Effects of the Sources of Reionization on 21cm Redshift Space Distortions

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Main Anisotropies along the LoS

• Redshift Space Distortions.

 Effect of the Finite Light Travel Time (Light Cone Effect).

Redshift Space Distortion







Observed power spectrum will be "anisotropic" (or LoS dependent)



 $P^{s}(k,\mu)$







Power spectrum in Redshift Space

$$P^{s}(k,\mu) = \sum_{l \text{ even}} \mathcal{P}_{l}(\mu) P_{l}^{s}(k)$$

$$P_{l}^{s}(k) = \frac{(2l+1)}{4\pi} \int \mathcal{P}_{l}(\mu) P^{s}(k) d\Omega$$

• Kaiser, N., 1987, MNRAS, 227, 1

o Hamilton, A. J. S., 1992, APJL, 385, L5

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$$P_0^s = \overline{\delta T_b}^2(z) \left[\frac{1}{5} P_{\rho_M \rho_M} + P_{\rho_{H \ I} \rho_{H \ I}} + \frac{2}{3} P_{\rho_{H \ I} \rho_M} \right]$$

$$P_2^s = 4 \,\overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_M \rho_M} + \frac{1}{3} P_{\rho_{H \ I} \rho_M} \right]$$

$$P_4^s = \frac{8}{35} \overline{\delta T_b}^2(z) P_{\rho_M \rho_M} \quad \text{when } T_{\text{S}} >> T_{\text{CMB}}$$

$$P_0^s = \overline{\delta T_b}^2(z) \left[\frac{1}{5} (P_{\rho_M \rho_M}) + P_{\rho_{H \ I} \rho_{H \ I}} + \frac{2}{3} P_{\rho_{H \ I} \rho_M} \right]$$

$$P_{2}^{s} = 4 \,\overline{\delta T_{b}}^{2}(z) \left[\frac{1}{7} P_{\rho_{M}\rho_{M}} + \frac{1}{3} P_{\rho_{H}} \right]$$

$$P_4^s = \frac{8}{35} \overline{\delta T_b}^2(z) P_{\rho_M \rho_M} \quad \text{when } T_{\text{S}} >> T_{\text{CMB}}$$

$$P_0^s = \overline{\delta T_b}^2(z) \left[\frac{1}{5} P_{\rho_M \rho_M} + P_{\rho_{H \ I} \rho_H} + \frac{2}{3} P_{\rho_{H \ I} \rho_M} \right]$$

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$$P_0^s = \overline{\delta T_b}^2(z) \left[\frac{1}{5} P_{\rho_M \rho_M} + P_{\rho_{H^{-1}} \rho_{H^{-1}}} + \frac{2}{3} P_{\rho_{H^{-1}} \rho_M} \right]$$
$$P_2^s = 4 \overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_M \rho_M} + \frac{1}{3} P_{\rho_{H^{-1}} \rho_M} \right]$$
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$$P_0^s = \overline{\delta T_b}^2(z) \left[\frac{1}{5} P_{\rho_M \rho_M} + P_{\rho_{H \ I} \rho_{H \ I}} + \frac{2}{3} P_{\rho_{H \ I} \rho_M} \right]$$

Quadrupole Moment

$$P_2^s = 4 \overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_M \rho_M} + \frac{1}{3} P_{\rho_{H^{-1}} \rho_M} \right]$$

What happens to the quadrupole moment under different reionization source models?

Majumdar, Bhardwaj & Choudhury, 2013, MNRAS, 434, 3
 Majumdar, Molloma, Datta et al. 2014, MNRAS, 443, 4

o Majumdar, Mellema, Datta et al., 2014, MNRAS, 443, 4

Effect of Spin Temperature Fluctuations

$$P_2^s = 4 \overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_M \rho_M} + \frac{1}{3} P_{\rho_{H I} \rho_M} \right]$$

$$P_2^s = 4 \,\overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_{\rm M},\rho_{\rm M}} + \frac{1}{3} P_{\rho_{\rm HI},\rho_{\rm M}} + \frac{1}{3} P_{\eta,\rho_{\rm M}} \right]$$

$$\eta(z, \mathbf{x}) = 1 - \frac{T_{\rm CMB}(z)}{T_{\rm S}(z, \mathbf{x})}$$

Effect of Spin Temperature Fluctuations

$$P_2^s = 4 \overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_M \rho_M} + \frac{1}{3} P_{\rho_{H \ I} \rho_M} \right]$$

$$P_2^s = 4 \overline{\delta T_b}^2(z) \left[\frac{1}{7} P_{\rho_{\mathrm{M}},\rho_{\mathrm{M}}} + \frac{1}{3} P_{\rho_{\mathrm{HI}},\rho_{\mathrm{M}}} + \frac{1}{3} P_{\eta,\rho_{\mathrm{M}}} \right]$$
$$\eta(z, \mathbf{x}) = 1 - \frac{T_{\mathrm{CMB}}(z)}{T_{\mathrm{S}}(z, \mathbf{x})}$$

Effect of Spin Temperature Fluctuations



Simulation

- DM distribution: P³M N-body simulation (from PRACE4LOFAR project, Dixon et al. in prep.)
- Box size: (714.28 Mpc)³
- \circ Minimum halo mass used: 2.0x10⁹ M_{sun}
- 21-cm brightness temperature fields: Excursion set based semi-numerical formalism (on a 600³ grid)
- All source models are tuned to have same reionization history (i.e. x_{HI} vs z)

Source Models

Reionization scenario	UV	UIB	SXR	PL n	Non-uniform recombination
Fiducial	100%	_	_	1.0	No
Clumping	100%	_	_	1.0	Yes
UIB dominated	20%	80%	_	1.0	No
SXR dominated	20%	_	80%	1.0	No
UV+SXR+UIB	50%	10%	40%	1.0	No
PL 2.0	_	_	_	2.0	No
PL 3.0	_	—	_	3.0	No

Source Mod	dels	Unifo	rm Ioniz	zation F	Background
Reionization scenario	UV	UB	SXR	PL n	Non-uniform recombination
Fiducial	100%	_	_	1.0	No
Clumping	100%	_	_	1.0	Yes
(UIB dominated)	20%	80%	_	1.0	No
SXR dominated	20%	_	80%	1.0	No
UV+SXR+UIB	50%	10%	40%	1.0	No
PL 2.0	_	_	_	2.0	No
PL 3.0	—	—	_	3.0	No

Soft X-ray Photons					
Reionization scenario	UV	UIB	SXR	PL n	Non-uniform recombination
Fiducial	100%	_	_	1.0	No
Clumping	100%	_	_	1.0	Yes
UIB dominated	20%	80%	_	1.0	No
SXR dominated	20%	_	80%	1.0	No
UV+SXR+UIB	50%	10%	40%	1.0	No
PL 2.0	_	_	_	2.0	No
PL 3.0	—	—	—	3.0	No

Source Models

Reionization scenario	UV	UIB	SXR	PL n	Non-uniform recombination
Fiducial	100%	_	_	1.0	No
Clumping	100%	_	_	1.0	Yes
UIB dominated	20%	80%	_	1.0	No
SXR dominated	20%	_	80%	1.0	No
UV+SXR+UIB	50%	10%	40%	10	No
PL 2.0	_	_	_	$\left(2.0\right)$	No
PL 3.0	_	_	—	3.0	No

No. of photons \propto (halo mass)ⁿ



Monopole Moment







Cross-Correlation

X=Aeîiθ Y=Beîiφ

 $P\downarrow XY = |A||B|\cos(\varphi - \theta)$





Light Cone Effect on the Quadrupole Moment



Robustness of the Quadrupole Moment & EoR History



Robustness of the Quadrupole Moment & EoR History



Robustness of the Quadrupole Moment & EoR History

$$\bar{x}_{\mathrm{H}\ \mathrm{I}}(z) = \begin{cases} 0 & \text{if } z < z_r \\ 1 - \exp[-\beta(z - z_r)] & \text{if } z \ge z_r \end{cases}$$

$$\mathcal{L} = \prod_{i} \frac{1}{\sqrt{2\pi\sigma}} \exp\left[\frac{\left(P_{2,i}^{s,\text{meas}} - P_{2,i}^{s,\text{true}}\right)^2}{2\sigma^2}\right]$$



Summary

- The quadrupole moment of the power spectrum provides a very robust estimate of the EoR 21-cm signal compared to the monopole moment.
- It is very insensitive to the properties of the reionization sources (or ionization topology) as long as their spatial distribution is even loosely correlated to the distribution of collapsed bound structures.
- It is also expected to be less sensitive to the spin temperature fluctuations compared to the monopole moment.
- This roboustness of the quadrupole moment can be used to extract the history of EoR very efficiently.

Thanks



Robustness of the Quadrupole Moment







Considering the model with quasi-linear approximations ----

$$P^{s}(k,\mu) = \overline{\delta T_{b}}^{2}(z) \left[P_{\rho_{H}}{}_{I}\rho_{H}{}_{I}(k) + 2\mu^{2} P_{\rho_{H}}{}_{I}\rho_{M}(k) + \mu^{4} P_{\rho_{M}}{}_{\rho_{M}}(k) \right]$$

when $T_S >> T_{CMB}$

Considering the model with quasi-linear approximations ----

$$P^{s}(k,\mu) = \overline{\delta T_{b}}^{2}(z) \left[P_{\rho_{H_{1}}\rho_{H_{1}}}(k) + 2\mu^{2} P_{\rho_{H_{1}}\rho_{M}}(k) + \mu^{4} P_{\rho_{M}\rho_{M}}(k) \right]$$

Power Spectrum of total Matter Density Field

Considering the model with quasi-linear approximations ----

$$P^{s}(k,\mu) = \overline{\delta T_{b}}^{2}(z) \left[P_{\rho_{H_{1}}\rho_{H_{1}}}(k) + 2\mu^{2} P_{\rho_{H_{1}}\rho_{M}}(k) + \mu^{4} P_{\rho_{M}\rho_{M}}(k) \right]$$

Power Spectrum of HI Density Field

Considering the model with quasi-linear approximations ----

$$P^{s}(k,\mu) = \overline{\delta T_{b}}^{2}(z) \left[P_{\rho_{H_{1}}\rho_{H_{1}}}(k) + 2\mu^{2} P_{\rho_{H_{1}}\rho_{M}}(k) + \mu^{4} P_{\rho_{M}\rho_{M}}(k) \right]$$

Cross-Power Spectrum of HI and Matter Density Fields

Considering the model with quasi-linear approximations ----

$$P^{s}(k,\mu) = \overline{\delta T_{b}}^{2}(z) \left[P_{\rho_{H_{1}}\rho_{H_{1}}}(k) + 2\mu^{2} P_{\rho_{H_{1}}\rho_{M}}(k) + \mu^{4} P_{\rho_{M}\rho_{M}}(k) \right]$$





Monopole Moment



Effect of this anisotropy on the 21-cm power spectrum will be significant.



 $P^{s}(k,\mu)$

- Bharadwaj & Ali, 2004, MNRAS, 352, 142
- o Bharadwaj & Ali, 2005, MNRAS, 356, 4
- o Barkana & Loeb, 2005, ApJL, 624, L65



