

# Frustrated magnetism, magnetization plateaux, and the mystery of $\text{SrCu}_2(\text{BO}_3)_2$

Solving a 15-year old puzzle



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Switzerland

# Collaborators



## Theorists

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K. Schmidt (Dortmund), J. Dorier (Lausanne)  
S. Manmana, A. Honecker (Göttingen)  
**P. Corboz (Amsterdam)**

## Experimentalists

M. Takigawa, Y. Matsuda (ISSP, Tokyo)  
C. Berthier, M. Horvatic, S. Krämer, I. Sheikin,  
F. Lévy-Bertrand (Grenoble)

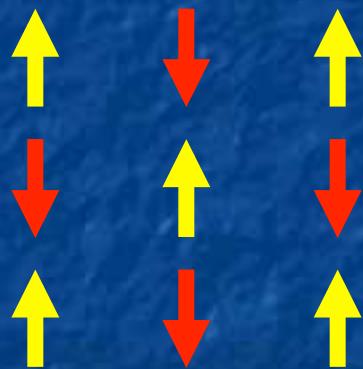
# Scope

- A few words about frustrated magnetism  
→ degeneracy, enhanced fluctuations, spin liquids
- Magnetization of frustrated magnets  
→ plateaus (1/3 on triangular lattice, etc.)
- The intriguing case of  $\text{SrCu}_2(\text{BO}_3)_2$   
→ plateaus at 1/8, 2/15, 1/6, 1/4, 1/3, 1/2 ?!
- Conclusions

# Heisenberg model

$$\mathcal{H} = \sum_{(i,j)} J_{ij} \vec{S}_i \cdot \vec{S}_j$$

Long-range order



Spin gap

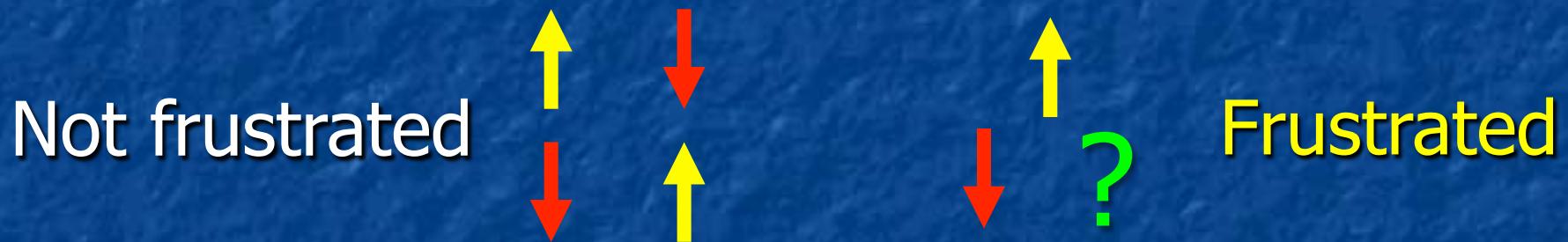


Singlet-triplet gap

Goldstone modes:  
Spin waves

$$\Delta = J' + O(J) \quad (J \ll J')$$

# Frustration



Antiferromagnetic coupling + odd loops



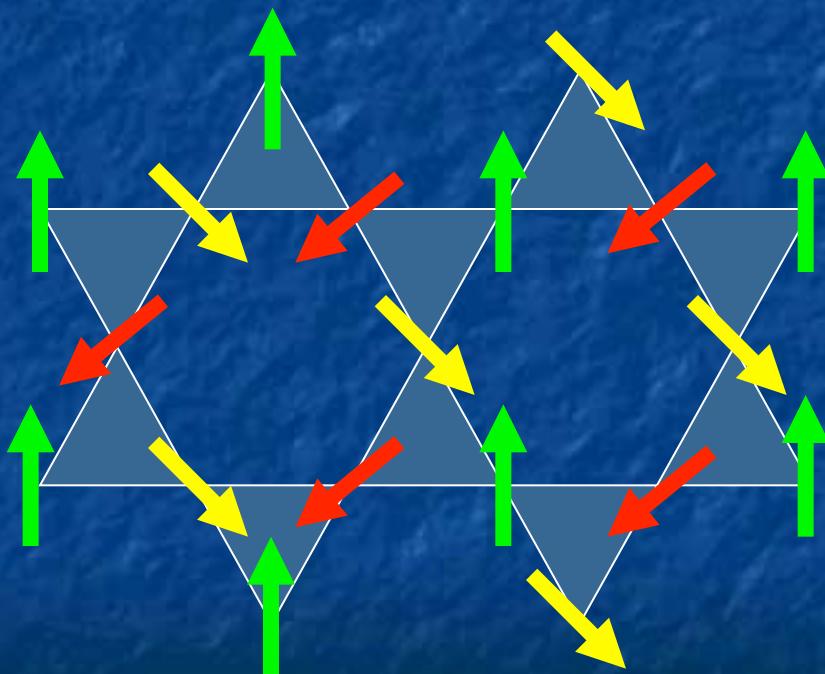
Competition between exchange paths = frustration

# Possible consequences of frustration

- Infinite degeneracy of **classical** ground state (kagome, pyrochlore,...)
- Enhanced **quantum fluctuations**
  - Resonating Valence Bond spin liquids
  - Algebraic spin liquids
- Magnetization plateaus

# Kagome antiferromagnet

Infinite degeneracy



RVB spin liquid?

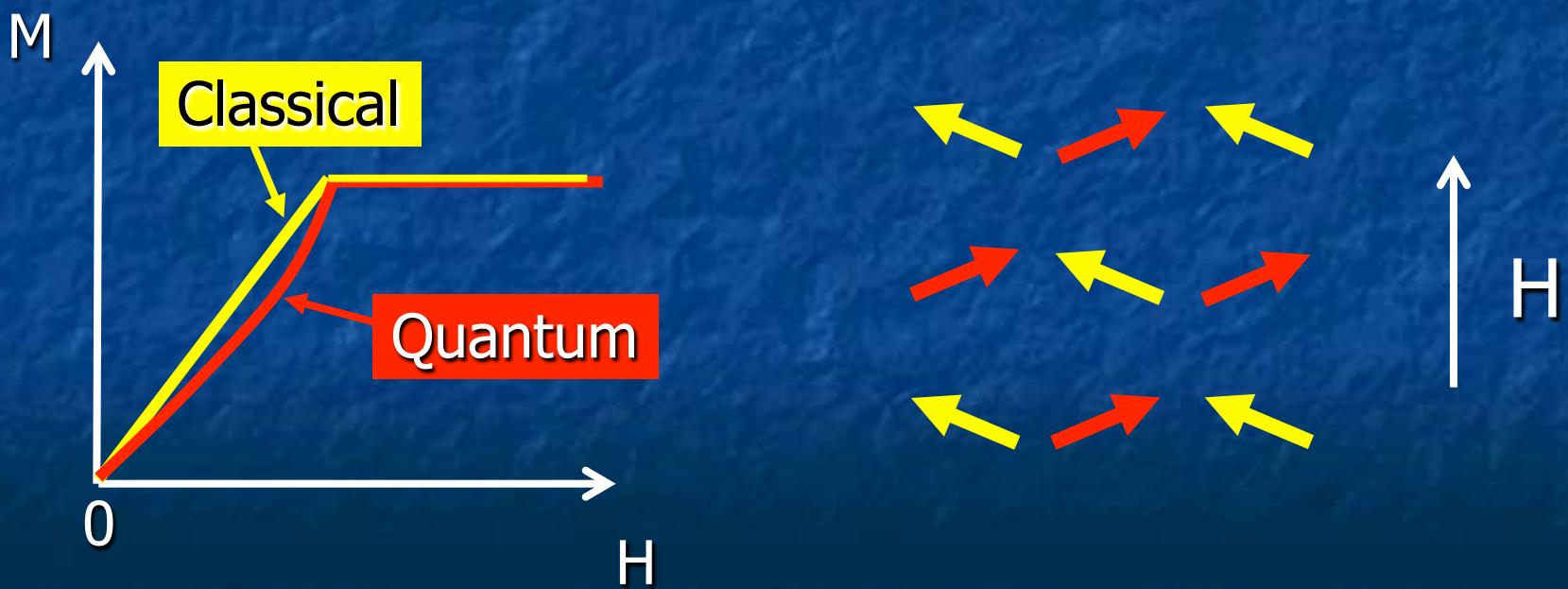
- Zeng, Elser, PRB 1995
- FM, PRL 1998
- Han, Huse White, Science 2011

Algebraic spin liquid?

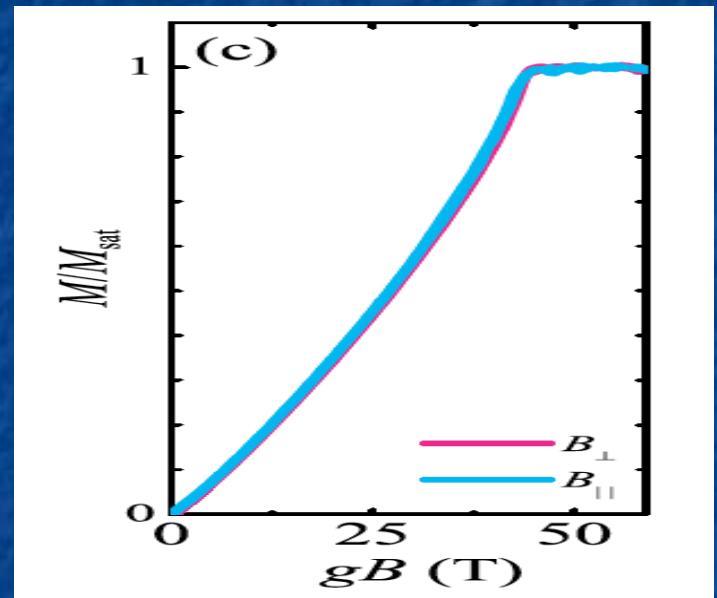
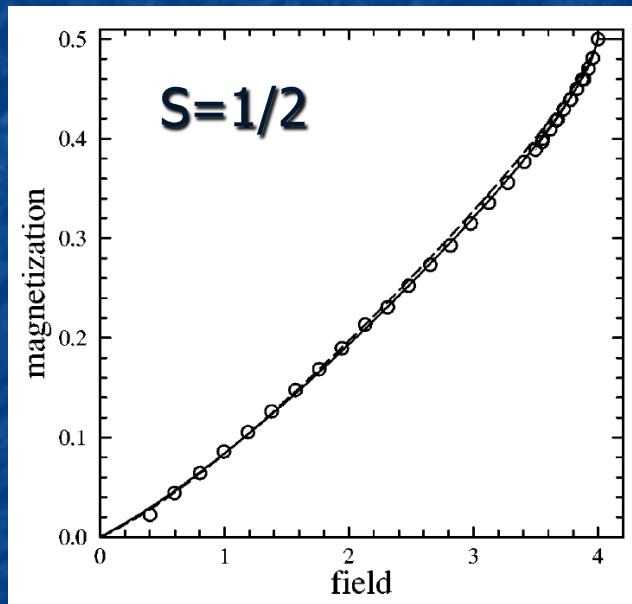
- Ran, Hermele, Lee, Wen, PRL 2007

# Magnetization of Néel AF

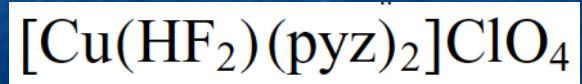
$$\mathcal{H} = J \sum_{\langle i,j \rangle} \vec{S}_i \cdot \vec{S}_j - g\mu_B H \sum_i S_i^z$$



# Square lattice antiferromagnet

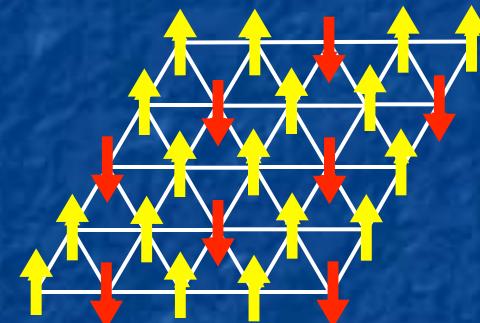
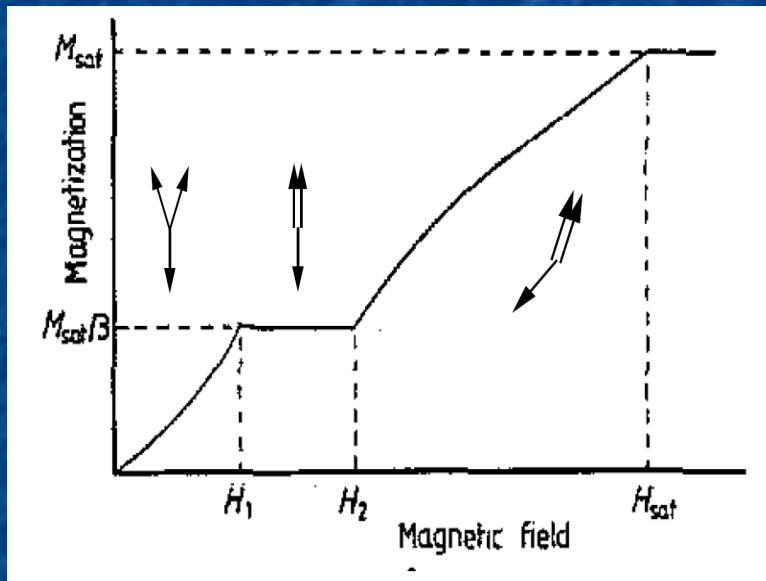


- Circles: QMC Igarashi, 1992
- Dashed and solid lines:  
1<sup>st</sup> and 2<sup>nd</sup> order spin-wave  
Zhitomirsky, Nikuni, 1998



Goddard et al, 2008

# 1/3 magnetization plateau

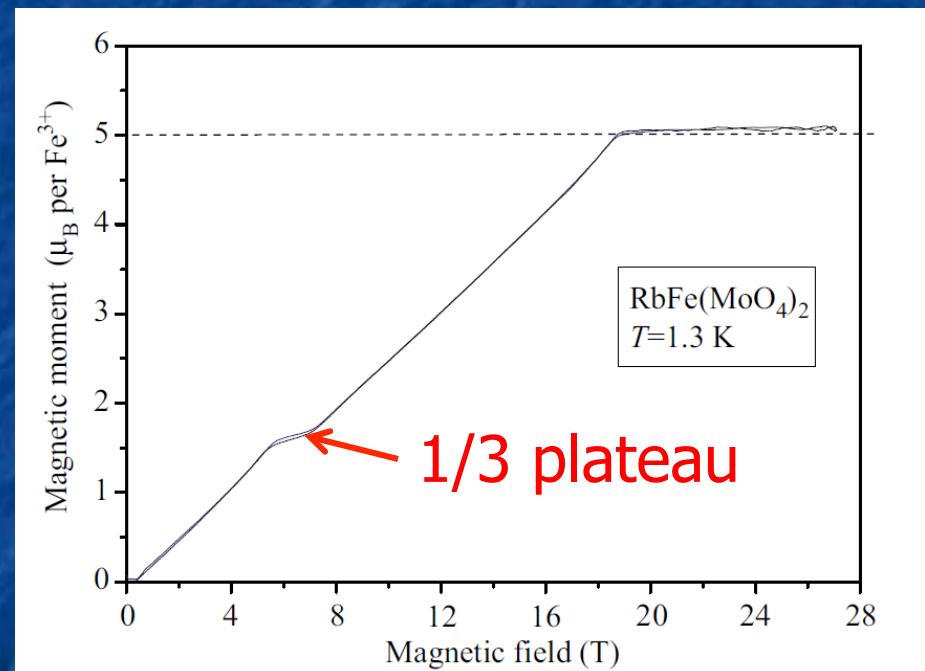
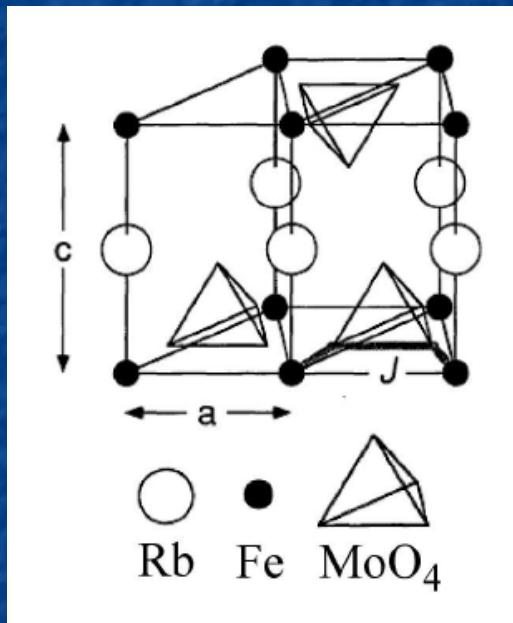


Chubukov and Golosov  
JPMC 1991

Collinear structure stabilized in a field range by zero-point fluctuations

# $\text{RbFe}(\text{MoO}_4)_2$

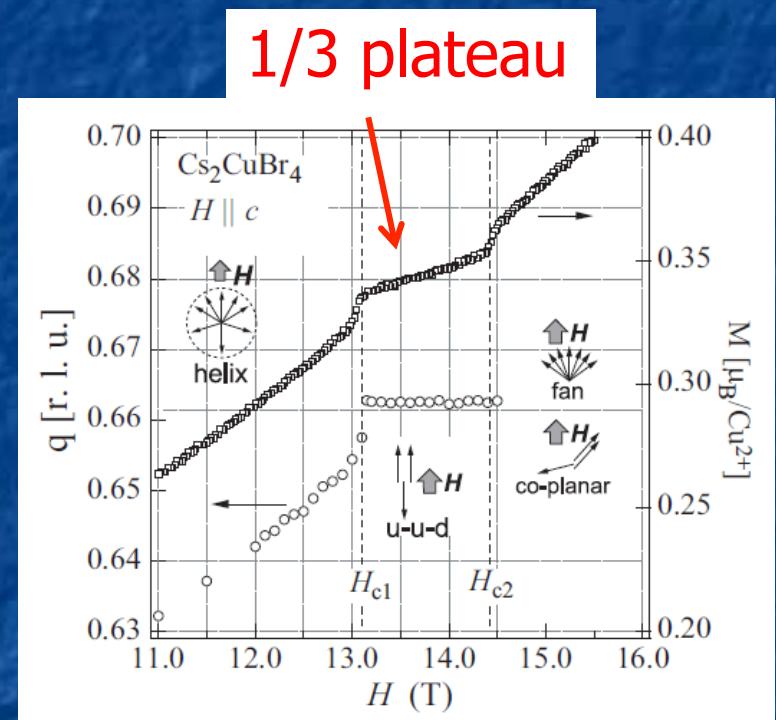
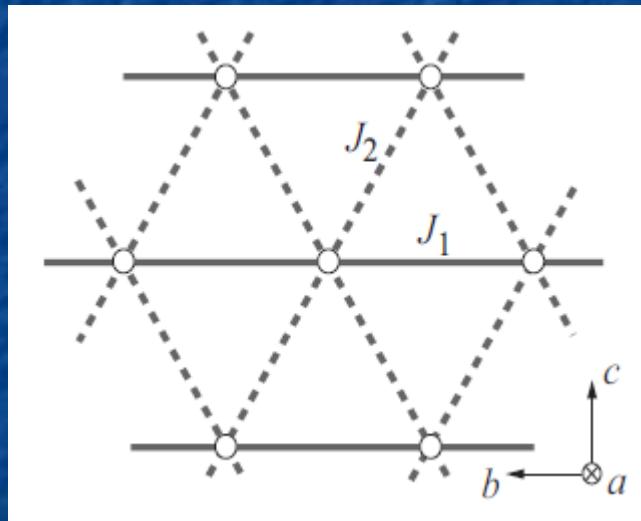
Smirnov et al, PRB 2007



Spin 5/2, triangular lattice

# $\text{Cs}_2\text{CuBr}_4$

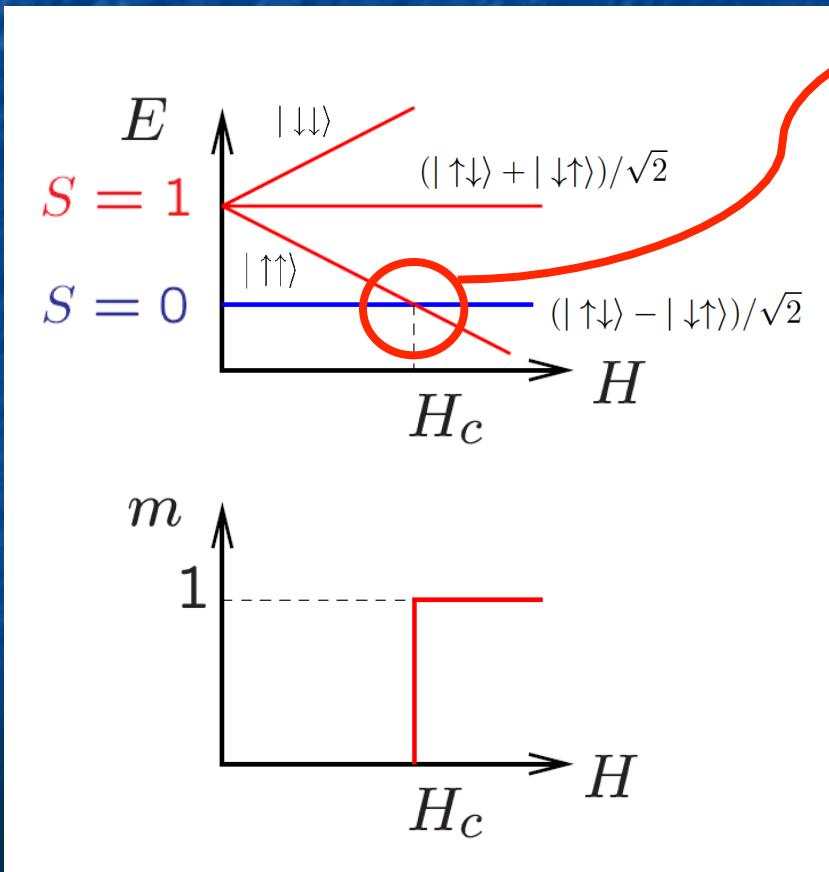
Ono et al,  
J. Phys. Cond. Mat. 2004



Spin 1/2, anisotropic triangular lattice

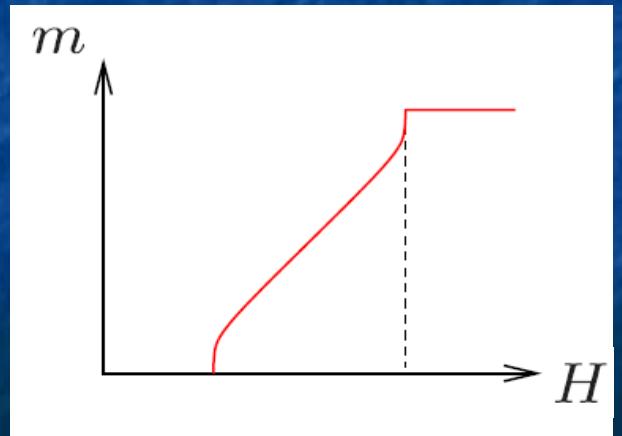
# Gapped systems (ladders,...)

Isolated dimers

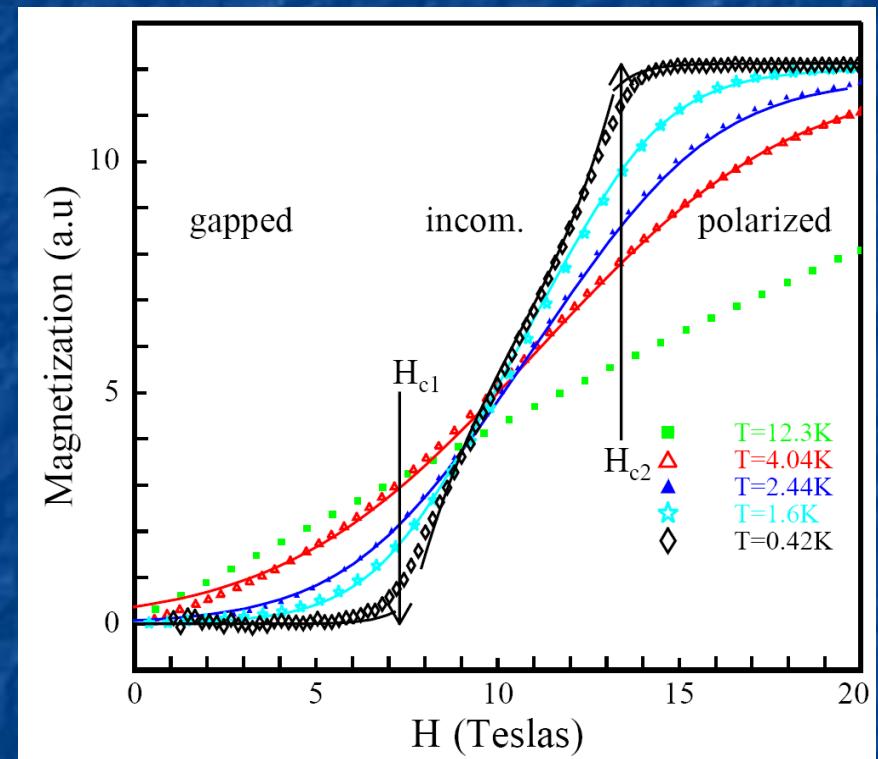
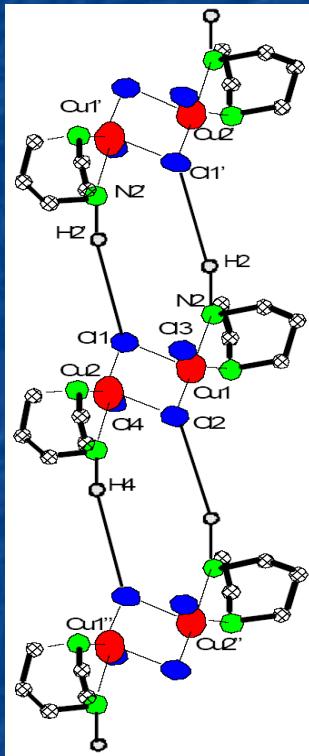


$S=0$ : empty site  
 $S_z=1$ : boson

Coupled dimers  
(ladders,...)



# Magnetization of CuHpCl

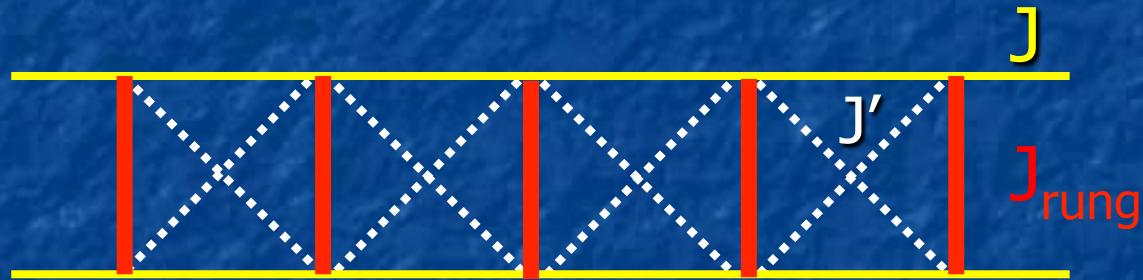


CuHpCl

Chaboussant et al, EPJB '98

No magnetization up to critical field  $H_{c1}$

# Frustrated ladder



$$H = -t \sum_{\langle i,j \rangle} (a_i^\dagger a_j + a_j^\dagger a_i) - \mu \sum_i n_i + V \sum_{\langle i,j \rangle} n_i n_j$$

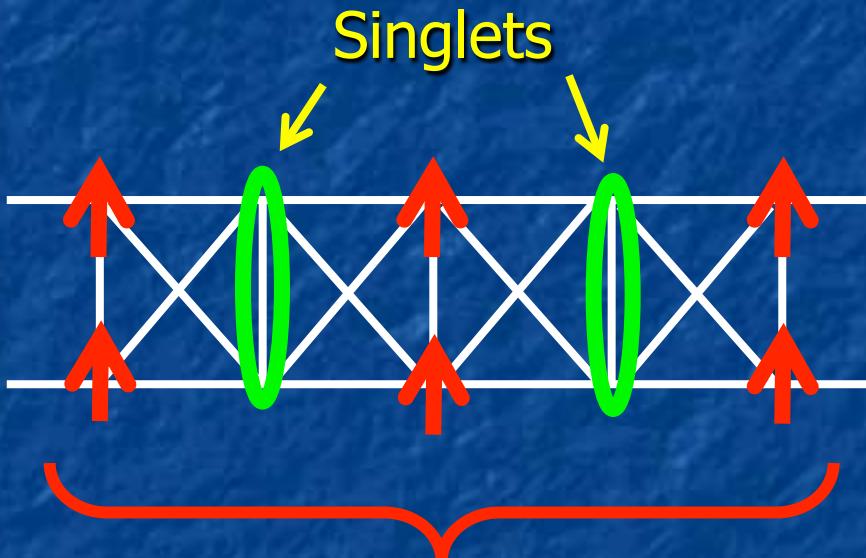
$$t = (J - J')/2$$

$$\mu = H - H_c$$

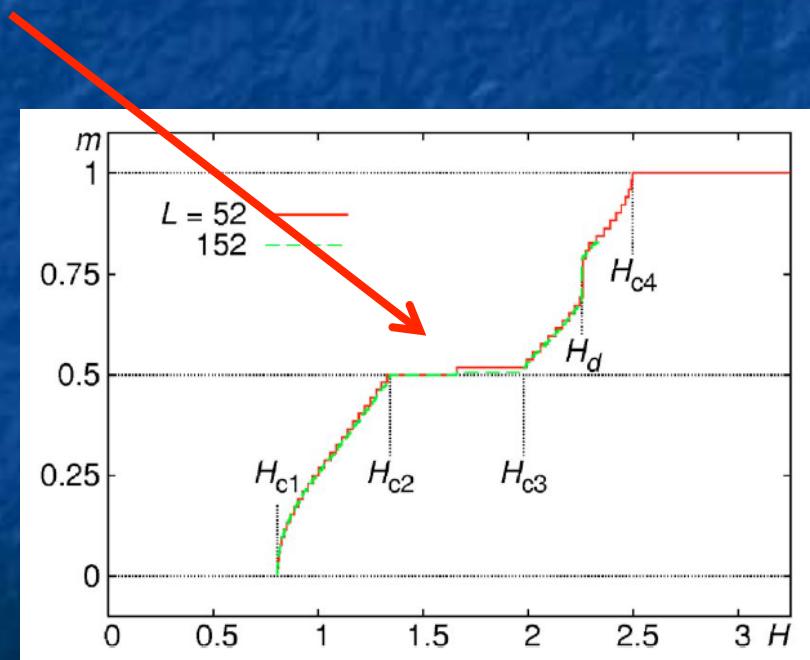
$$V = (J + J')/2$$

- Mott insulator for  $V > 2t$  Ovchinnikov 1971
- 1/2 plateau for  $1/3 < J'/J < 3$  Totsuka PRB 98; FM EPJB 98

# Frustrated ladder



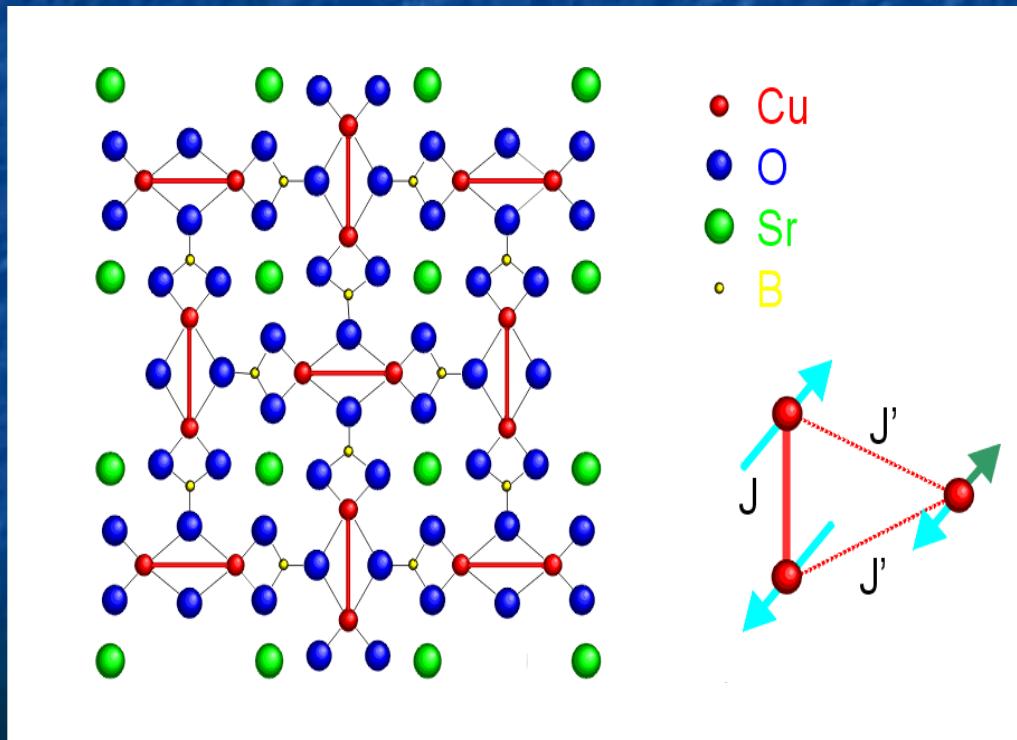
Crystal of triplets



Fouet et al, 2006

# $\text{SrCu}_2(\text{BO}_3)_2$

Smith and Keszler, JSSC 1991

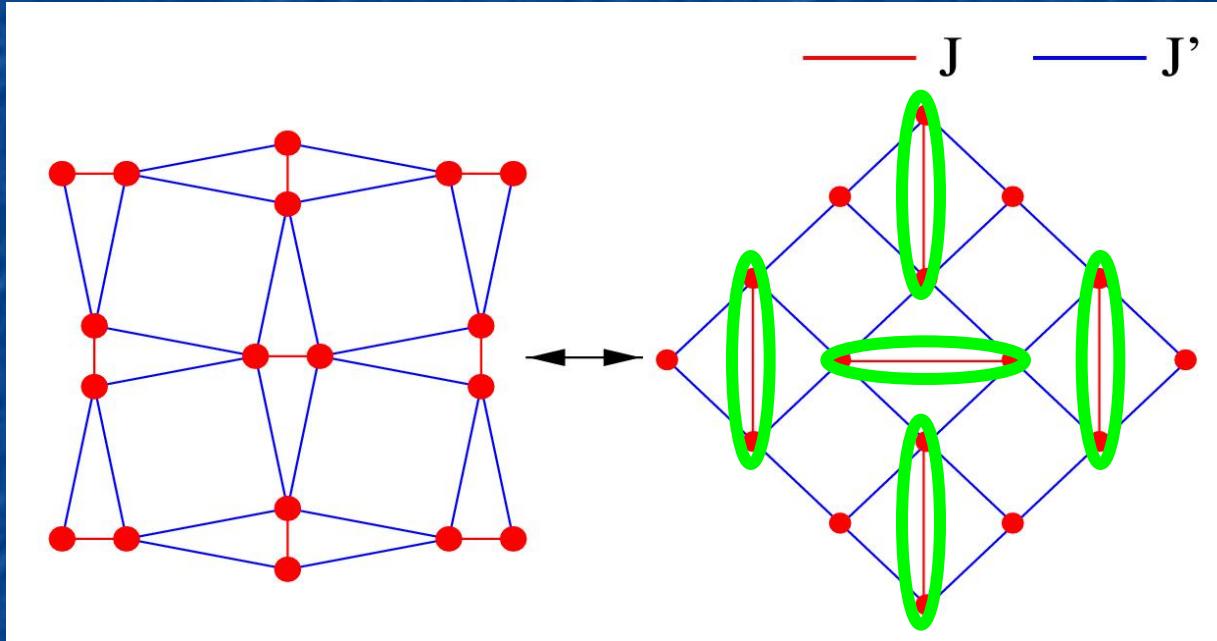


$\text{Cu}^{2+} \rightarrow \text{Spin } 1/2$

$J \approx 85 \text{ K}$

$J'/J \approx 0.65$

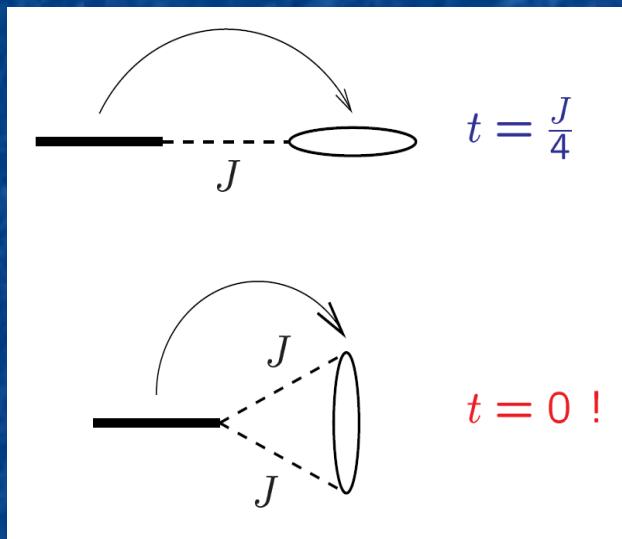
# Shastry-Sutherland model



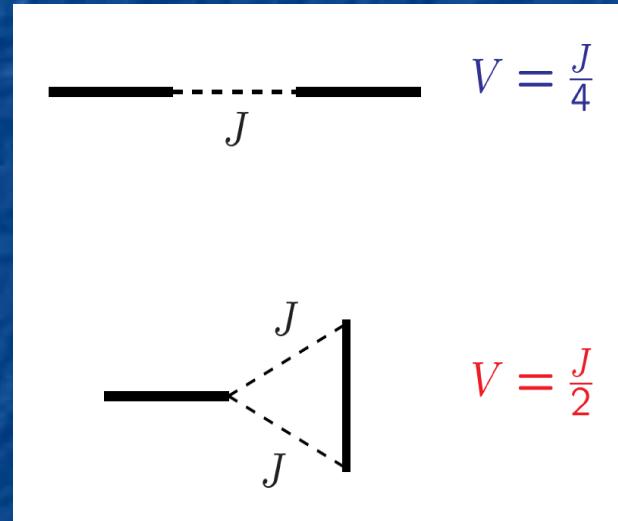
Ground-state = Product of singlets on  $J$ -bonds  
Shastry and Sutherland, '81  
→ Spin gap, and plateau at 0

# Effect of frustration

## Triplet Hopping



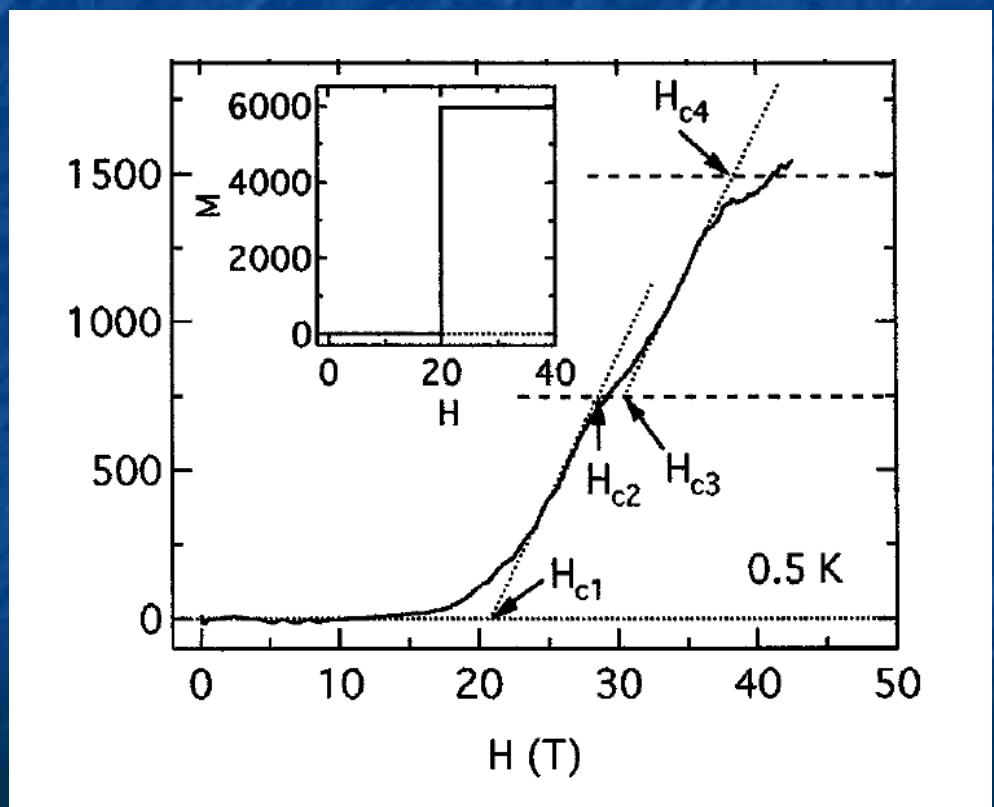
## Triplet Repulsion



- Kinetic energy << potential energy
- Long-range repulsion
  - Crystals of triplets with high commensurability
  - Magnetization plateaux

## Exact Dimer Ground State and Quantized Magnetization Plateaus in the Two-Dimensional Spin System $\text{SrCu}_2(\text{BO}_3)_2$

H. Kageyama,<sup>1,2,\*</sup> K. Yoshimura,<sup>1,3,†</sup> R. Stern,<sup>3</sup> N. V. Marshnikov,<sup>2</sup> K. Onizuka,<sup>2</sup> M. Kato,<sup>1</sup> K. Kosuge,<sup>1</sup>  
C. P. Slichter,<sup>3</sup> T. Goto,<sup>2</sup> and Y. Ueda<sup>2</sup>

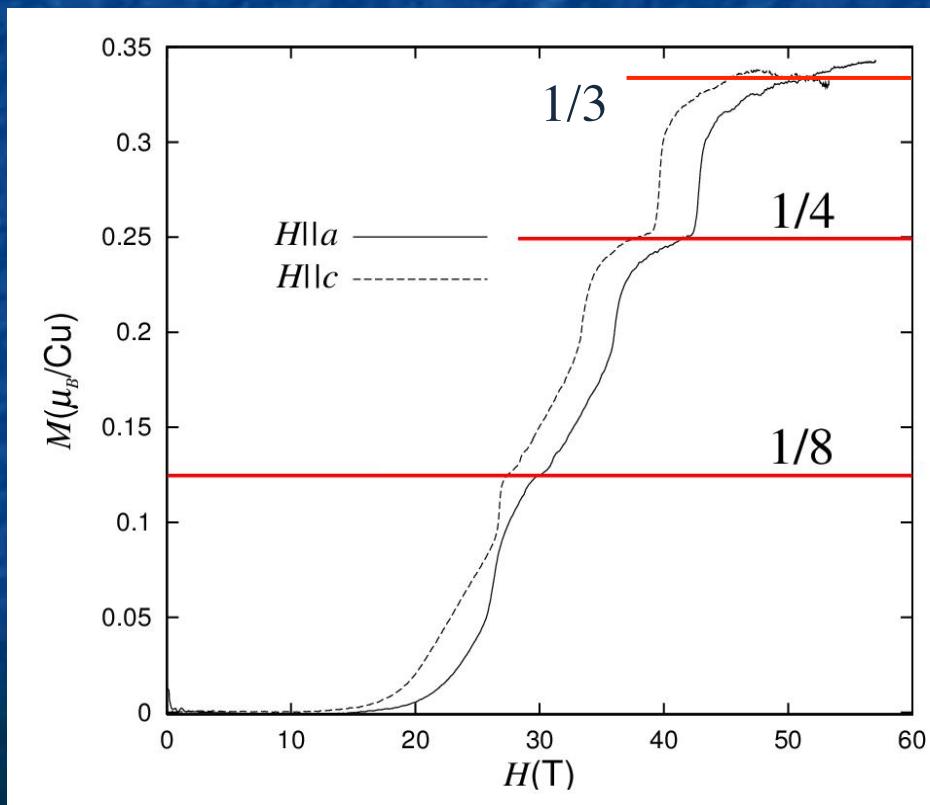


Anomalies

- $M=0$
- $M=1/8$
- $M=1/4$

## 1/3 Magnetization Plateau in $\text{SrCu}_2(\text{BO}_3)_2$ - Stripe Order of Excited Triplets -

Kenzo ONIZUKA, Hiroshi KAGEYAMA\*, Yasuo NARUMI<sup>1,2</sup>,  
Koichi KINDO<sup>2,1</sup>, Yutaka UEDA and Tsuneaki GOTO



Plateaus

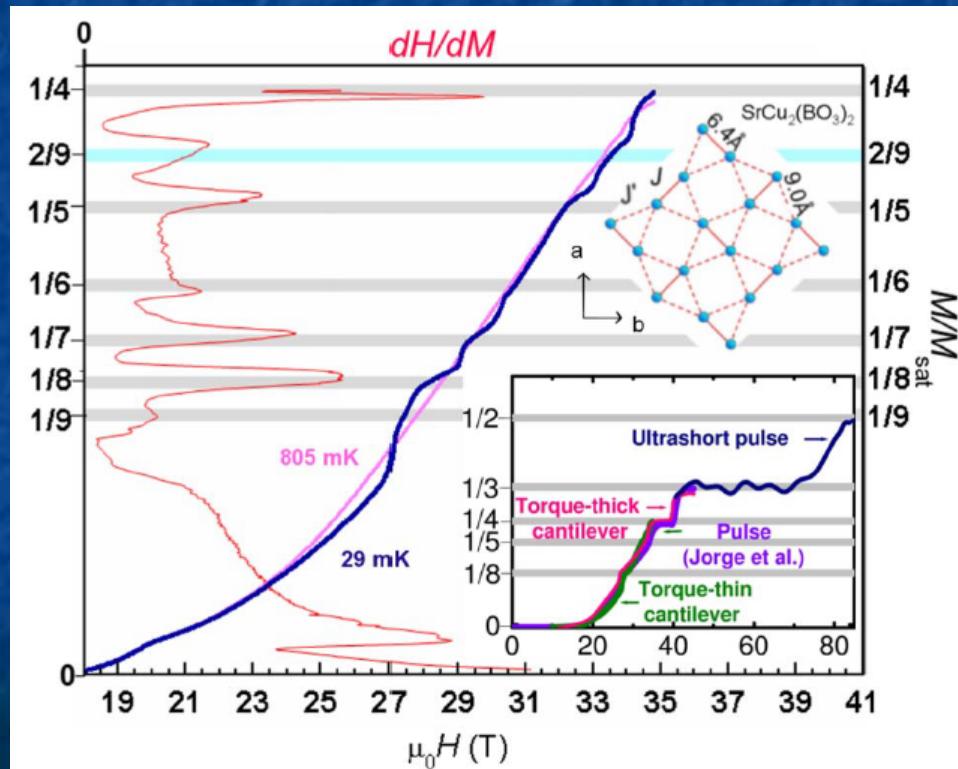
- $M=0$
- $M=1/8$
- $M=1/4$
- $M=1/3$

# Fractalization drives crystalline states in a frustrated spin system

Suchitra E. Sebastian<sup>a,1</sup>, N. Harrison<sup>b</sup>, P. Sengupta<sup>c</sup>, C. D. Batista<sup>c</sup>, S. Francoual<sup>b</sup>, E. Palm<sup>d</sup>, T. Murphy<sup>d</sup>, N. Marcano<sup>a</sup>, H. A. Dabkowska<sup>e</sup>, and B. D. Gaulin<sup>e</sup>

[www.pnas.org/cgi/doi/10.1073/pnas.0804320105](http://www.pnas.org/cgi/doi/10.1073/pnas.0804320105)

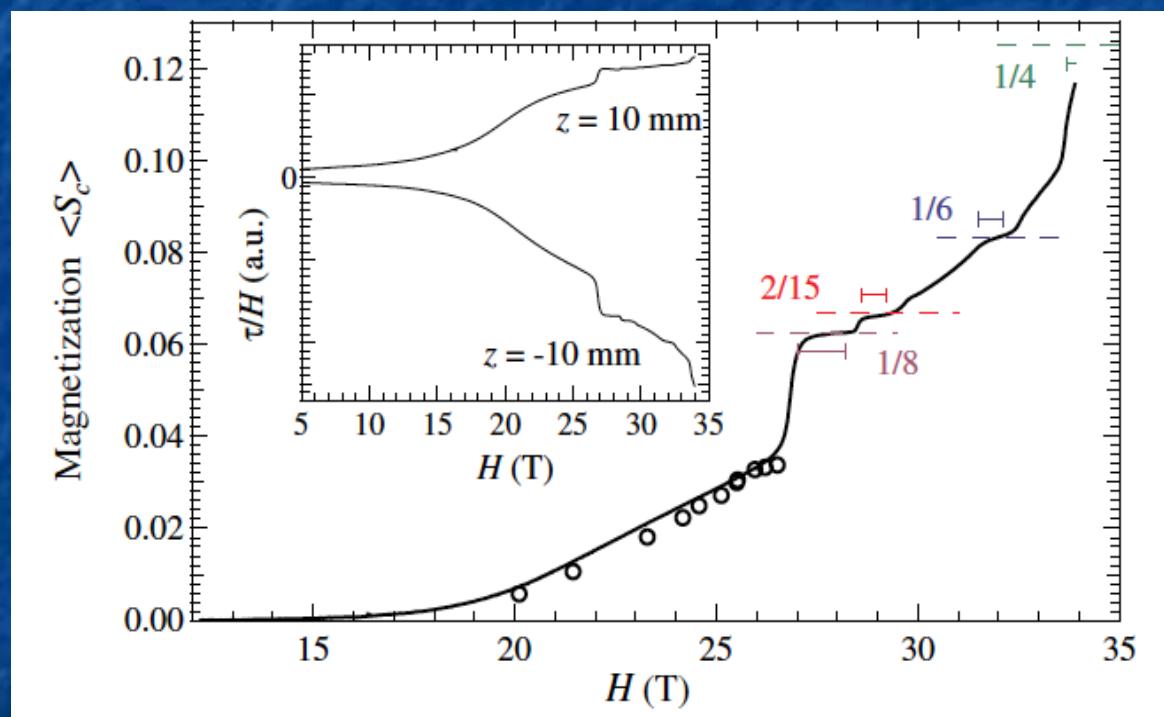
PNAS | December 23, 2008 | vol. 105 | no. 51 | 20157–20160



1/9, 1/8, 1/7, 1/6,  
1/5, 2/9, 1/4

# Incomplete Devil's Staircase in the Magnetization Curve of $\text{SrCu}_2(\text{BO}_3)_2$

M. Takigawa,<sup>1,\*</sup> M. Horvatić,<sup>2</sup> T. Waki,<sup>3</sup> S. Krämer,<sup>2</sup> C. Berthier,<sup>2</sup> F. Lévy-Bertrand,<sup>2,†</sup> I. Sheikin,<sup>2</sup> H. Kageyama,<sup>4</sup> Y. Ueda,<sup>1</sup> and F. Mila<sup>5</sup>



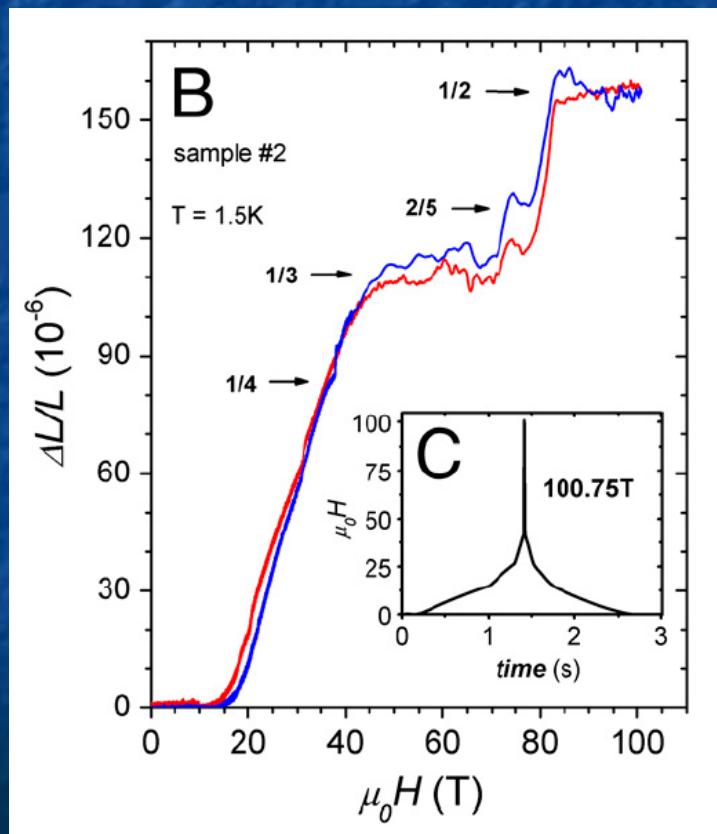
1/8, 2/15, 1/6, 1/4, ...

# Magnetostriction and magnetic texture to 100.75 Tesla in frustrated $\text{SrCu}_2(\text{BO}_3)_2$

Marcelo Jaime<sup>a,b,1</sup>, Ramzy Daou<sup>c,d</sup>, Scott A. Crooker<sup>a,b</sup>, Franziska Weickert<sup>b</sup>, Atsuko Uchida<sup>a,b</sup>, Adrian E. Feiguin<sup>e</sup>, Cristian D. Batista<sup>f</sup>, Hanna A. Dabkowska<sup>g</sup>, and Bruce D. Gaulin<sup>g,h</sup>

12404–12407 | PNAS | July 31, 2012 | vol. 109 | no. 31

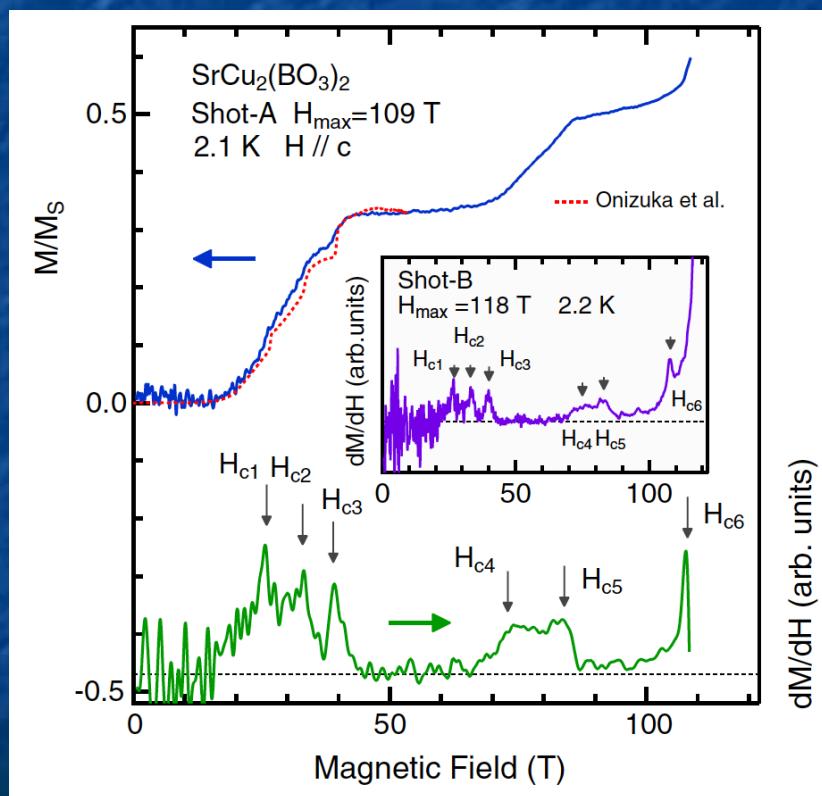
[www.pnas.org/cgi/doi/10.1073/pnas.1200743109](http://www.pnas.org/cgi/doi/10.1073/pnas.1200743109)



...,  $1/4$ ,  $1/3$ ,  $2/5$ ,  $1/2$

# Magnetization of $\text{SrCu}_2(\text{BO}_3)_2$ in Ultrahigh Magnetic Fields up to 118 T

Y. H. Matsuda,<sup>1,\*</sup> N. Abe,<sup>1</sup> S. Takeyama,<sup>1</sup> H. Kageyama,<sup>2</sup> P. Corboz,<sup>3</sup> A. Honecker,<sup>4,5</sup> S. R. Manmana,<sup>4</sup> G. R. Foltin,<sup>6</sup> K. P. Schmidt,<sup>6</sup> and F. Mila<sup>7</sup>



...1/3, 1/2

# Critical summary I

0, 1/9, 1/8, 2/15, 1/7, 1/6, 1/5, 2/9, 1/4, 1/3, 2/5, 1/2

Too many suspects!

# Critical summary II

Only pulsed field

Only magnetostriction

$0, 1\cancel{9}, 1/8, 2/15, 1\cancel{7}, 1/6, 1\cancel{5}, 2\cancel{9}, 1/4, 1/3, 2\cancel{5}, 1/2$

$0, 1/8, 2/15, 1/6, 1/4, 1/3, 1/2$

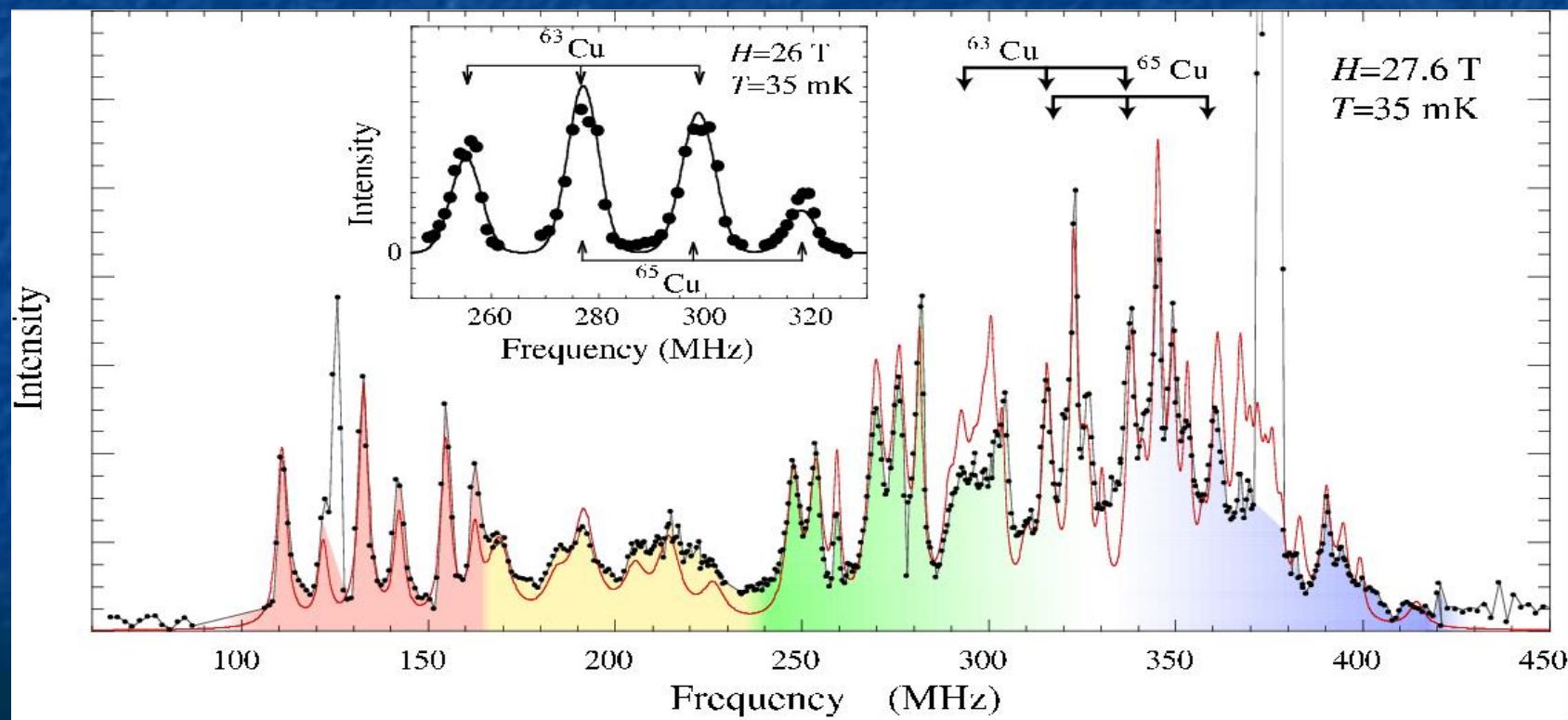
Quantum Hall Effect or broken spatial symmetry?

# Magnetic Superstructure in the Two-Dimensional Quantum Antiferromagnet $\text{SrCu}_2(\text{BO}_3)_2$

K. Kodama,<sup>1</sup> M. Takigawa,<sup>1\*</sup> M. Horvatić,<sup>2</sup> C. Berthier,<sup>2,3</sup>  
H. Kageyama,<sup>1</sup> Y. Ueda,<sup>1</sup> S. Miyahara,<sup>1,4</sup> F. Becca,<sup>4</sup> F. Mila<sup>4</sup>

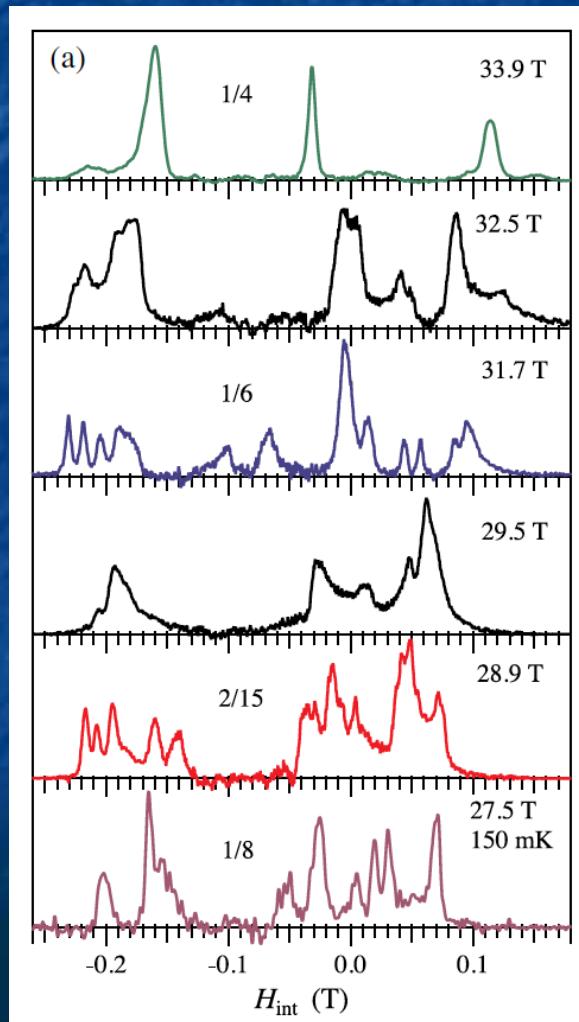
www.sciencemag.org SCIENCE VOL 298 11 OCTOBER 2002

Broken symmetry  
in 1/8 plateau



## Incomplete Devil's Staircase in the Magnetization Curve of $\text{SrCu}_2(\text{BO}_3)_2$

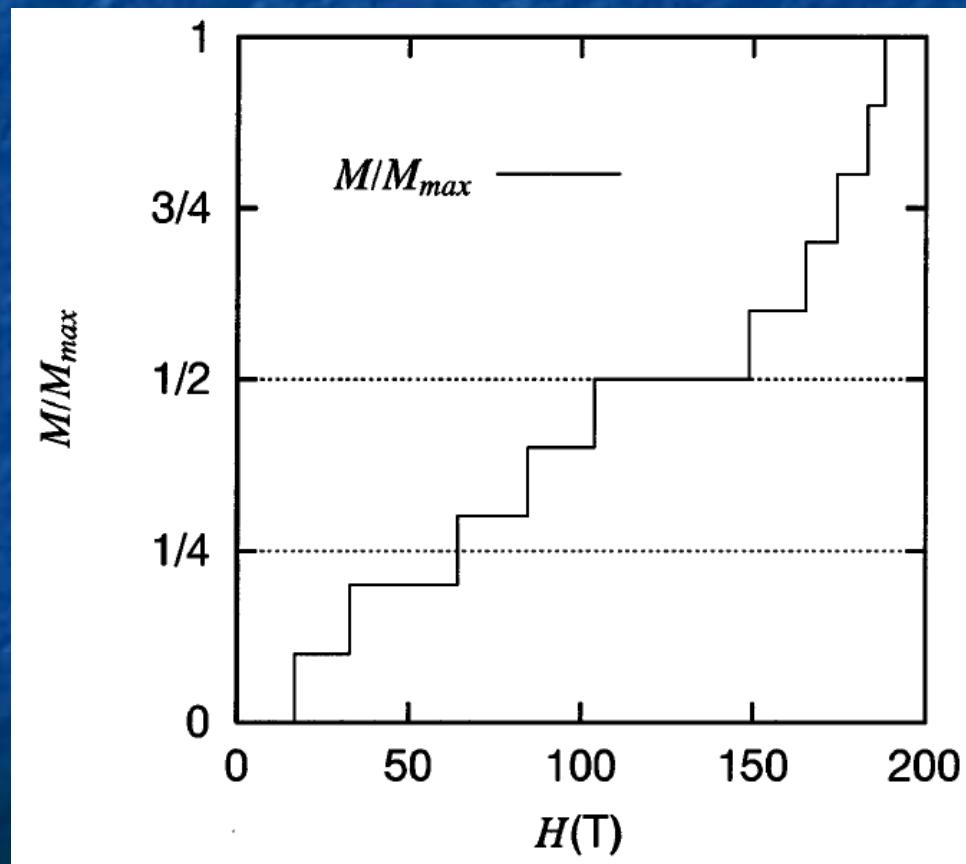
M. Takigawa,<sup>1,\*</sup> M. Horvatić,<sup>2</sup> T. Waki,<sup>3</sup> S. Krämer,<sup>2</sup> C. Berthier,<sup>2</sup> F. Lévy-Bertrand,<sup>2,†</sup> I. Sheikin,<sup>2</sup> H. Kageyama,<sup>4</sup> Y. Ueda,<sup>1</sup> and F. Mila<sup>5</sup>



Broken symmetry in  
 $1/8$ ,  $2/15$ ,  $1/6$  and  $1/4$  plateaus

# Exact Dimer Ground State of the Two Dimensional Heisenberg Spin System $\text{SrCu}_2(\text{BO}_3)_2$

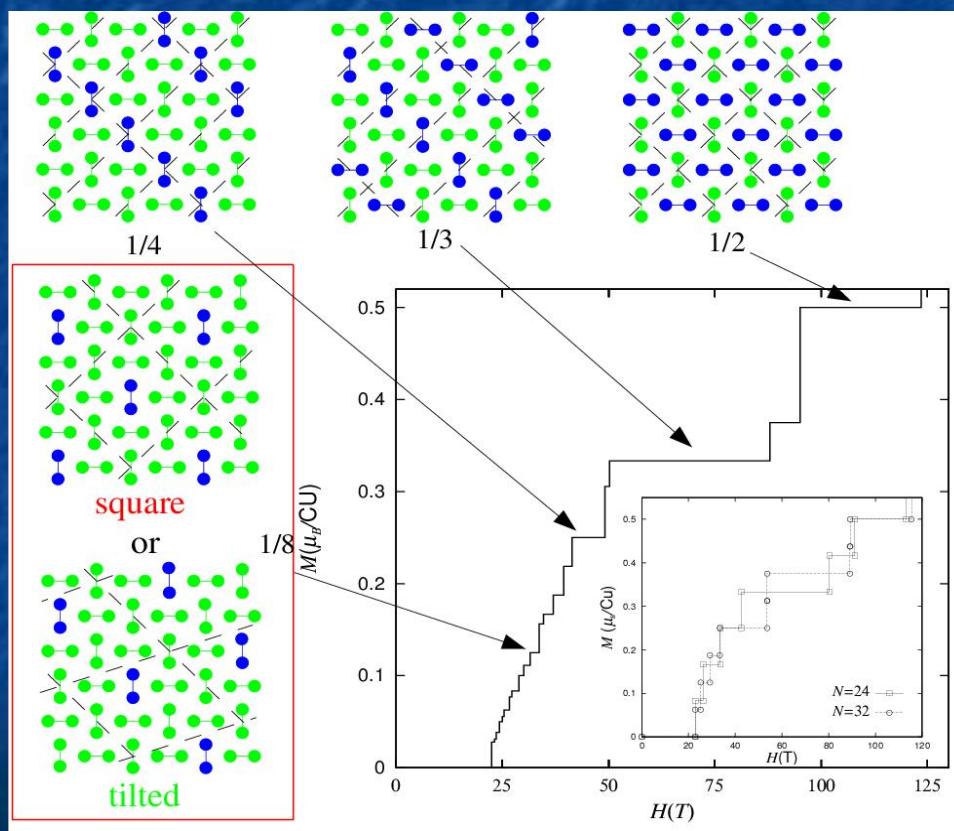
Shin Miyahara and Kazuo Ueda



- 20-site cluster  
→ finite-size effects

# Superstructures at magnetization plateaus in $\text{SrCu}_2(\text{BO}_3)_2$

Shin Miyahara and Kazuo Ueda



- Simple ansatz for long-range triplet-triplet interaction
- Many plateaus

# Magnetization plateaus of the Shastry-Sutherland model for $\text{SrCu}_2(\text{BO}_3)_2$ : Spin-density wave, supersolid, and bound states

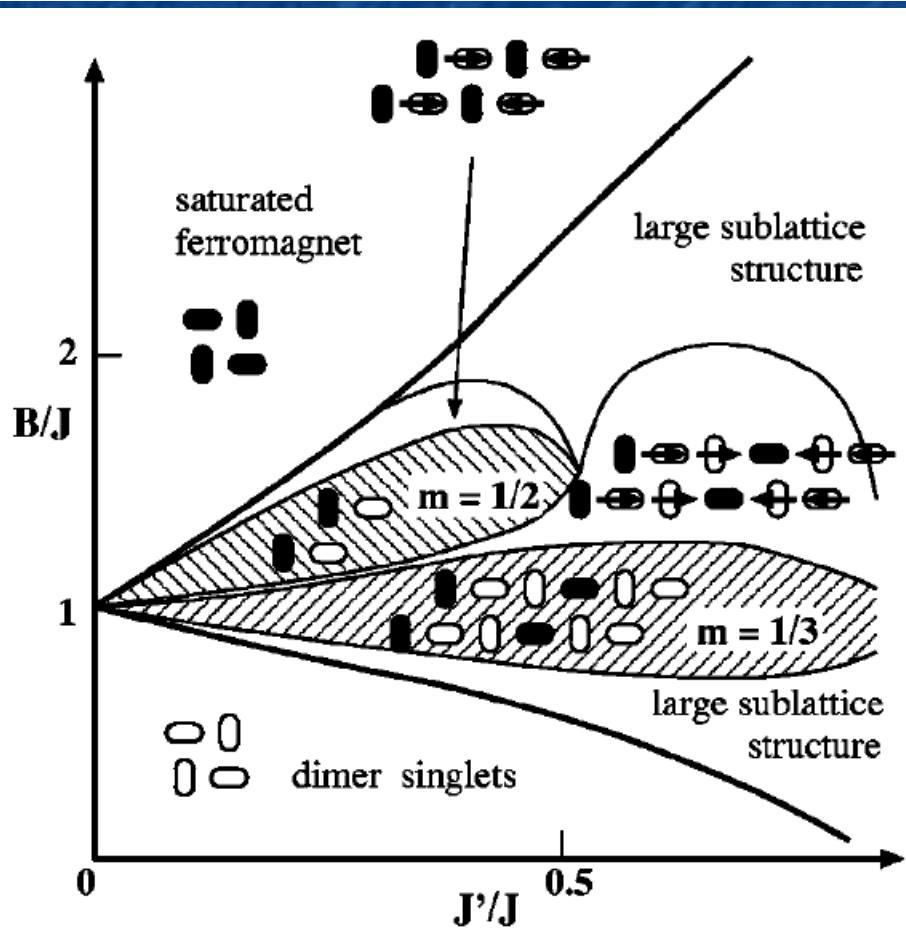
Tsutomu Momoi\*

*Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138*

Keisuke Totsuka

*Department of Physics, Kyushu University, Hakozaki, Higashi-ku, Fukuoka-shi 812-8581, Japan*

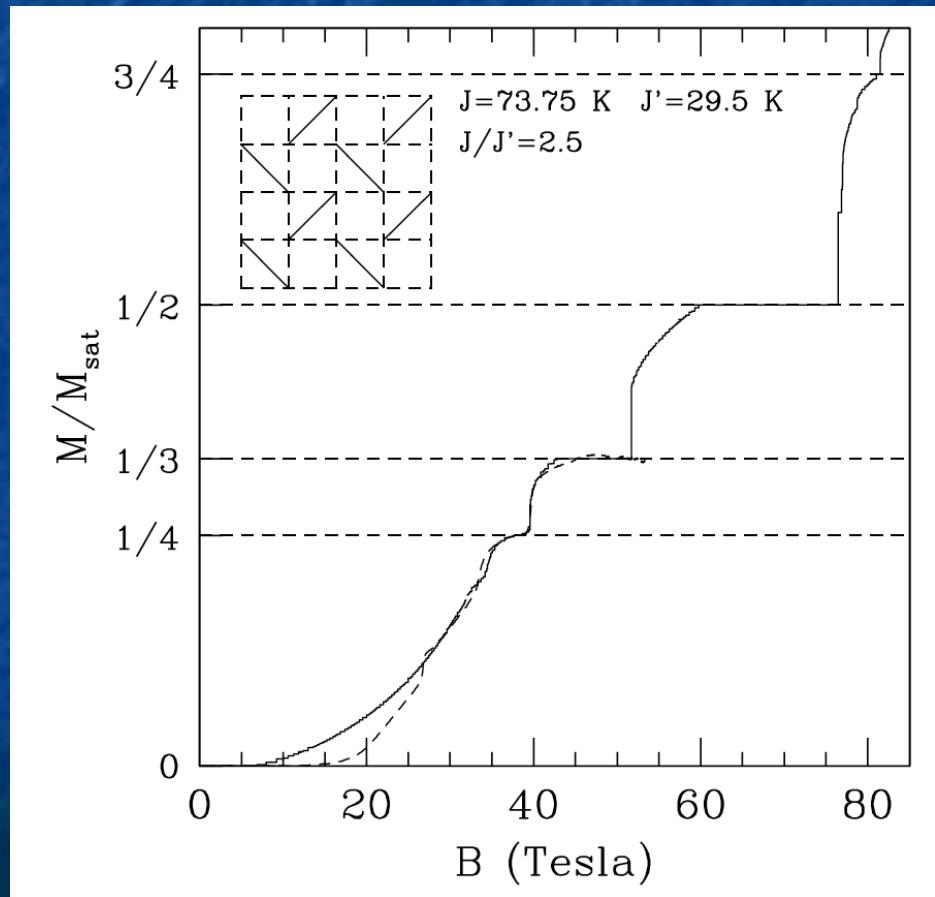
(Received 1 June 2000)



- Effective model to third order in  $(J'/J)^3$
- Only short-range triplet-triplet interaction
- Only  $1/3$  and  $1/2$  plateaus
- Spin-supersolids

## Magnetization Plateaus of $\text{SrCu}_2(\text{BO}_3)_2$ from a Chern-Simons Theory

G. Misguich,<sup>1</sup> Th. Jolicoeur,<sup>2</sup> and S. M. Girvin<sup>3,4</sup>

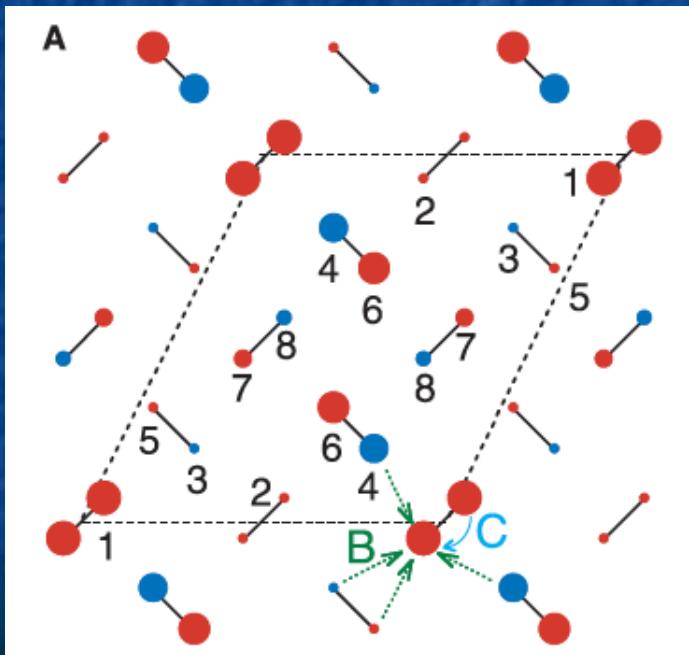


- No breaking of spatial symmetry
- Incompatible with NMR

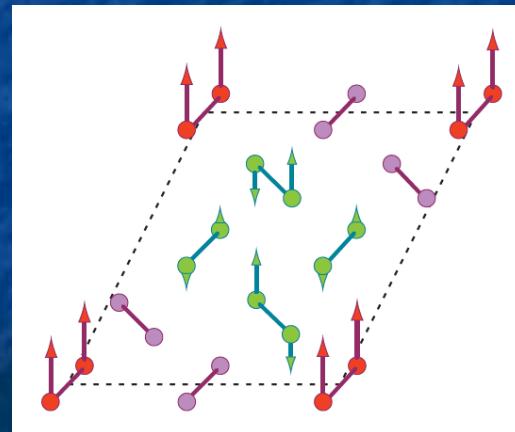
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www.sciencemag.org SCIENCE VOL 298 11 OCTOBER 2002



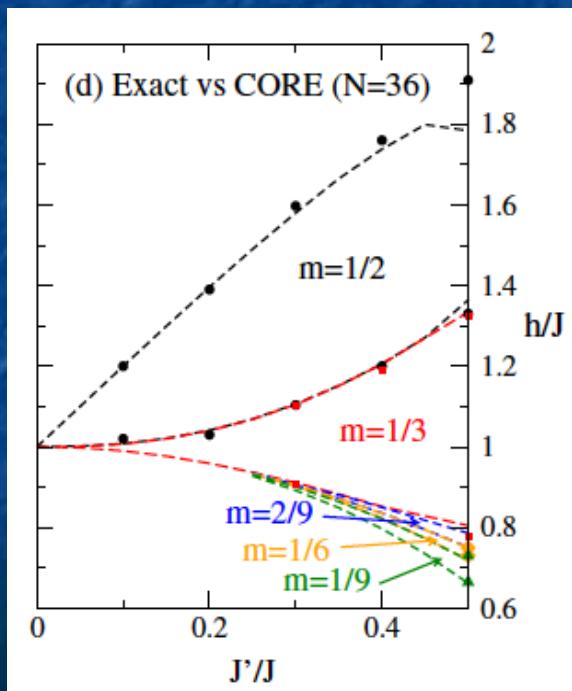
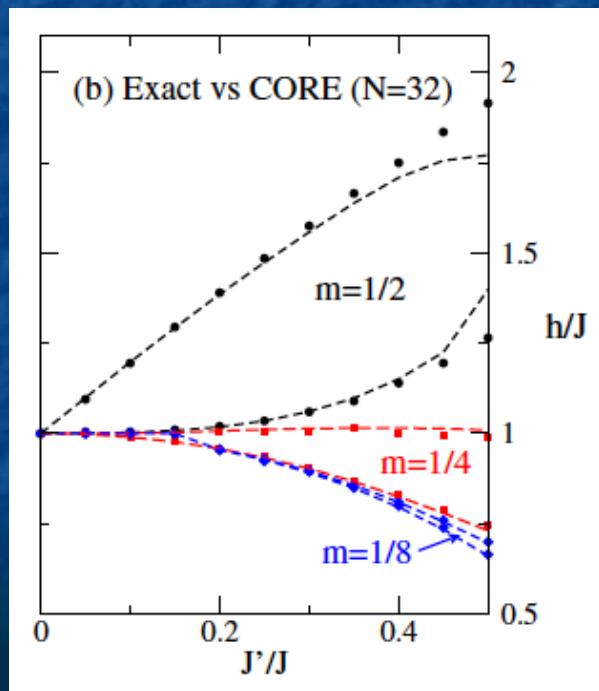
1/8 plateau  
= crystal of (dressed) triplets



## Effective Theory of Magnetization Plateaux in the Shastry-Sutherland Lattice

A. Abendschein<sup>1,2</sup> and S. Capponi<sup>1,2,\*</sup>

Effective model with CORE (Contractor Renormalization)  
Exact diagonalization on small clusters

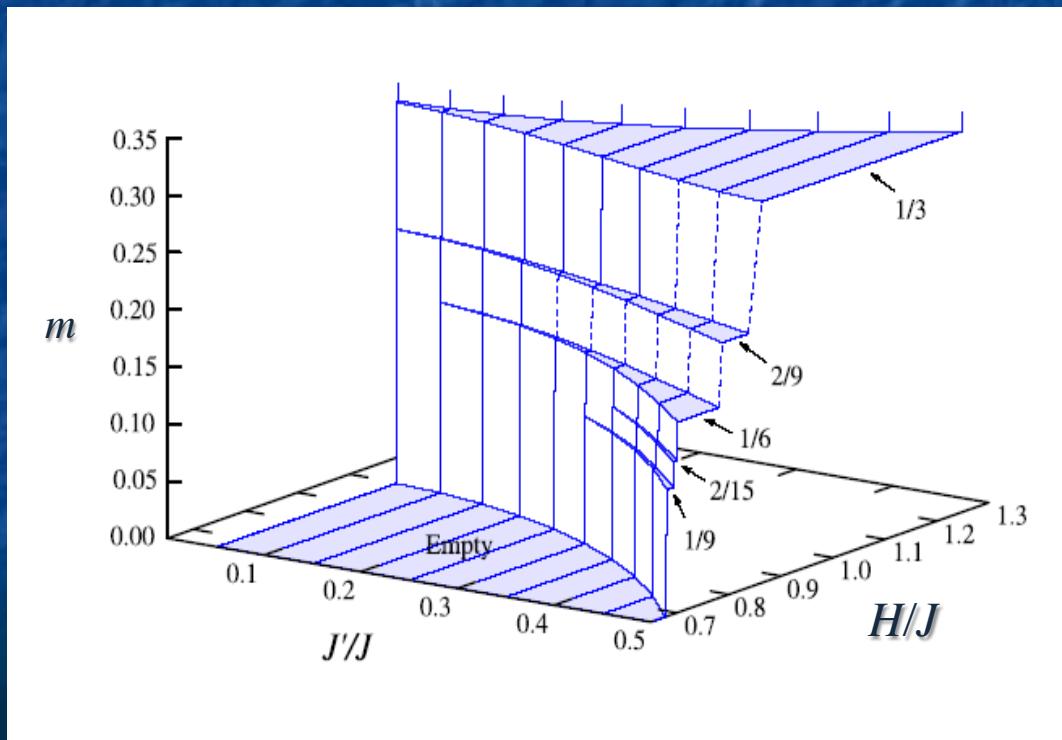


- Large finite-size effects
- No  $2/15$  plateau
- No information on plateau structures

## Theory of Magnetization Plateaux in the Shastry-Sutherland Model

J. Dorier,<sup>1</sup> K. P. Schmidt,<sup>2,\*</sup> and F. Mila<sup>1</sup>

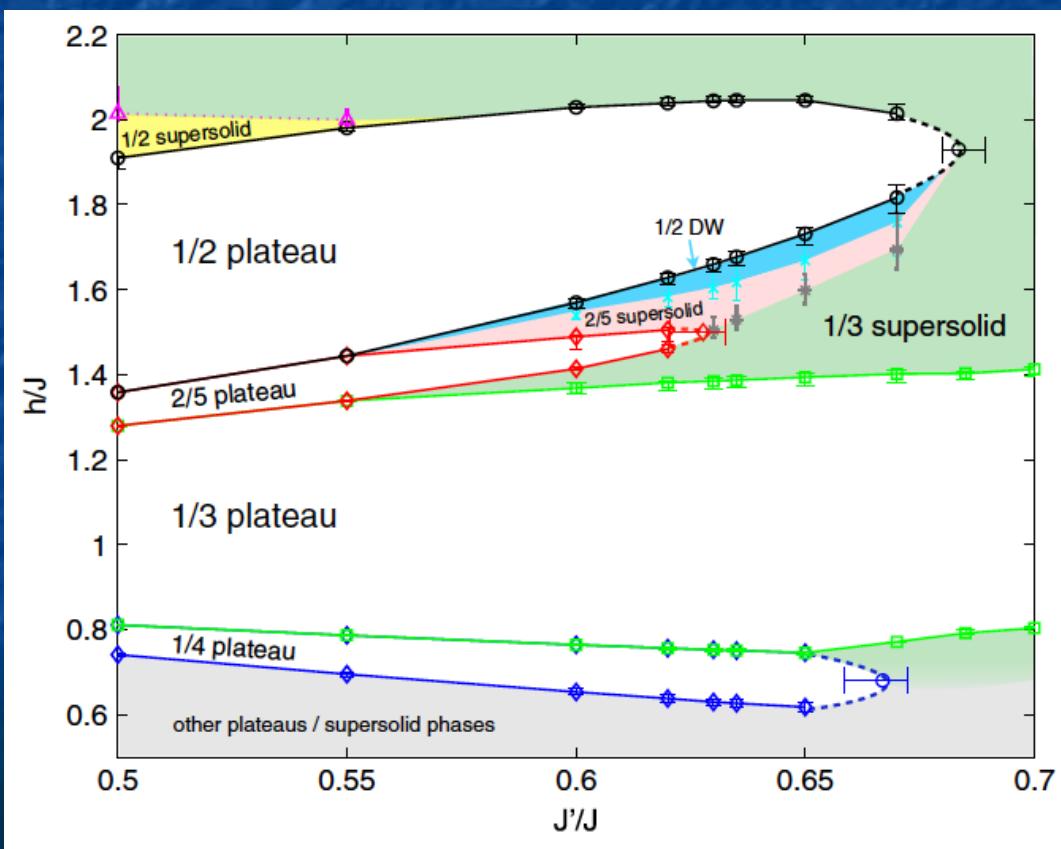
Long-range triplet-triplet interaction with high-order perturbation theory in  $J'/J$



Triplet crystals:  
not the right sequence  
(at least up to  $J'/J=0.5$ )

# Magnetization of $\text{SrCu}_2(\text{BO}_3)_2$ in Ultrahigh Magnetic Fields up to 118 T

Y. H. Matsuda,<sup>1,\*</sup> N. Abe,<sup>1</sup> S. Takeyama,<sup>1</sup> H. Kageyama,<sup>2</sup> P. Corboz,<sup>3</sup> A. Honecker,<sup>4,5</sup> S. R. Manmana,<sup>4</sup> G. R. Foltin,<sup>6</sup> K. P. Schmidt,<sup>6</sup> and F. Mila<sup>7</sup>



iPEPS

A variational approach  
derived from quantum  
information



Verstraete & Cirac, 2004

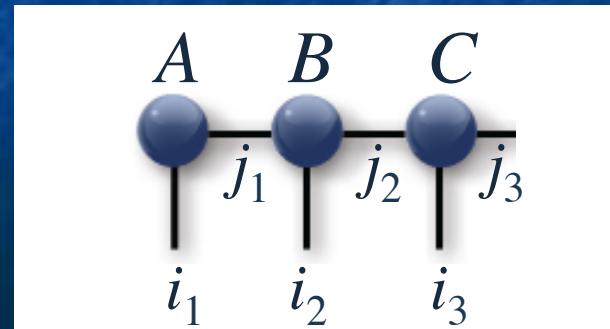
# Tensor network ansatz

$$|\psi\rangle = \sum_{i_1 \dots i_N} c_{i_1 \dots i_N} |i_1\rangle \otimes \dots \otimes |i_N\rangle$$

$c_{i_1 \dots i_N} \simeq \text{trace over a product of tensors}$

Example: Matrix product state in 1D

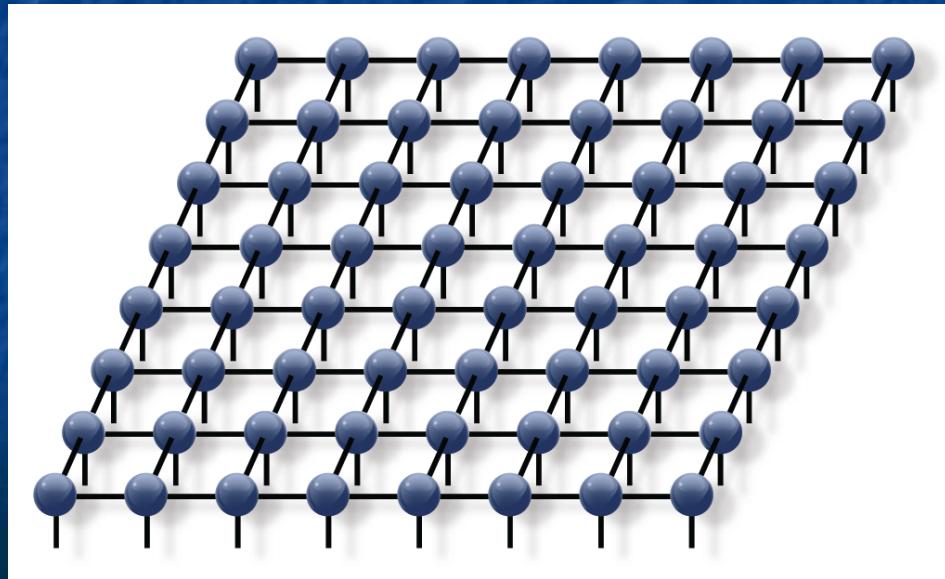
$$c_{i_1 i_2 i_3 \dots} \simeq \sum_{j_1 j_2 \dots} A_{i_1}^{j_1} B_{i_2}^{j_1 j_2} C_{i_3}^{j_2 j_3} \dots$$



# Generalization to 2D

PEPS = product of entangled pair states

Verstraete and Cirac, 2004



$A_i^{j_1 j_2 j_3 j_4}$  = rank-5 tensor

$j_1, j_2, j_3, j_4 = 1, \dots, D$

# Variational approach

- PEPS: minimize the energy w.r.t. tensor elements
- Other schemes: renormalization (MERA,...)
- Advantage:  $\text{dim}=\text{pol}(D,N)$ , not  $\exp(N)$
- Why can it work?  
→ reproduces the ‘area law’ for the entanglement entropy in the GS of a local Hamiltonian

$$S = -\text{tr} (\rho_A \log \rho_A) \sim \partial A$$

- How large should D be? It depends...

# Results on frustrated magnets

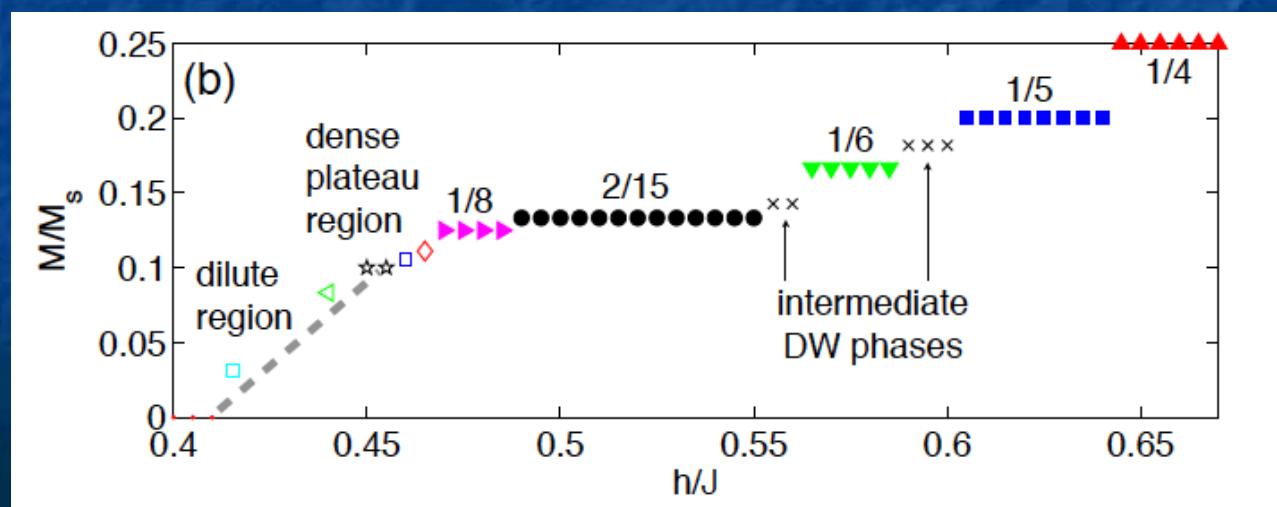
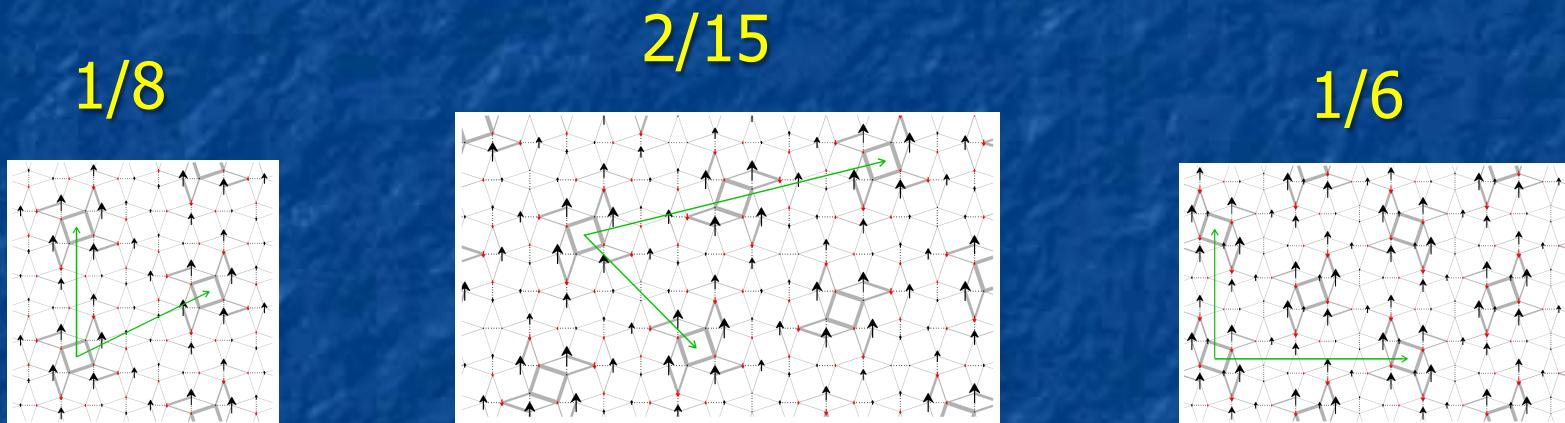
- Kagome: difficult
- $J_1$ - $J_2$  on square lattice: intermediate phase around  $J_2/J_1=1/2$ , but no consensus on its nature
  - Plaquette?  
Yu and Kao, PRB 2012
  - Topological spin liquid?  
Mezzacapo, PRB 2012; Wang, Gu, Verstraete, Wen 2012
- Checkerboard: intermediate plaquette phase, in agreement with previous results  
Chan, Han, Duan, PRB 2011

# Shastry-Sutherland

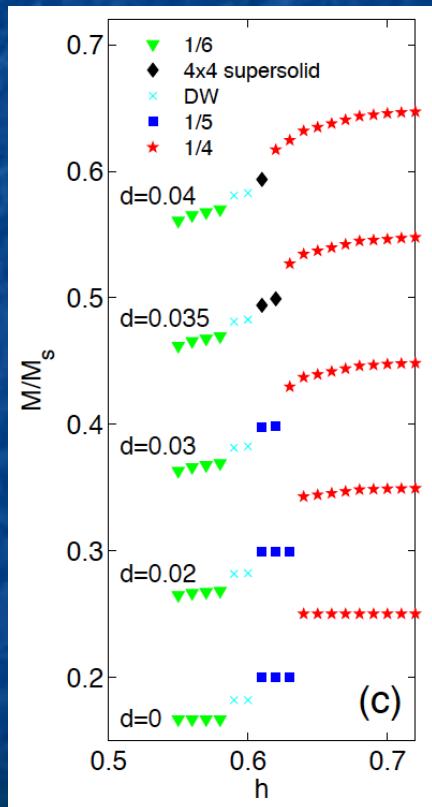
- Isacsson and Syljuasen, PRE 2006  
→ D=2, no intermediate phase in zero field
- Lou, Suzuki, Harada, Kawashima, arxiv 2012  
→ MERA  
→ intermediate plaquette phase  
→ plateaux and supersolid phases in a field
- P. Corboz and FM, PRB 2013, PRL 2014  
→ plateau sequence of  $\text{SrCu}_2(\text{BO}_3)_2$

# Crystals of Bound States in the Magnetization Plateaus of the Shastry-Sutherland Model

Philippe Corboz<sup>1</sup> and Frédéric Mila<sup>2</sup>



# Why no 1/5 plateau?

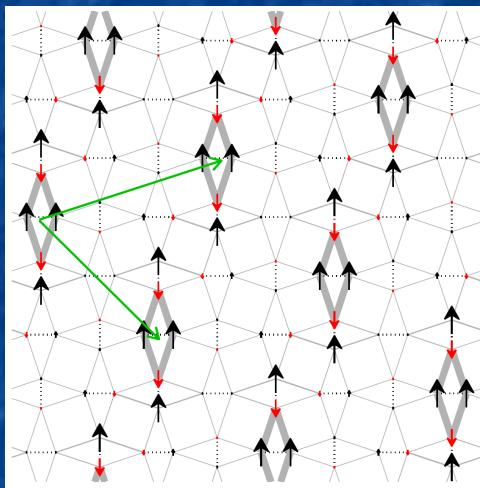


Killed by Dzyaloshinskii-Moriya  
Interaction

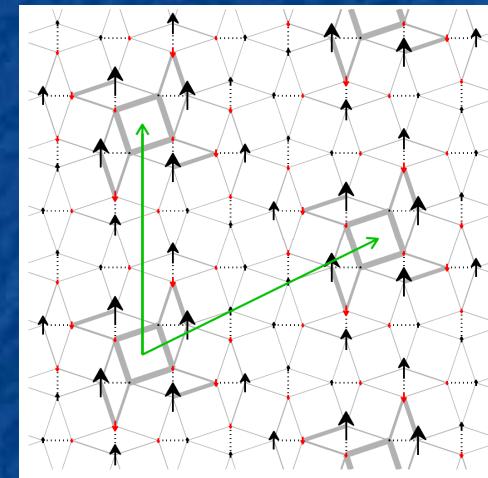
P. Corboz, FM, PRL 2014

# Triplets versus bound-states

Triplet crystal



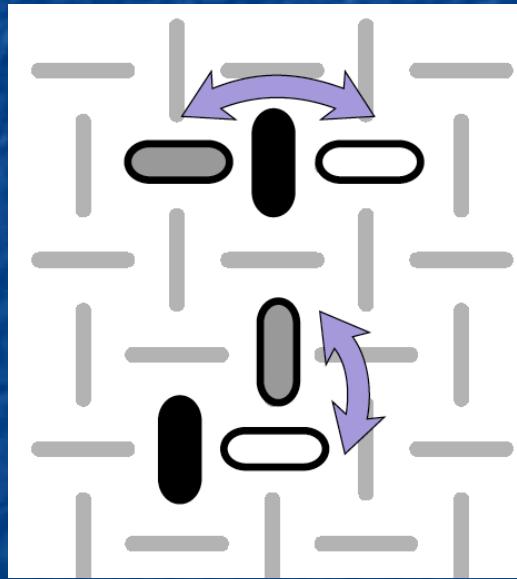
Bound state crystal



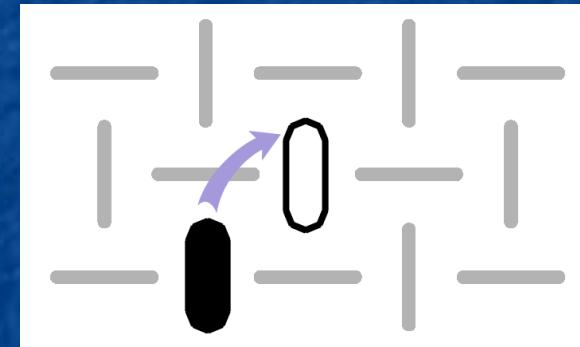
- Why are bound state crystals favoured?
- Why did it take 15 years to identify them?

# Why are bound states favoured?

Correlated hopping



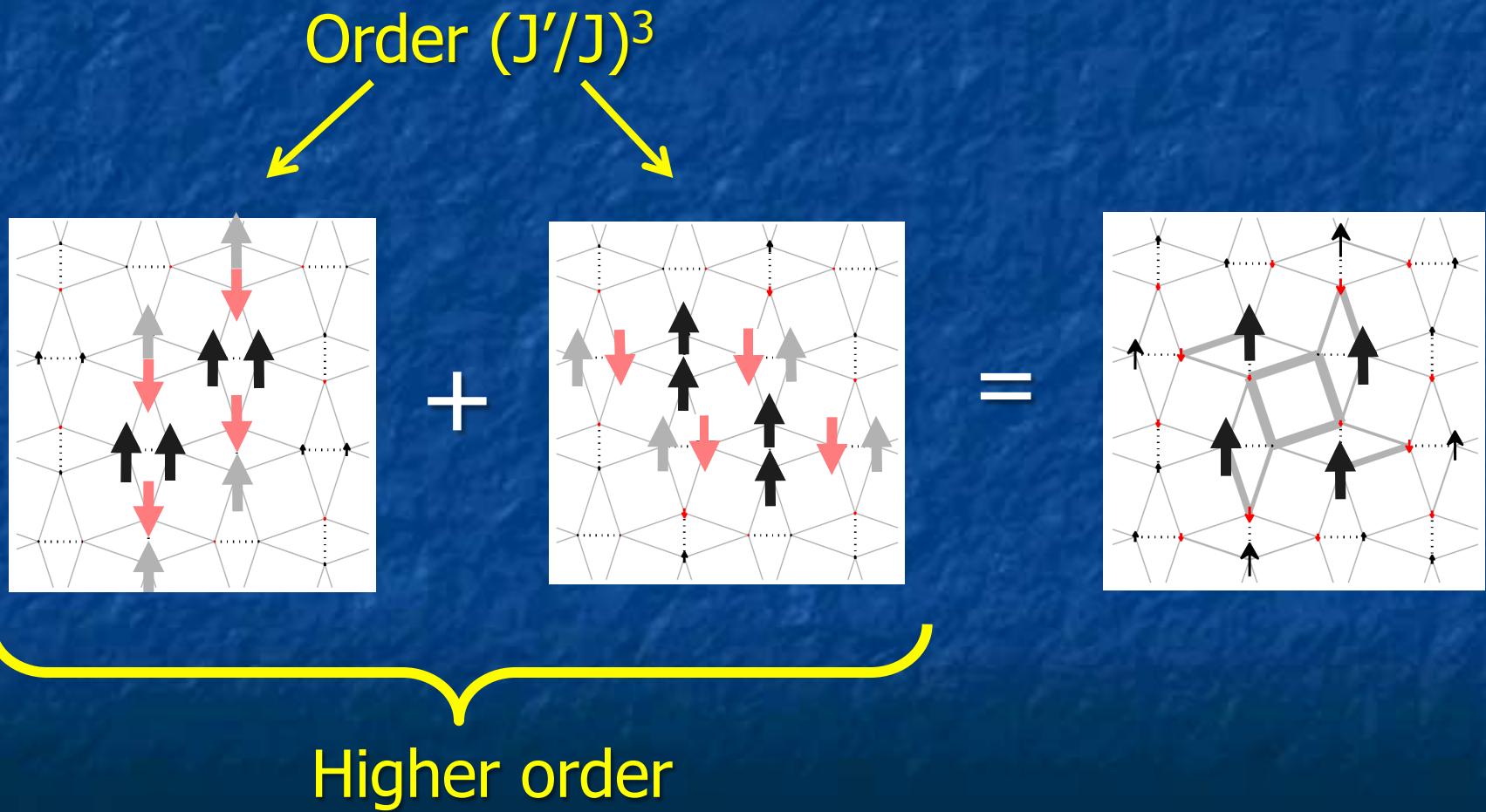
Single particle hopping



$$\alpha (J'/J)^6$$

$\alpha (J'/J)^2 \rightarrow$  gain in kinetic energy

# Boundstates as pinwheels



# Magnetization plateaus of the Shastry-Sutherland model for $\text{SrCu}_2(\text{BO}_3)_2$ : Spin-density wave, supersolid, and bound states

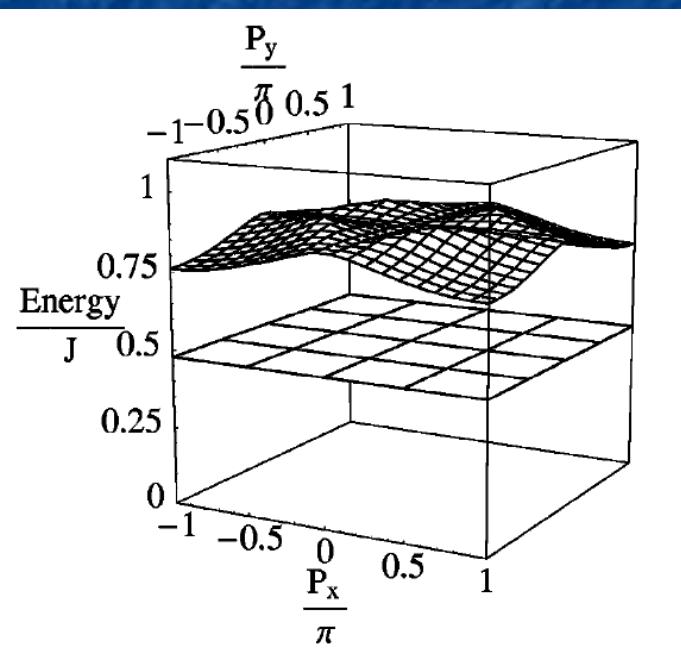
Tsutomu Momoi\*

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*Department of Physics, Kyushu University, Hakozaki, Higashi-ku, Fukuoka-shi 812-8581, Japan*

(Received 1 June 2000)



Bound state dispersion:  
minimum at the zone corner

# Magnetization plateaus of the Shastry-Sutherland model for $\text{SrCu}_2(\text{BO}_3)_2$ : Spin-density wave, supersolid, and bound states

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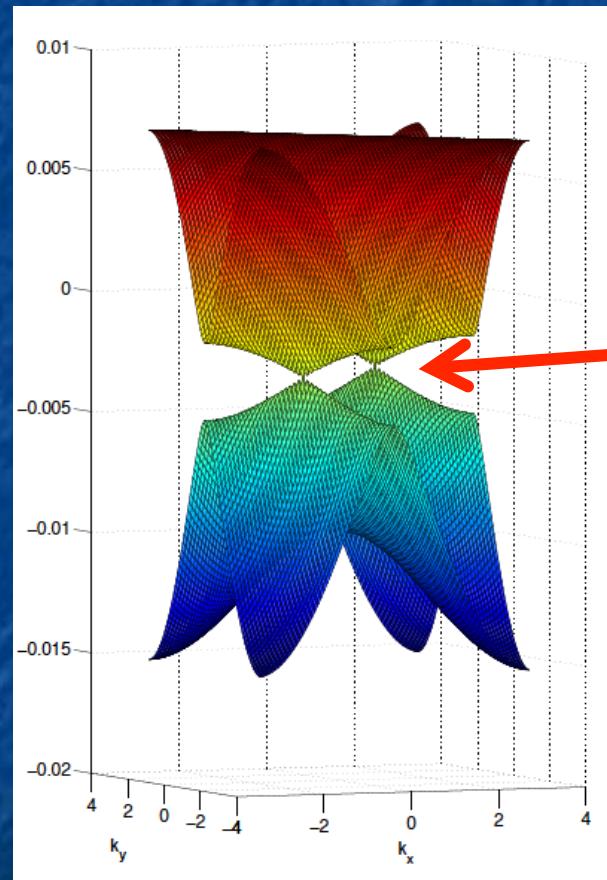
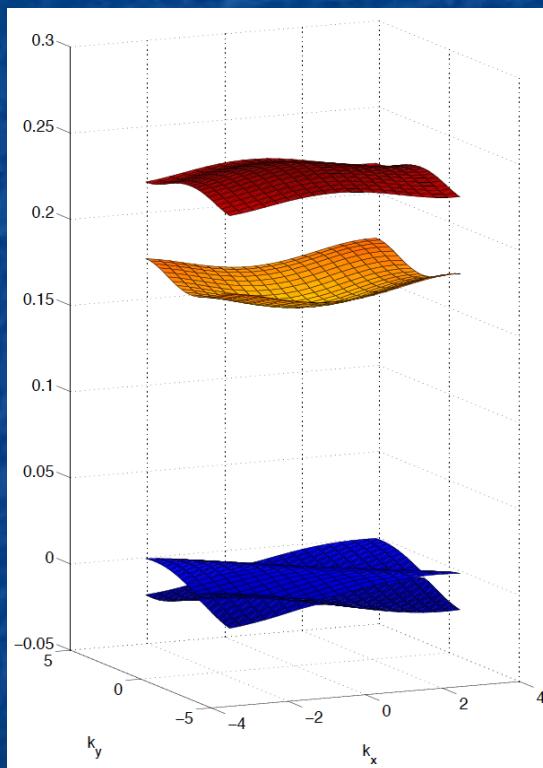
*Department of Physics, Kyushu University, Hakozaki, Higashi-ku, Fukuoka-shi 812-8581, Japan*

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If bound states tend to crystallize, bound states loose kinetic energy and then they will be unbound by repulsive interaction between triplets.

Is that really true?

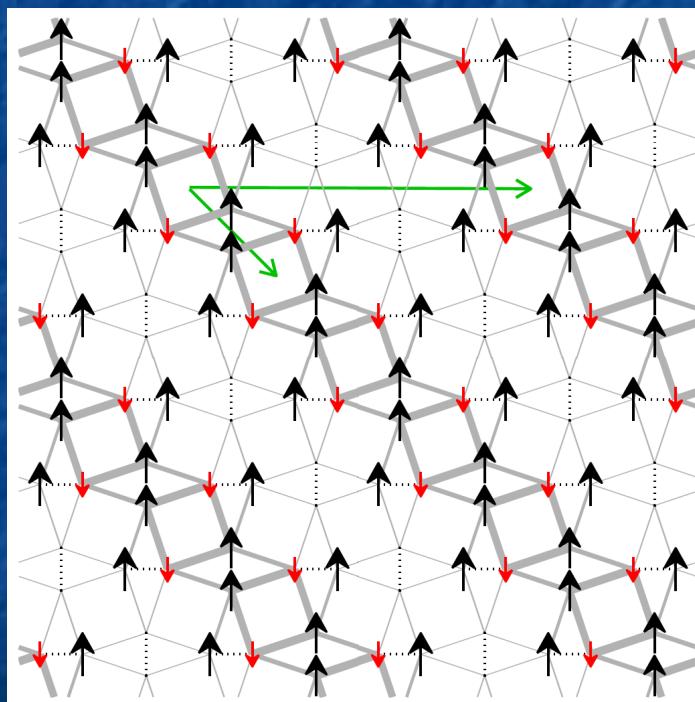
# Localized bound state



Negative  
reference  
energy!

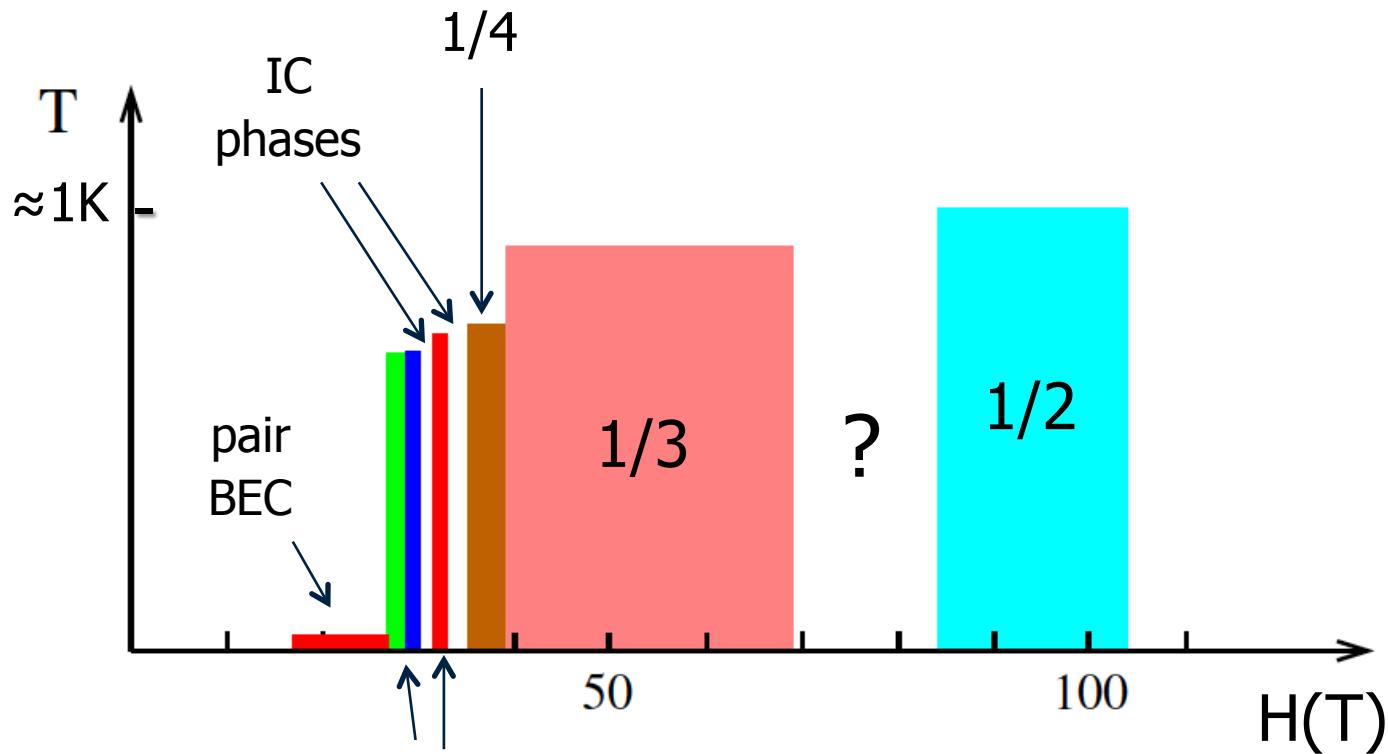
# Dense triplets as bound states

1/4 plateau



1/4 , 1/3 and 1/2 plateaux:  
can be seen both as crystals  
of triplets and as crystals of  
bound states.

# Tentative phase diagram



Pair BEC: Bendjama, Kumar, FM, PRL 2005;  $T_c$ : Takigawa et al, PRL 2008

# Conclusions on $\text{SrCu}_2(\text{BO}_3)_2$

- Natural plateau sequence
  - $1/8, 2/15, 1/6$ : crystals of bound states
  - new predictions for their symmetry
- In progress
  - revisit effective triplet model
  - compare predictions with NMR
- Still to be done
  - X-ray or neutron in high field ( $>27 \text{ T}$ )?

# Conclusions

- Magnetization of frustrated magnets
  - very rich
  - reveals several fundamental aspects of strongly correlated systems

