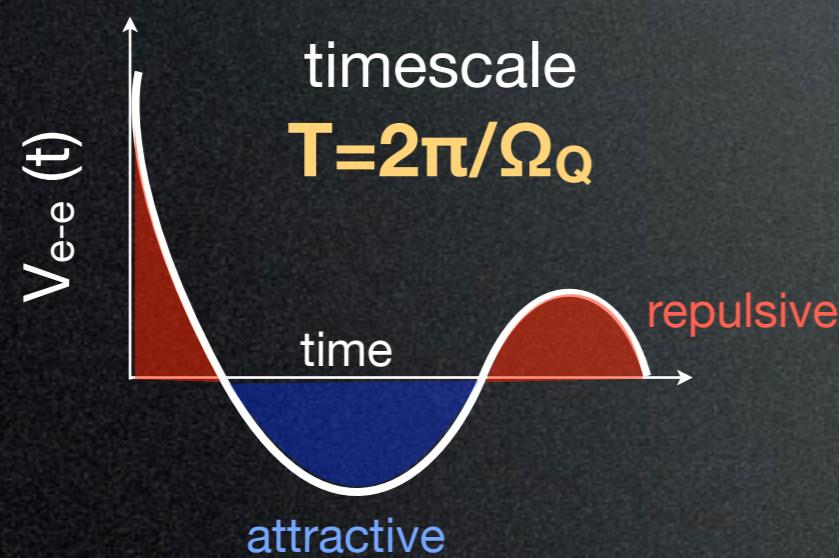
$$\hat{H} = \sum_{\mathbf{k}} \epsilon_{\mathbf{k}} \hat{c}_{\mathbf{k}}^\dagger \hat{c}_{\mathbf{k}} + \sum_{\mathbf{k}, \mathbf{Q}} M_Q \hat{c}_{\mathbf{k}+\mathbf{Q}}^\dagger c_{\mathbf{k}} (\hat{b}_{\mathbf{Q}} + \hat{b}_{-\mathbf{Q}}^\dagger)$$

glue: lattice distortion



space

$$\lambda = 2\pi/|Q|$$



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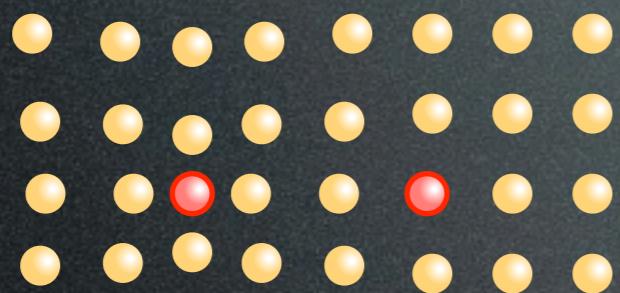
retarded vs instantaneous interaction

BCS

$$\hat{H} = \sum_{\mathbf{k}} \epsilon_{\mathbf{k}} \hat{c}_{\mathbf{k}}^\dagger \hat{c}_{\mathbf{k}} + \sum_{\mathbf{k}, \mathbf{Q}} M_Q \hat{c}_{\mathbf{k}+\mathbf{Q}}^\dagger c_{\mathbf{k}} (\hat{b}_{\mathbf{Q}} + \hat{b}_{-\mathbf{Q}}^\dagger)$$

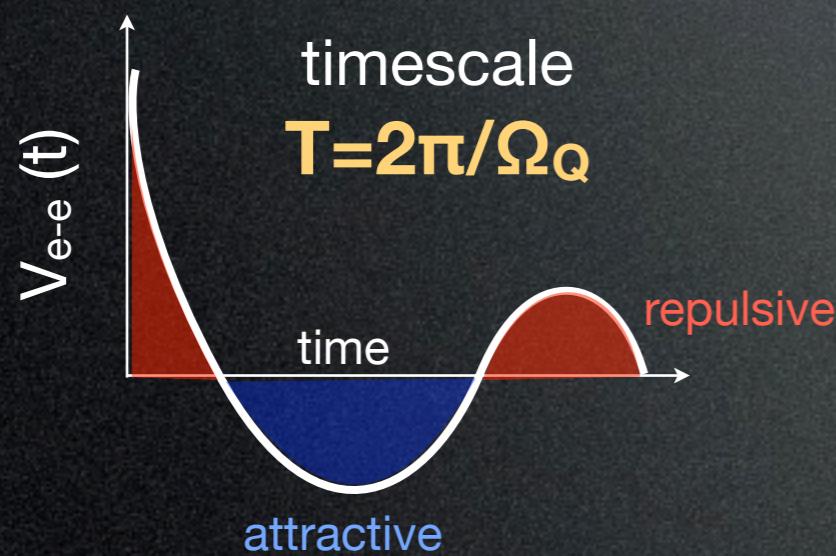
boson

glue: lattice distortion



space

$$\lambda = 2\pi/|Q|$$



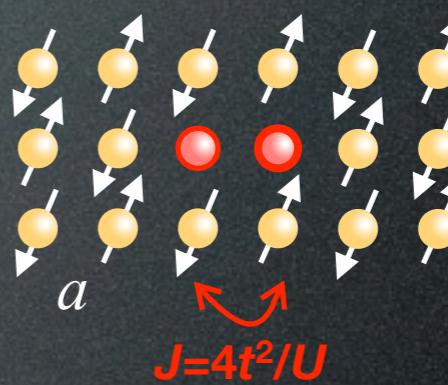
Hubbard model

$$\hat{H} = - \sum_{i,j,\sigma} (t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + c.c.) + U \sum_i \hat{n}_{i,\uparrow} \hat{n}_{i,\downarrow} - \mu \sum_i \hat{n}_i$$

space

d-wave: other glue?

$$\lambda = a \ (\approx 4 \text{ \AA})$$



timescales

$$\hbar/t_{ij} \ (\approx 2 \text{ fs})$$

vs

$$\hbar/J \ (\approx 6 \text{ fs})$$

vs

$$\hbar/U \ (\approx 0.3 \text{ fs})$$



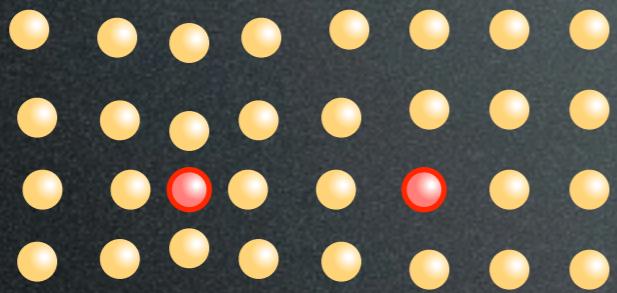
retarded vs instantaneous interaction

BCS

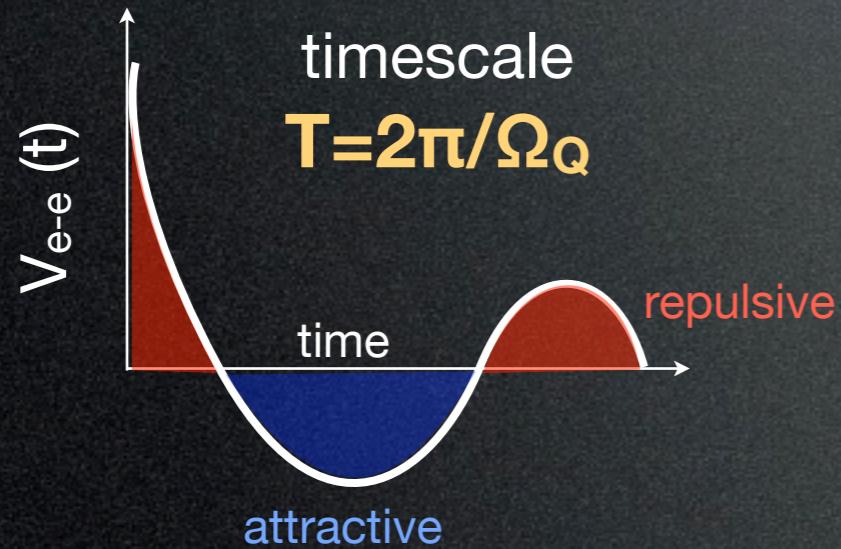
$$\hat{H} = \sum_{\mathbf{k}} \epsilon_{\mathbf{k}} \hat{c}_{\mathbf{k}}^\dagger \hat{c}_{\mathbf{k}} + \sum_{\mathbf{k}, \mathbf{Q}} M_Q \hat{c}_{\mathbf{k}+\mathbf{Q}}^\dagger c_{\mathbf{k}} (\hat{b}_{\mathbf{Q}} + \hat{b}_{-\mathbf{Q}}^\dagger)$$

boson

glue: lattice distortion



space
 $\lambda=2\pi/|Q|$



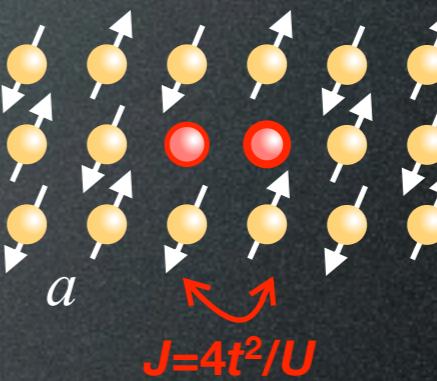
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space

d-wave: other glue?

$\lambda=a$ (≈ 4 Å)



\hbar/t_{ij} (≈ 2 fs)

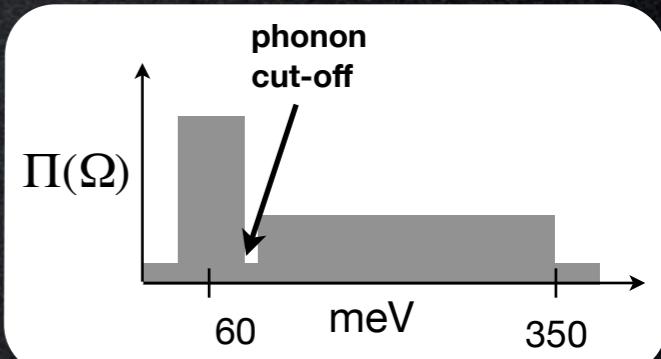
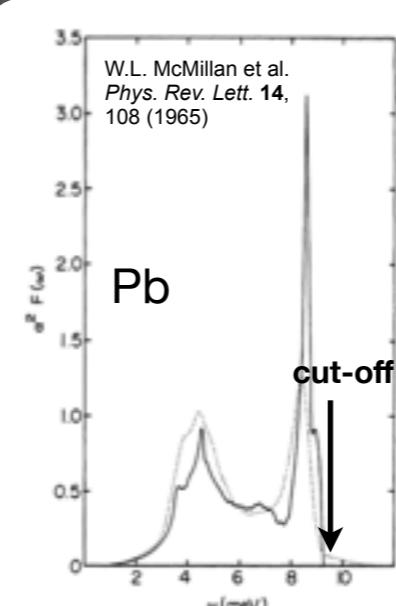
vs

\hbar/J (≈ 6 fs)

vs

\hbar/U (≈ 0.3 fs)

$$\Pi(\Omega) = \alpha^2 F(\Omega) + I^2 \chi(\Omega)$$



J. Carbotte et al. Rep. Prog. Phys. **74**, 066501 (2011)



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Single-colour pump-probe experiments and electron-phonon coupling

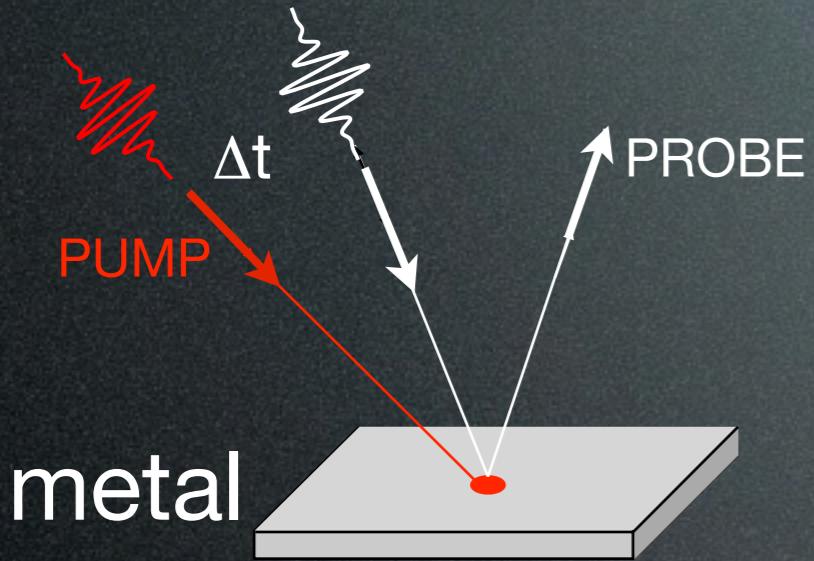


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non-equilibrium reflectivity and e-ph coupling

pump probe on metals



2-temperature model

$$\frac{\partial T_e}{\partial t} = \frac{G(\Pi_{lat}, T_{lat}, T_e)}{\gamma_e T_e} + \frac{p}{\gamma_e T_e}$$

$$\frac{\partial T_{lat}}{\partial t} = -\frac{G(\Pi_{lat}, T_{lat}, T_e)}{C_{lat}}$$

PUMP

G/C

determines the
dynamics in the
time domain

P.B. Allen, *Phys. Rev Lett.* **59** 1460 (1987)

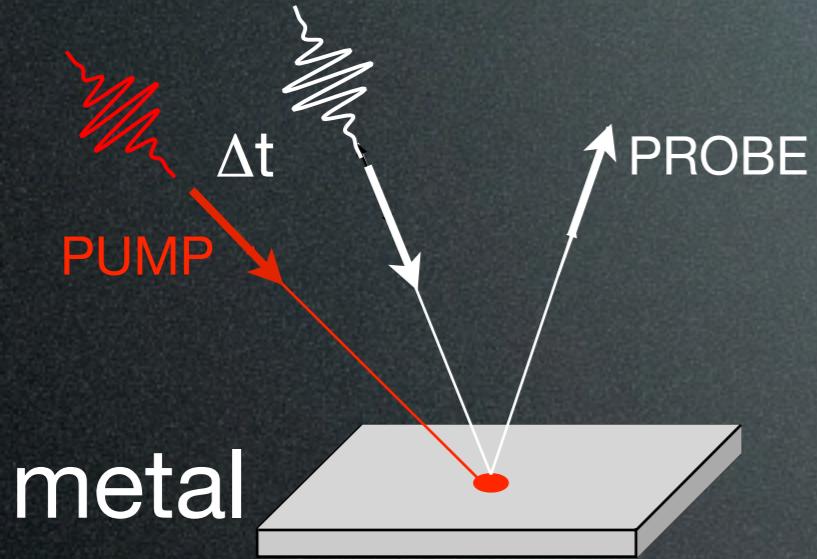


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$$\frac{\partial T_e}{\partial t} = \frac{G(\Pi_{lat}, T_{lat}, T_e)}{\gamma_e T_e} + \frac{p}{\gamma_e T_e}$$

G/C

$$\frac{\partial T_{lat}}{\partial t} = -\frac{G(\Pi_{lat}, T_{lat}, T_e)}{C_{lat}}$$

determines the dynamics in the time domain

$$G = \frac{3\gamma_e}{\pi\hbar k_B^2} \int_0^\infty d\Omega \Pi_{lat}(\Omega) [N(\Omega, T_{lat}) - N(\Omega, T_e)]$$

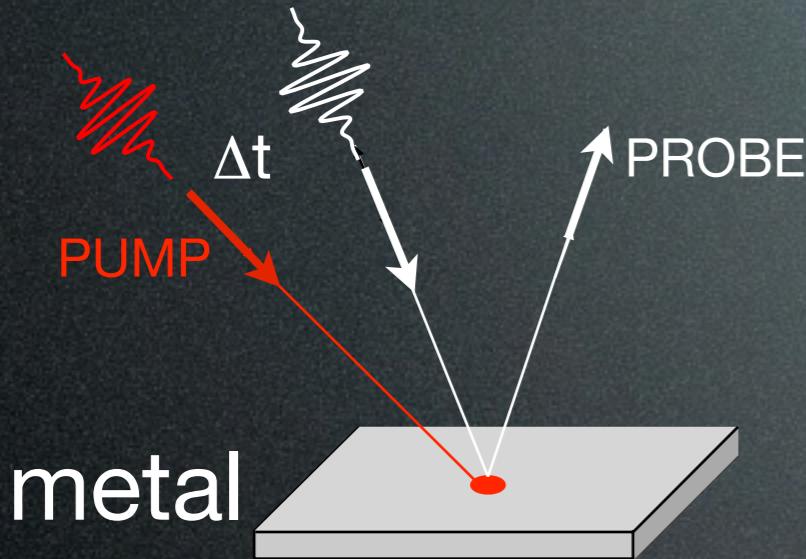


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non-equilibrium reflectivity and e-ph coupling

pump probe on metals



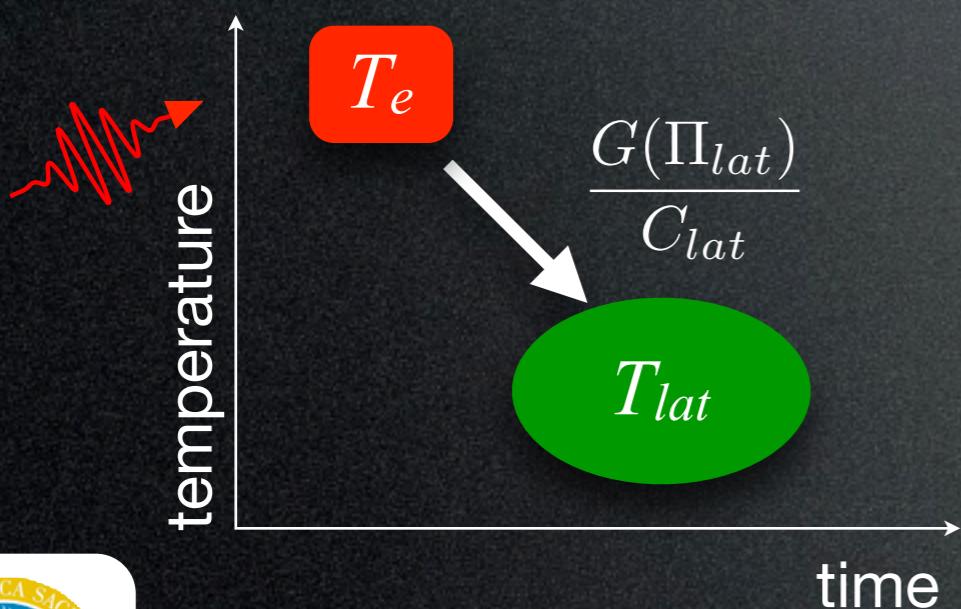
2-temperature model

$$\gamma_e T_e \frac{\partial T_e}{\partial t} = -G \cdot (T_e - T_{lat}) + p(t)$$

$$C_{lat} \frac{\partial T_{lat}}{\partial t} = G \cdot (T_e - T_{lat})$$

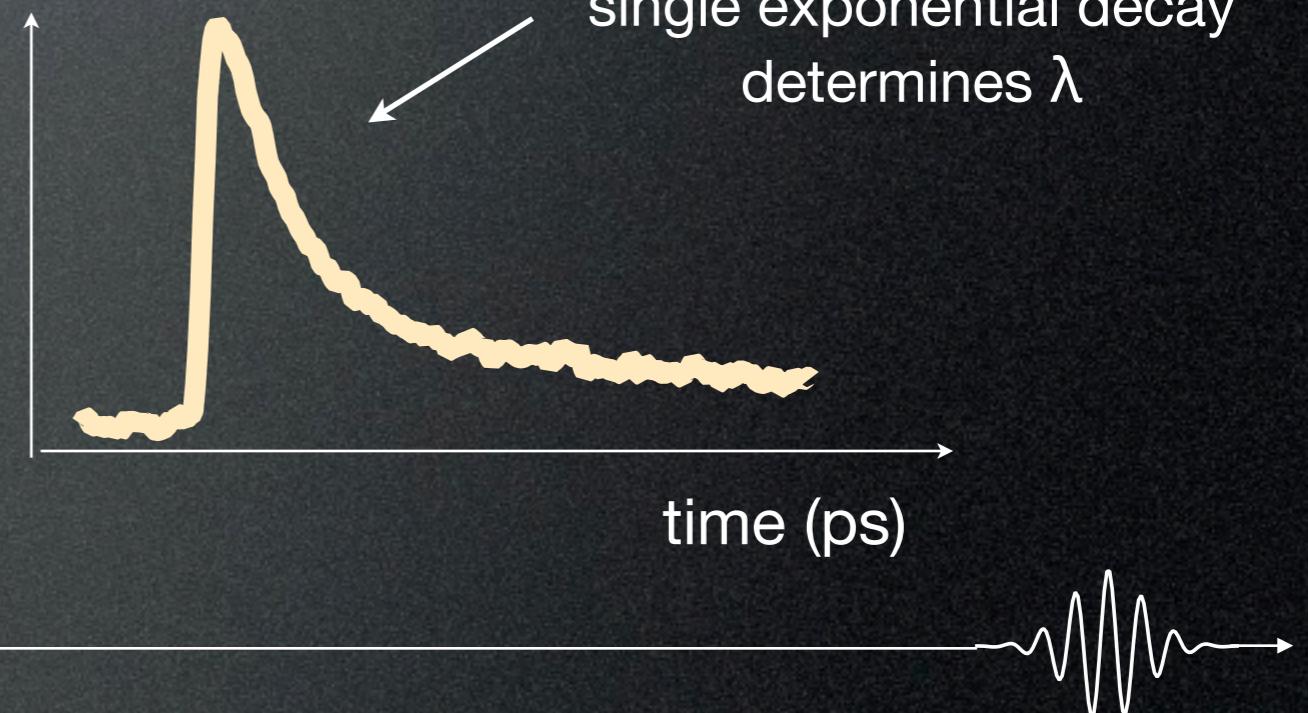
PUMP
G/C
determines the dynamics in the time domain

hierarchy in the dynamics



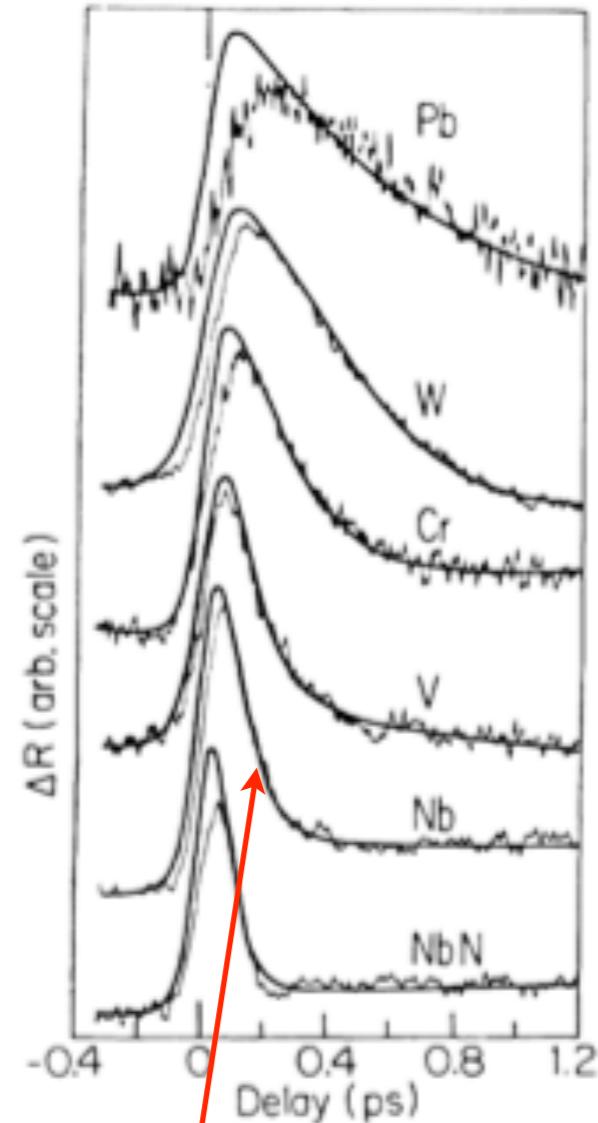
$$\delta R/R \propto \delta T_e/T_e$$

in conventional metals:
single exponential decay
determines λ



non-equilibrium spectroscopy on metals

electron-phonon coupling in metals



single exponential decay

	$T_e(0)$ (K) ^a	$\lambda_{\text{exp}}\langle\omega^2\rangle$ (meV ²)	$\langle\omega^2\rangle$ (meV ²)	λ_{exp}	λ_{int}
Cu	590	29 ± 4	377 ^b	0.08 ± 0.01	0.10 ^b
Au	650	23 ± 4	178 ^c	0.13 ± 0.02	0.15 ^c
Cr	716	128 ± 15	987 ^d	0.13 ± 0.02	...
W	1200	112 ± 15	425 ^e	0.26 ± 0.04	0.26 ^e
V	700	280 ± 20	352 ^f	0.80 ± 0.06	0.82 ^f
Nb	790	320 ± 30	275 ^g	1.16 ± 0.11	1.04 ^g
Ti	820	350 ± 30	601 ^g	0.58 ± 0.05	0.54 ^g
Pb	570	45 ± 5	31 ^h	1.45 ± 0.16	1.55 ^h
NbN	1070	640 ± 40	673 ^j	0.95 ± 0.06	1.46 ^j
V ₃ Ga	1110	370 ± 60	448 ^k	0.83 ± 0.13	1.12 ^k

S.D. Brorson et al. *Phys. Rev. Lett.* **64**, 2172 (1990)

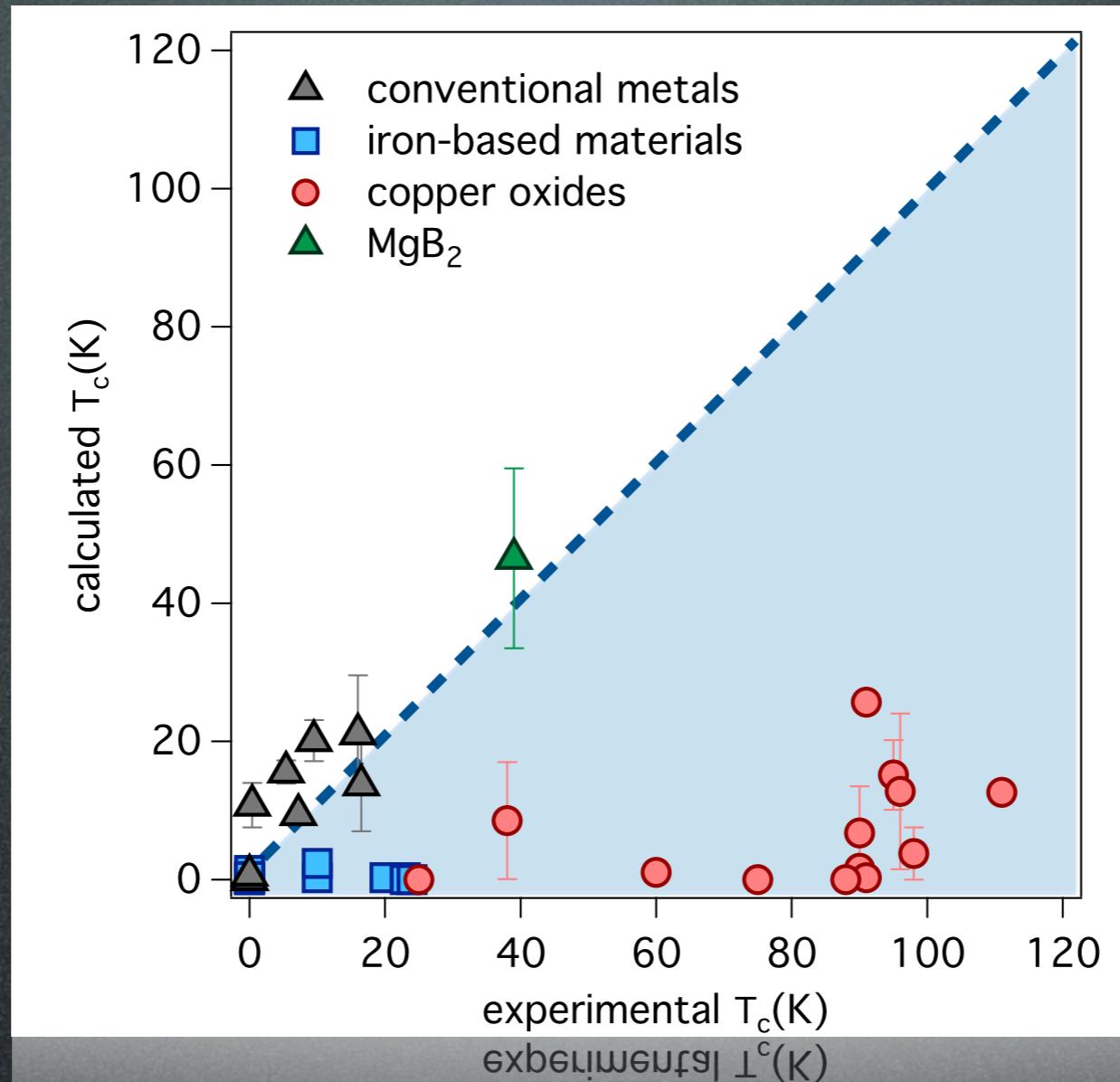
$$\tau_{e-lat} = \frac{\pi k_B^2 T_e}{3\hbar\lambda\langle\Omega^2\rangle}$$

Eliashberg coupling $\lambda = 2 \int_0^\infty \alpha^2 F(\Omega)/\Omega d\Omega$



electron-phonon coupling in superconductors

electron-phonon coupling
 λ obtained from time-
 resolved techniques



CUPRATES:
 the glue is not in the
 electron-phonon
 coupling!!

μ^* : effective Coulomb repulsion
 $g=1$ for s-wave BCS superconductors
 $g<1$ for d-wave superconductors

$$T_c = 0.83 \tilde{\Omega} e^{-\frac{1.04(1+\lambda)}{g(\lambda - \mu^*(1+0.62\lambda))}}$$

P.B. Allen and R.C. Dynes, *Phys. Rev. B* **12**, 905 (1975)
 A. J. Millis, S. Sachdev, C. M. Varma, *Phys. Rev. B* **37**, 4975 (1988).

$\tilde{\Omega}$: frequency log-average



Beyond single-colour experiments to access the ultrafast dynamics



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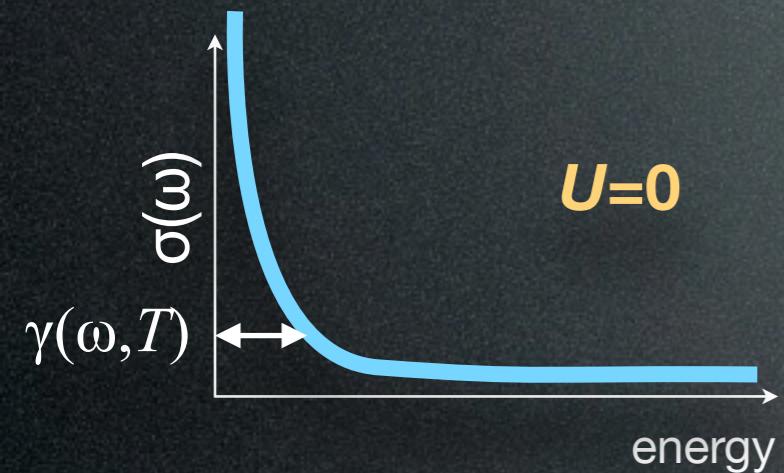


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Extended Drude Model

equilibrium optical conductivity

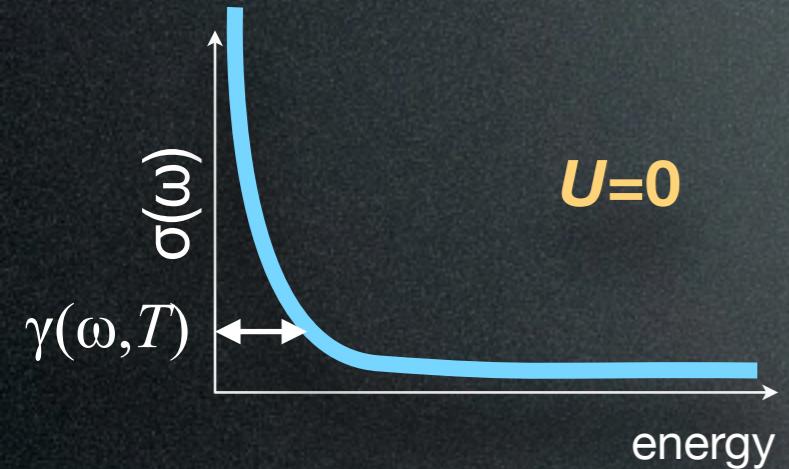


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Extended Drude Model

equilibrium optical conductivity



Drude dielectric function $\epsilon_D(\omega) = 1 - \frac{\omega_p^2}{\omega(\omega + i\tau^{-1})}$ τ constant

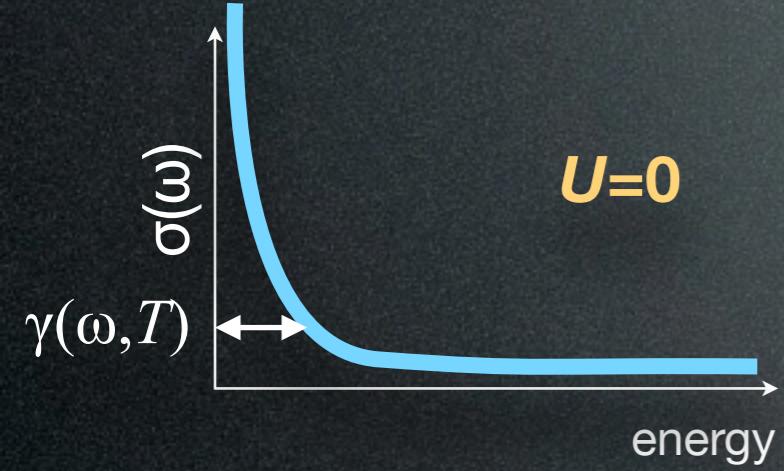


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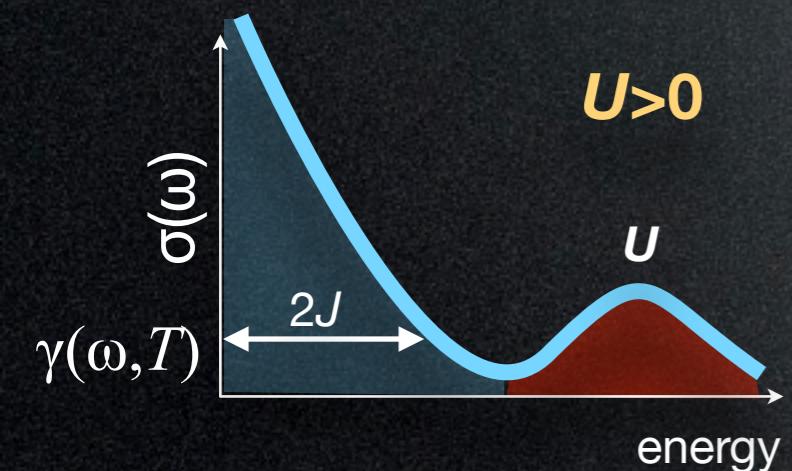
Extended Drude Model

equilibrium optical conductivity



Drude dielectric function $\epsilon_D(\omega) = 1 - \frac{\omega_p^2}{\omega(\omega + i\tau^{-1})}$ τ constant

Extended Drude dielectric function $\epsilon_D(\omega, T) = 1 - \frac{\omega_p^2}{\omega(\omega + M(\omega, T))}$ Memory function



Fermi-Dirac distributions

$$\frac{M(\omega)}{\omega} = \left\{ \int \frac{f(\xi) - f(\xi + \omega)}{\omega + \Sigma^*(\xi) - \Sigma(\xi + \omega)} d\xi \right\}^{-1} - 1$$

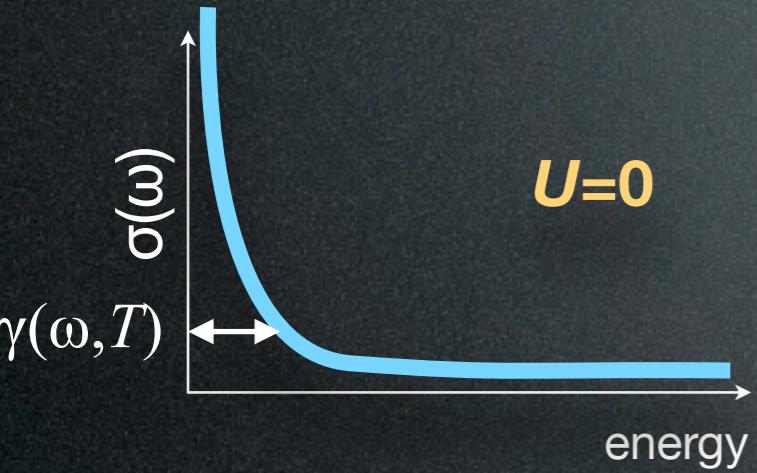


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Extended Drude Model

equilibrium optical conductivity

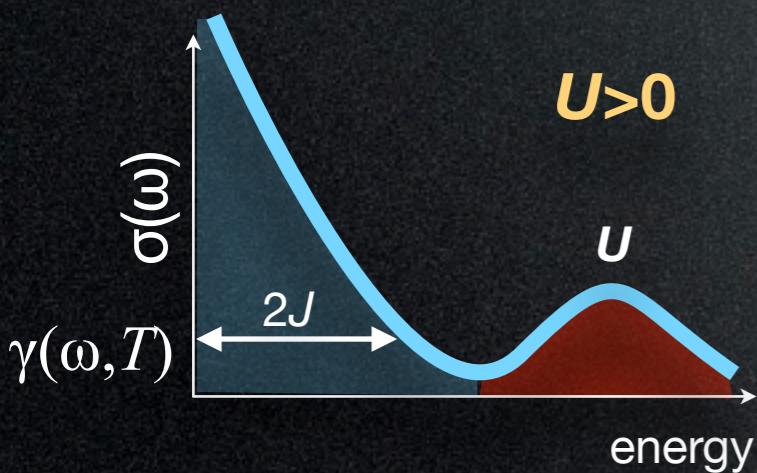


Drude dielectric function

$$\epsilon_D(\omega) = 1 - \frac{\omega_p^2}{\omega(\omega + i\tau^{-1})} \quad \tau \text{ constant}$$

Extended Drude dielectric function

$$\epsilon_D(\omega, T) = 1 - \frac{\omega_p^2}{\omega(\omega + M(\omega, T))} \quad \text{Memory function}$$



Fermi-Dirac distributions

$$\frac{M(\omega)}{\omega} = \left\{ \int \frac{f(\xi) - f(\xi + \omega)}{\omega + \Sigma^*(\xi) - \Sigma(\xi + \omega)} d\xi \right\}^{-1} - 1$$

single particle self-energy

$$\Sigma(\xi) = \int_0^\infty d\Omega L(\xi, \Omega; T) \Pi(\Omega)$$

bosonic function
 $\alpha^2 F(\Omega)$ in conventional metals

T-dep. kernel function
 $L(\xi, \Omega; T_e) + L(\xi, \Omega; T_b)$

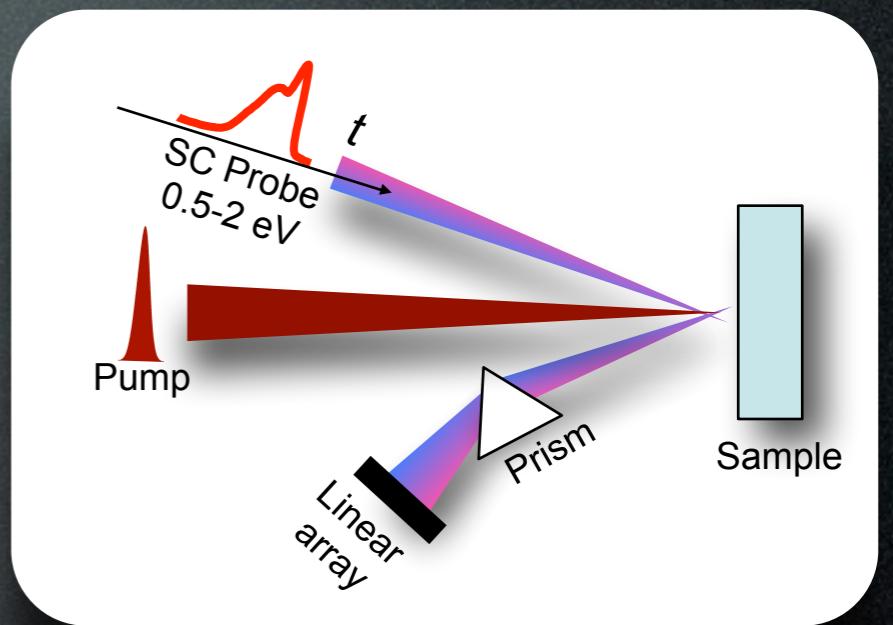
$$\Pi(\Omega) = \alpha^2 F(\Omega) + I^2 \chi(\Omega)$$



The “glue” problem in time-resolved optics

A) Is it really possible to describe the electron dynamics of doped cuprates in terms of charge carriers interacting with bosonic fluctuations?
(energy resolution)

B) On which timescale the charge carriers can exchange energy with bosons?
(temporal resolution)



$$\frac{\delta R}{R}(\omega, t) = \frac{R_{exc}(\omega, t) - R_{eq}(\omega)}{R_{eq}(\omega)}$$



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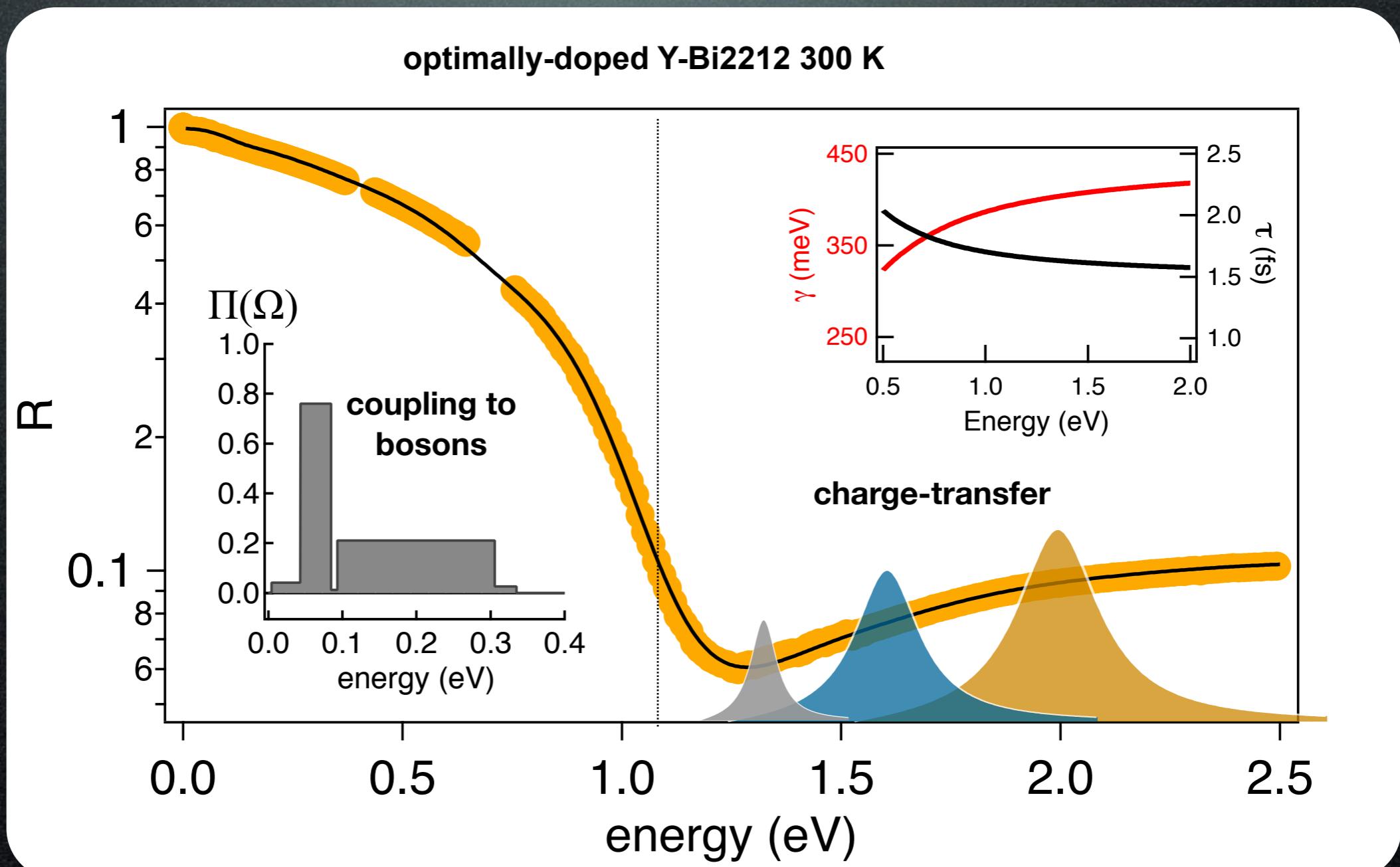


Equilibrium optical properties of cuprates

scattering rate

$$ImM(\omega, T) = \hbar/\tau(\omega, T)$$

D.N. Basov et al.
Rev. Mod. Phys. **85**, 471 (2011)



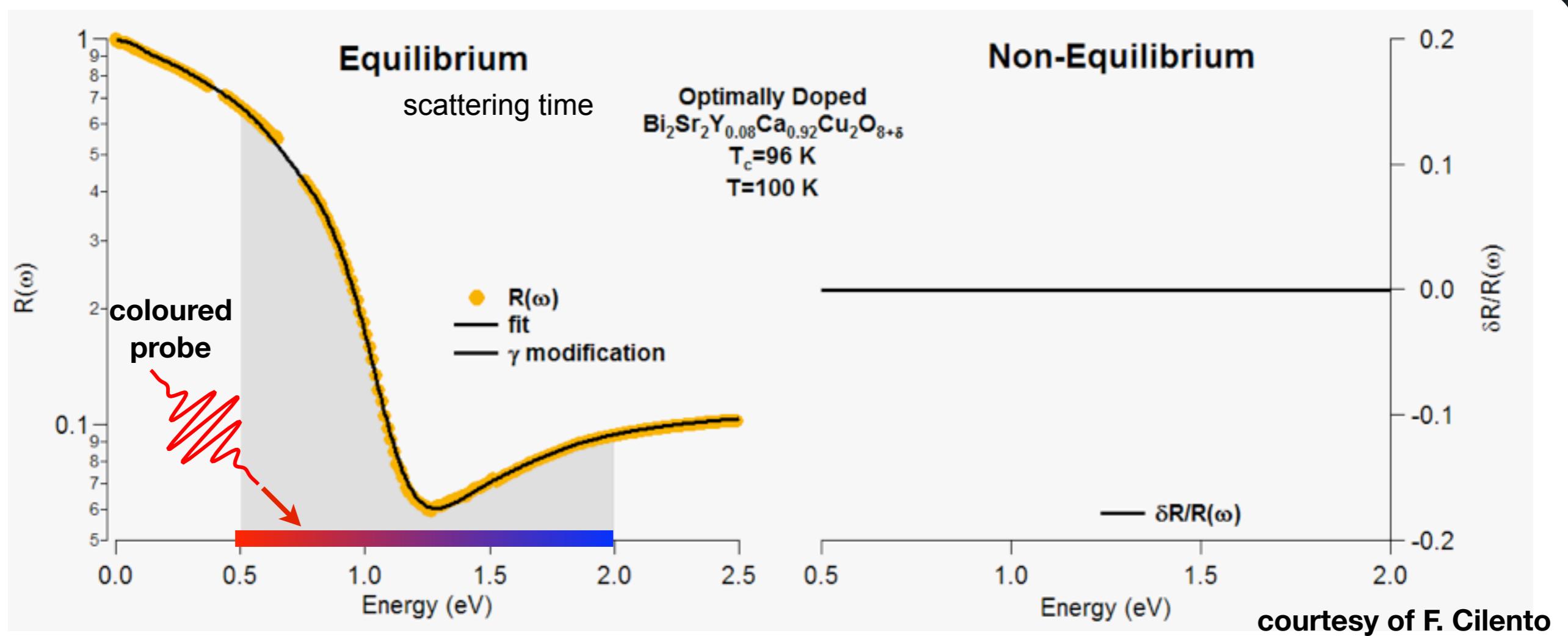
E. van Heumen et al., *Phys. Rev. B* **79**, 184512 (2009)

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NON-equilibrium optical properties of cuprates

effective e-bos scenario: change of scattering rate $\gamma(\omega, T)$



isosbestic point at 1.1 eV

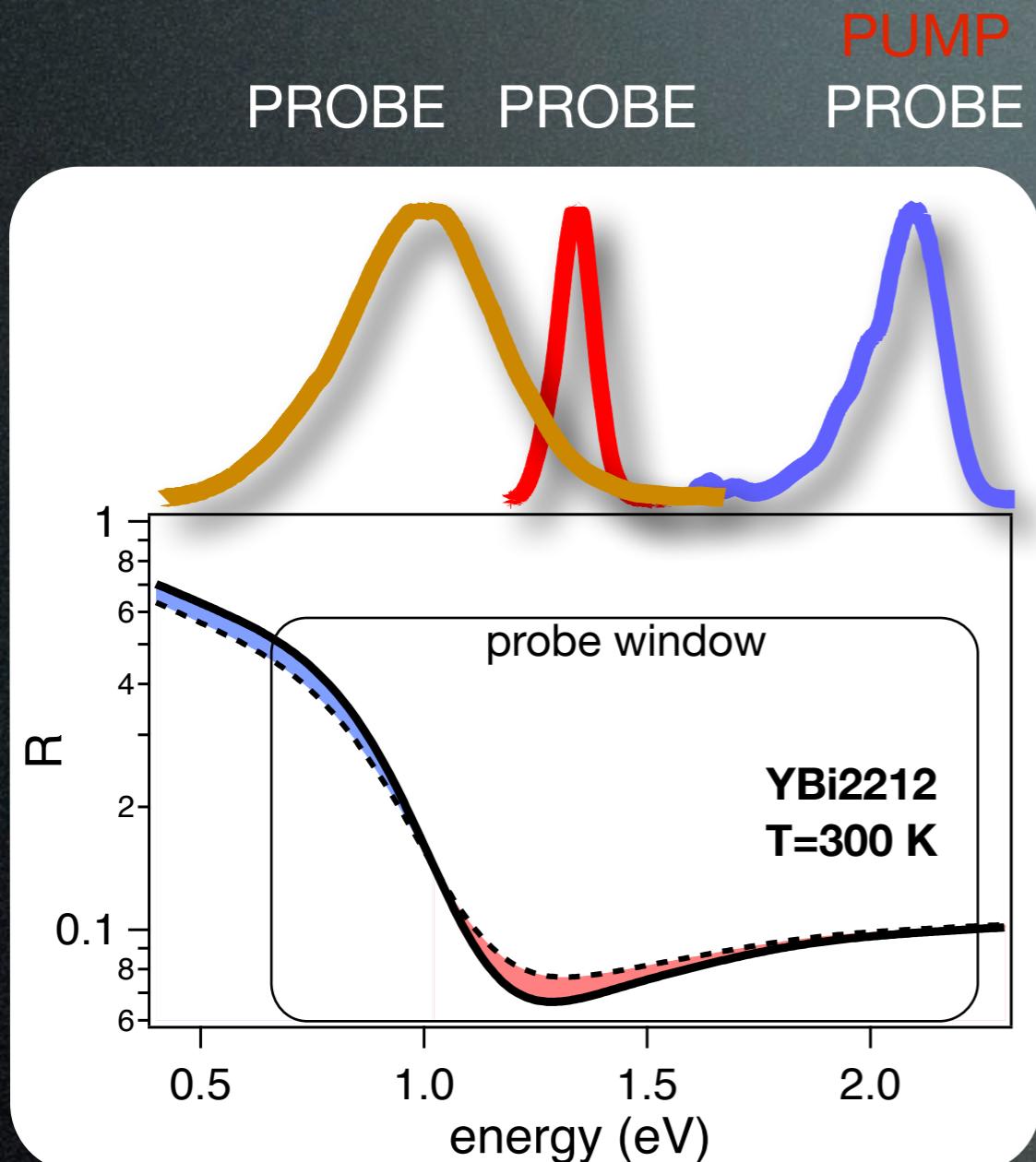
$$\delta R(\omega, \gamma) = \frac{\partial R}{\partial \gamma}(\omega) \delta \gamma$$

$$\frac{\partial R}{\partial \gamma}(1.1) = 0$$

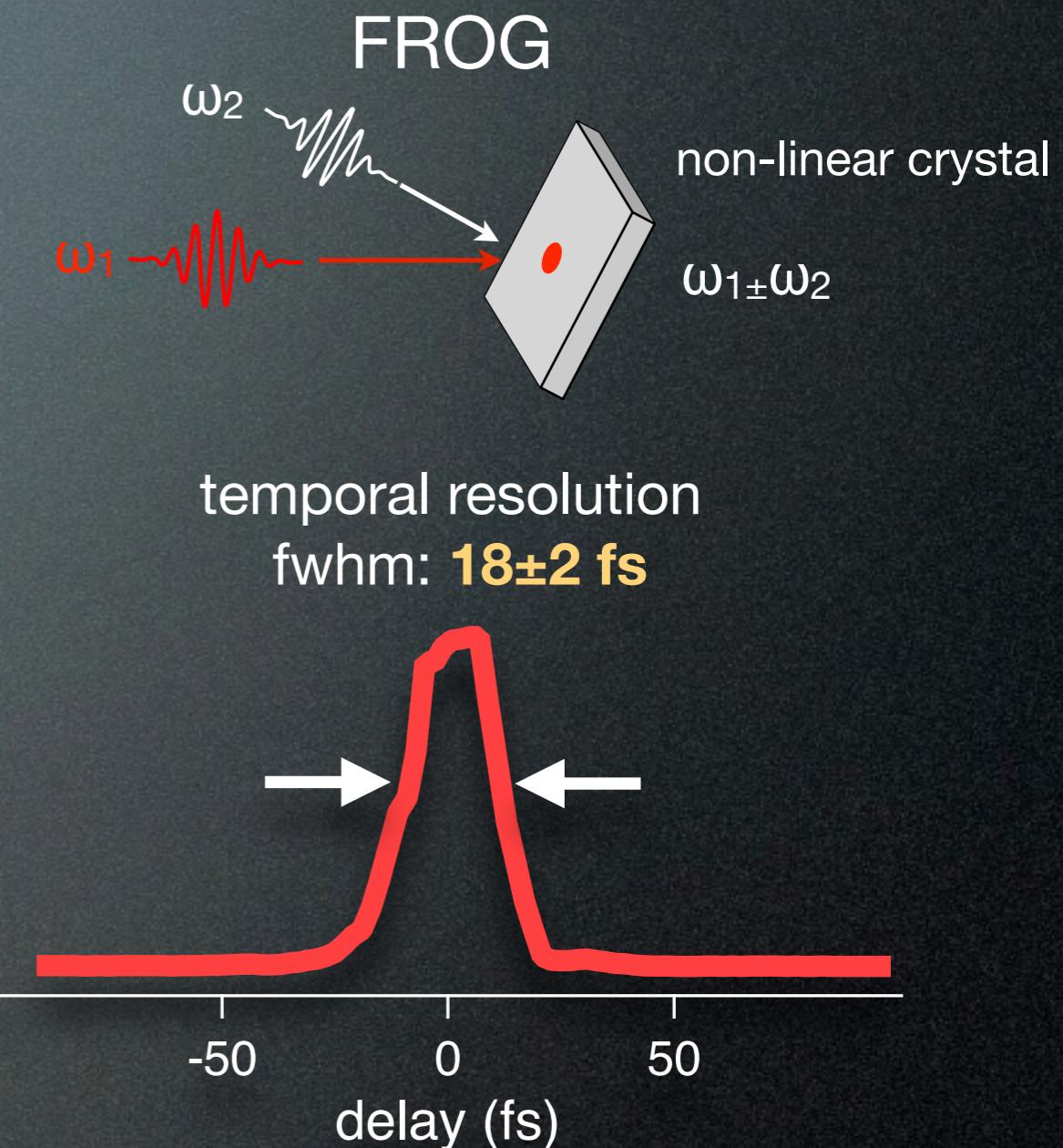
M. Greger et al. *Phys. Rev. B* **87**, 195140 (2013)



ultra high temporal resolution

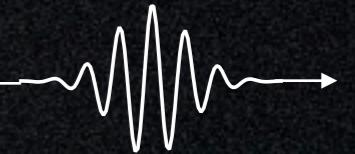


G. Cerullo's group (Politecnico of Milan)

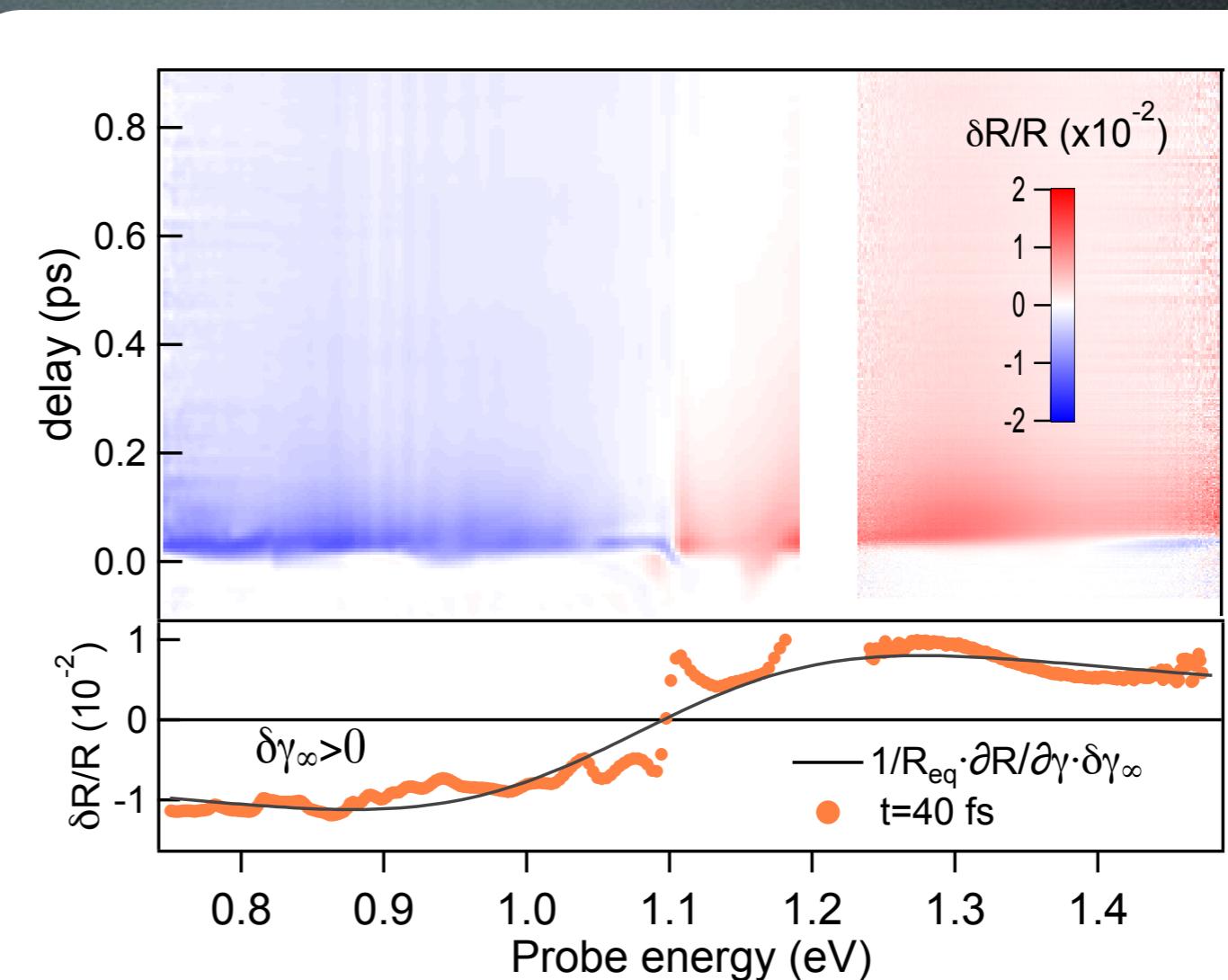
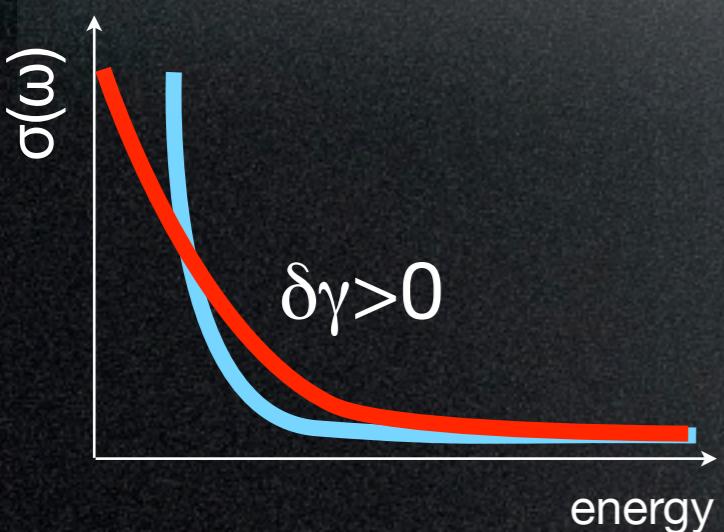
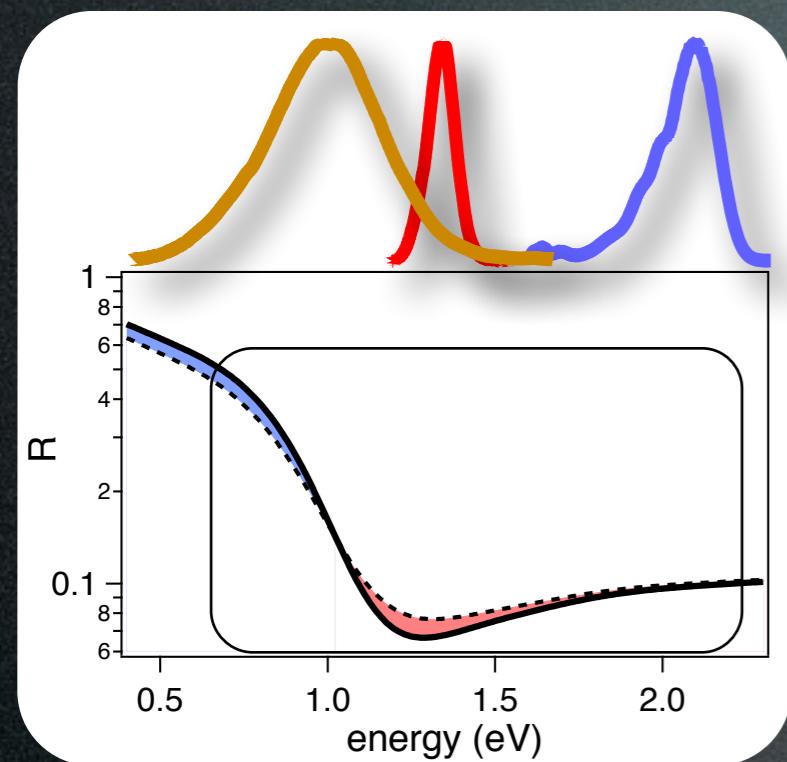


optimally doped $\text{Bi}_2\text{Sr}_2\text{Y}_{0.08}\text{Ca}_{0.92}\text{Cu}_2\text{O}_{8+\delta}$ (YBi2212)
 $T_c=96$ K

C. Giannetti



results on YBi2212 $T=300$ K

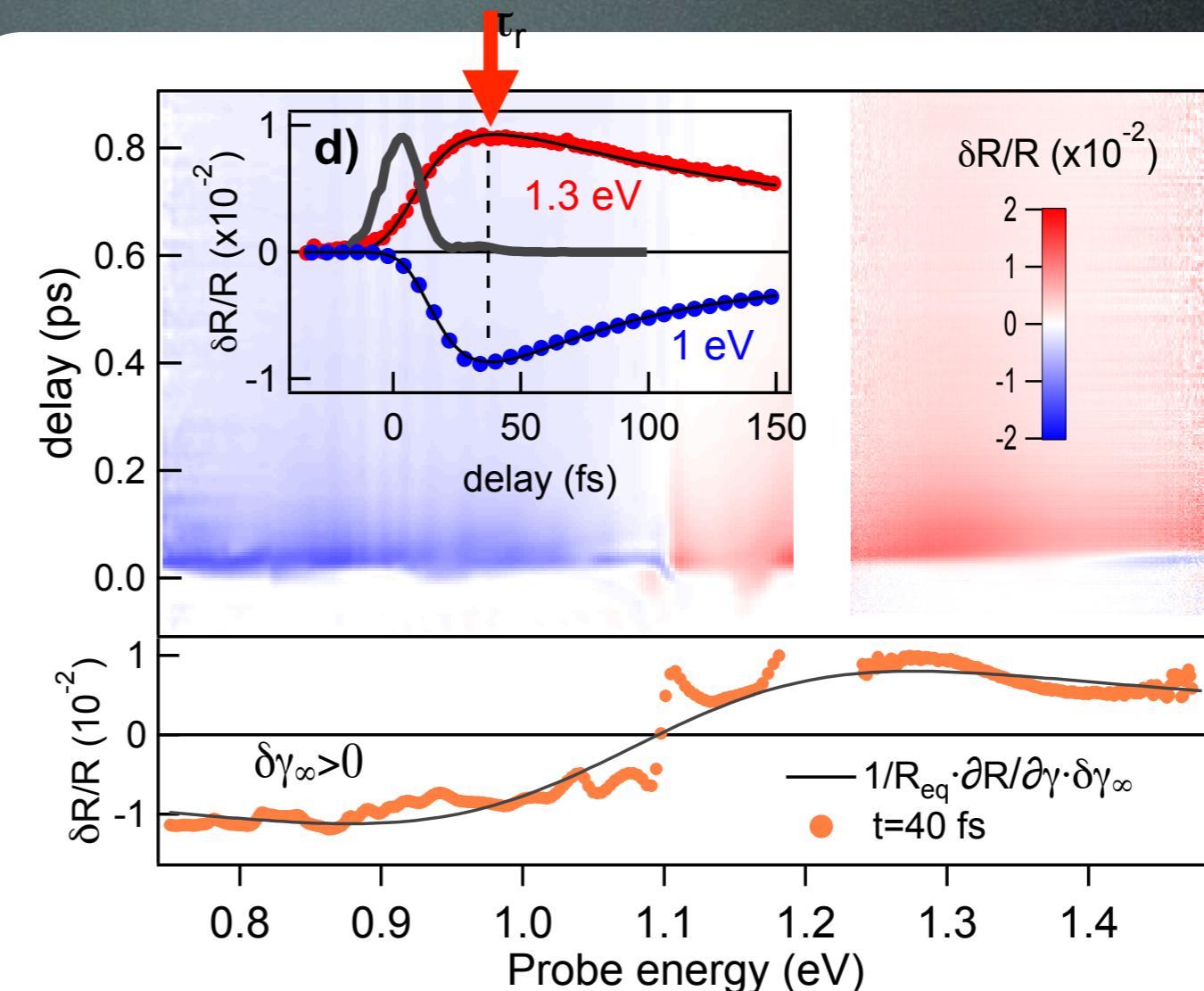
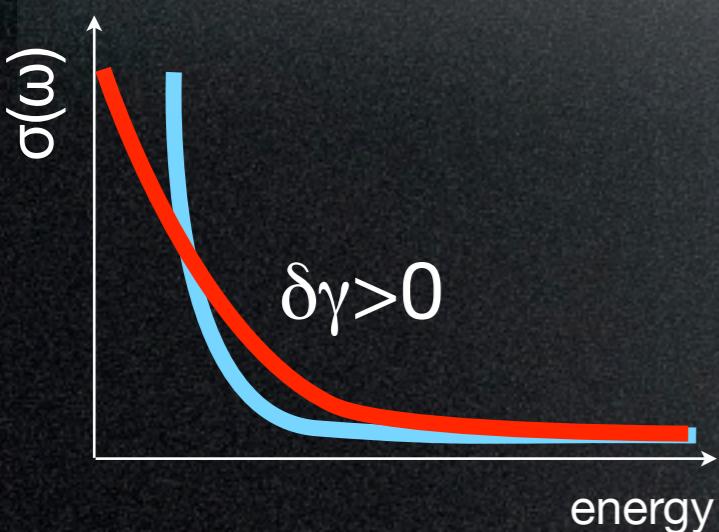
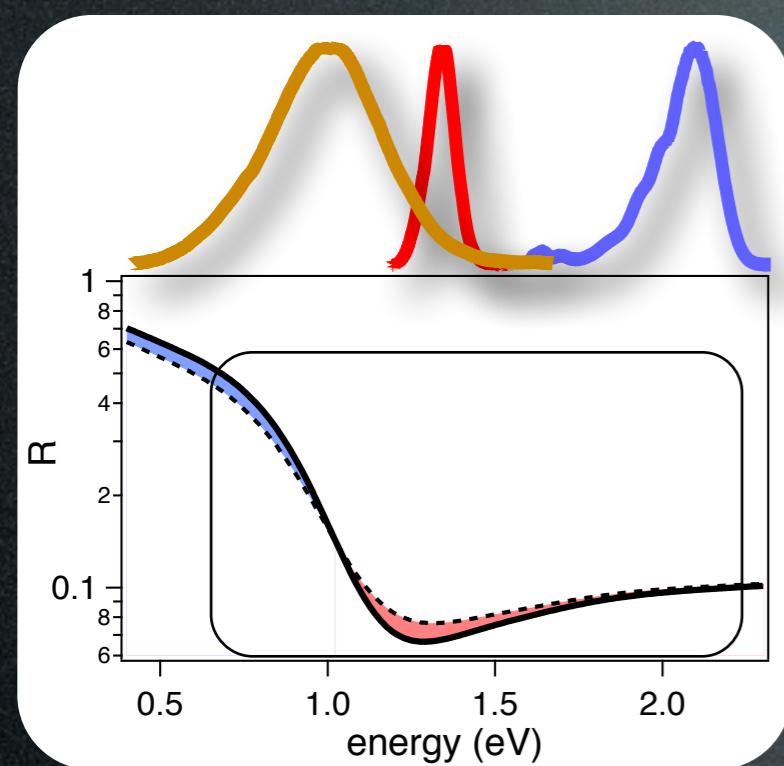


no variation at the energy scale $U!$ $\rightarrow \delta R(\omega) \propto \delta T_b$

- A) In 10-20 fs it is really possible to describe the dynamics of doped cuprates in terms of charge carriers interacting with effective bosonic fluctuations



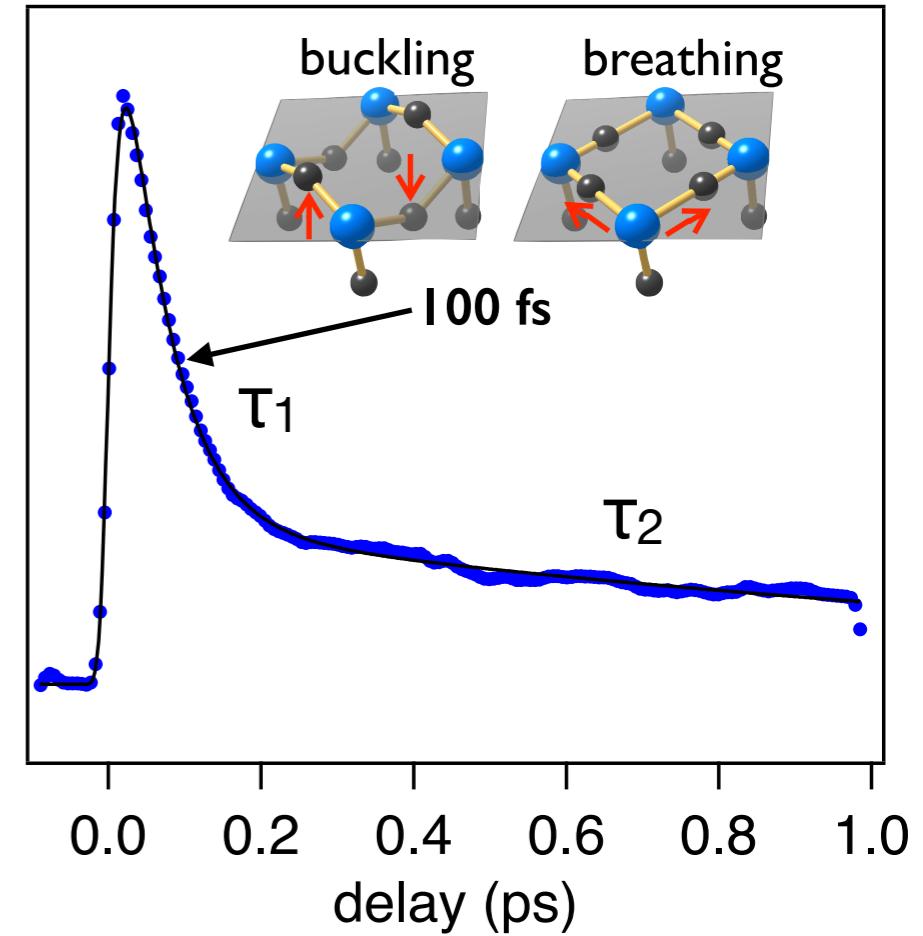
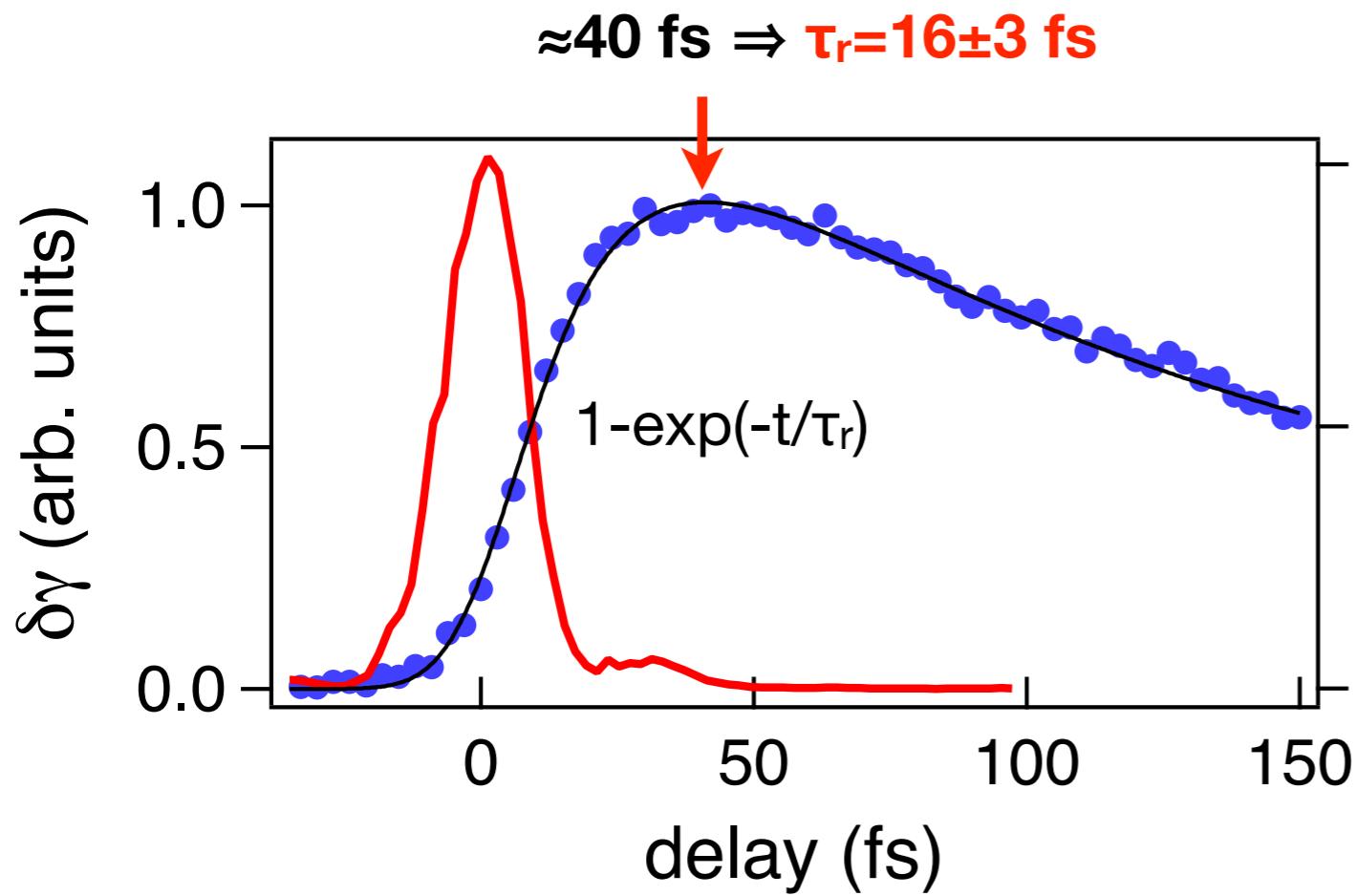
results on YBi2212 $T=300$ K



A) In 10-20 fs it is really possible to describe the dynamics of doped cuprates in terms of charge carriers interacting with effective bosonic fluctuations



retarded e-boson interaction

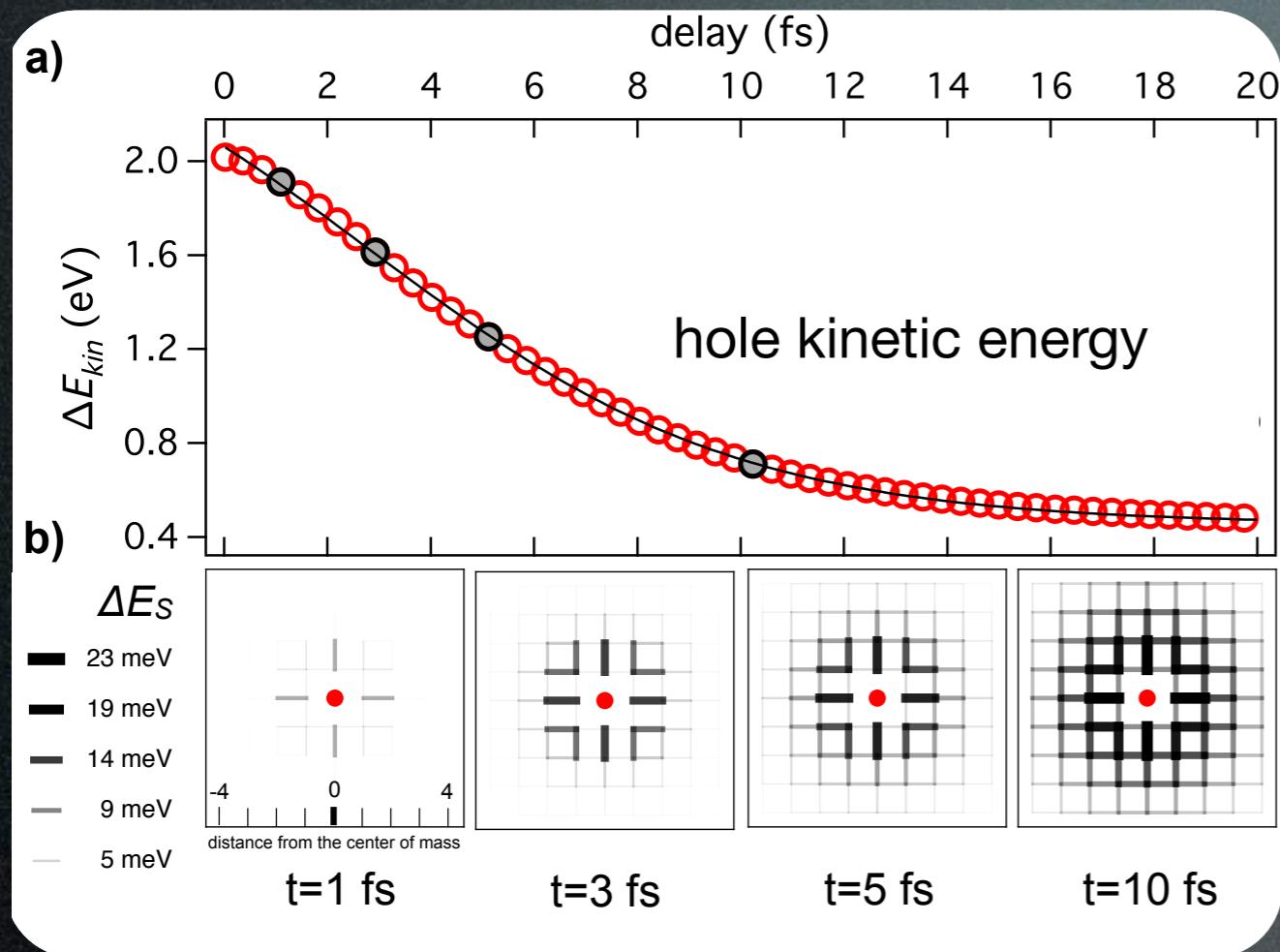


$$\delta R(\omega) \propto \delta T_b$$

B) In 16 fs photoexcited carriers can exchange energy with bosons ✓



non-equilibrium t-J model

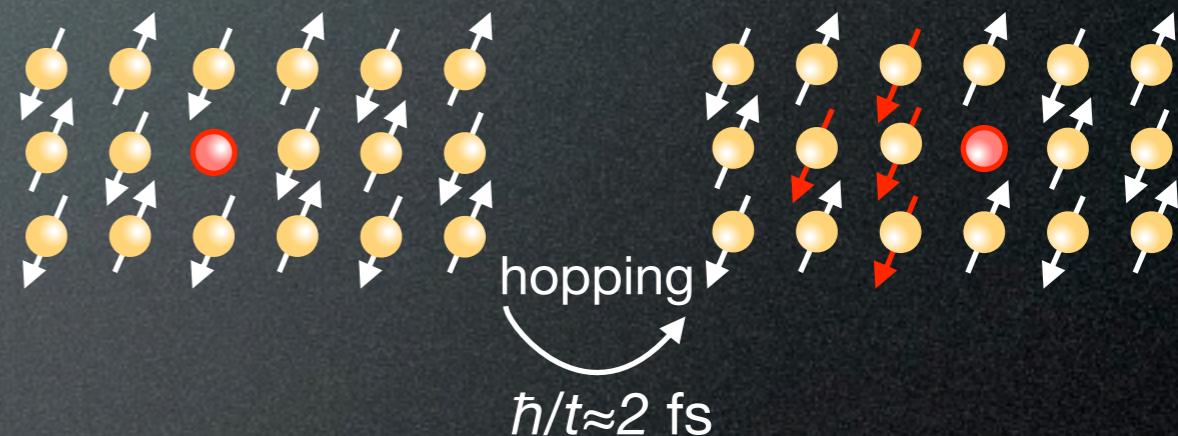


energy transfer to AF background
→ time-dep. Schr. equation

D. Golez, M. Mierzejewski, J. Bonca, L. Vidmar

Ultrafast coupling to high-energy short-range AF spin fluctuations:

- 10 fs coupling to short-range AF spin background

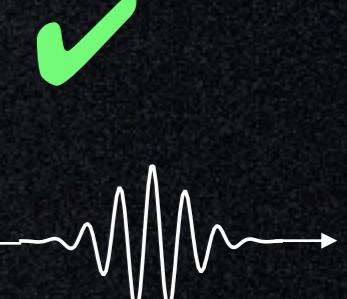


$$H = -t_h \sum_{\langle \mathbf{lj} \rangle, \sigma} (c_{\mathbf{l},\sigma}^\dagger \tilde{c}_{\mathbf{j},\sigma} + \text{h.c.}) + J \sum_{\langle \mathbf{lj} \rangle} \mathbf{S}_{\mathbf{l}} \cdot \mathbf{S}_{\mathbf{j}}$$

$t=360 \text{ meV}$ $J=120 \text{ meV}$

B) In 16 fs photoexcited carriers can exchange energy with bosons ✓

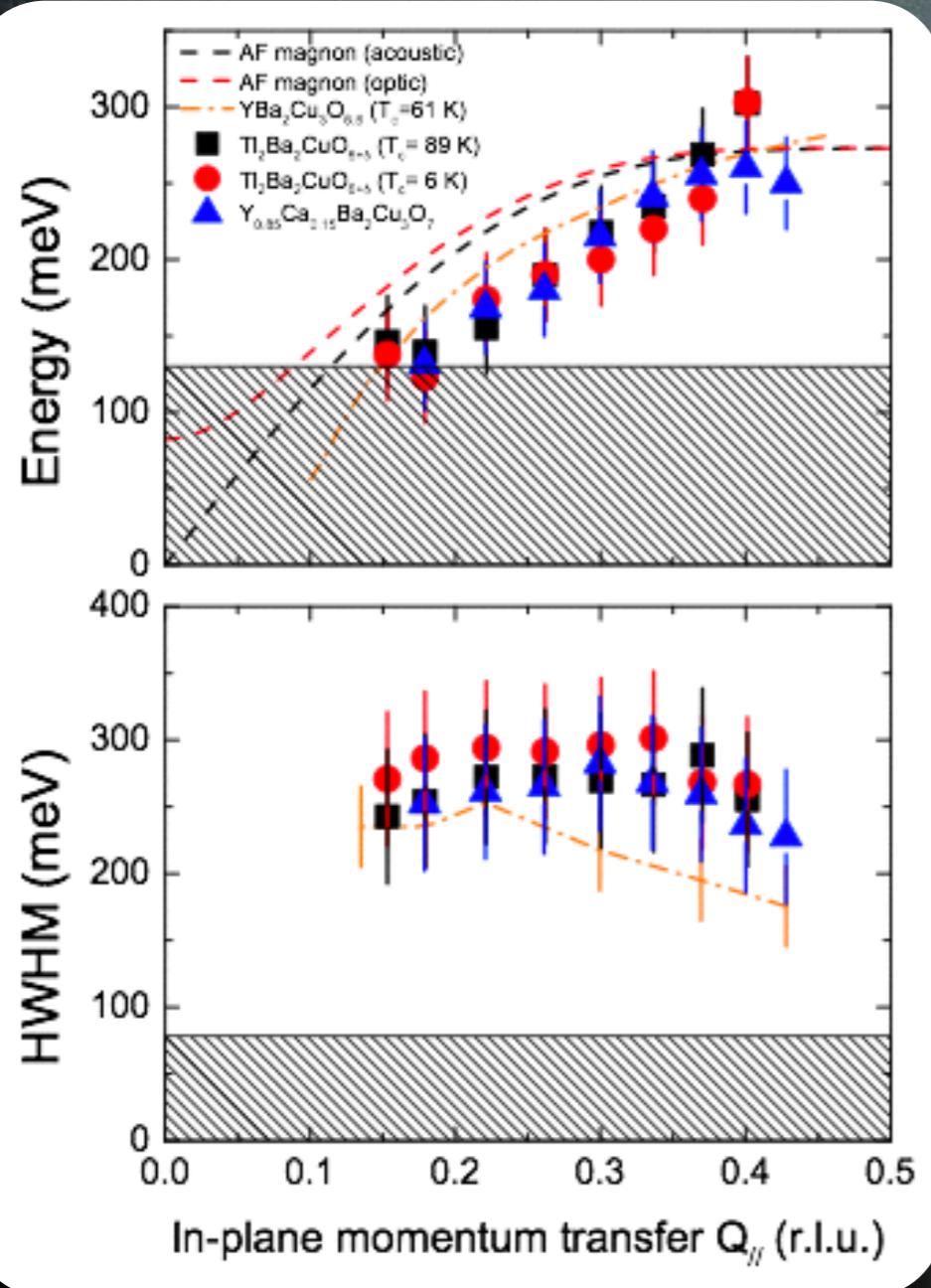
in agreement with non-eq. DMFT (M. Eckstein & P. Werner *arXiv:1410.3956*)



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short-range antiferromagnetic fluctuations

RIXS: Dispersion of paramagnons in cuprates at all dopings



Ultrafast coupling to high-energy short-range AF spin fluctuations:

- ~2 fs optical scattering rate
- 2-3 fs the inverse of paramagnon width (2-300 meV)

$$\tau = -\frac{\hbar}{\gamma} \quad \hbar = 658 \text{ meV fs}$$

M. Le Tacon et al, *Nature Physics* **7**, 725 (2011)
M. Dean et al, *Nature Materials* **12**, 1019 (2013)

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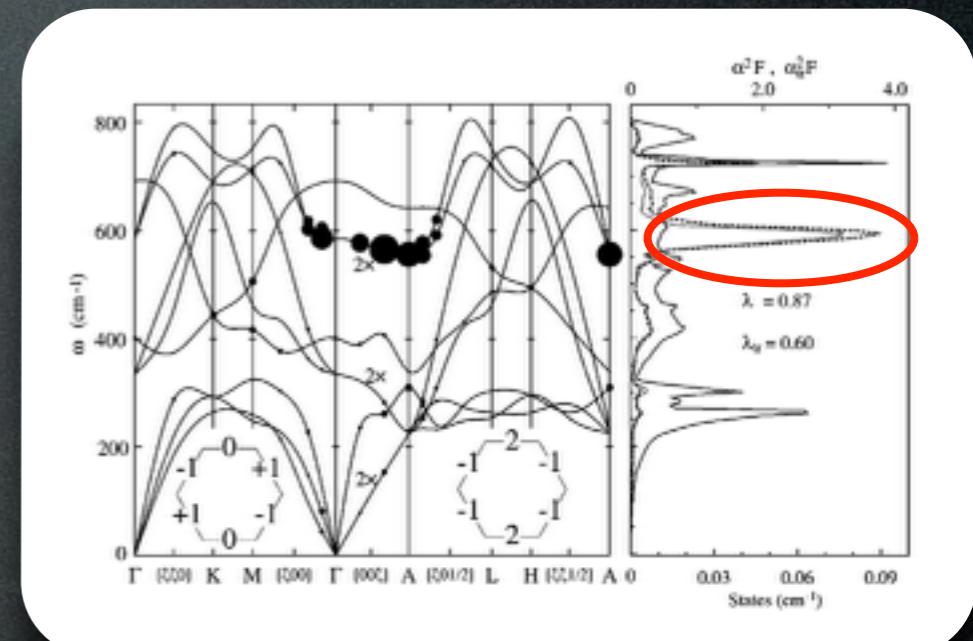
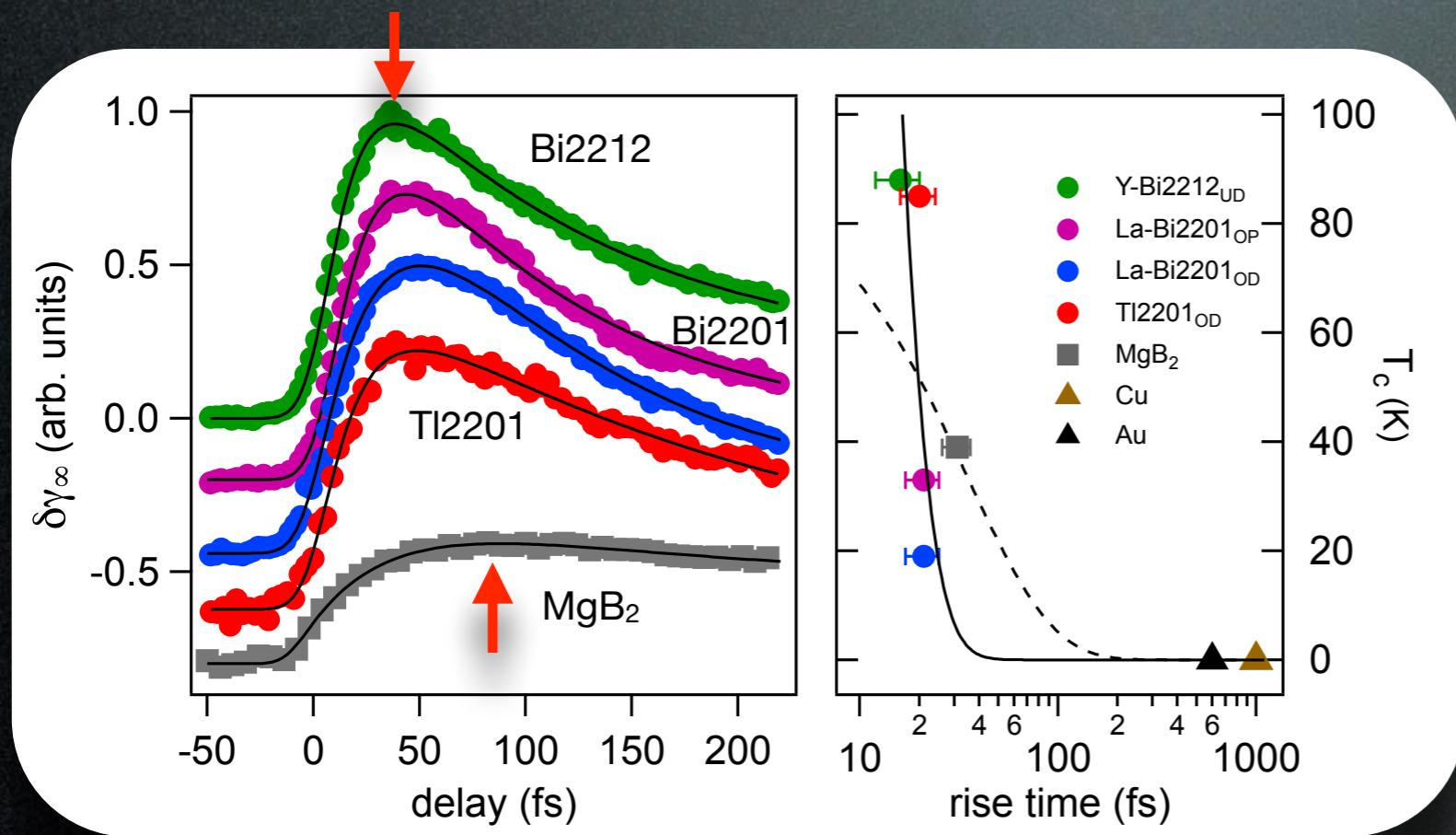
comparing to conventional superconductors (MgB_2)

MgB_2 : conventional BCS system with $T_c=40$ K

$$\approx 100 \text{ fs} \Rightarrow \tau_r = 32 \pm 3 \text{ fs}$$

coupling with optical mode at
70 meV

$$\tau_r \sim 2\pi/\Omega_Q$$



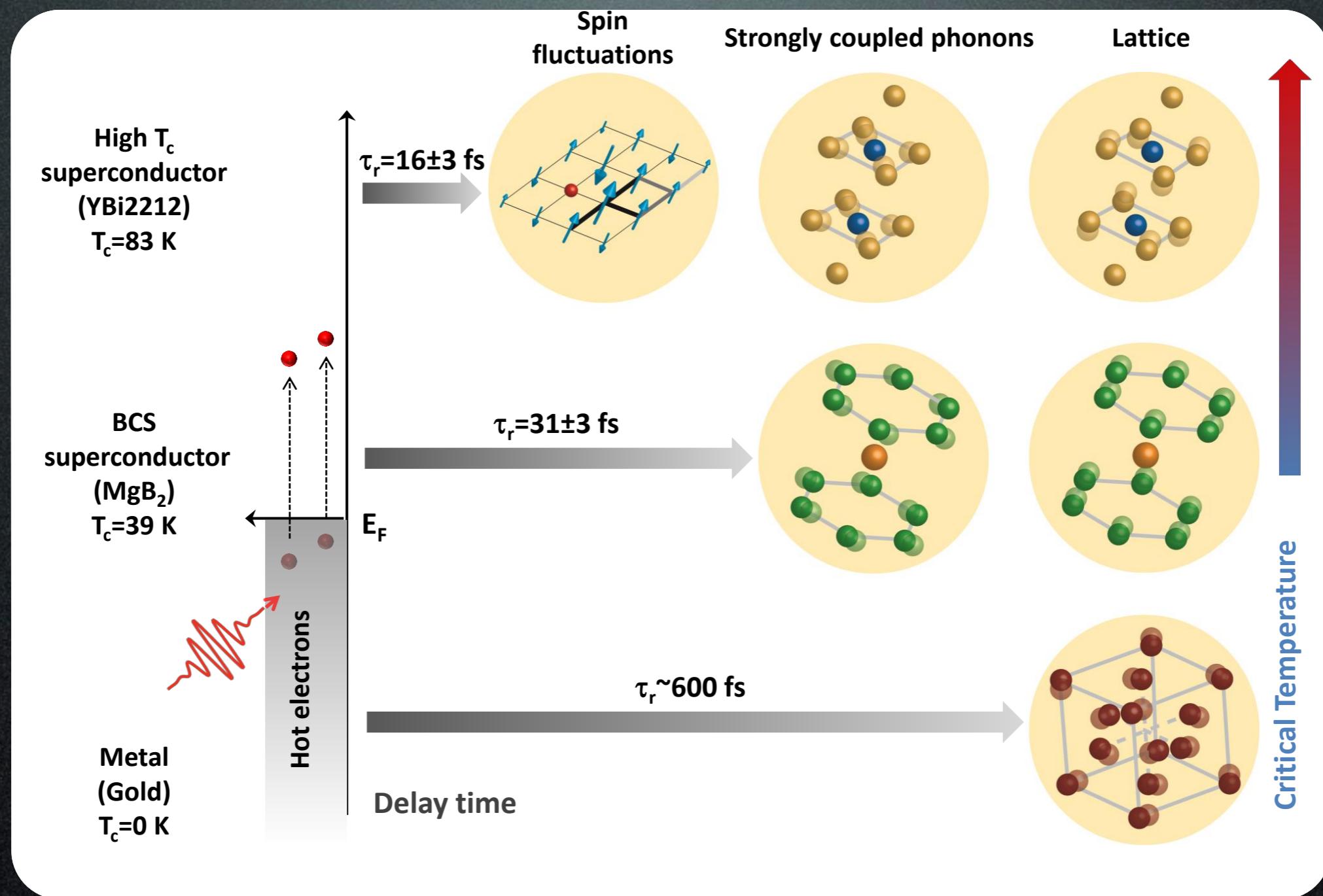
Y. Kong et al., *Phys. Rev. B* **64**, 020501(R) (2001)

P.B. Allen and R.C. Dynes, *Phys. Rev. B* **12**, 905 (1975)
A.J. Millis et al., *Phys. Rev. B* **37**, 4975 (1988).

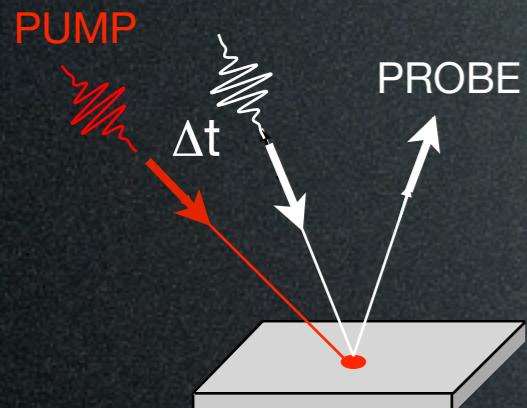
$$T_c = 0.83\tilde{\Omega} e^{-\frac{1.04(1+\lambda)}{g(\lambda - \mu^*(1+0.62\lambda))}}$$



Summary



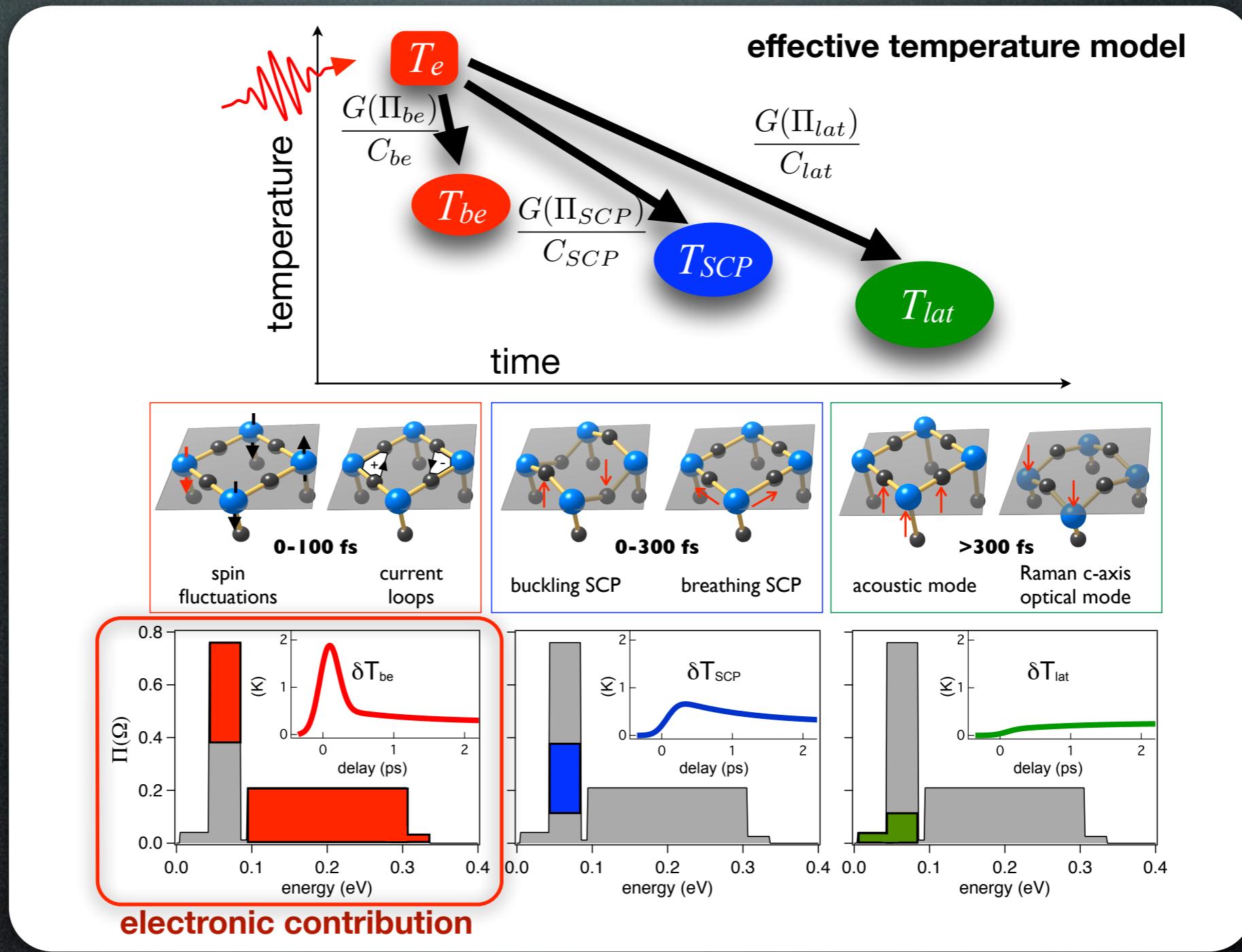
electron-boson coupling in CUPRATES



$$\Pi(\Omega) = \alpha^2 F(\Omega) + I^2 \chi(\Omega)$$

E. van Heumen et al., *Phys. Rev. B* **79**, 184512 (2009)

J. Carbotte et al. *Rep. Prog. Phys.* **74**, 066501 (2011)



S. Dal Conte et al., *Science* **335**, 1600 (2012)

in agreement with the glue extracted from the Hubbard model

E. Gull & A.J. Millis *arXiv:1407.0704* (2014) → see talk: E. Gull



Summary

- Non-equilibrium spectroscopy for disentangling different degrees of freedom
- Snapshots of the electron-boson coupling in correlated materials
- Coupling of the charge-carriers with AF fluctuations on the 15 fs timescale
- Effective glue in optimally and overdoped cuprates



C. Giannetti



People and Collaborations

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- Ultrafast optics group (Università degli Studi di Trieste)

F. Cilento, G. Coslovich, D. Fausti, F. Parmigiani

- Ultrafast optics group (Politecnico di Milano)

S. Dal Conte, D. Brida, G. Cerullo

- Equilibrium spectroscopies

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M. Greven (University of Minnesota & Stanford University)

L. Chauviere, B. Keimer (MPG-UBC center for QM)

- Non-equilibrium models of correlated materials

L. Vidmar (LMU Munich), M. Mierzejewski (Katowice), D. Golez, J. Bonca (Ljubljana)

M. Capone, M. Fabrizio (SISSA, Trieste)

- Equilibrium optical properties of HTSC

D. van der Marel (Université de Genève)



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Thank you!



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