Frustration and Entanglement

ICTP international workshop on Current Trends in Frustrated Magnetism SPS, JNU



Ujjwal Sen HRI, Allahabad

Outline



- 1. Understanding entanglement
- 2. Entanglement in many-body physics
- 3. What is frustration?
- 4. Characterizing "classical" frustration in q systems
- 5. Frustration and Entanglement
 - I. Area Law
 - II. Genuine multipartite entanglement
- 6. End remarks

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Understanding entanglement LOCC paradigm in quantum info

 If the state is shared between two or more parties, the parties would only be able to act locally.
 Allowed operations: LOCC.



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Not this!!





What do we mean by LOCC?



• Alice makes a measurement and communicates her result to Bob (say, by a phone call).





What do we mean by LOCC?



- Alice makes a measurement and communicates her result to Bob (say, by a phone call).
- Then depending on her result, Bob will make his measurement and communicate his result to Alice.
- And so on.





• Quantum states that can be prepared by $LOCC \rightarrow Separable$ states.

• Otherwise \rightarrow Entangled states.





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- How do they look like?





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- Separable *pure* states: products over pure states of individual systems.



- Quantum states that can be prepared by $LOCC \rightarrow Separable states.$
- How do they look like? Mathematically?
- Separable states: mixtures of products over pure states of individual systems.



Which "entanglements" can we *compute*? Circa 2000

• Nielsen, Preskill, Wootters et al.









Which "entanglements" can we *compute*? Circa 2000

• Nielsen, Preskill, Wootters et al.

Idea of using entanglement-like concepts in quantum many-body phenomena was put forward.



- Nielsen, Preskill, Wootters et al.
- Osborne and Nielsen, QIP'02, PRA'02
- Osterloh, Amico, Falci, Fazio, Nature'02





Which "entanglements" can we *compute*?

To see the behavior of entanglement in real systems, it is *not* sufficient

to understand an entanglement measure conceptually.



Which "entanglements" can we *compute*?

To see the behavior of entanglement in real systems, it is *not* sufficient to understand an entanglement measure conceptually. We must also be able to *compute* it for the states of the real systems.



Which "entanglements" can we *compute*?

• Bipartite states.



- Bipartite states.
- For mixed two-party states, only entanglement of formation of two-qubit states.



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For mixed two party states only

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Modulo certain additivity problems.



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Modulo certain additivity problems.



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- For mixed two-party states, only entanglement of formation of two-qubit states.
- In higher dimensions, logarithmic negativity can be calculated. But it cannot detect bound entanglement.



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Logneg of a two-party state is $log_2(2N + 1)$.



Which "entanglements" can we *compute*?

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Logneg of a two-party state is $log_2(2N + 1)$.

N = sum of mod of negative eigenvalues in partial transpose of state.



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- For **pure** two-party states, local von Neumann entropy is a "good" measure of entanglement, and is computable.



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- For **mixed** two-party states, only entanglement of formation of two-qubit states.
- For **pure** two-party states, local von Neumann entropy is a "good" measure of entanglement, and is computable.

Possible in arbitrary dimensions.



Which "entanglements" can we *compute*?

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This sets the stage for the QI – many-body interface.

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Indeed, two of the main directions of study are 1. EoF of reduced densities of spin-1/2 ground states 2. Scaling of local entropy in ground state partitions

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Understanding entanglement

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• Many notions available.

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a. Geometric measure

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a. Geometric measure

Wei, Goldbart, PRA'03 Balsone, DellAnno, DeSiene, Illuminatti, PRA'08 +

- Many *notions* available.
- However, not all r computable.

a. Geometric measureb. Global measure

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Meyer, Wallach, JMP'02

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- a. Geometric measure
- b. Global measure
- c. Generalized geometric measure

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A. Sen(De), US, PRA'10

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S Sachdev, QPT



• Transitions at zero temperature.





- Transitions at zero temperature.
- Implying, transition not temp. driven.





- Transitions at zero temperature.
- Implying, transition not temp. driven.
- Driven by system parameter, like a magnetic field.





Typical situation:

• H = H(int) + a H(field)





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- H = H(int) + a H(field)
- Ground state of H





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- "a" can be changed.

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Typical situation:

- H = H(int) + a H(field)
- Ground state of H \leftarrow guarantees T=0
- GS depends on "a".
- "a" can be changed.
- Nonanalyticity appears in some physical quantity as "a" is changed.

S Sachdev, QPT













The reduced state is a two-qubit state.

Spin-1/2 Chain



The prescription:



*The prescription:*1. Find ground state of spin-1/2 system



*The prescription:*1. Find ground state of spin-1/2 system2. Remove all spins except two NNs



The prescription:

- 1. Find ground state of spin-1/2 system
- 2. Remove all spins except two NNs
- 3. Find EoF of resulting two-site density



The prescription:

- 1. Find ground state of spin-1/2 system
- 2. Remove all spins except two NNs
- 3. Find EoF of resulting two-site density
- 4. Investigate it wrt the relevant system parameter





 $\Sigma J [(1 + \gamma) S_x^i S_x^{i+1} (1 - \gamma) S_v^i S_v^{i+1}] - a S_z^i$





$$\Sigma J [(1 + \gamma) S_x^i S_x^{i+1} (1 - \gamma) S_y^i S_y^{i+1}] - a S_z^i$$

S are half of Pauli matrices.





$$\Sigma J [(1 + \gamma) S_x^i S_x^{i+1} (1 - \gamma) S_y^i S_y^{i+1}] - a S_z^i$$

Quantum phase transition at h=1.



For $\gamma = 1$: Transverse Ising Model.



$$\Sigma J [(1 + \gamma) S_x^i S_x^{i+1} (1 - \gamma) S_y^i S_y^{i+1}] - a S_z^i$$

Quantum phase transition at h=1.

Entanglement in many-body physics



Linking QI with concepts in quantum statistical mechanics and quantum phase transitions.

Near QPT in 1D transverse Ising model, 2-site entanglement remains short ranged, while 2-site

Entanglement, however, does show signs of criticality.



Entanglement in many-body physics







The prescription:

- 1. Find ground state of spin-1/2 system
- 2. Remove all spins except two NNs
- 3. Find EoF of resulting two-site density
- 4. Investigate it wrt the relevant system parameter


Entanglement in many-body physics Two-site densities

Why *ground* state?

The prescription:

- 1. Find ground state of spin-1/2 system
- 2. Remove all spins except two NNs
- 3. Find EoF of resulting two-site density
- 4. Investigate it wrt the relevant system parameter



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Entanglement in many-body physics Two-site densities

The prescription:

Find ground s
Remove all s
Find EoF of r
Investigate it

Guarantees that there are no thermal effects.

Why ground state?



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Entanglement in many-body physics Two-site densities

The prescription:

Find ground s
Remove all s
Find EoF of r
Investigate it

Thermal states, time-evolved states also considered.

Why ground state?



Entanglement in many-body physics Two-site densities

The prescription:

1. Find ground state of spin-1/2 system

- 2. Remove all spins except two NNs
- 3. Find EoF of resulting two-site density
- 4. Investigate it wrt the relevant system parameter

Why NN?



Entanglement in many-body physics Two-site densities

The prescription:

Find ground s
Remove all s
Find EoF of r
Investigate it

Why NN?

In many instances, but NOT all, NNN and so on have little to no entanglement. Entanglement in many-body physics Multiparty entanglement

Multiparty entanglement detects QPT

- a. Geometric measure
- b. Global measure
- c. Generalized geometric measure (GGM)

Entanglement in many-body physics Geometric measure detects QPT



Wei, Das, Mukhopadhyay, Vishveshwara, Goldbart, PRA'05

Entanglement in many-body physics

Global measure of multipartite entanglement detects QPT



deOliviera, Rigolin, deOliviera, PRA'06

Entanglement in many-body physics GGM detects QPT



Entanglement in many-body physics GGM –1/2 GGM detects QPT



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From a classical perspective



From a classical perspective

Consider an Ising model:

 $\mathcal{H} = J \sum \sigma_i \sigma_j; \quad J > 0$



From a classical perspective

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 $\mathcal{H} = J \sum \sigma_i \sigma_j; \quad J > 0$





From a classical perspective

Failure to have spin configuration to minimize individual interaction terms



From a *quantum* perspective



From a quantum perspective

Draw a parallel



From a quantum perspective

Classical frustration: spin configuration

Quantum frustration: GSs of two terms not same

$$\mathcal{H} = \mathcal{H}_{loc} + \mathcal{H}_{int}$$

Dawson and Nielsen, PRA 69, 052316 (2004)



Classical spin configuration



Cannot get optimal spin configuration





Quantum non-commutativity



Cannot get optimal spin \longleftrightarrow GSs of two terms not same configuration

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Characterizing "classical" frustration in q systems

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Classical frustration

Characterizing "classical" frustration in q systems





Sen(De), US, Dziarmaga, Sanpera, Lewenstein, PRL'08 Jindal, Rane, Dhar, Sen(De), US, PRA'14



• Given H, $|\Gamma\rangle$,



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replace one-body, two-body etc. in H by Ising ones, i.e. by σ_{i}^{z} or $\sigma_{i}^{z}\sigma_{i}^{z}$ etc.



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 $Find \; H^{\rm I}$

• Given H, $|\Gamma\rangle$,

replace one-body, two-body etc. in H by Ising ones, i.e. by σ_{i}^{z} or $\sigma_{i}^{z}\sigma_{j}^{z}$ etc. Find H^I

Frustrated Non-Frustrated $H^{I} = \sum_{k} H^{k}_{f} + \sum_{l} H^{l}_{nf}$ $\Phi = avg \quad \frac{\sum_{k} \langle \Gamma | H^{k}_{f} | \Gamma \rangle}{\sum_{l} |\langle \Gamma | H^{l}_{nf} | \Gamma \rangle|}$

Ising model: $H=J \Sigma \sigma_i^z \sigma_j^z$ with J>0

 $H^{I} = \sum_{k} H^{k}{}_{f} + \sum_{l} H^{l}{}_{nf}$

Characterizing "classical" frustration in q systems *Frustration degree* Ising model: $H=J \Sigma \sigma_{i}^{z} \sigma_{i}^{z}$

 $H^{I} = \sum_{k} H^{k}{}_{f} + \sum_{l} H^{l}{}_{nf}$ 3 $\sigma_1^z \sigma_2^z | \downarrow \rangle \rightarrow 1$

with J>0
Characterizing "classical" frustration in q systems *Frustration degree*

Ising model: $H=J \Sigma \sigma_i^z \sigma_j^z$ with J>0





Characterizing "classical" frustration in q systems *Frustration degree* Ising model: $H=J \Sigma \sigma_{i}^{z} \sigma_{i}^{z}$ with J>0 $H^{I} = \sum_{k} H^{k}{}_{f} + \sum_{l} H^{l}{}_{nf}$

 $\begin{array}{ccc} \sigma_{1}^{z} \sigma_{2}^{z} |\downarrow\rangle & \rightarrow 1 \\ \sigma_{2}^{z} \sigma_{3}^{z} |\downarrow\rangle & \rightarrow -1 \end{array}$

3

Characterizing "classical" frustration in q systems *Frustration degree* Ising model: $H=J \Sigma \sigma^{z}_{i} \sigma^{z}_{i}$ with J>0 $H^{I} = \sum_{k} H^{k}{}_{f} + \sum_{l} H^{l}{}_{nf}$

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We r talking abt interacting systems.





Would be true (trivially) if ...





Boundary particles are *pure* entangled states.





Boundary particles are *pure* entangled states. Plus no long-range entangled pairs.





Typical situation is far from being such.





Typical situation is far from being such. Usually intricately multiparty quantum correlated.





$$\mathrm{S}(\rho_\mathrm{L})\sim\mathrm{L}^{\mathrm{d}\text{-}1}$$

L: characteristic length of A





$$S(\rho_L) \sim L^{d-1}$$

L: characteristic length of A





Block entanglement: $E(|\Psi\rangle_{L:N-L})$

$$E(|\Psi\rangle_{L:N-L}) = S(\rho_L)$$





$$E(|\Psi\rangle_{L:N-L}) = S(\rho_L) \sim L^{d-1} = constant$$

away from criticality









$$E(|\Psi\rangle_{L:N-L}) = S(\rho_L) \sim \ln L$$

at criticality





at co



Lot of progress in different directions.

Lot of progress in different directions. A case study: Frustrated systems



Main Thesis

Highly frustrated systems do not follow area law



Main Thesis

Highly frustrated systems do not follow area law

while

Weakly frustrated systems follow same area law as nonfrustrated systems away from criticality

Sen(De), US, Dziarmaga, Sanpera, Lewenstein, PRL'08 Jindal, Rane, Dhar, Sen(De), US, PRA'14



Area Law for frustrated systems

- 1. Long range Ising model
- 2. Majumdar Ghosh model
- 3. Shastry-Sutherland model
- 4. Ising chain with NN interactions

Cooling/Quenching Method



 $|\Phi\rangle_{in} \equiv |\psi\rangle_1 \otimes |\psi\rangle_2 \otimes |\psi\rangle_3 \otimes ... \otimes |\psi\rangle_N$

Cooling/Quenching Method



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> Project $|\Phi\rangle_{in}$ onto the ground state space of the model.

$$|\Phi\rangle_{\rm f} = (\sum |\Gamma\rangle_i \langle \Gamma|) |\Phi\rangle_{\rm in}$$

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Project $|\Phi\rangle_{in}$ onto the ground state space of the model.

$$|\Phi\rangle_{\rm f} = (\sum |\Gamma\rangle_i \langle \Gamma|) |\Phi\rangle_{\rm in}$$

 \succ Calculate $E_{N/2:N/2}(|\Phi\rangle_f)$.

Solution Maximize $E_{N/2:N/2}(|\Phi\rangle_f)$ over all choices of the initial state.



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After quenching:

 $|\psi\rangle$ = superposition of all vectors with m | 0 \rangle s and m | 1 \rangle s





After quenching:

 $|\psi\rangle$ = superposition of all vectors with m | 0 \rangle s and m | 1 \rangle s

$$E_{k:2m-k} = \frac{1}{2} \log k$$





After quenching:

 $|\psi\rangle$ = superposition of all vectors with m | 0 \rangle s and m | 1 \rangle s

 $E_{k:2m-k} = \frac{1}{2} \log divergence}$









► Possible area law: $k^{1-1/d}$ with $d \rightarrow \infty$





▶ Possible area law: $k^{1-1/d}$ with $d \rightarrow \infty$

Note: Effect due to frustration. Not due to long-range interactions.





▶ Possible area law: $k^{1-1/d}$ with $d \rightarrow \infty$

Note: Effect due to frustration. Not due to long-range interactions. Ising with J<0 : constant block entanglement.



Area Law for frustrated systems

1. Long range Ising model

2. Majumdar Ghosh model

3. Shastry-Sutherland model

4. Ising chain with NN interactions



$H=J_1 \Sigma \sigma_i \sigma_{i+1} + J_2 \Sigma \sigma_i \sigma_{i+2} \text{ with } J_1, J_2 > 0; J_2 = J_1/2$



 $H=J_{1} \Sigma \sigma_{i} \sigma_{i+1} + J_{2} \Sigma \sigma_{i} \sigma_{i+2}$

 $\Phi \approx 1/2$









After quenching:





After quenching: $E \ge 2$ (even) or 1 (odd)





After quenching:





After quenching:

Numerically, E = 2.3 for 8 spins









for frustrated systems









Area Law









Area law

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Main Thesis

Highly frustrated systems do not follow any area law

Highly frustrated systems r near-maximally genuine multi-party entangled

While

Weakly frustrated systems do not have a similar definite behavior regarding genuine multi-party entanglement.

Weakly frustrated systems follow the same area law as nonfrustrated systems away from criticality.

Sen(De), US, Dziarmaga, Sanpera, Lewenstein, PRL'08 Jindal, Rane, Dhar, Sen(De), US, PRA'14 Frustrated systems: Area law and Genuine multiparty entanglement



In C O N C L U S I O N





NO Area law

High genuine multiparty entanglement





Area law

No definite genuine multiparty entanglement



More work done

- Adv. Phys. 56, 243 (2007)
- Rev. Mod. Phys. 80, 517 (2008)







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References r incomplete!