

Phase Space Methods for the Analysis and Simulation of CDM Dynamics

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Abel, Hahn, Kaehler (2012), MNRAS
Kaehler, Hahn, Abel (2012), IEEE TVCG
Hahn, Abel, Kaehler (2013), MNRAS
Angulo, Hahn, Abel (2013), MNRAS
Hahn, Angulo, Abel (2014), MNRAS subm.
Hahn & Angulo (2015), MNRAS subm.

What is Dark Matter?

microscopic

proton = 1GeV, WIMP 100GeV? $\rightarrow 10^{21}/g$



continuum limit

cold (or at most lukewarm)

e.g. thermally produced at very early times, cooled since then



$v_{\text{thermal}} \ll v_{\text{bulk}}$

negligible cross-section

weak-scale or even weaker



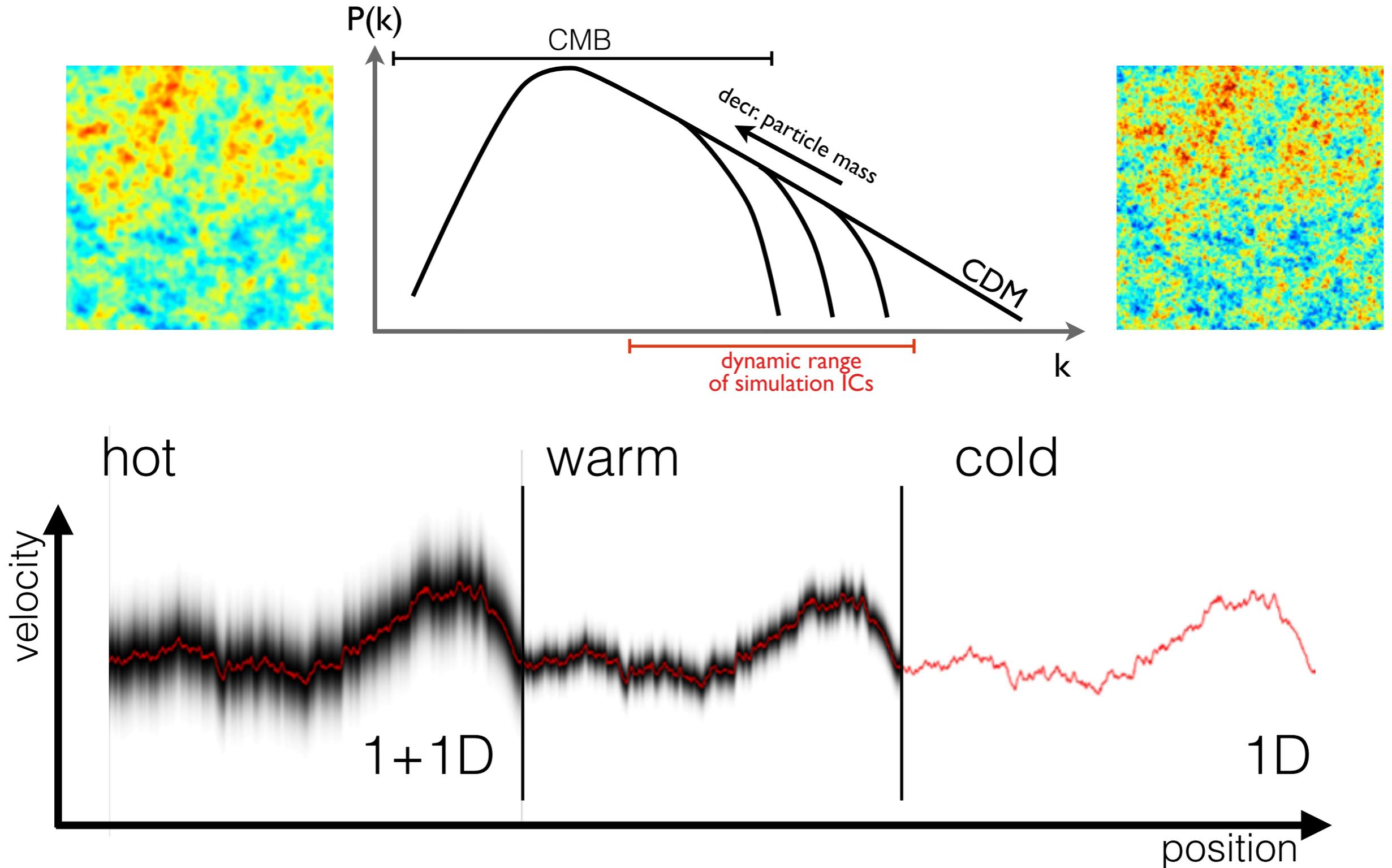
$\sigma_{\text{DM}} \ll \sigma_{\text{em}}$

collisionless

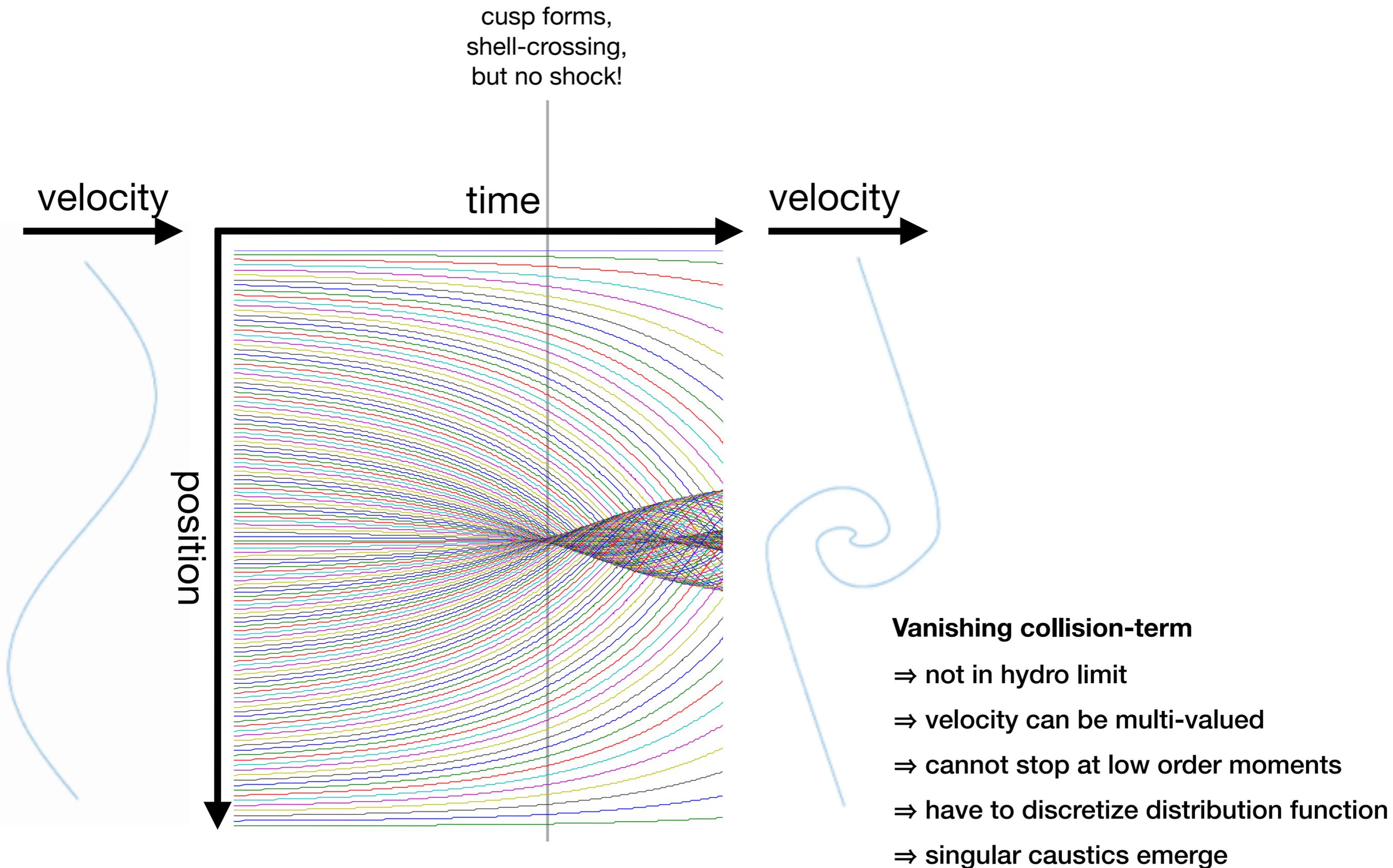
...and also the dominant gravitating component (~80%)

at first order, structure formation is well described by assuming all matter is dark matter

Dark Matter - properties on small scales



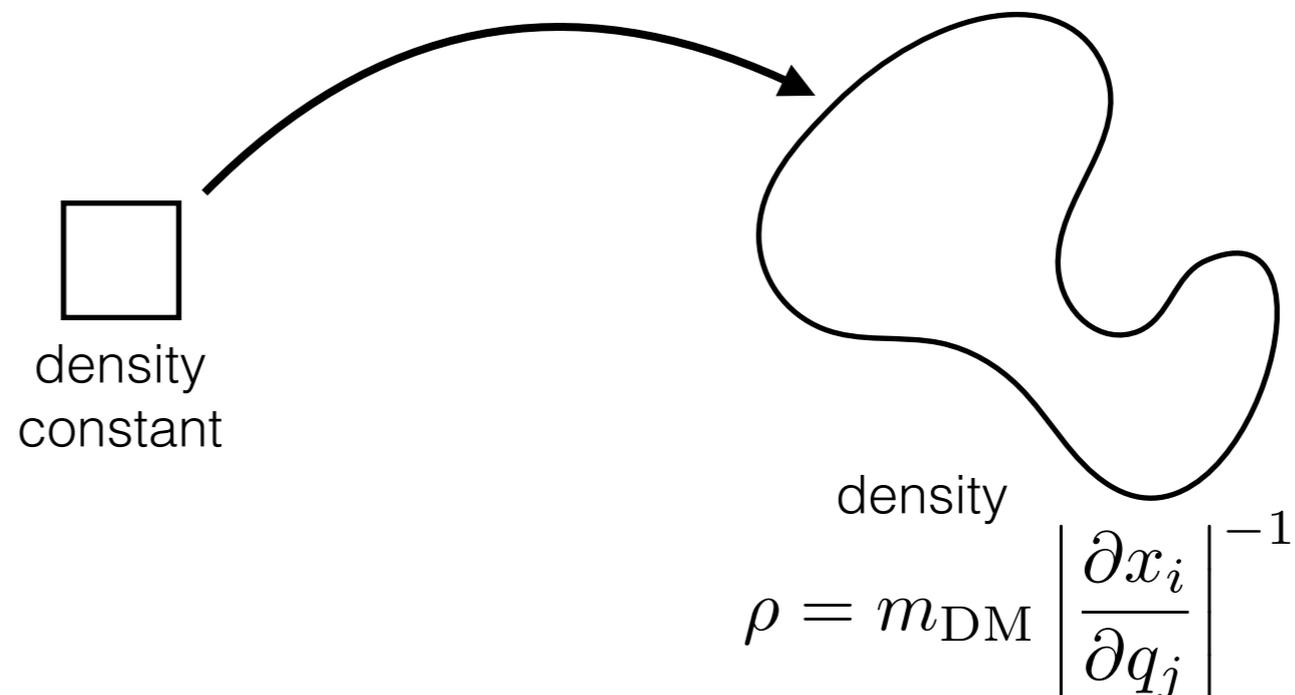
1D behaviour under self-gravity



Dark Matter - fluid flow

Lagrangian description, evolution of fluid element

$$\mathcal{Q} \subset \mathbb{R}^3 \rightarrow \mathbb{R}^6 : \mathbf{q} \mapsto (\mathbf{x}_{\mathbf{q}}(t), \mathbf{v}_{\mathbf{q}}(t))$$



For DM, motion of any point \mathbf{q} depends only on gravity

$$(\dot{\mathbf{x}}_{\mathbf{q}}, \dot{\mathbf{v}}_{\mathbf{q}}) = (\mathbf{v}_{\mathbf{q}}, -\nabla\phi)$$

unlike hydro, no internal temperature, entropy, pressure

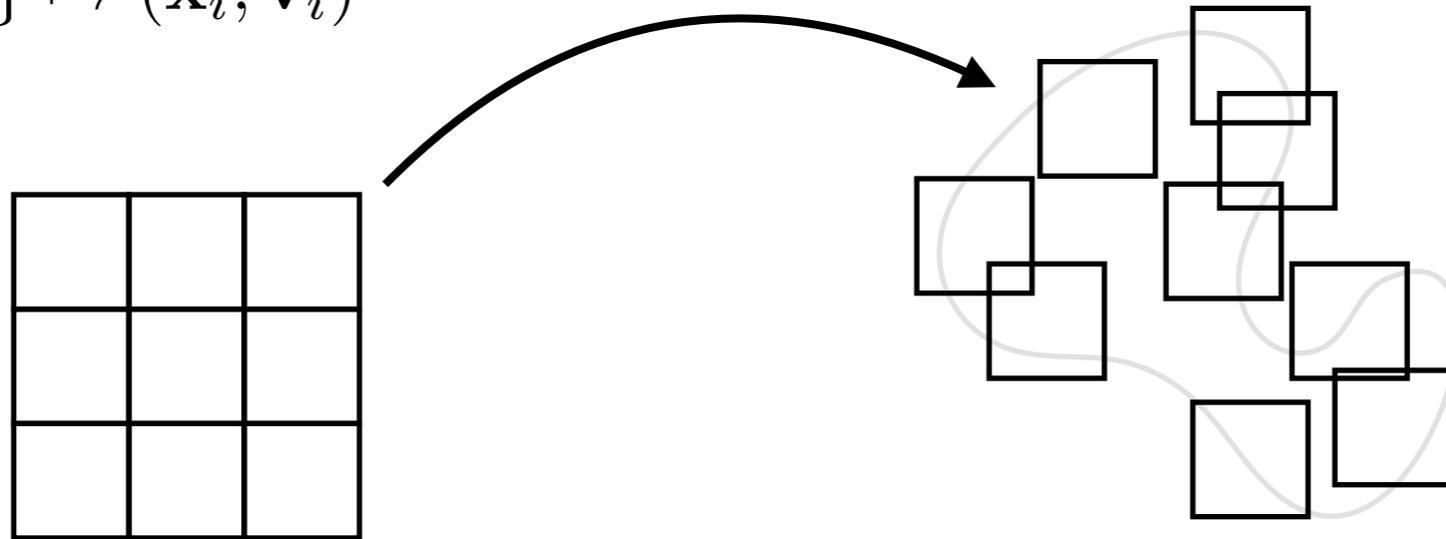
So the quest is to solve Poisson's equation

$$\Delta\phi = 4\pi G\rho$$

N-body vs. continuum approximation

The N-body approximation:

$$i \in \{1 \dots N\} \mapsto (\mathbf{x}_i, \mathbf{v}_i)$$



\Rightarrow EoM are just Hamiltonian N-body eq. (method of characteristics)

for small N, density field is poorly estimated,

$$\rho = m_p \sum \delta_D(x - x_i) \otimes W$$

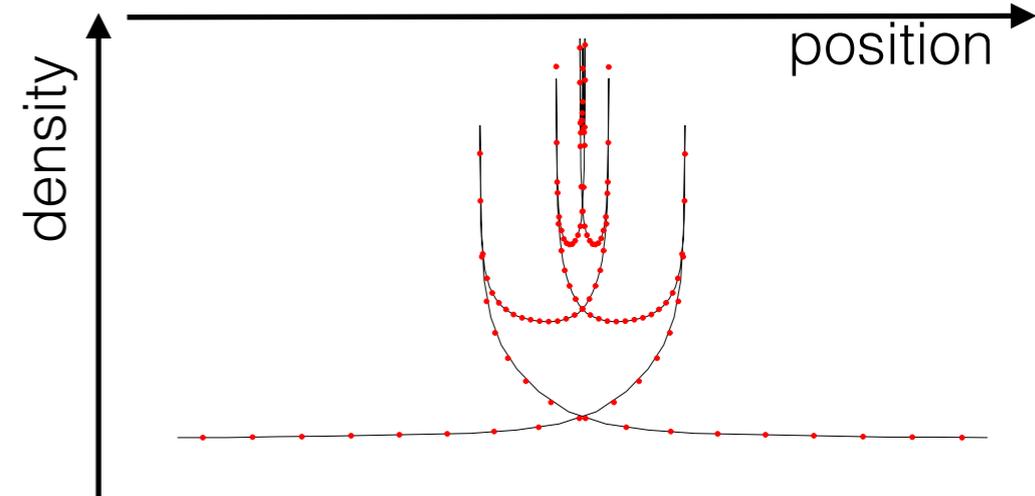
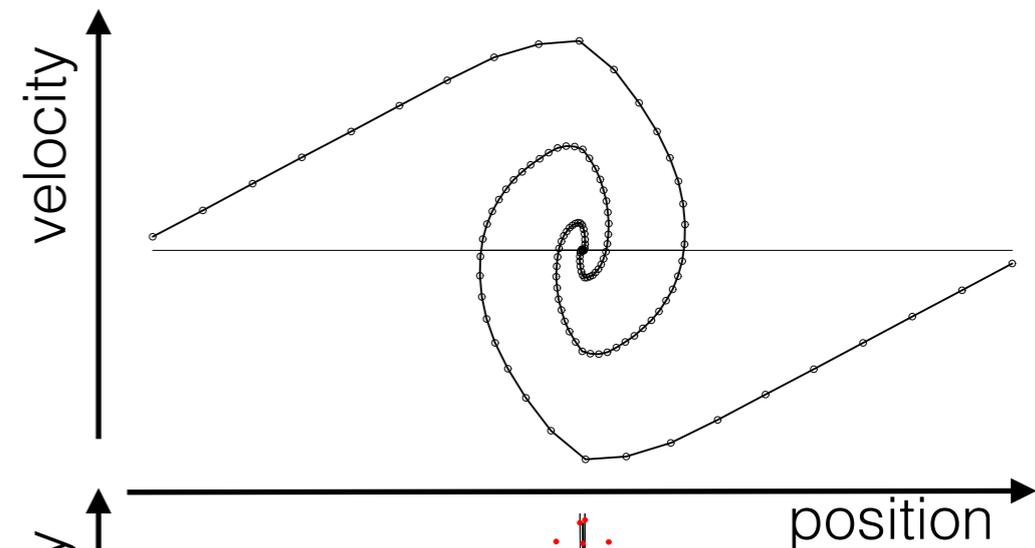
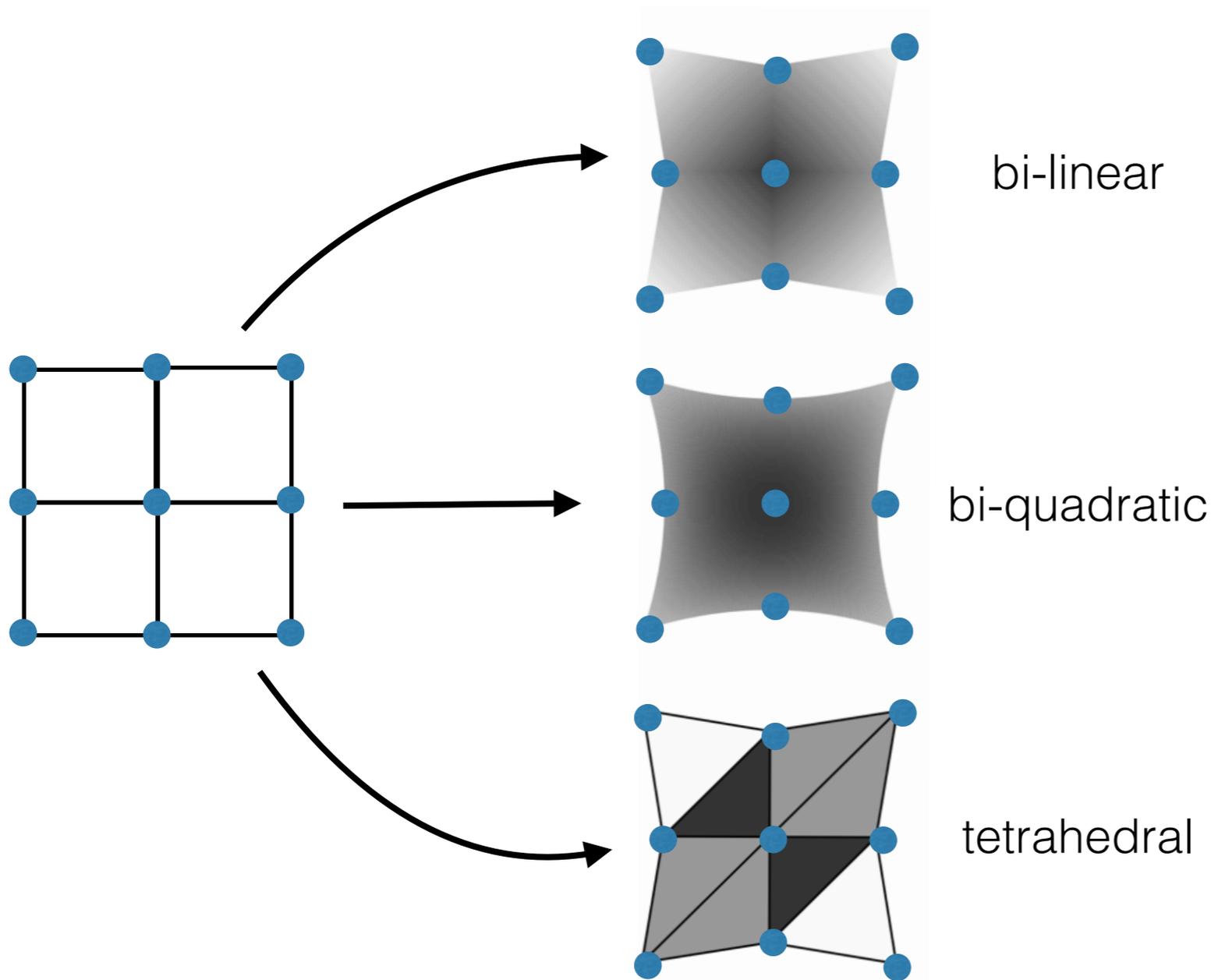
continuum structure is given up, but ‘easy’ to solve for forces

hope that as $N \rightarrow$ very large numbers, approach collisionless continuum

Lagrangian elements

Define little piecewise maps:

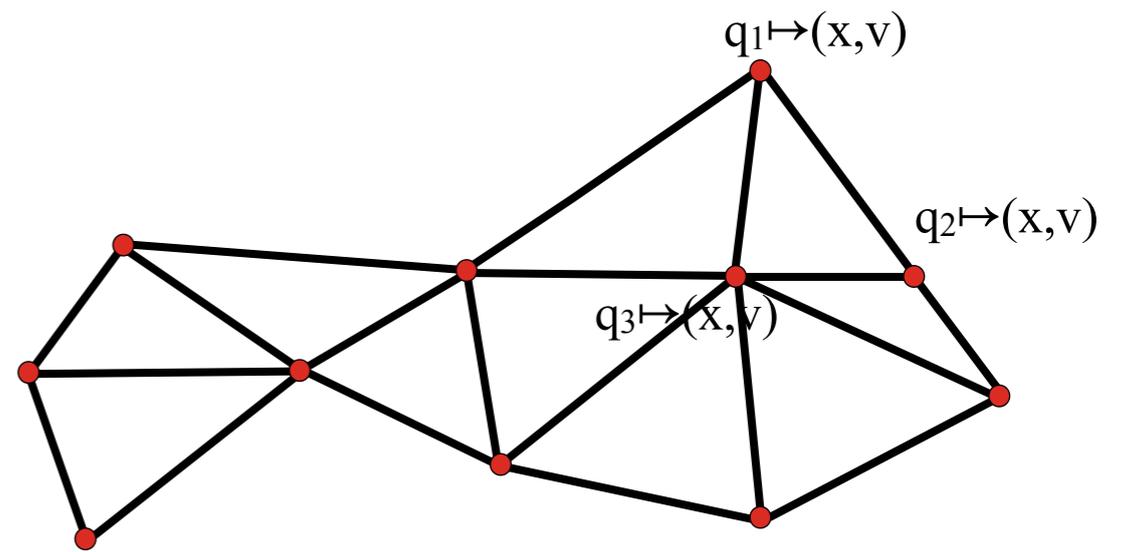
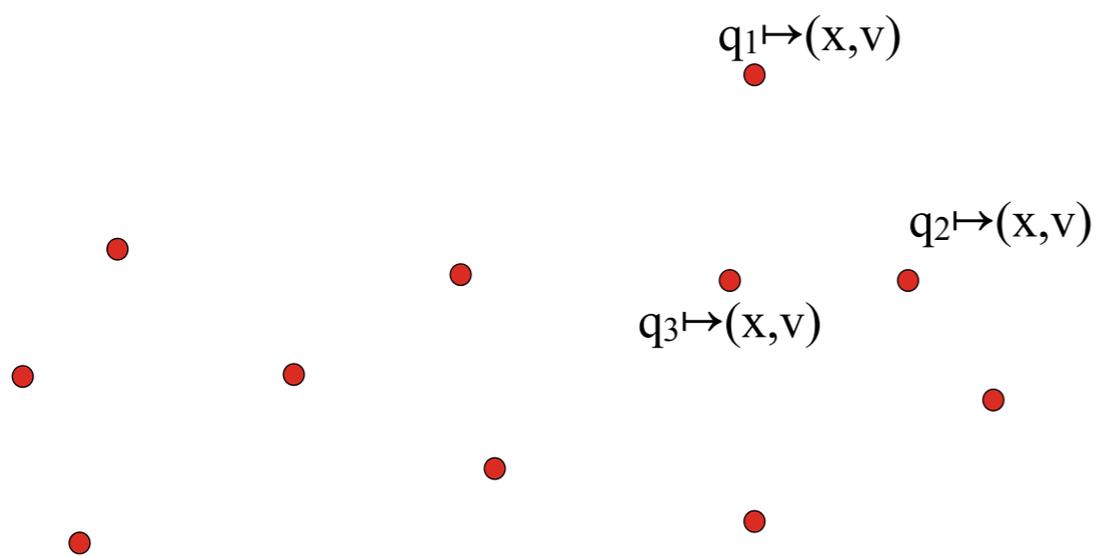
$$Q_i \subset \mathbb{R}^3 \rightarrow \mathbb{R}^6 : \mathbf{q} \mapsto (\mathbf{x}_{\mathbf{q}}(t), \mathbf{v}_{\mathbf{q}}(t))$$



$$\rho = m_{\text{DM}} \left| \frac{\partial x_i}{\partial q_j} \right|^{-1}$$

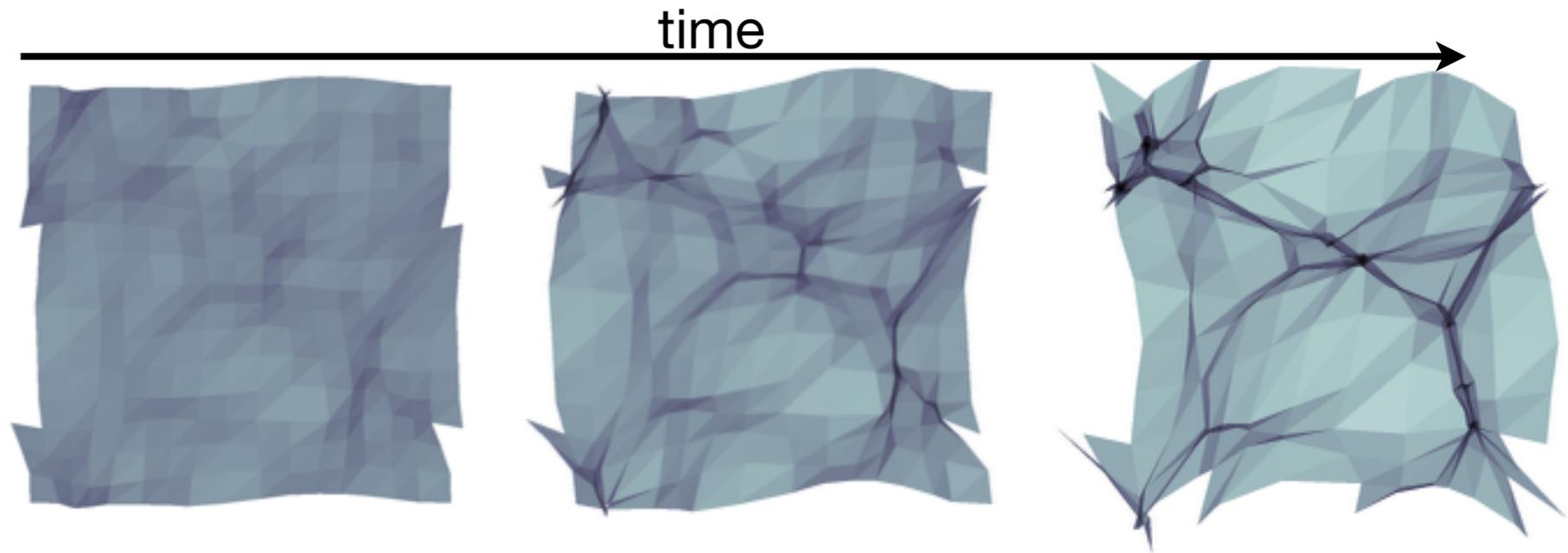
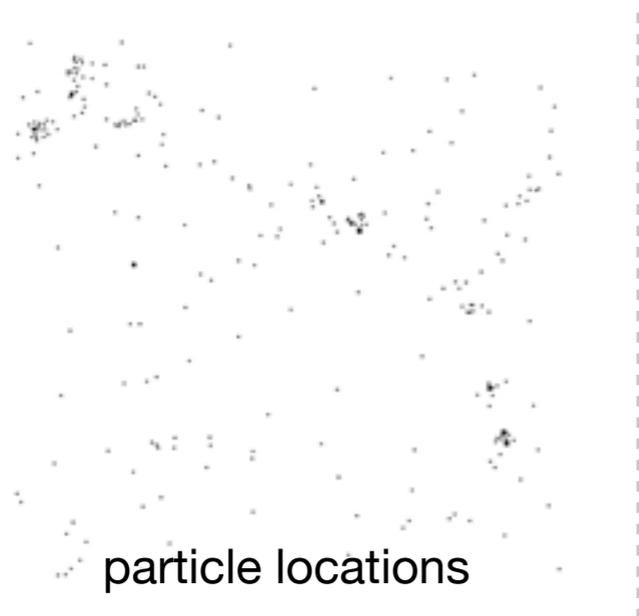
cost:
truncation error
in EoM!

Describing the density field

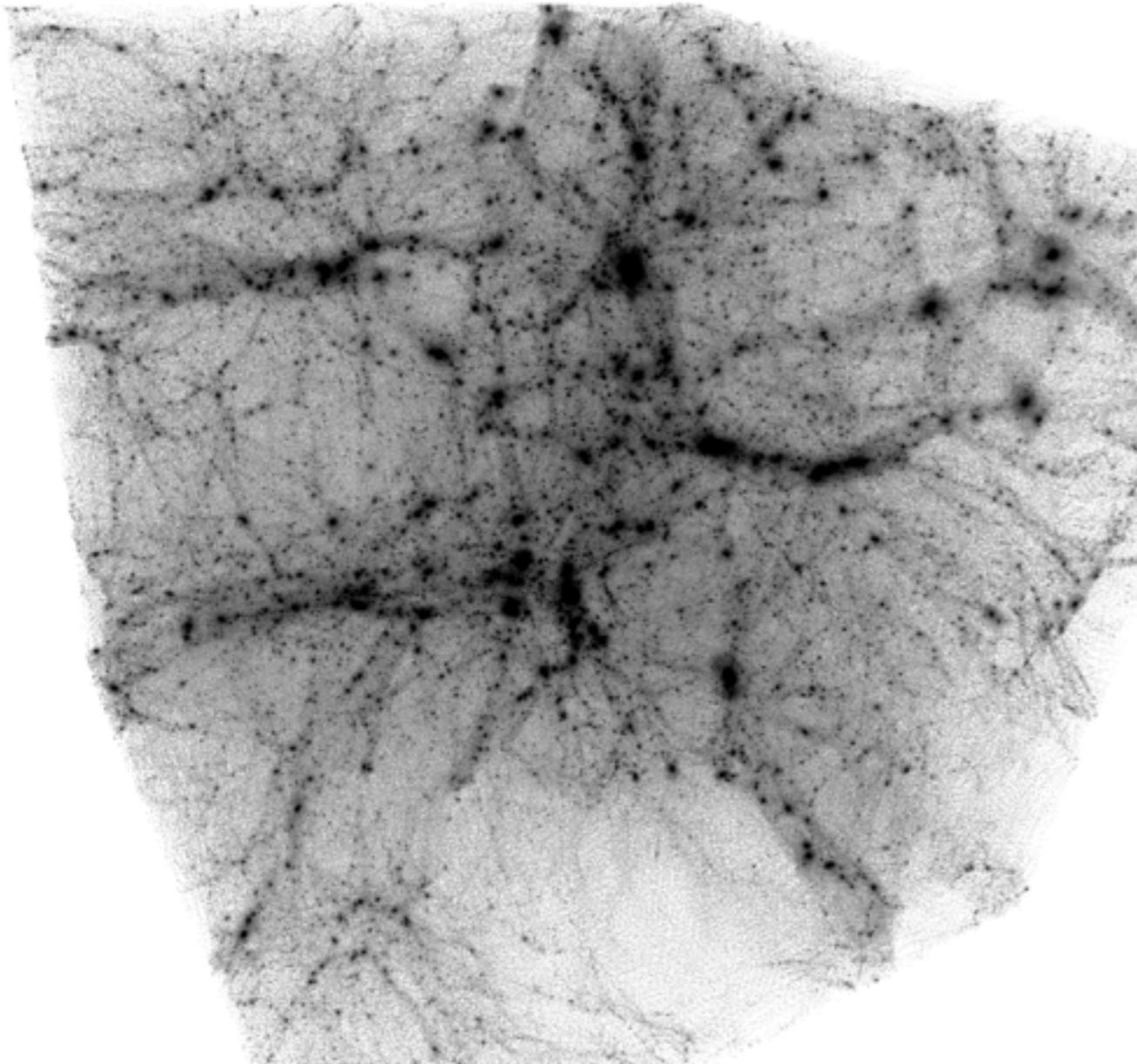


$$\rho = m_p \sum \delta_D(x - x_i) \otimes W$$

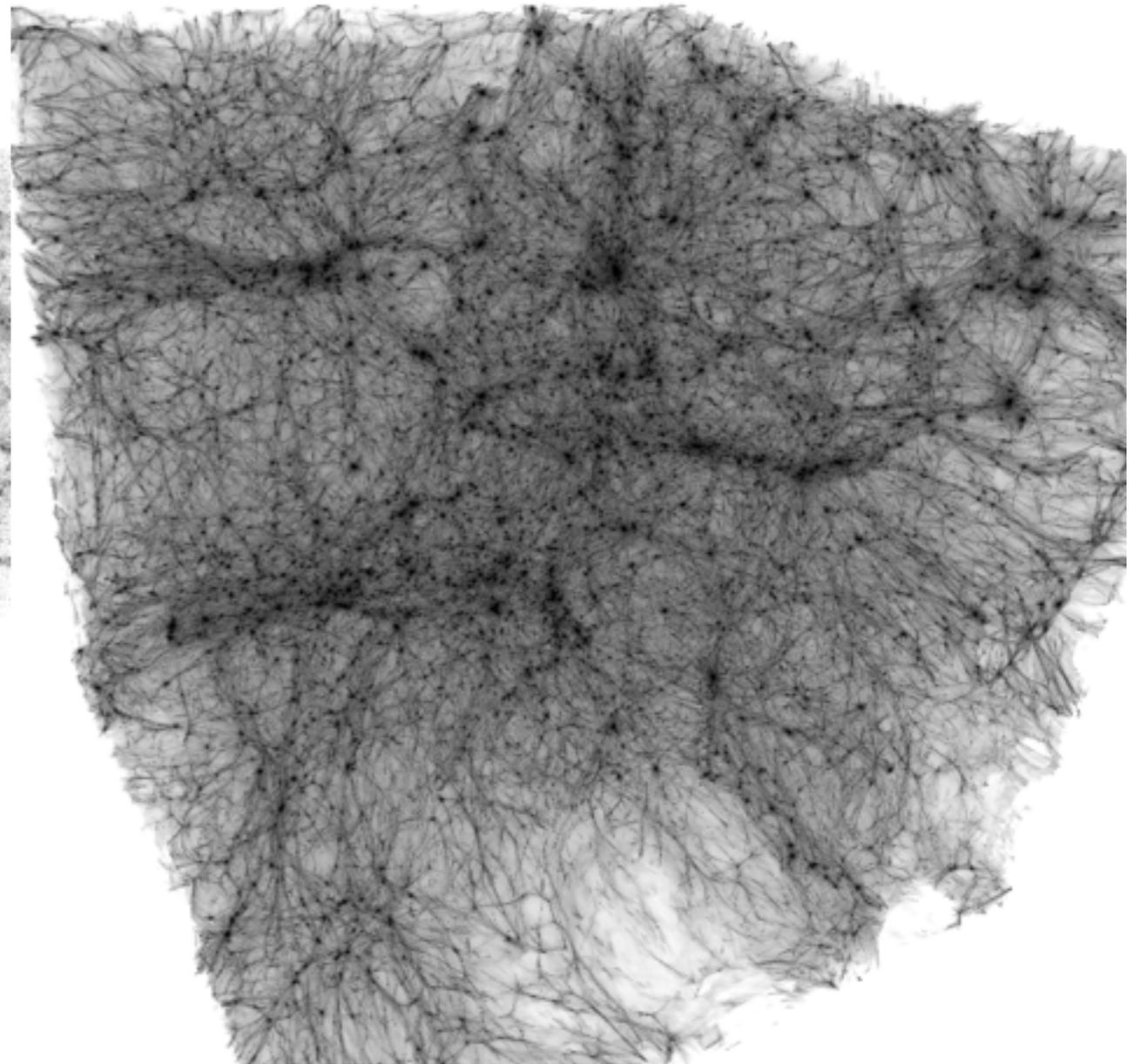
$$\rho = m_p \sum_{\text{streams}} \left| \det \frac{\partial x_i}{\partial q_j} \right|^{-1}$$



Three dimensions



rendering points for particles.

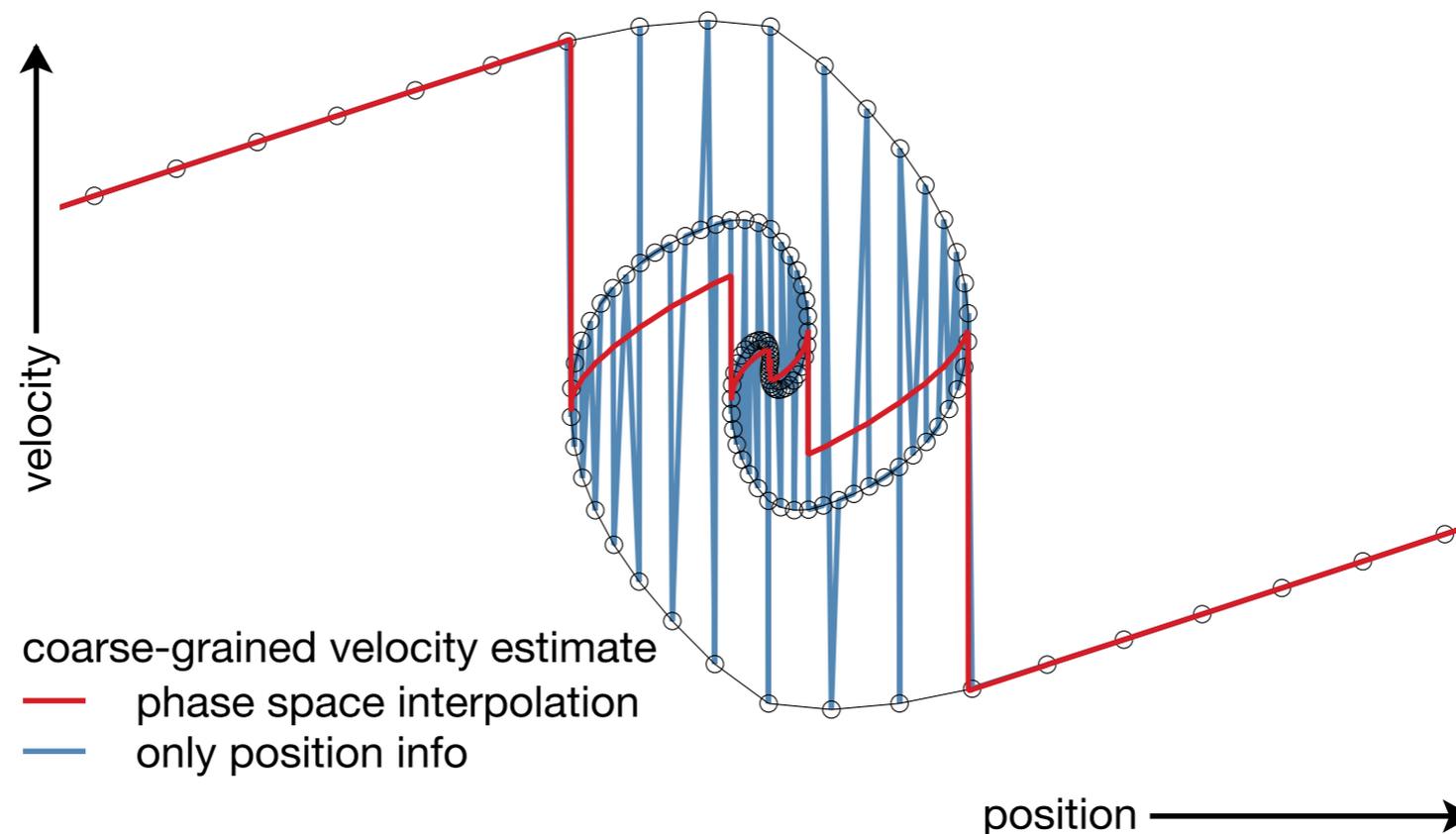


rendering tetrahedral phase space cells.

Same simulation data! (Abel, Hahn, Kaehler 2012)

Problem: How to measure the bulk velocity field?

- Interpolate between neighbouring N-body particles
- “neighbouring” in phase space, not configuration space
- account for averaging over streams (“coarse-graining”)



- Coarse-grained bulk velocity field:

$$\langle \mathbf{v} \rangle \equiv \frac{\int_{\mathbb{R}^3} \mathbf{v} f(\mathbf{x}, \mathbf{v}) d^3v}{\int_{\mathbb{R}^3} f(\mathbf{x}, \mathbf{v}) d^3v} = \frac{\sum_{s \in S} \mathbf{v}_s(\mathbf{x}) \rho_s(\mathbf{x})}{\sum_{s \in S} \rho_s(\mathbf{x})}$$

- result is discontinuous across caustics

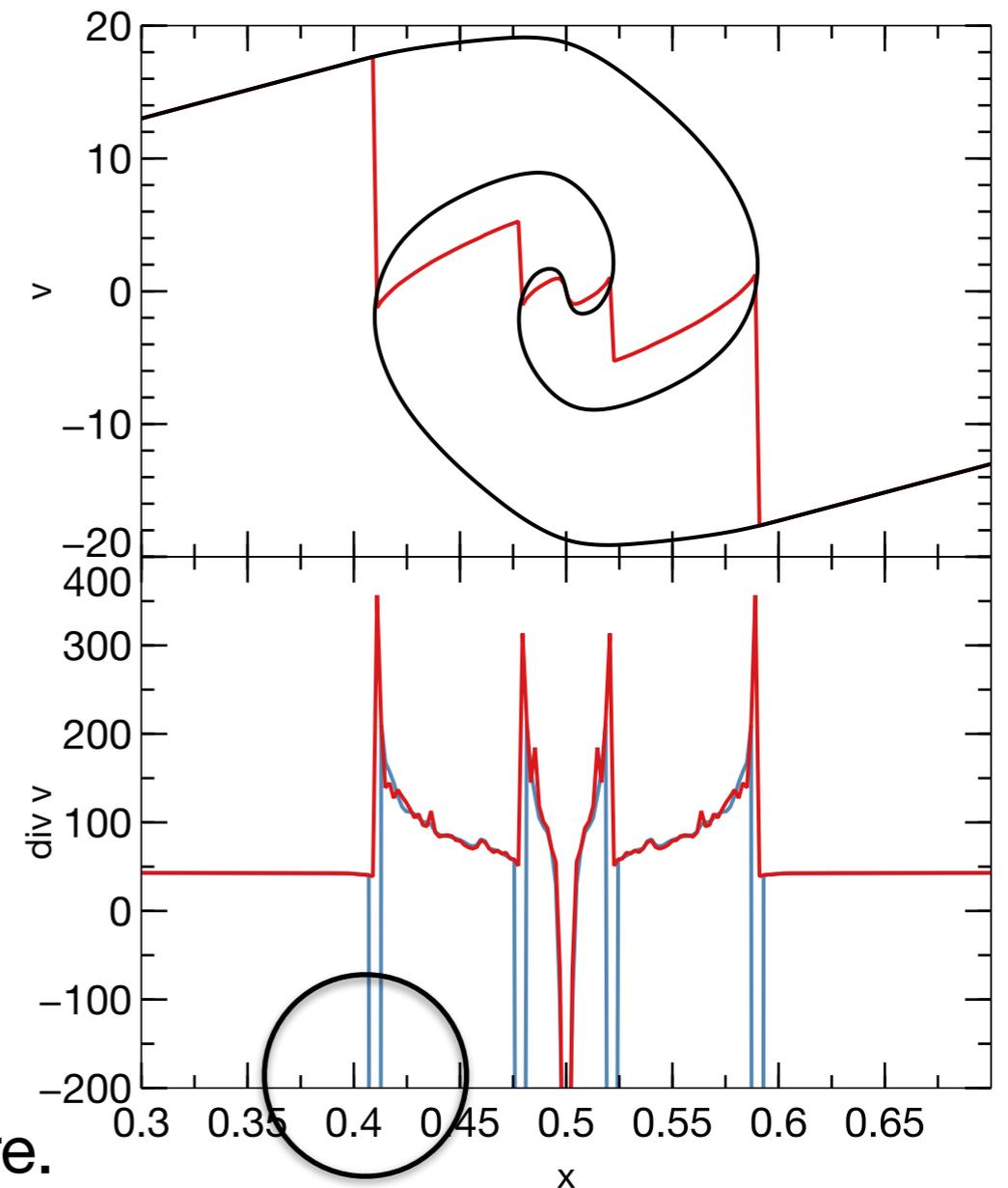
Derivatives of the bulk velocity field

- **Discontinuities make ordinary derivatives ill-defined without coarse-graining!**
- **Away from discontinuities:**
Need to explicitly evaluate action of derivative on **projected** field:

$$\nabla \cdot \langle \mathbf{v} \rangle = \langle (\nabla \log \rho) \cdot (\mathbf{v} - \langle \mathbf{v} \rangle) \rangle + \langle \nabla \cdot \mathbf{v} \rangle$$

$$\nabla \times \langle \mathbf{v} \rangle = \langle (\nabla \log \rho) \times (\mathbf{v} - \langle \mathbf{v} \rangle) \rangle + \langle \nabla \times \mathbf{v} \rangle$$

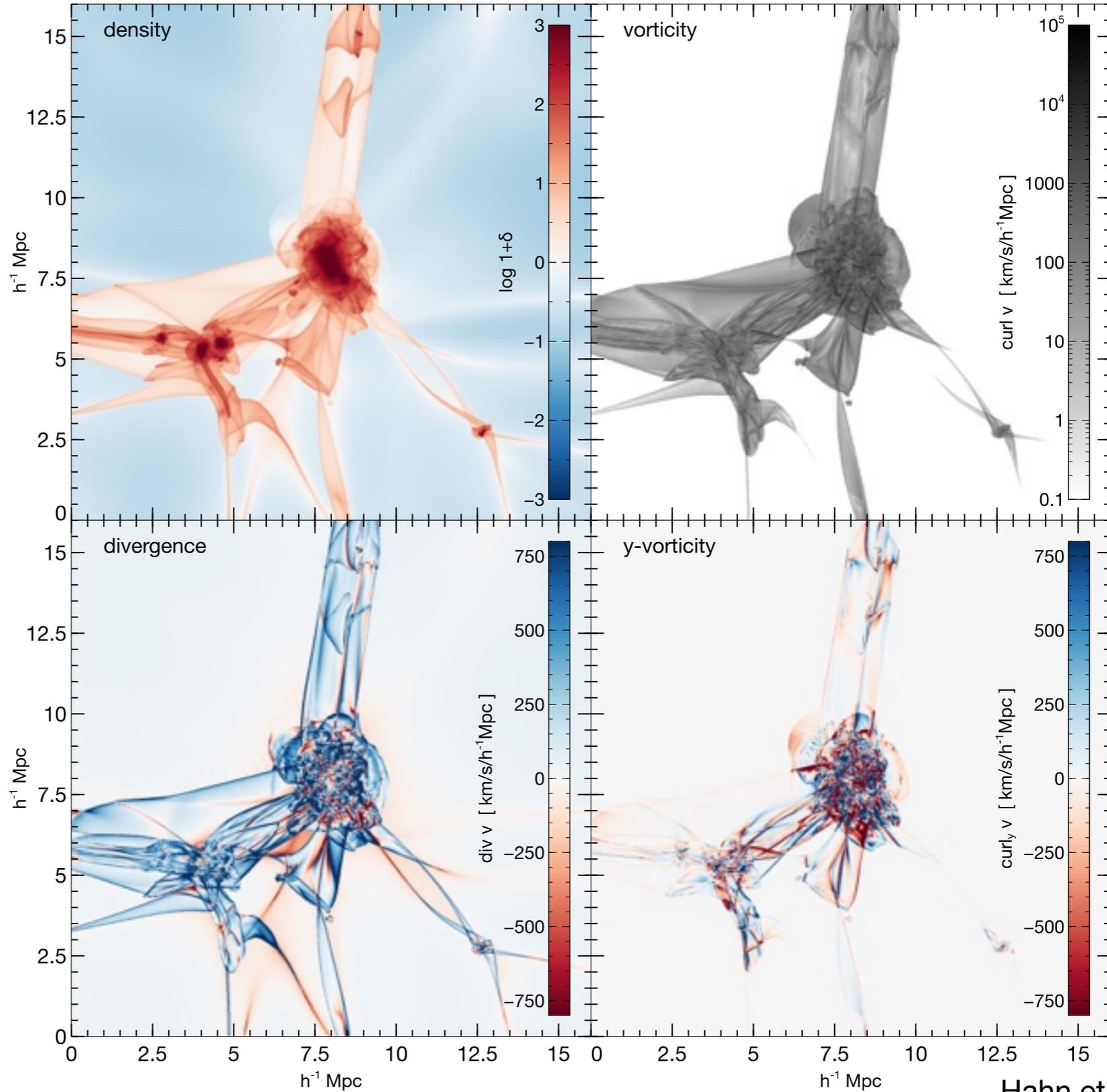
- **Vorticity for std. gravity pure multi-stream phenomenon!!**
- **At discontinuities:**
Derivatives are singular, but have finite measure.



**compressive singularities
at caustics**

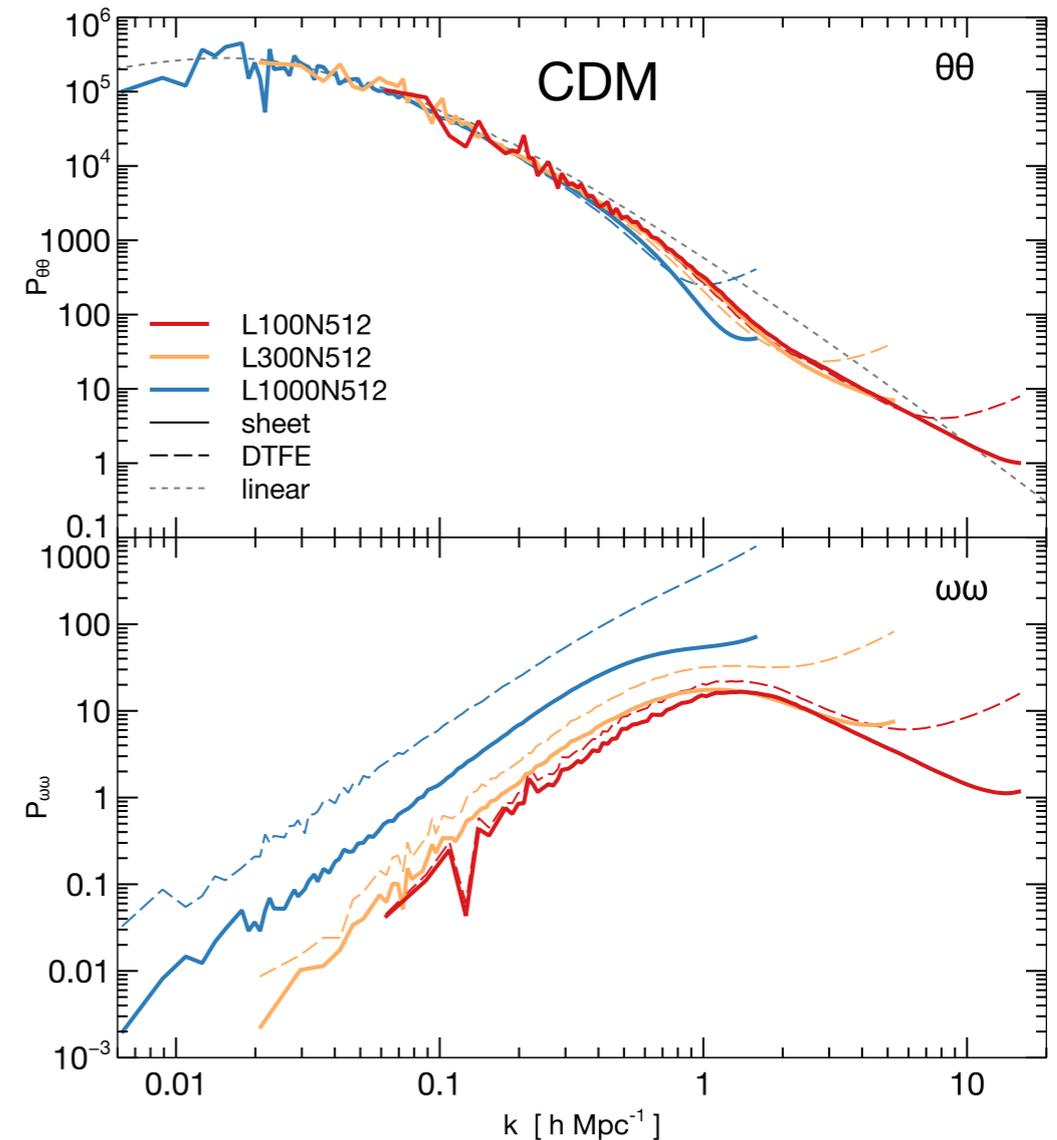
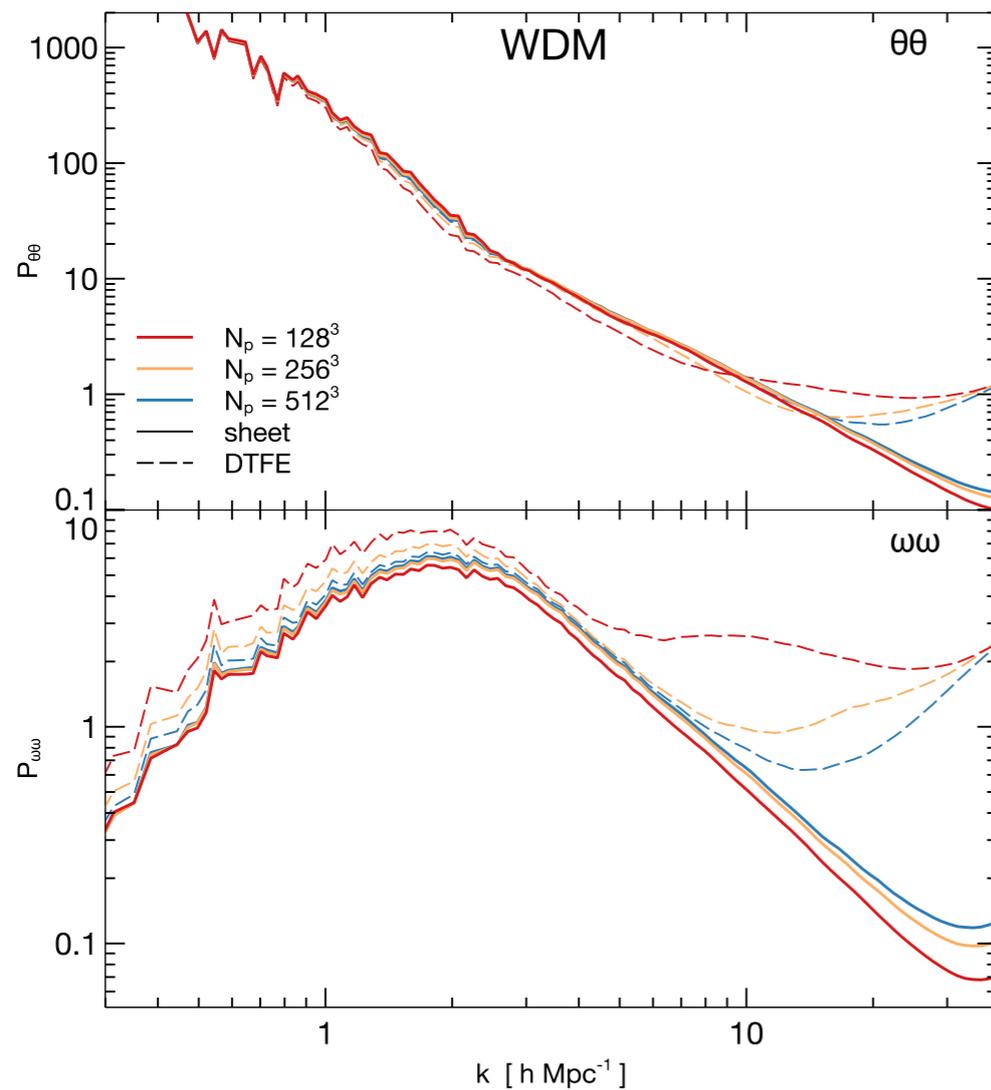
Properties of the cosmic velocity field II

Dark matter
fluid mechanics!

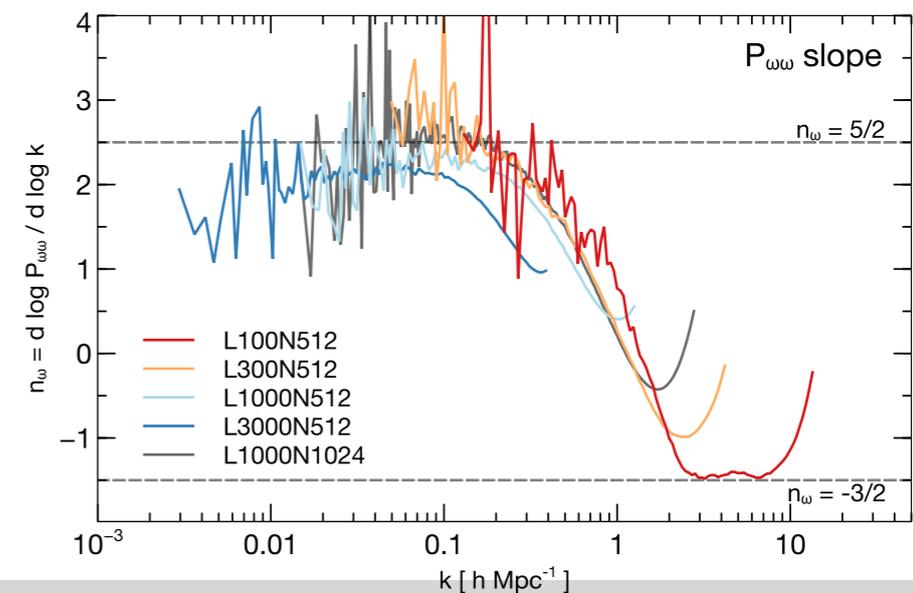


Hahn et al. 2014a

Spectral properties of the cosmic velocity field I



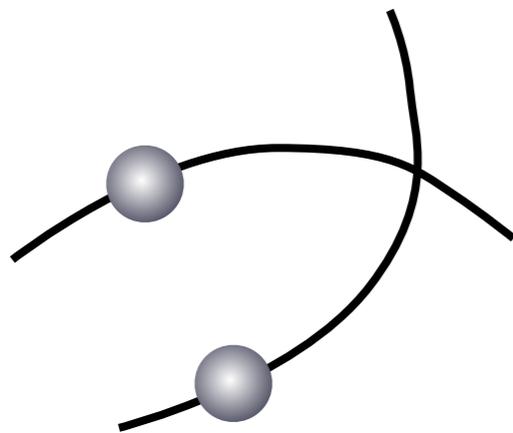
- Faster convergence (for WDM: convergence!)
- Better small scale properties



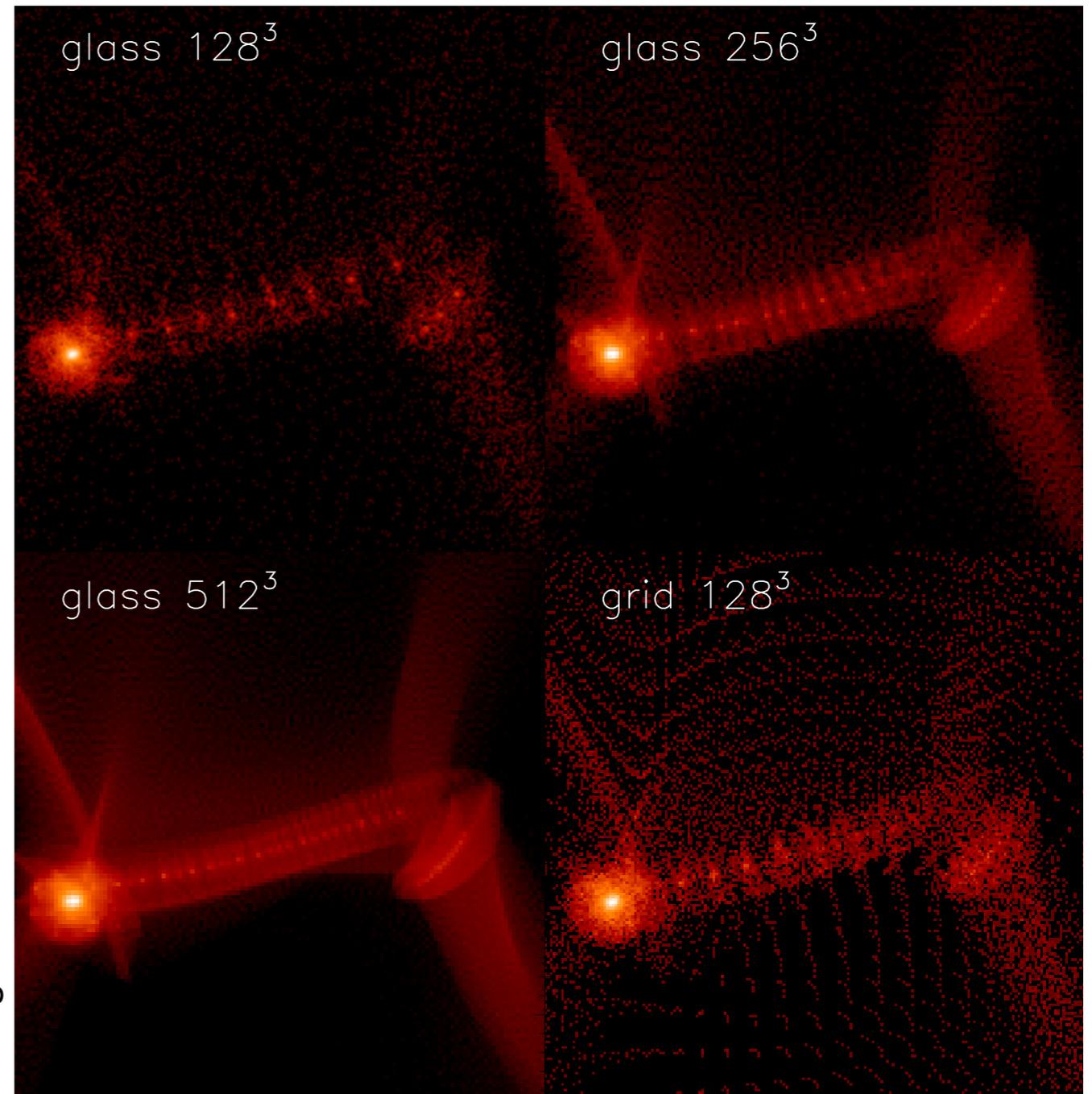
Problems of the N-body method: WDM

Main Problem: two-body effects, directly related to force softening

Scattering



Clumping/
Fragmentation



Most obvious for non-CDM simulations!

(e.g. Centrella&Melott 1983, Melott&Shandarin 1989, Wang&White 2007)

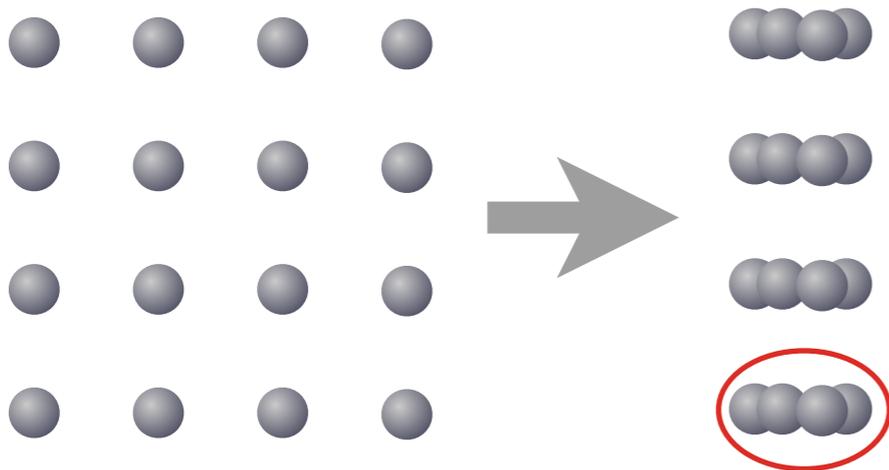
Improving on N-body....

N-body

$$\dot{\mathbf{x}}_i = \frac{1}{m a^2} \mathbf{p}_i \quad \text{and} \quad \dot{\mathbf{p}}_i = -m \nabla_x \phi|_{\mathbf{x}_i}$$

point-wise and Hamiltonian

need softening,
no knowledge what it
should be (empirical)



Lagrangian phase-space element

$$\dot{\mathbf{x}}_{\mathbf{q}} = \mathbf{v}_{\mathbf{q}}, \quad \text{and} \quad \dot{\mathbf{v}}_{\mathbf{q}} = - \nabla_x \phi|_{\mathbf{x}_{\mathbf{q}}}, \quad \text{with } \mathbf{q} \in \mathcal{Q}$$

continuum structure (diff w.r.t. \mathbf{q}), approx by

$$P_k = \{ \pi(\mathbf{q}) \mid \pi(\mathbf{q}) = \sum_{\alpha, \beta, \gamma=0}^k a_{\alpha\beta\gamma} q_0^\alpha q_1^\beta q_2^\gamma \}$$

-> EoM for polynomial coefficients

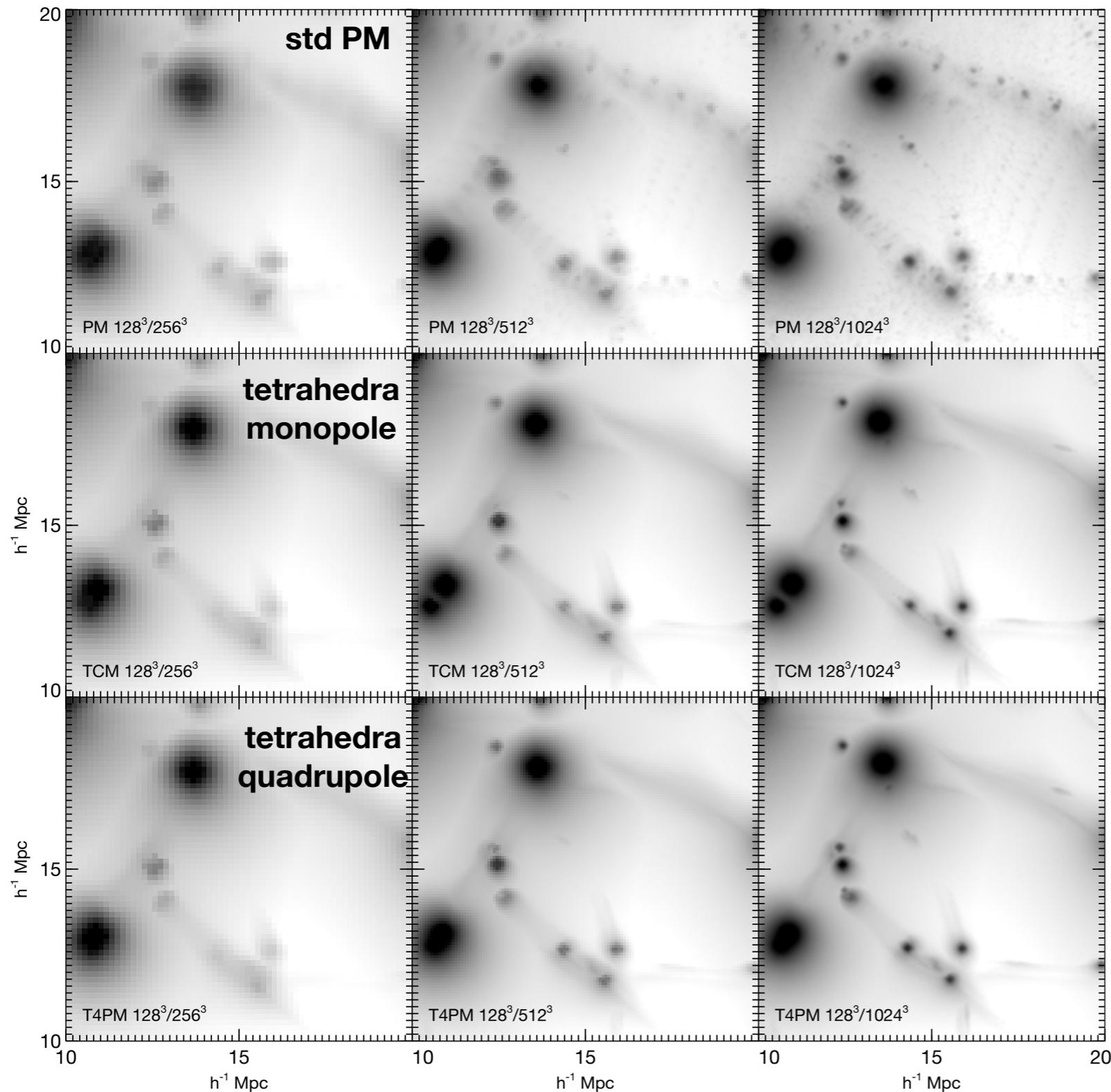
$$\dot{\mathbf{x}}_{\alpha\beta\gamma} = \mathbf{v}_{\alpha\beta\gamma}, \quad \dot{\mathbf{v}}_{\alpha\beta\gamma} = -J^{-1} \mathbf{f}_{\alpha\beta\gamma}$$

explicit truncation error:

$$\Delta \dot{\mathbf{v}} = -J^{-1} \sum_{\alpha, \beta, \gamma=k+1}^{\infty} \mathbf{f}_{\alpha\beta\gamma} q_0^\alpha q_1^\beta q_2^\gamma$$

Using tets for simulations: 300eV toy WDM problem

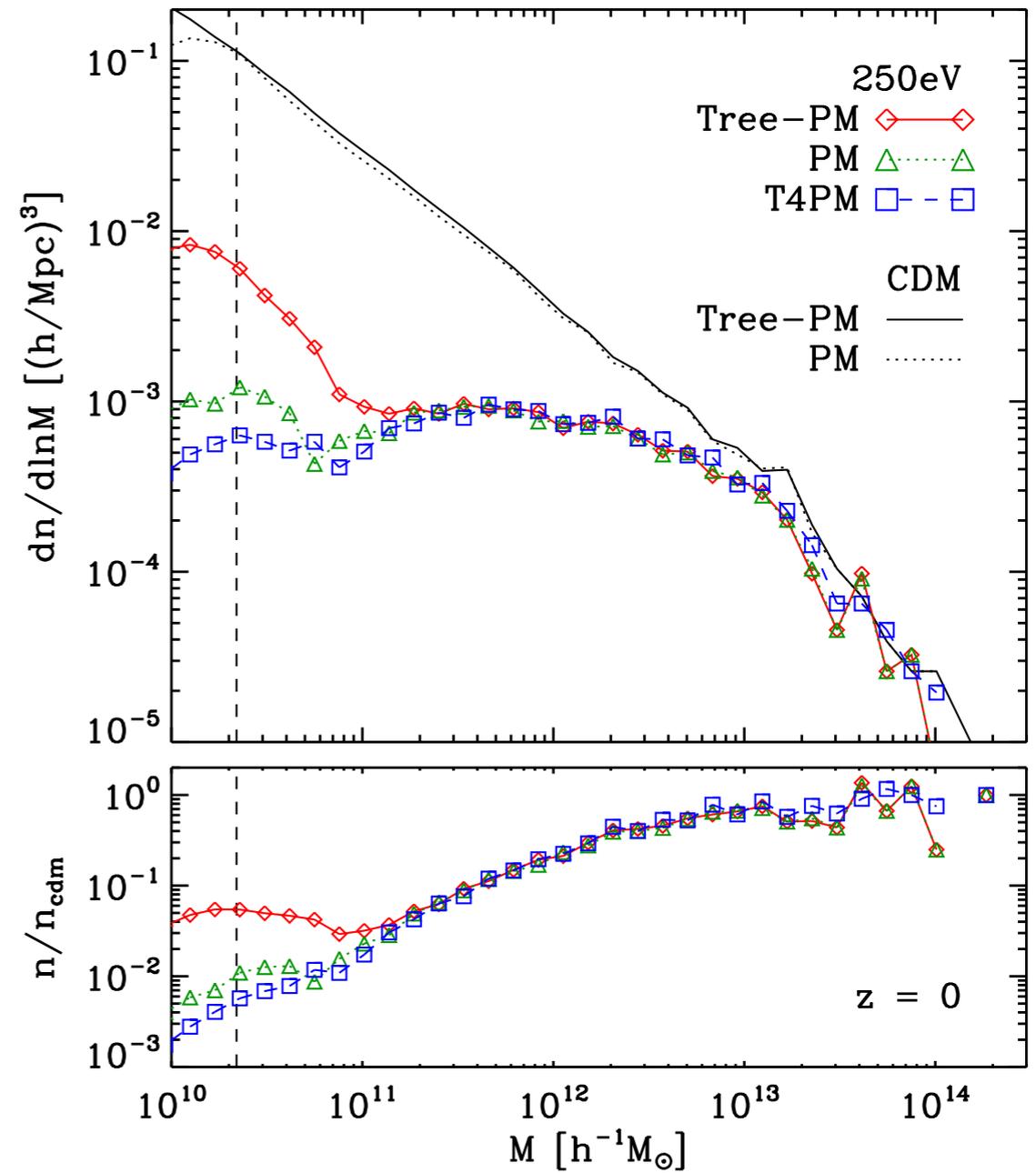
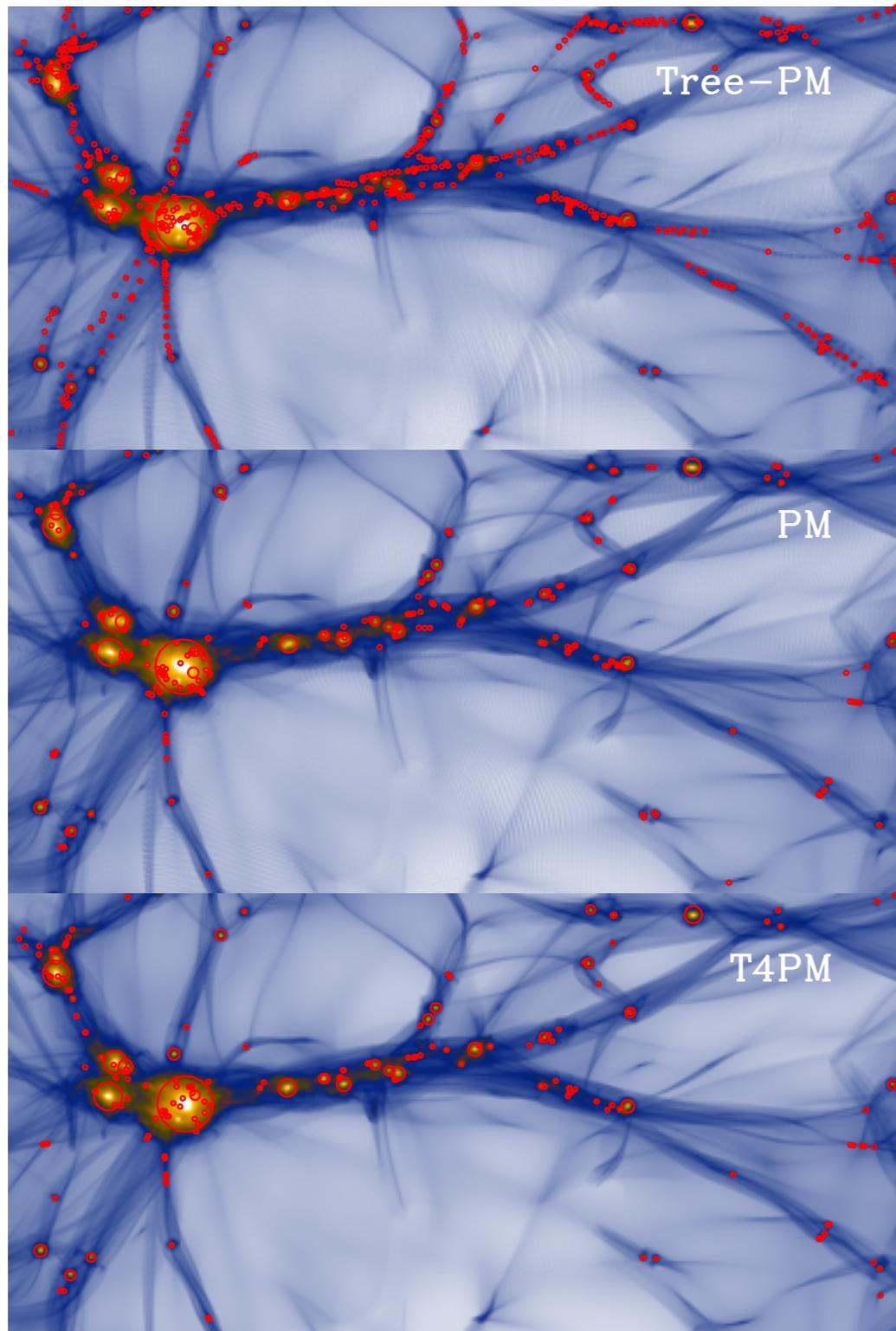
fixed mass resolution, varying force resolution:



force res. \rightarrow
features become sharper
fragmentation appears

sheet tessellation
based method cures
artificial fragmentation

First determination of WDM halo mass function!



Angulo, Hahn & Abel 2013

Limitations - diffusion/loss of energy cons.

Mixing - (phase or chaotic)

need increasingly larger number of elements to trace the sheet surface



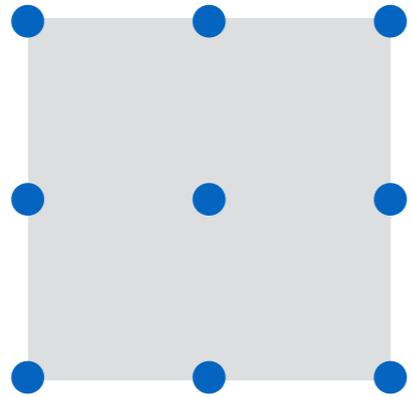
hi-res N-body



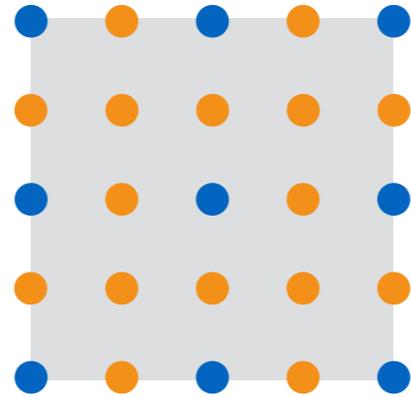
tesselated cube orbiting
in non-harmonic potential

Need adaptive refinement

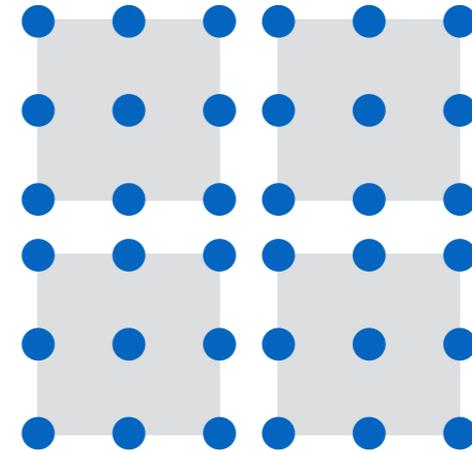
adaptive refinement:



a. element is flagged for refinement

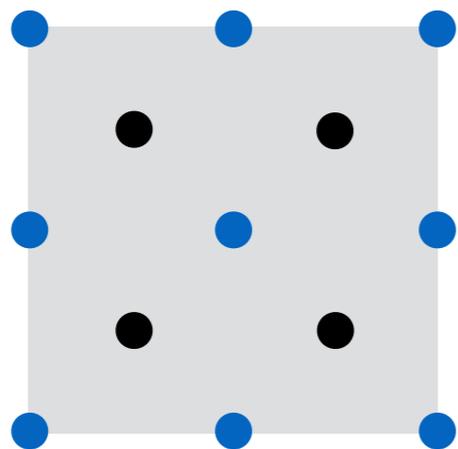


b. positions and velocities are determined at mid-points

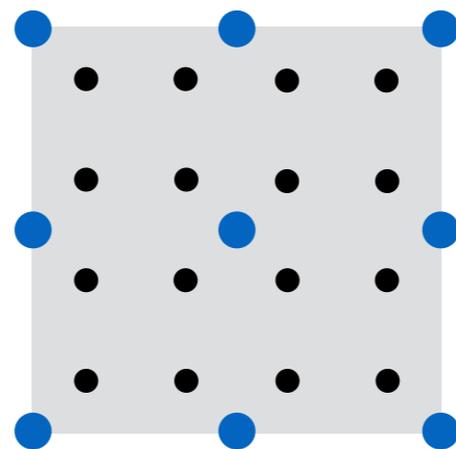


c. new elements are created using the mid-point values

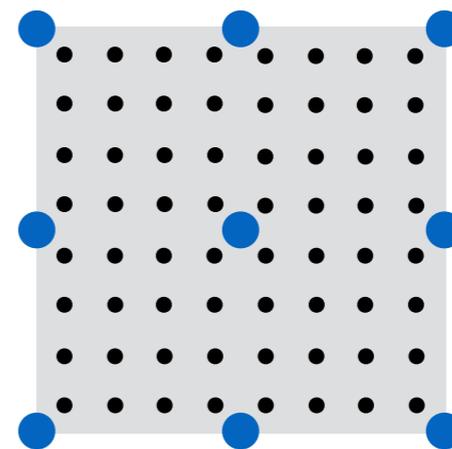
approximate element mass distribution by recursively deposited 'mass carrier particles' (these are not *active*, *i.e.* no degrees of freedom)



m=1



m=2



m=3

Hahn & Angulo 2015

refinement + higher order!



hi-res N-body



tesselated cube orbiting
in non-harmonic potential



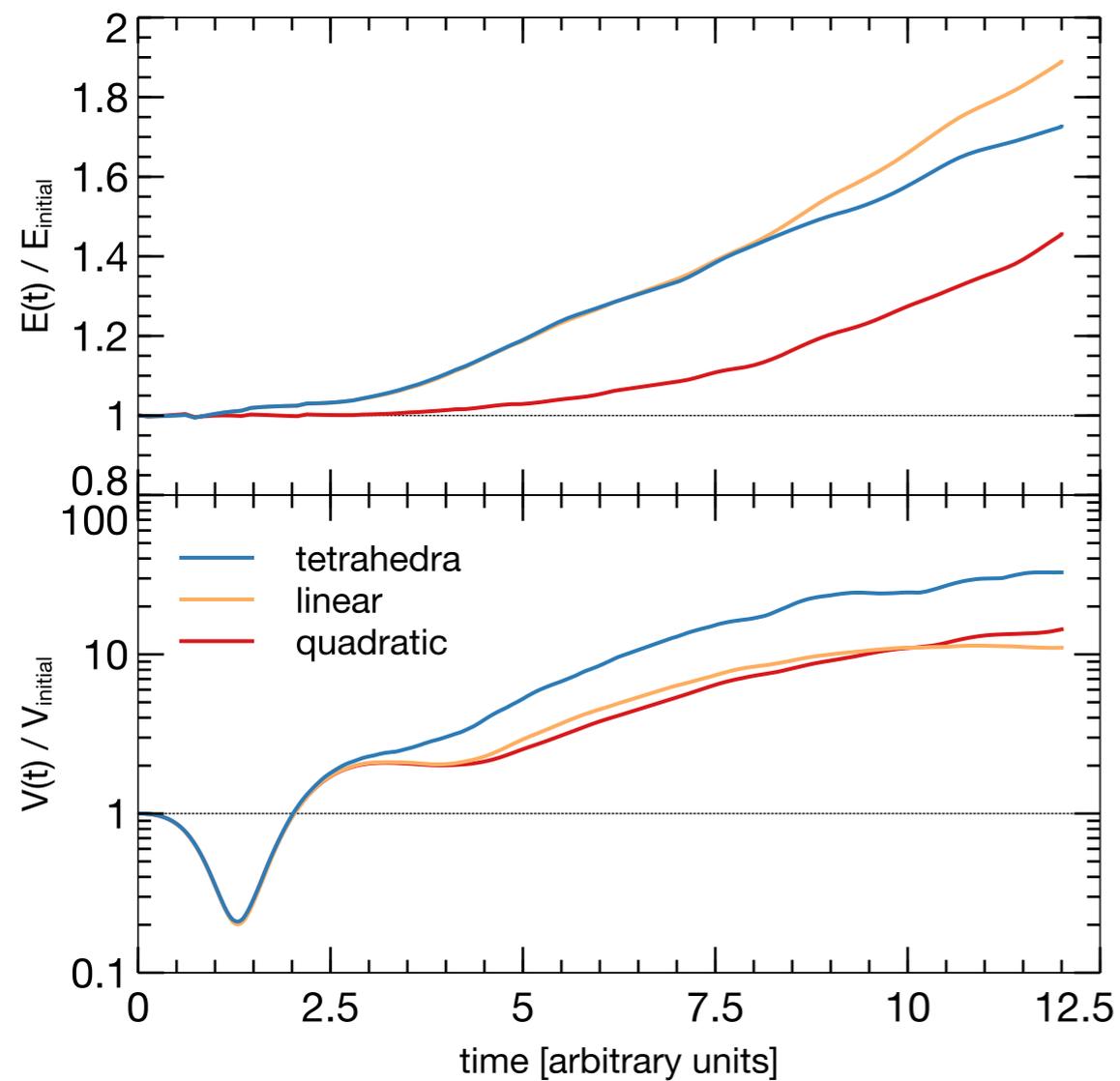
adaptively refined tri-quadratic
phase-space element

first alternative to N-body in highly non-linear regime!
+ able to track fine-grained phase space

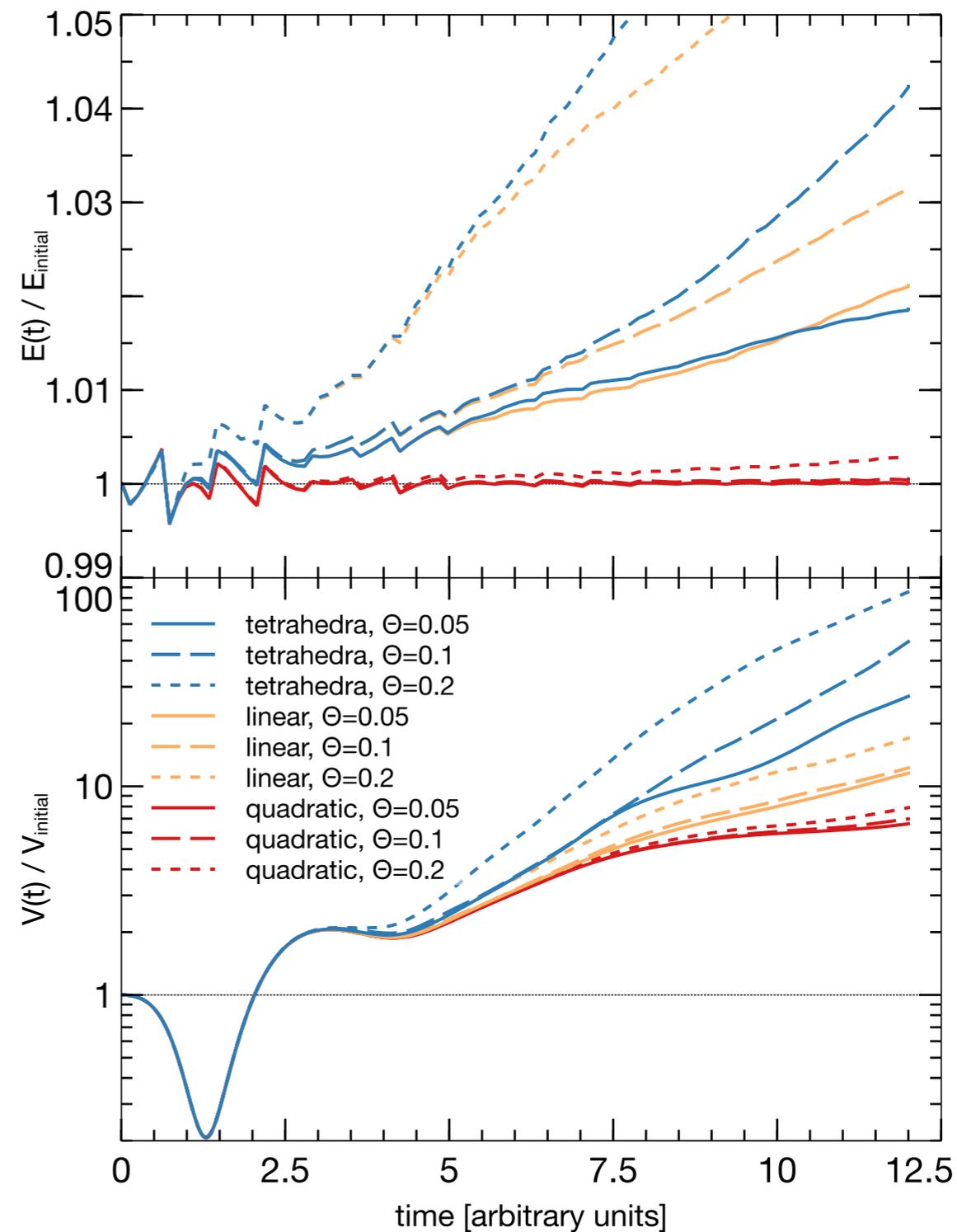
Hahn & Angulo 2015

Orbit test

no refinement

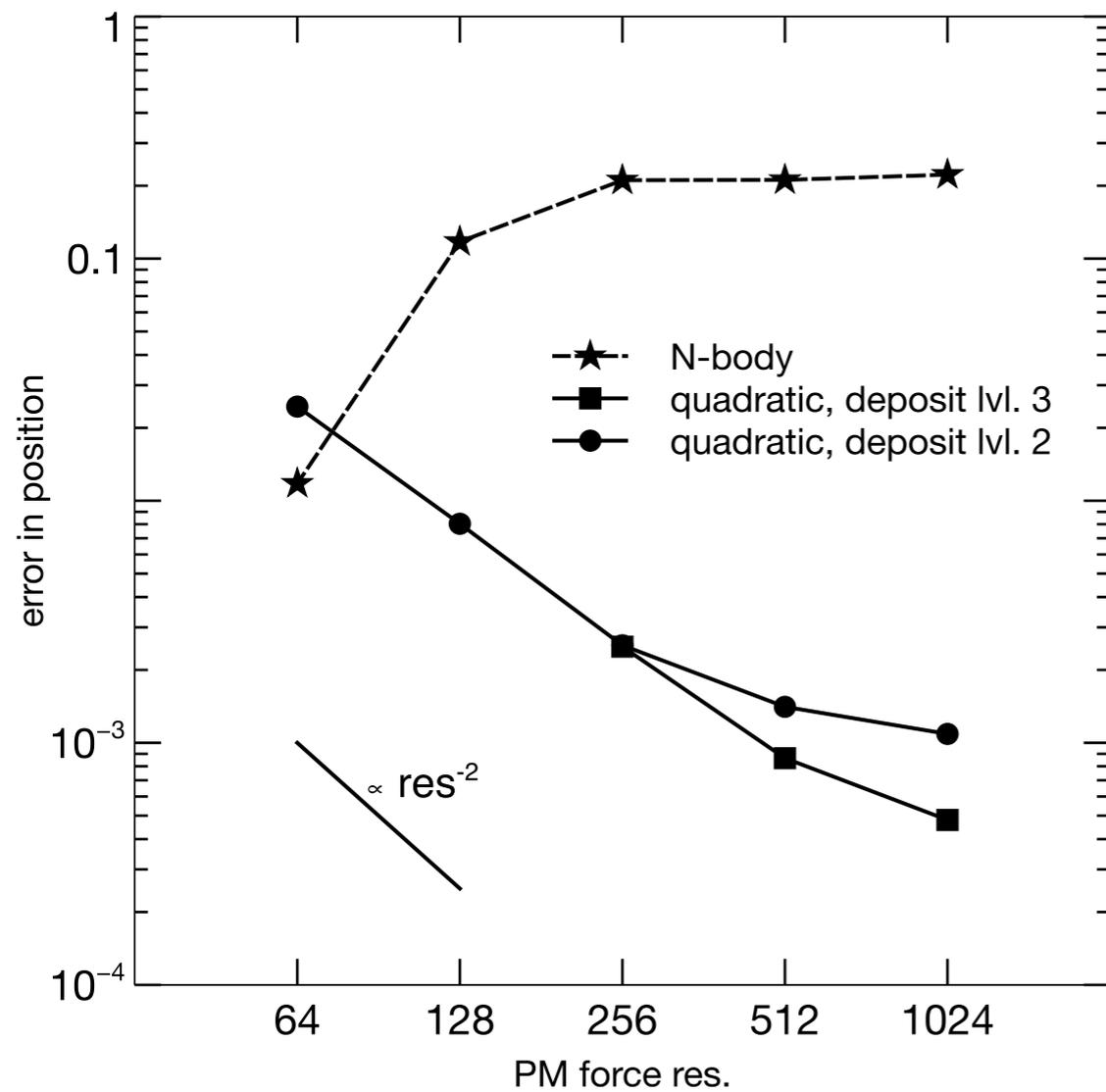


refinement

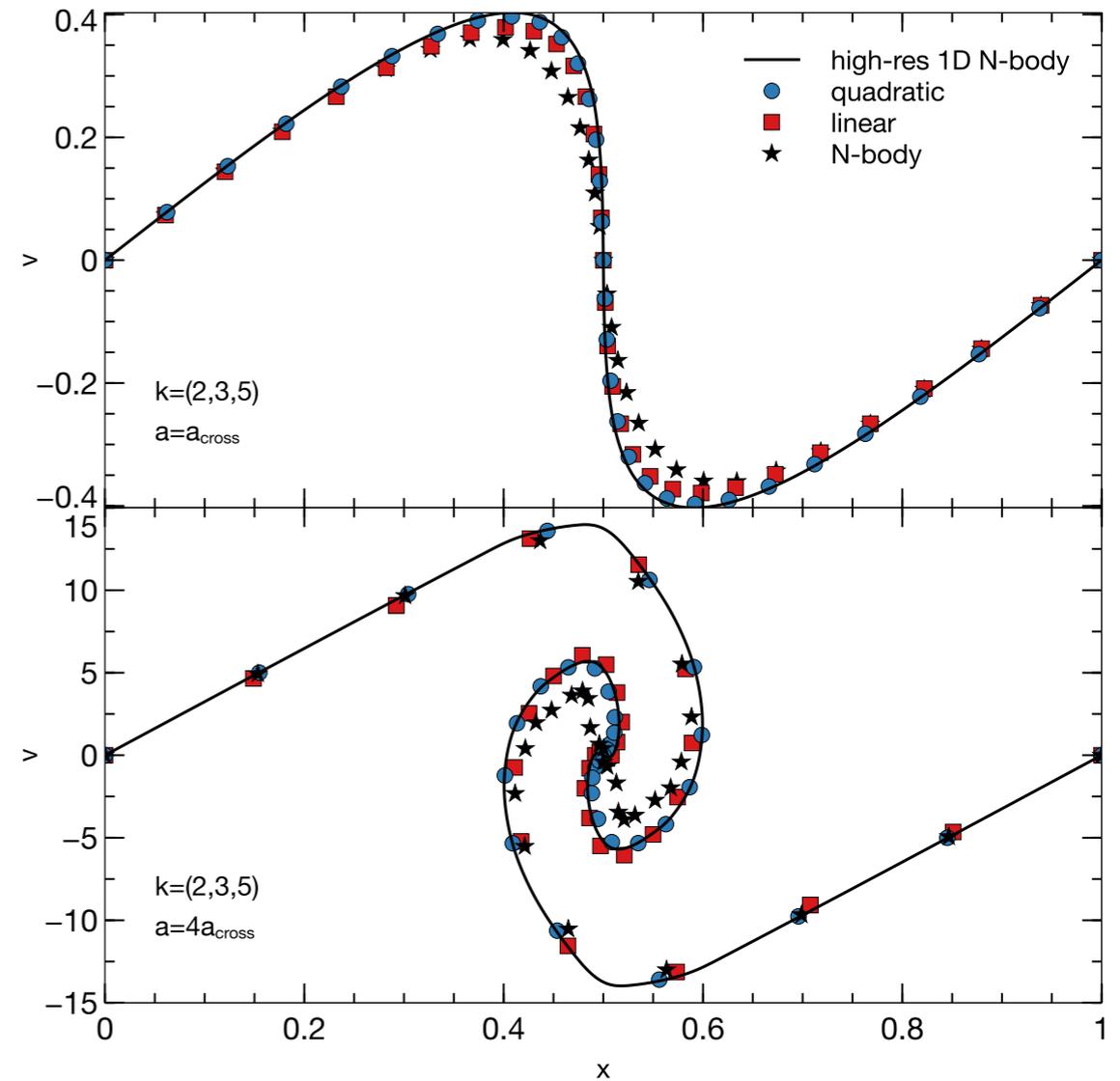


Self-gravitating tests 1D

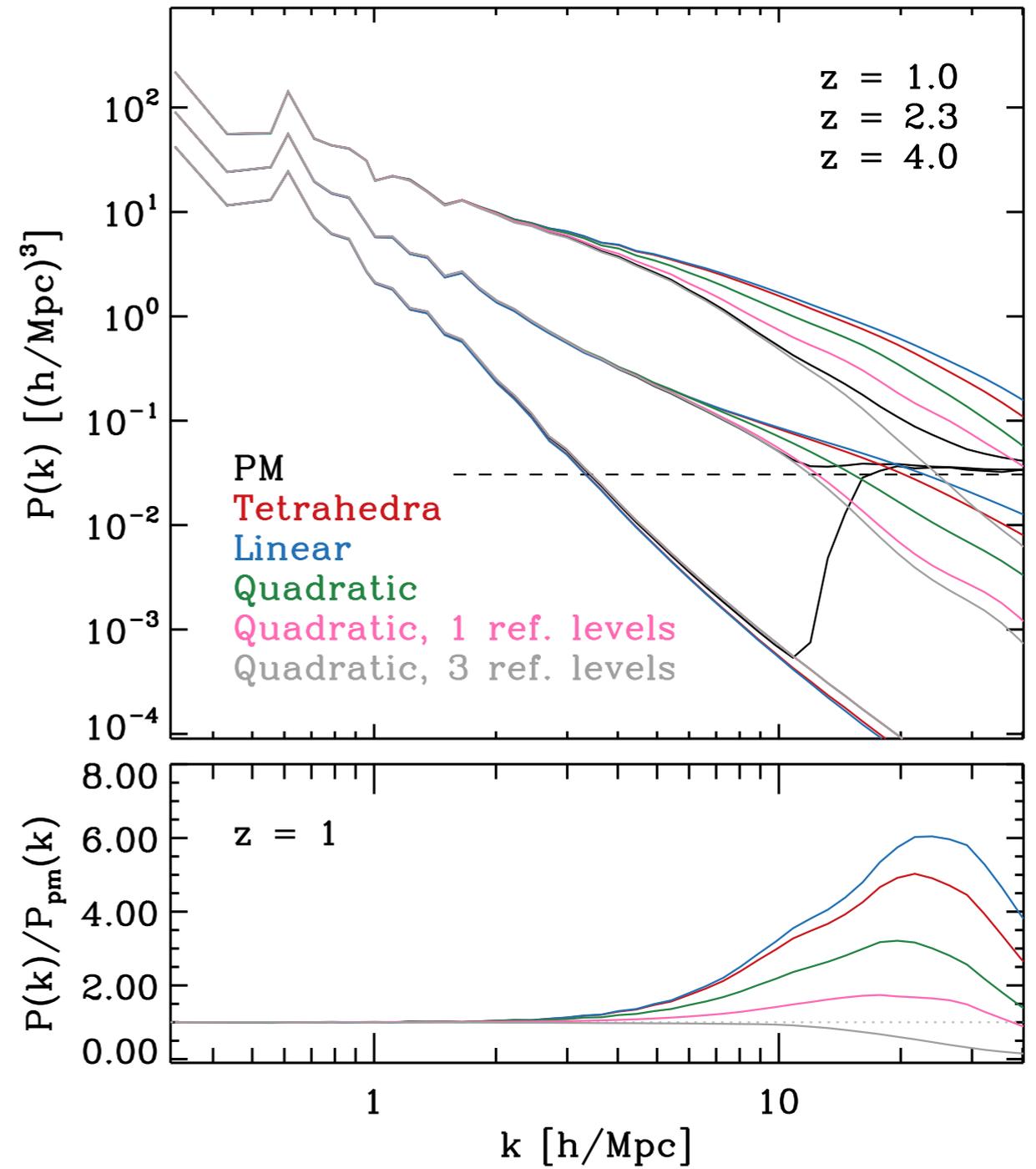
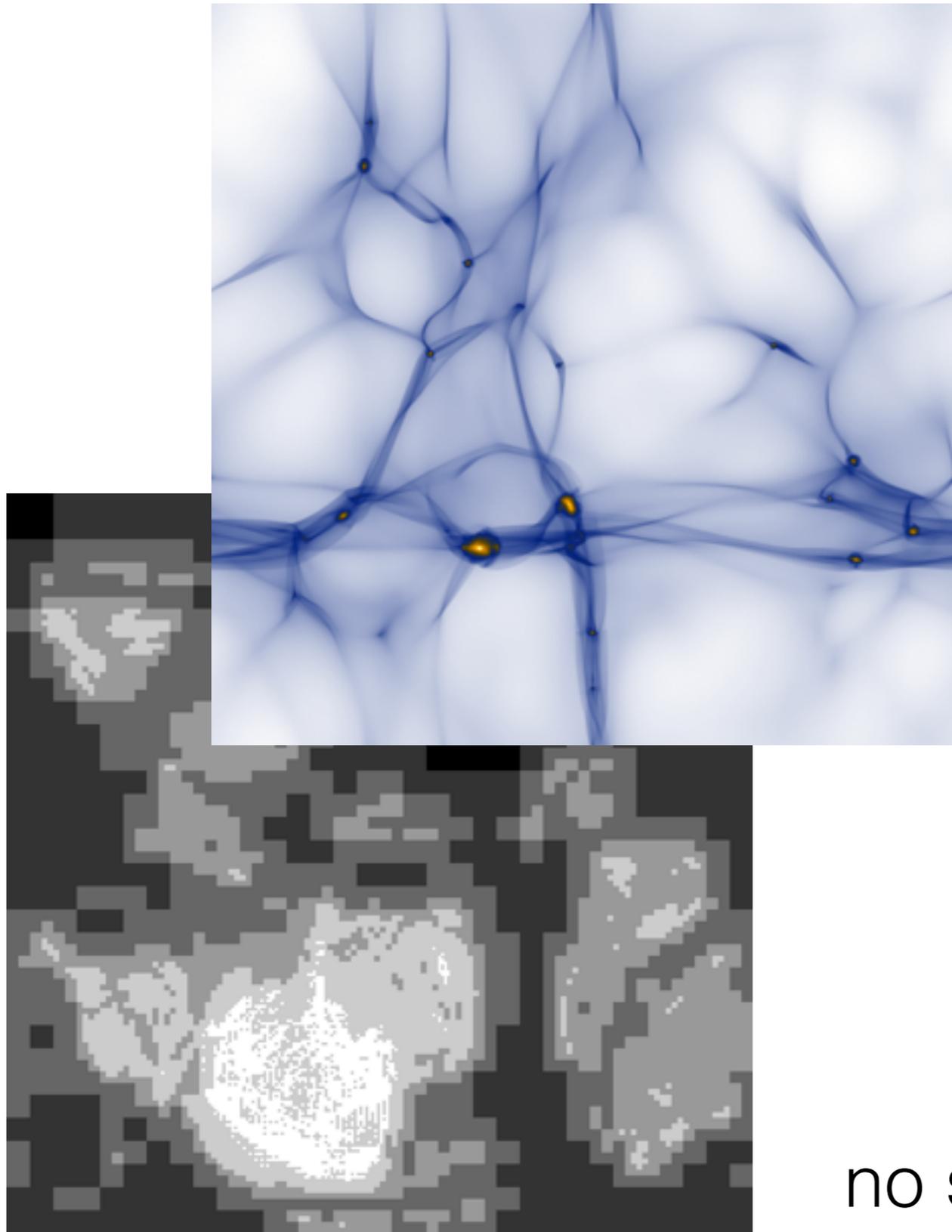
32³ particle plane wave,
axis aligned



32³ particle plane wave,
oblique



let's go cosmological



no shotnoise!!!

$a = 0.015625000$

Conclusions

- Lagrangian elements can give new insights into existing simulations (density/velocity fields, multi-stream analysis,...)
- Provide also self-consistent simulation technique.
(functional when using high-order and adaptive refinement)
- Solves two-body and fragmentation problems of N-body
- First methodological test of N-body in deeply non-linear regime
- Stay tuned for halo properties...