ICTP international workshop on
Current Trends in Frustrated Magnetism
SPS, JNU


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## Outline

1. Understanding entanglement
2. Entanglement in many-body physics
3. What is frustration?
4. Characterizing "classical" frustration in q systems
5. Frustration and Entanglement
I. Area Law
II. Genuine multipartite entanglement
6. End remarks

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## Understanding entanglement

## LOCC paradigm in quantum info

- If the state is shared between two or more parties, the parties would only be able to act locally.
Allowed operations: LOCC.


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- What do we mean by LOCC?


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- What do we mean by LOCC?


Not this!!

## Understanding entanglement



What do we mean by LOCC?


- Alice makes a measurement and communicates her result to Bob (say, by a phone call).


## Understanding entanglement



What do we mean by LOCC?


- Alice makes a measurement and communicates her result to Bob (say, by a phone call).
- Then depending on her result, Bob will make his measurement and communicate his result to Alice.
- And so on.


## Understanding entanglement

## Separable and Entangled states

- Quantum states that can be prepared by LOCC $\rightarrow$ Separable states.
- Otherwise $\rightarrow$ Entangled states.


## Understanding entanglement

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- How do they look like?


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- Separable pure states: products over pure states of individual systems.


## Understanding entanglement

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- Quantum states that can be prepared by LOCC $\rightarrow$ Separable states.
- How do they look like? Mathematically?
- Separable states: mixtures of products over pure states of individual systems.


## Understanding entanglement

Which "entanglements" can we compute?

## Circa 2000

- Nielsen, Preskill, Wootters et al.


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Idea of using entanglement-like concepts in quantum many-body phenomena was put forward.

## Understanding entanglement

Which "entanglements" can we compute?

## Circa 2000

- Nielsen, Preskill, Wootters et al.
- Osborne and Nielsen, QIP’02, PRA'02
- Osterloh, Amico, Falci, Fazio, Nature'02


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To see the behavior of entanglement in real systems, it is not sufficient
to understand an entanglement measure conceptually.

## Understanding entanglement

Which "entanglements" can we compute?

To see the behavior of entanglement in real systems, it is not sufficient
to understand an entanglement measure conceptually. We must also be able to compute it for the states of the real systems.

## Understanding entanglement

## Which "entanglements" can we compute?

- Bipartite states.


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- For mixed two-party states, only entanglement of formation of two-qubit states.


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Ene mivad twa nortur ctatac anly

Entanglement of formation of a two-party state is the number of singlets that r required to create the state by LOCC.

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Modulo certain additivity problems.

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## Understanding entanglement

Which "entanglements" can we compute?

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- For mixed two-party states, only entanglement of formation of two-qubit states.
- In higher dimensions, logarithmic negativity can be calculated. But it cannot detect bound entanglement.


## Understanding entanglement

## Which "entanglements" can we compute?

- Bipartite states.

Enr mivad trun nnotry atotac anly

$$
\begin{aligned}
& \text { Logneg of a two-party state is } \\
& \qquad \log _{2}(2 N+1) .
\end{aligned}
$$

## Understanding entanglement

Which "entanglements" can we compute?

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Far mivad trun narter atatac anly

Logneg of a two-party state is $\log _{2}(2 N+1)$.
$\mathrm{N}=$ sum of mod of negative eigenvalues in partial transpose of state.

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- For mixed two-party states, only entanglement of formation of two-qubit states.
- For pure two-party states, local von Neumann entropy is a "good" measure of entanglement, and is computable.

Possible in arbitrary dimensions.

## Understanding entanglement

## Which "entanglements" can we compute?

- Bipartite states.

This sets the stage for the
QI - many-body interface.

## Understanding entanglement

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Indeed, two of the main directions of study are

## 1. EoF of reduced densities

of spin-1/2 ground states
2. Scaling of local entropy
in ground state partitions

1 Ussivic 11 arvillaly unininsivis.

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## Understanding entanglement Multiparty entanglement

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a. Geometric measure


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Wei, Goldbart, PRA’03
Balsone, DellAnno, DeSiene, Illuminatti, PRA'08 + ......

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b. Global measure


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Meyer, Wallach, JMP’02 $+\ldots .$.

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c. Generalized geometric measure


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A. $\operatorname{Sen}(\mathrm{De}), \mathrm{US}, \mathrm{PRA}^{\prime} 10$


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Entanglement in many-body physics Quantum Phase Transitions

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- Transitions at zero temperature.


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- Implying, transition not temp. driven.


# Entanglement in many-body physics 

## Quantum Phase Transitions

- Transitions at zero temperature.
- Implying, transition not temp. driven.
- Driven by system parameter, like a magnetic field.


# Entanglement in many-body physics Quantum Phase Transitions 

## Typical situation:

- $\mathrm{H}=\mathrm{H}(\mathrm{int})+\mathrm{a} \mathrm{H}($ field $)$


# Entanglement in many-body physics Quantum Phase Transitions 

## Typical situation:

- $\mathrm{H}=\mathrm{H}(\mathrm{int})+\mathrm{a} \mathrm{H}(f i e l d)$
- Ground state of H


# Entanglement in many-body physics Quantum Phase Transitions 

## Typical situation:

- H = H(int) + a H(field)
- Ground state of $\mathrm{H} \leftarrow$ guarantees $\mathrm{T}=0$


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- GS depends on " a ".
- "a" can be changed.


# Entanglement in many-body physics Quantum Phase Transitions 

## Typical situation:

- H = H(int) + a H(field)
- Ground state of $\mathrm{H} \leftarrow$ guarantees $\mathrm{T}=0$
- GS depends on " a ".
- "a" can be changed.
- Nonanalyticity appears in some physical quantity as "a" is changed.

Entanglement in many-body physics Two-site densities


Entanglement in many-body physics Two-site densities


# Entanglement in many-body physics 

## Two-site densities



The reduced state is a two-qubit state.
Spin-1/2 Chain

Entanglement in many-body physics

## Two-site densities

## The prescription:

Entanglement in many-body physics

## Two-site densities

The prescription:

1. Find ground state of spin-1/2 system

Entanglement in many-body physics

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1. Find ground state of spin- $1 / 2$ system
2. Remove all spins except two NNs

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4. Investigate it wrt the relevant system parameter

# Entanglement in many-body physics Quantum XY spin model 



# Entanglement in many-body physics 

## Quantum XY spin model



$$
\Sigma \mathrm{J}\left[(1+\gamma) \mathrm{S}_{\mathrm{x}}^{\mathrm{i}} \mathrm{~S}_{\mathrm{x}}^{\mathrm{i}+1}+(1-\gamma) \mathrm{S}_{\mathrm{y}}^{\mathrm{i}} \mathrm{~S}_{\mathrm{y}}^{\mathrm{i}+1}\right]-\mathrm{a} \mathrm{~S}_{\mathrm{z}}^{\mathrm{i}}
$$

S are half of Pauli matrices.

# Entanglement in many-body physics Quantum XY spin model 



$$
\Sigma \mathrm{J}\left[(1+\gamma) \mathrm{S}_{\mathrm{x}}^{\mathrm{i}} \mathrm{~S}_{\mathrm{x}}^{\mathrm{i}+1}+(1-\gamma) \mathrm{S}_{\mathrm{y}}^{\mathrm{i}} \mathrm{~S}_{\mathrm{y}}^{\mathrm{i}+1}\right]-\mathrm{a} \mathrm{~S}_{\mathrm{z}}^{\mathrm{i}}
$$

Quantum phase transition at $\mathrm{h}=1$.

# Entanglement in many-body physics Quantum XY spin model 

## For $\gamma=1$ : Transverse Ising Model.

$$
\Sigma \mathrm{J}\left[(1+\gamma) \mathrm{S}_{\mathrm{x}}^{\mathrm{i}} \mathrm{~S}_{\mathrm{x}}^{\mathrm{i}+1}+(1-\gamma) \mathrm{S}_{\mathrm{y}}^{\mathrm{i}} \mathrm{~S}_{\mathrm{y}}^{\mathrm{i}+1}\right]-\mathrm{a} \mathrm{~S}_{\mathrm{z}}^{\mathrm{i}}
$$

Quantum phase transition at $\mathrm{h}=1$.

## Entanglement in many-body physics

Linking QI with concepts in quantum statistical mechanics and quantum phase transitions.

Near QPT in 1D transverse Ising model, 2-site entanglement remains short ranged, while 2-site correlation length diverges.

Entanglement, however, does show signs of criticality.


Osterloh, Amico, Falci, \& Fazio,
Nature 2002; Osborne \& Nielsen,
Phys. Rev. A 2002.

## Entanglement in many-body physics



# Entanglement in many-body physics 

## Two-site densities

The prescription:

1. Find ground state of spin- $1 / 2$ system
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# Entanglement in many-body physics 

## Two-site densities

## Why ground state?

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# Entanglement in many-body physics 

## Two-site densities

## Why ground state?

The prescription:

1. Find ground s
2. Remove all s]
3. Find EoF of $r$
4. Investigate it

Guarantees that there are no thermal effects.

# Entanglement in many-body physics 

## Two-site densities

## Why ground state?

The prescription:

1. Find grounds
2. Remove all s]
3. Find EoF of $r$
4. Investigate it

## Thermal states,

 time-evolved states also considered.
# Entanglement in many-body physics 

## Two-site densities

## Why NN?

The prescription:

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2. Remove all spins except two NNs
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# Entanglement in many-body physics 

## Two-site densities

## Why $N N ?$

The prescription:

1. Find ground s
2. Remove all s]

In many instances, but NOT all,
NNN and so on
have little to no entanglement.

# Entanglement in many-body physics 

 Multiparty entanglementMultiparty entanglement detects QPT
a. Geometric measure
b. Global measure
c. Generalized geometric measure (GGM)

## Entanglement in many-body physics Geometric measure detects QPT



Wei, Das, Mukhopadhyay, Vishveshwara, Goldbart, PRA’05

## Entanglement in many-body physics

Global measure of multipartite entanglement detects QPT

deOliviera, Rigolin, deOliviera, PRA’06

## Entanglement in many-body physics GGM detects QPT



## Entanglement in many-body physics

## GGM $-1 / 2$

## GGM detects QPT



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## Entanglement in many-body physics

## GGM $-1 / 2$ GGM detects QPT

## darivativa

## anisotropy

Blue dashes $\rightarrow 1$ (Ising)
Pink circles $\rightarrow 0.8$
Green dots $\rightarrow 0.2$
A. Sen(De), US, 1002.1253


## Entanglement in many-body physics

## GGM $-1 / 2$ GGM detects QPT



J_1-J_2 models in 1D \& 2D:
darive
A. Biswas, R. Prabhu, A. Sen(De), US, PRA'14
anisouvpy
Blue dashes $\rightarrow 1$ (Ising)
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A. $\operatorname{Sen}(\mathrm{De}), \mathrm{US}$, 1002.1253

## Outline

1. Understanding entanglement
2. Entanglement in many-body physics
3. What is frustration?
4. Characterizing "classical" frustration in q systems
5. Frustration and Entanglement
I. Area Law
II. Genuine multipartite entanglement
6. End remarks

## What is frustration?

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From a classical perspective

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Consider an Ising model:
$\mathcal{H}=J \sum \sigma_{i} \sigma_{j} ; J>0$

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Consider an Ising model:


$$
\mathcal{H}=J \sum \sigma_{i} \sigma_{j} ; J>0
$$

## What is frustration?

## From a classical perspective

Failure to have spin configuration to minimize individual interaction terms

## What is frustration?

## From a quantum perspective

## What is frustration?

## From a quantum perspective

## Draw a parallel

## What is frustration?

## From a quantum perspective

Classical frustration: spin configuration

Quantum frustration: GSs of two terms not same


$$
\mathcal{H}=\mathcal{H}_{\text {loc }}+\mathcal{H}_{\text {int }}
$$

## What is frustration?

## Classical <br> spin configuration



Cannot get optimal spin configuration

## What is frustration?

## Classical <br> spin configuration





Quantum non-commutativity

Cannot get optimal spin configuration

GSs of two terms not same

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## Characterizing "classical" frustration in q systems

# Characterizing "classical" frustration in $q$ systems 

Classical frustration

## Characterizing "classical" frustration in $q$ systems

## Classical frustration

"Frustration degree"

Sen(De), US, Dziarmaga, Sanpera, Lewenstein, PRL'08 Jindal, Rane, Dhar, Sen(De), US, PRA'14

# Characterizing "classical" frustration in q systems Frustration degree 

- Given $\mathrm{H},|\Gamma\rangle$,


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- Given $H,|\Gamma\rangle$,
replace one-body, two-body etc. in H by Ising ones, i.e. by $\sigma_{i}^{z}$ or $\sigma_{i}^{z} \sigma_{j}^{z}$ etc.


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Find $\mathrm{H}^{\mathrm{I}}$

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Find $\mathrm{H}^{\mathrm{I}}$


$$
\Phi=\operatorname{avg} \frac{\sum_{\mathrm{k}}\langle\Gamma| \mathrm{H}_{\mathrm{f}}^{\mathrm{k}}|\Gamma\rangle}{\left.\sum_{\mathrm{l}}\left|\langle\Gamma| \mathrm{H}_{\mathrm{nf}}^{\mathrm{t}}\right| \Gamma\right\rangle \mid}
$$

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Ising model: $\mathrm{H}=\mathrm{J} \Sigma \sigma_{\mathrm{i}}{ }_{\mathrm{i}}^{\mathrm{z}_{\mathrm{j}}} \quad$ with $\mathrm{J}>0$



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Reduced entropy S would depend on the surface of separation between A and B.


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## We r talking abt interacting systems.



Reduced entropy S would depend on the surface of separation between A and B.

Would be true (trivially) if ...


Reduced entropy $S$ would depend on the surface of separation between A and B.

Boundary particles are pure entangled states.


Reduced entropy S would depend on the surface of separation between A and B.

Boundary particles are pure entangled states. Plus no long-range entangled pairs.


Reduced entropy S would depend on the surface of separation between A and B.

## Typical situation is far from being such.



Reduced entropy S would depend on the surface of separation between A and B.

## Typical situation is far from being such.

Usually intricately multiparty quantum correlated.


Reduced entropy S would depend on the surface of separation between A and B.

$$
\mathrm{S}\left(\rho_{\mathrm{L}}\right) \sim \mathrm{L}^{\mathrm{d}-1}
$$

L: characteristic length of A


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$$
\mathrm{S}\left(\rho_{\mathrm{L}}\right) \sim \mathrm{L}^{\mathrm{d}-1}
$$

L: characteristic length of $A$

## Area Law:1D



$$
\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)=\mathrm{S}\left(\rho_{\mathrm{L}}\right)
$$

## Area Law:1D



$$
\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)=\mathrm{S}\left(\rho_{\mathrm{L}}\right) \sim \mathrm{L}^{\mathrm{d}-1} \equiv \mathrm{constant}
$$

away from criticality

## Area Law:1D



$$
\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)=\mathrm{S}\left(\rho_{\mathrm{L}}\right) \sim \mathrm{L}^{\mathrm{d}-1} \equiv \mathrm{c} \rho \text { tant }
$$

independent of block-size

## Area Law:1D



Block entanglement: $\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)$

$$
\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)=\mathrm{S}\left(\rho_{\mathrm{L}}\right) \sim \ln \mathrm{L}
$$

at criticality

## Area Law:1D



Block entanglement: $\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)$

$$
\underset{\underset{\text { atcDe }}{\mathrm{E}\left(|\Psi\rangle_{\mathrm{L}: \mathrm{N}-\mathrm{L}}\right)=\mathrm{S}\left(\rho_{\mathrm{L}}\right)} \sim \log \text { divergence }}{\sim \ln \mathrm{L}}
$$

## Lot of progress in different directions.

## Lot of progress in different directions. A case study:

 Frustrated systems
## Main Thesis

$>$ Highly frustrated systems do not follow area law

## Main Thesis

$>$ Highly frustrated systems do not follow area law

while

$>$ Weakly frustrated systems follow same area law as nonfrustrated systems away from criticality

Sen(De), US, Dziarmaga, Sanpera, Lewenstein, PRL'08
Jindal, Rane, Dhar, Sen(De), US, PRA'14

## Area Law

 for frustrated systems1. Long range Ising model
2. Majumdar Ghosh model
3. Shastry-Sutherland model
4. Is ing chain with $\mathfrak{N J V}$ interactions

## Coosing/Quenching Method

$>$ Initial state:

$$
|\Phi\rangle_{\mathrm{in}} \equiv|\psi\rangle_{1} \otimes|\psi\rangle_{2} \otimes|\psi\rangle_{3} \otimes \ldots \otimes|\psi\rangle_{\mathrm{N}}
$$

## Cooling/Quenching Method

$>$ Initial state:

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|\Phi\rangle_{\text {in }} \equiv|\psi\rangle_{1} \otimes|\psi\rangle_{2} \otimes|\psi\rangle_{3} \otimes \ldots \otimes|\psi\rangle_{\mathrm{N}}
$$

$>$ Project $|\Phi\rangle_{\text {in }}$ onto the ground state space of the model.

$$
|\Phi\rangle_{\mathrm{f}}=\left(\sum|\Gamma\rangle_{\mathrm{i}}\langle\Gamma|\right)|\Phi\rangle_{\mathrm{in}}
$$

## Cooling/Quenching Method

> Initial state:

$$
|\Phi\rangle_{\text {in }} \equiv|\psi\rangle_{1} \otimes|\psi\rangle_{2} \otimes|\psi\rangle_{3} \otimes \ldots \otimes|\psi\rangle_{\mathrm{N}}
$$

$>$ Project $|\Phi\rangle_{\text {in }}$ onto the ground state space of the model.

$$
|\Phi\rangle_{\mathrm{f}}=\left(\sum|\Gamma\rangle_{\mathrm{i}}\langle\Gamma|\right)|\Phi\rangle_{\mathrm{in}}
$$

$>$ Calculate $\mathrm{E}_{\mathrm{N} / 2: \mathrm{N} / 2}\left(|\Phi\rangle_{\mathrm{f}}\right)$.

## Cooling/Quenching Method

> Initial state:

$$
|\Phi\rangle_{\text {in }} \equiv|\psi\rangle_{1} \otimes|\psi\rangle_{2} \otimes|\psi\rangle_{3} \otimes \ldots \otimes|\psi\rangle_{\mathrm{N}}
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$>$ Project $|\Phi\rangle_{\text {in }}$ onto the ground state space of the model.

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|\Phi\rangle_{\mathrm{f}}=\left(\sum|\Gamma\rangle_{\mathrm{i}}\langle\Gamma|\right)|\Phi\rangle_{\mathrm{in}}
$$

$>$ Calculate $\mathrm{E}_{\mathrm{N} / 2: \mathrm{N} / 2}\left(|\Phi\rangle_{\mathrm{f}}\right)$.
$>$ Maximize $\mathrm{E}_{\mathrm{N} / 2: \mathrm{N} / 2}\left(|\Phi\rangle_{\mathrm{f}}\right)$ over all choices of the initial state.

# for frustrated systems 

1. Long range Ising model
2. Majumdar Ghosh model
3. Shastry-Sutherland model
4. Ising chain with $\mathfrak{N V V}$ interactions

## Long range Ising model

 $\mathrm{H}=\mathrm{J} \Sigma \sigma^{\mathbf{Z}}{ }_{i} \sigma^{\mathbf{Z}}{ }_{j} \quad$ with $\mathrm{J}>0$
## Long range Ising model

 $\mathrm{H}=\mathrm{J} \Sigma \sigma^{\mathrm{Z}}{ }_{i} \sigma^{\mathrm{Z}}{ }_{j} \quad$ with $\mathrm{J}>0 \quad \Phi \approx 1$Long range Ising model

$$
\mathrm{H}=\mathrm{J} \Sigma \sigma_{i}^{\mathrm{Z}} \sigma^{\mathrm{Z}}{ }_{j} \quad \text { with } \mathrm{J}>0
$$

After quenching:
$|\psi\rangle=$ superposition of all vectors with $\mathrm{m}|0\rangle \mathrm{s}$ and $\mathrm{m}|1\rangle \mathrm{s}$

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After quenching:
$|\psi\rangle=$ superposition of all vectors with
$\mathrm{m}|0\rangle \mathrm{s}$ and $\mathrm{m}|1\rangle \mathrm{s}$

$$
E_{k: 2 m-k}=1 / 2 \log k
$$

## Long range Ising model

$$
\mathrm{H}=\mathrm{J} \Sigma \sigma_{i}^{\mathrm{Z}} \sigma^{\mathrm{Z}}{ }_{j} \quad \text { with } \mathrm{J}>0
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## Area law Clear departure from area law

$>$ Long range Ising model: "Infinite" dimensions

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Note: Effect due to frustration.
Not due to long-range interactions.

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$>$ Long range Ising model: "Infinite" dimensions
$>$ Possible area law: $\mathrm{k}^{1-1 / \mathrm{d}}$ with $\mathrm{d} \rightarrow \infty$

Note: Effect due to frustration.
Not due to long-range interactions.
Ising with $\mathrm{J}<0$ : constant block entanglement.

## Area Law

## for frustrated systems

1. Long range Ising model
2. Majumdar Ghosh model
3. Shastry-Sutherland model
4. Ising chain with $\mathfrak{N N V}$ interactions

## Majumdar-Ghosh mode[

$\mathrm{H}=\mathrm{J}_{1} \Sigma \sigma_{i} \sigma_{i+1}+\mathrm{J}_{2} \Sigma \sigma_{i} \sigma_{i+2}$ with $\mathrm{J}_{1}, \mathrm{~J}_{2}>0 ; \mathrm{J}_{2}=\mathrm{J}_{1} / 2$

## Majumdar-Ghosh mode[

$$
\mathrm{H}=\mathrm{J}_{1} \Sigma \sigma_{i} \sigma_{i+1}+\mathrm{J}_{2} \Sigma \sigma_{i} \sigma_{i+2}
$$

$$
\Phi \approx 1 / 2
$$

## Majumdar-Ghosk mode[

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Ground state:
(1)


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Ground state:


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Ground state:
(1)


After quenching:

$$
E \geq 2 \quad \text { (even) or } 1
$$

(odd)

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Ground state:
(1)


After quenching:
$E \geq 2 \quad($ even $)$ or 1
$E \leq \log 5($ even $)$ or
$\log 3($ odd $)$

## Majumdar-Ghosh model

$\mathrm{H}=\mathrm{J}_{1} \Sigma \sigma_{i} \sigma_{i+1}+\mathrm{J}_{2} \Sigma \sigma_{i} \sigma_{i+2}$ with $\mathrm{J}_{1}, \mathrm{~J}_{2}>0 ; \mathrm{J}_{2}=\mathrm{J}_{1} / 2$
Ground state:
(1)


After quenching:

$$
\begin{aligned}
& E \geq 2 \quad(\text { even }) \text { or } 1 \\
& E \leq \log 5(\text { even }) \text { or } \\
& \log 3(\text { odd })
\end{aligned}
$$

Numerically, $\mathrm{E}=2.3$ for 8 spins

## Majumdar-Ghosh model

$\mathrm{H}=\mathrm{J}_{1} \Sigma \sigma_{i} \sigma_{i+1}+\mathrm{J}_{2} \Sigma \sigma_{i} \sigma_{i+2}$ with $\mathrm{J}_{1}, \mathrm{~J}_{2}>0 ; \mathrm{J}_{2}=\mathrm{J}_{1} / 2$
Ground state:
(1)


After quenching:

$$
\begin{aligned}
& \mathrm{E} \geq 2 \quad \text { (even) } \\
& \mathrm{E} \leq \log 5 \text { a }
\end{aligned}
$$

Numerically, $\mathrm{E}=2.3$ for \& spons


## Area Law

 for frustrated systems

NO Area law


Area law

## Outline

1. Understanding entanglement
2. Entanglement in many-body physics
3. What is frustration?
4. Characterizing "classical" frustration in q systems

## 5. Frustration and Entanglement

I. Area Law
II. Genuine multipartite entanglement
6. End remarks

## Main Thesis

$>$ Highly frustrated systems do not follow any area law
Highly frustrated systems r near-maximally genuine multi-party entangled

## While

Weakly frustrated systems do not have a similar definite behavior regarding genuine multi-party entanglement.
$>$ Weakly frustrated systems follow the same area law as nonfrustrated systems away from criticality.

Sen(De), US, Dziarmaga, Sanpera, Lewenstein, PRL'08 Jindal, Rane, Dhar, Sen(De), US, PRA '14

Frustrated systems: Area law and Genuine multiparty entangโement


NO Area law
High genuine multiparty entanglement


Area law
No definite genuine multiparty entanglement

## More work done

- Adv. Phys. 56, 243 (2007)
- Rev. Mod. Phys. 80, 517 (2008)


## Thank you!




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References r incomplete!

