

# 7th International Workshop on Multiple Partonic Interactions at the LHC (Trieste, Italy, November 22-27, 2015)

CMS Experiment at LHC, CERN  
Data recorded: Thu Sep 13 05:21:23 2012 CEST  
Run/Event: 202792 / 1737666483  
Lumi section: 918  
Orbit/Crossing: 240400935 / 1986

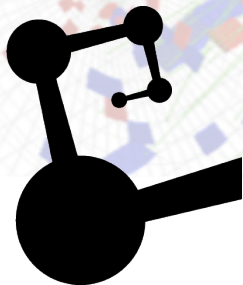
## Jet effects in high multiplicity pp events

**Antonio Ortiz**

**November 25, 2015**

Work in collaboration with: Gyula Bencedi,  
Héctor Bello and Satyajit Jena

Instituto de  
Ciencias  
Nucleares  
UNAM



# Outline

- ☐ Introduction
- ☐ Tools
- ☐ Particle production as a function of the event multiplicity and hardness
- ☐ Energy dependence
- ☐ Summary



# INTRODUCTION



# Introduction

- ❑ Small systems (like those produced in pp and p-Pb collisions) have attracted the attention of the heavy ion community because:
  - ❑ In high multiplicity events, sQGP-like signatures have been found (flow & long range azimuthal correlations)
  - ❑ The origin of such effects is still unknown
  - ❑ More differential studies are needed



# Introduction

A hydro-inspired model (Blast-Wave):

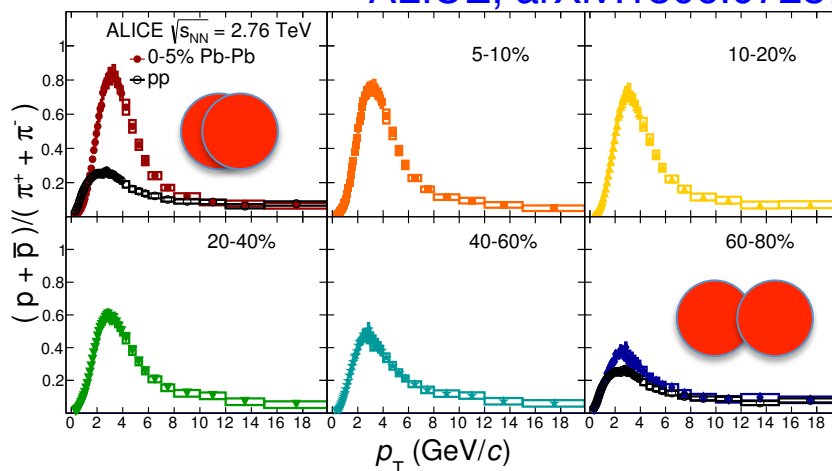
$$\frac{1}{p_T} \frac{dN}{dp_T} \propto \int_0^R r dr m_T I_0 \left( \frac{p_T \sinh \rho}{T_{\text{kin}}} \right) K_1 \left( \frac{m_T \cosh \rho}{T_{\text{kin}}} \right)$$

$$\rho = \tanh^{-1} \beta_T = \tanh^{-1} \left( \left( \frac{r}{R} \right)^n \beta_s \right)$$

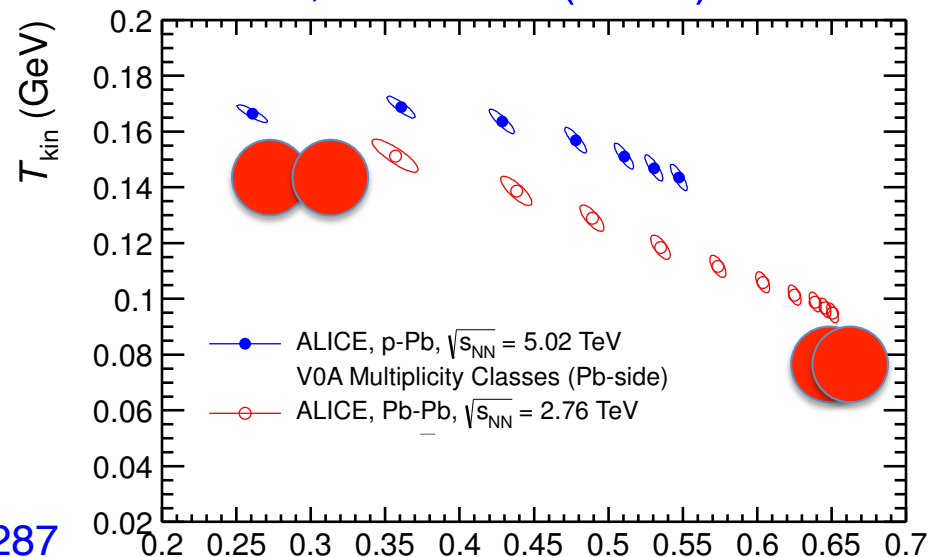
Describes the  $p_T$  spectra of identified hadrons in:

□ p-Pb and Pb-Pb data

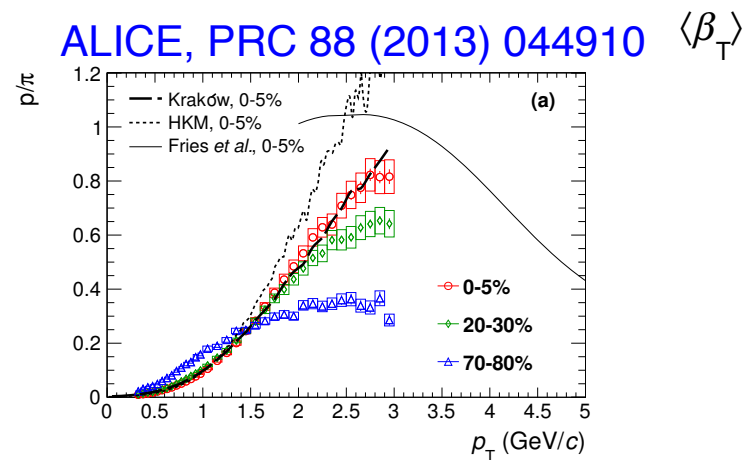
ALICE, arXiv:1506.07287



ALICE, PLB 728 (2014) 25-38



ALICE, PRC 88 (2013) 044910



# Introduction

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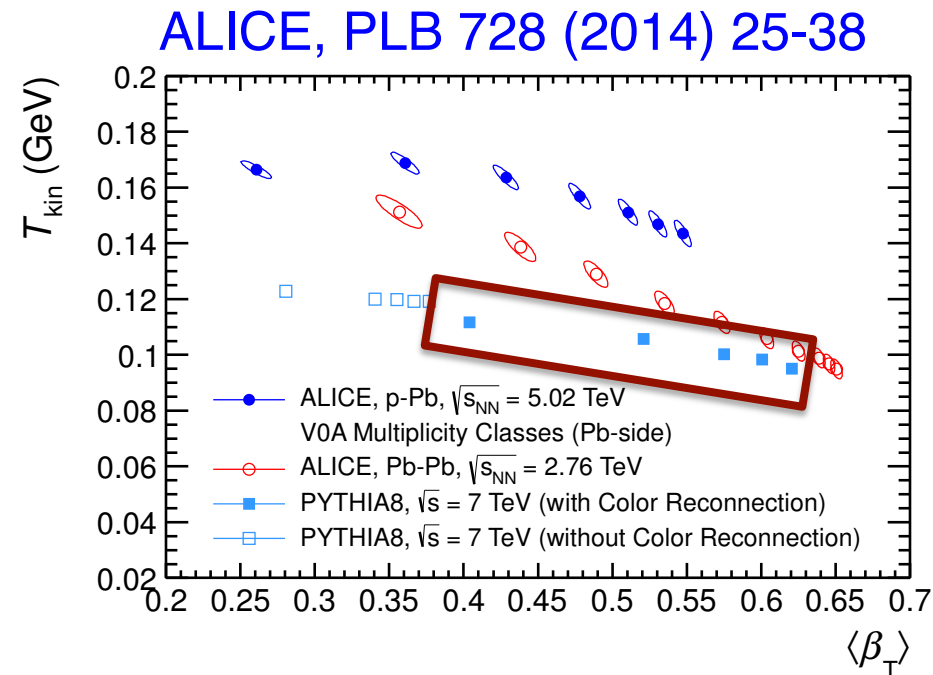
Describes the  $p_T$  spectra of identified hadrons in:

□ p-Pb and Pb-Pb data

□ Also the  $p_T$  distributions generated with Pythia (where no hydrodynamical evolution is assumed)

□ It has been discussed that color reconnection (CR) produces radial flow-like patterns due to boosted strings

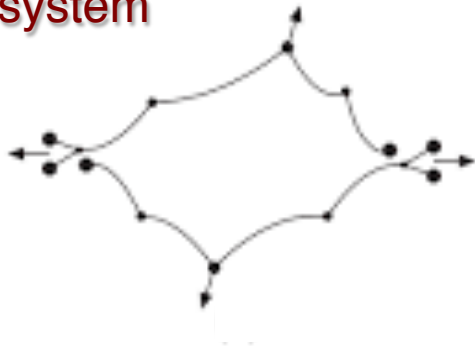
G. Paic, E. Cuautle, P. Christiansen, I. Maldonado and A. O., PRL 111 (2013) 042001



# Introduction

\* Figure taken from: G. Gustafson, [Acta Phys. Polon. B40, 1981 \(2009\)](#)

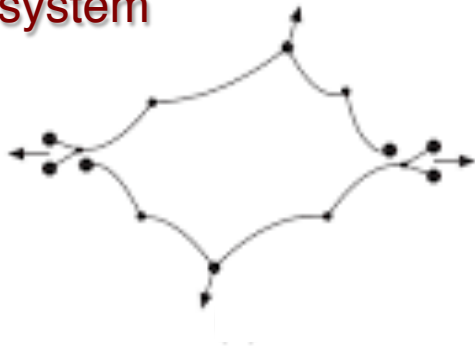
1<sup>st</sup> partonic  
system



# Introduction

\* Figure taken from: G. Gustafson, [Acta Phys. Polon. B40, 1981 \(2009\)](#)

1<sup>st</sup> partonic  
system



+2<sup>nd</sup> partonic  
system

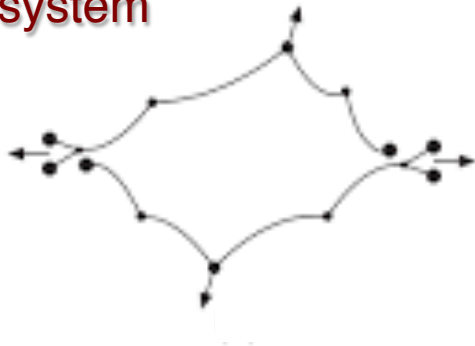




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1<sup>st</sup> partonic  
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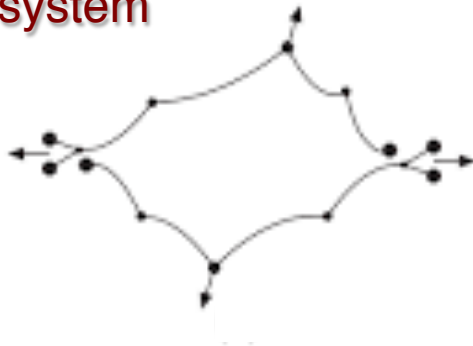
When CR is activated



# Introduction

\* Figure taken from: G. Gustafson, [Acta Phys. Polon. B40, 1981 \(2009\)](#)

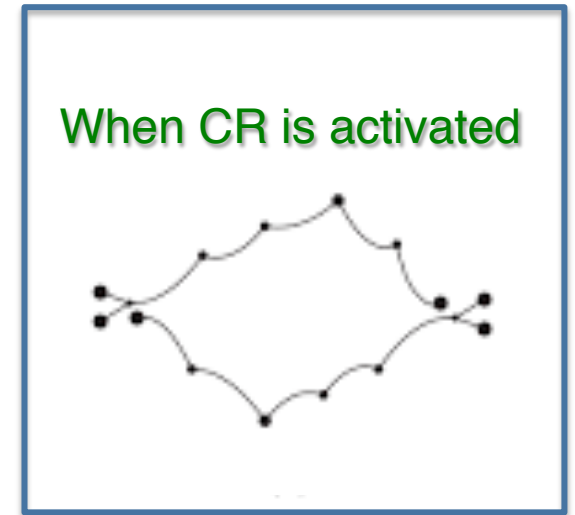
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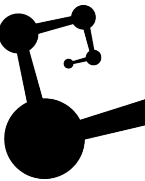
When CR is activated



This was the focus of this work:  
[PRL 111 \(2013\) 042001](#)

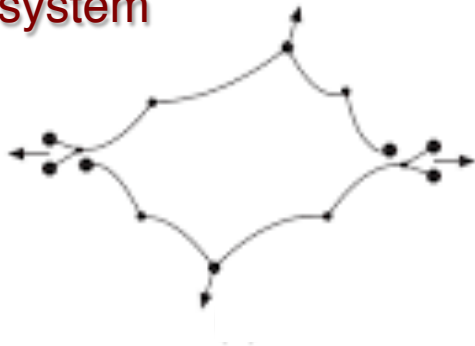
- The more  $N_{\text{MPI}}$  the higher the flow-like effect

# Introduction



\* Figure taken from: G. Gustafson, [Acta Phys. Polon. B40, 1981 \(2009\)](#)

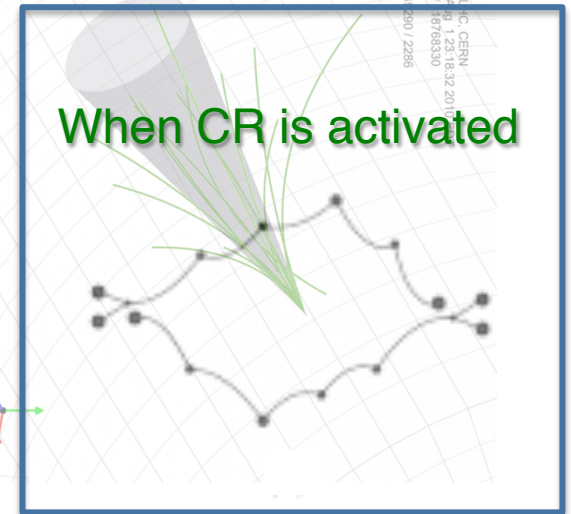
1<sup>st</sup> partonic  
system



+2<sup>nd</sup> partonic  
system



When CR is activated



Due to the large  $N_{\text{MPI}}$  a high  $p_T$  jet in the event is expected (high probability):

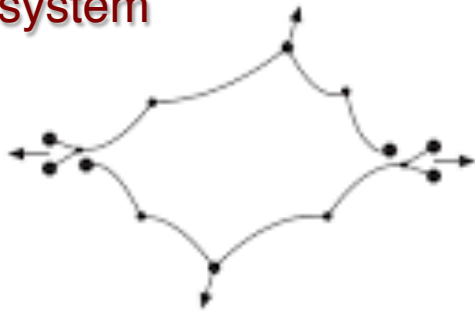
- ☐ Can we quantify the effects of the high  $p_T$  jets?
- ☐ I would expect a higher boost with increasing the parton  $p_T$

# Introduction



\* Figure taken from: G. Gustafson, *Acta Phys. Polon. B40*, 1981 (2009)

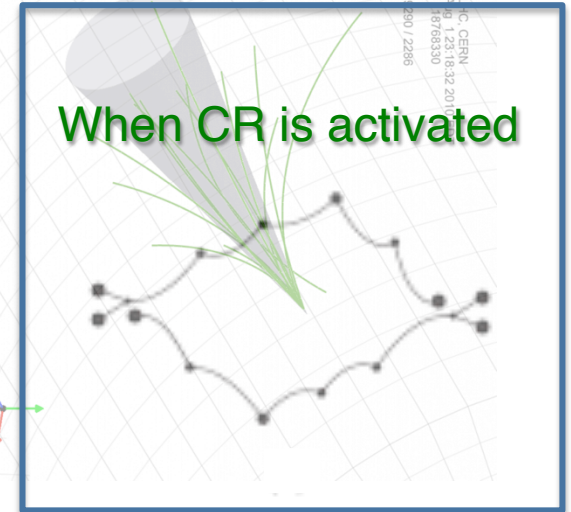
1<sup>st</sup> partonic  
system



+2<sup>nd</sup> partonic  
system



When CR is activated



In the CR model used in the tune Monash 2013 (Mo2013), an MPI system with a scale  $p_T$  of the hard interaction (normally  $2 \rightarrow 2$ ) can be merged with one of a harder scale with a probability that is:

$$P(p_T) = \frac{(RR \times p_{T0})^2}{(RR \times p_{T0})^2 + p_T^2}$$

Reconnection Range ( $RR$ ): 0-10

Tune Monash 2013:  $RR \times p_{T0} \approx 3$

<http://home.thep.lu.se/~torbjorn/pythia82html/Welcome.html>



# Introduction

\* Figure taken from: G. Gustafson, [Acta Phys. Polon. B40, 1981 \(2009\)](#)

**This work:** study the properties of the pp events as a function of their multiplicity (z =  $dN/d\eta / \langle dN/d\eta \rangle$ ) & their jet content (leading jet  $p_T$ )

in CR is activated

multiplicity ( $z = dN/d\eta / \langle dN/d\eta \rangle$ )

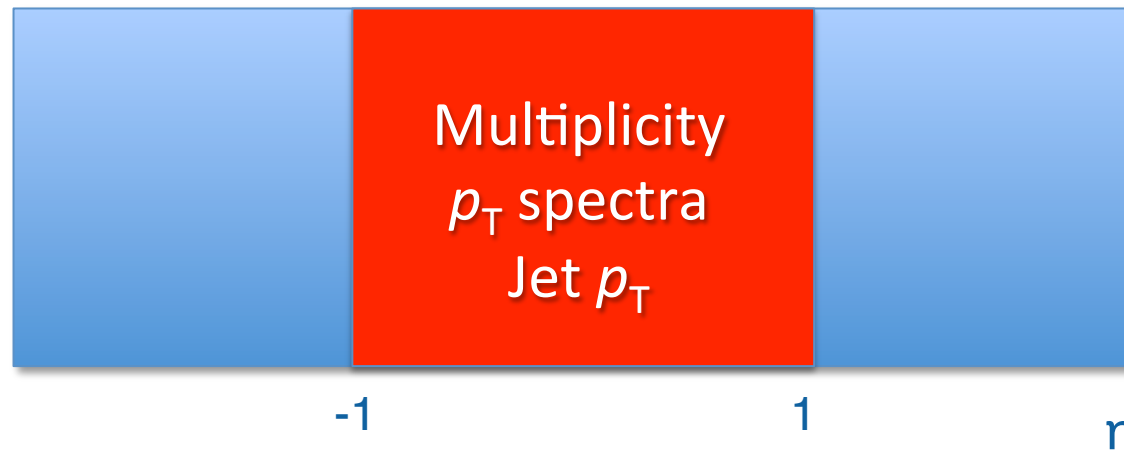
& their jet content (leading jet  $p_T$ )

# Tools

- ❑ Generator: **Pythia 8.212**, [T. Sjöstrand et. al, CPC191 \(2005\) 159](#)
  - ❑ Tune Monash 2013, [P. Skands, EPJC74 \(2014\) 8, 3024](#)
  - ❑ 900M events
    - ❑ 7 TeV (reference), 0.9 TeV, 2.76 TeV and 13 TeV
- ❑ Jet Finder: **FastJet 3.1.3**, [M. Cacciari et al., EPJC72\(2012\)1896](#)
  - ❑ Anti- $k_T$  algorithm
  - ❑  $R=0.4$
  - ❑  $p_T^{\min} = 5 \text{ GeV}$
  - ❑ Visible particles (Pythia definition) are considered for the jet reconstruction

# INCLUSIVE PARTICLE PRODUCTION AS A FUNCTION OF THE EVENT MULTIPLICITY AND HARDNESS





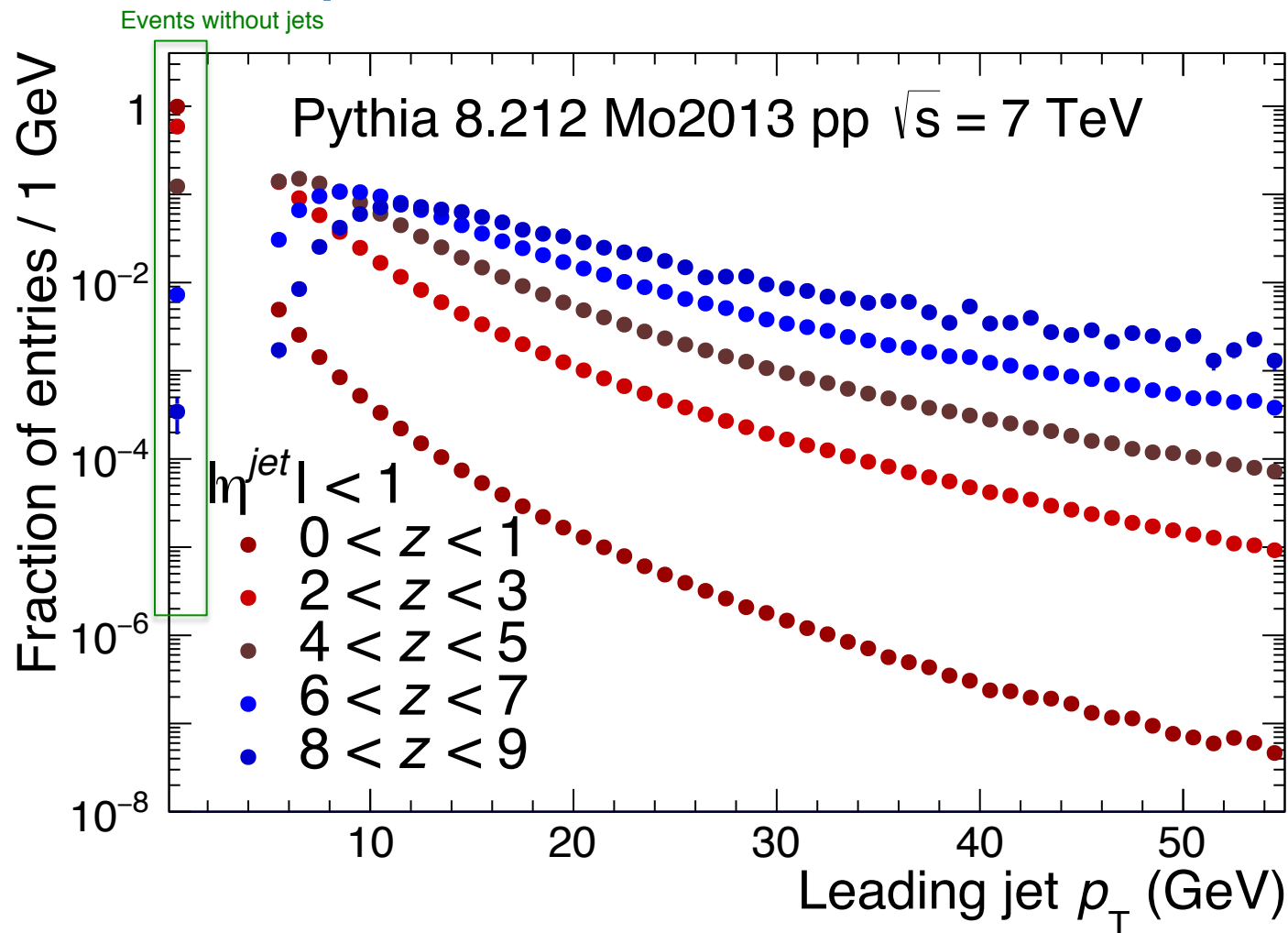
\* The underlying event contribution to the jet  $p_T$  was not studied, because we are only interested in the event classification

# INCLUSIVE PARTICLE PRODUCTION AS A FUNCTION OF THE EVENT MULTIPLICITY AND HARDNESS





# $p_T^{\text{jet}}$ vs. multiplicity

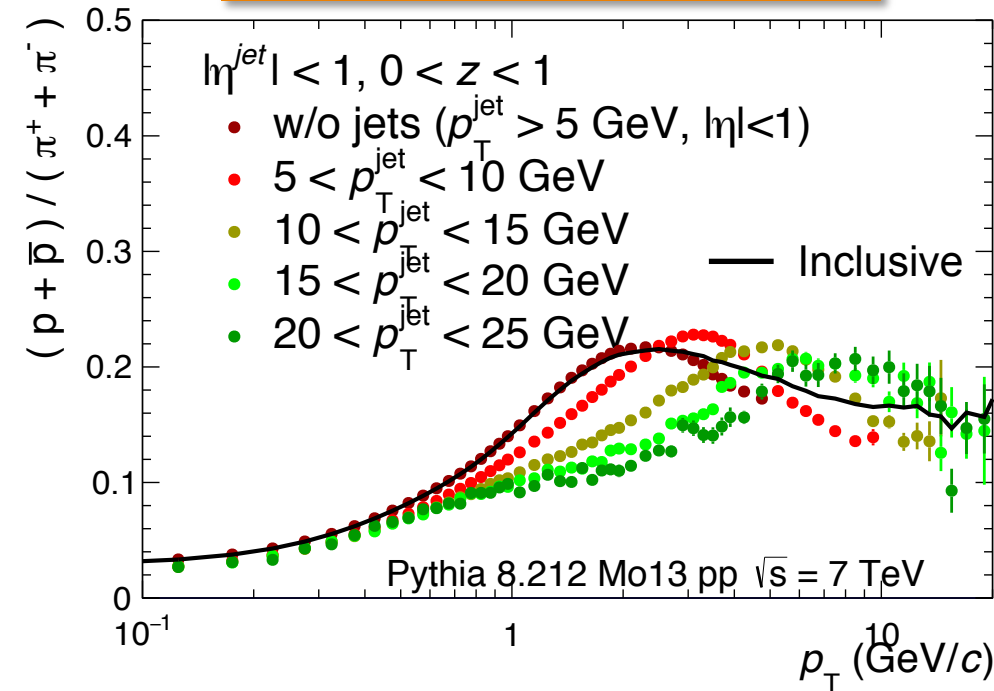


$\langle z \rangle$	$\langle p_T^{\text{jet}} \rangle$ (GeV)
0.5	7.09
1.5	7.49
2.5	7.83
3.5	8.48
4.5	9.55
5.5	11.1
6.5	13.2
7.5	15.85

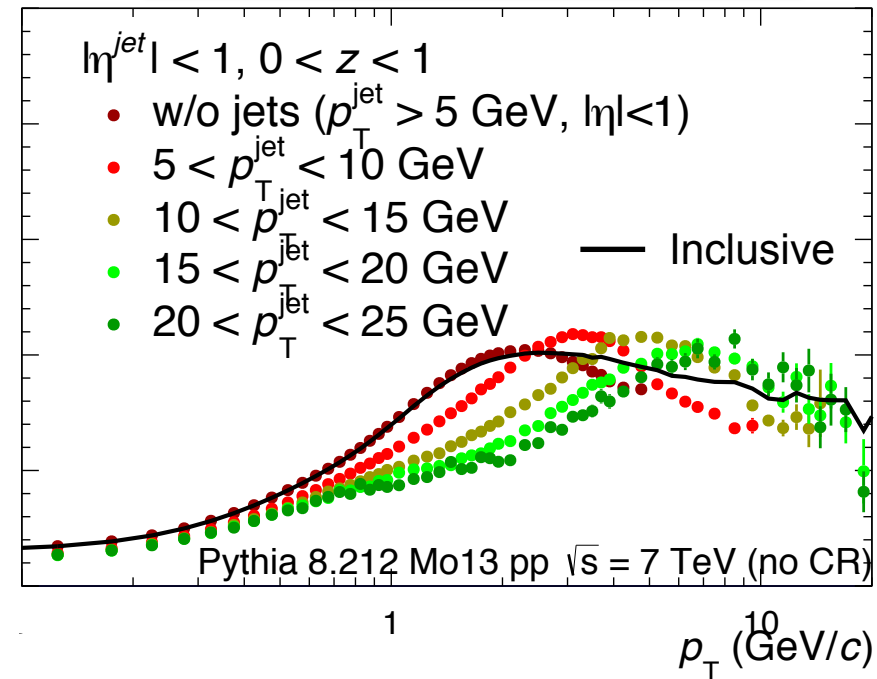
The higher the event multiplicity the higher the average  $p_T^{\text{jet}}$

# $p/\pi$ vs. $p_T$ (low multiplicity)

With Color Reconnection

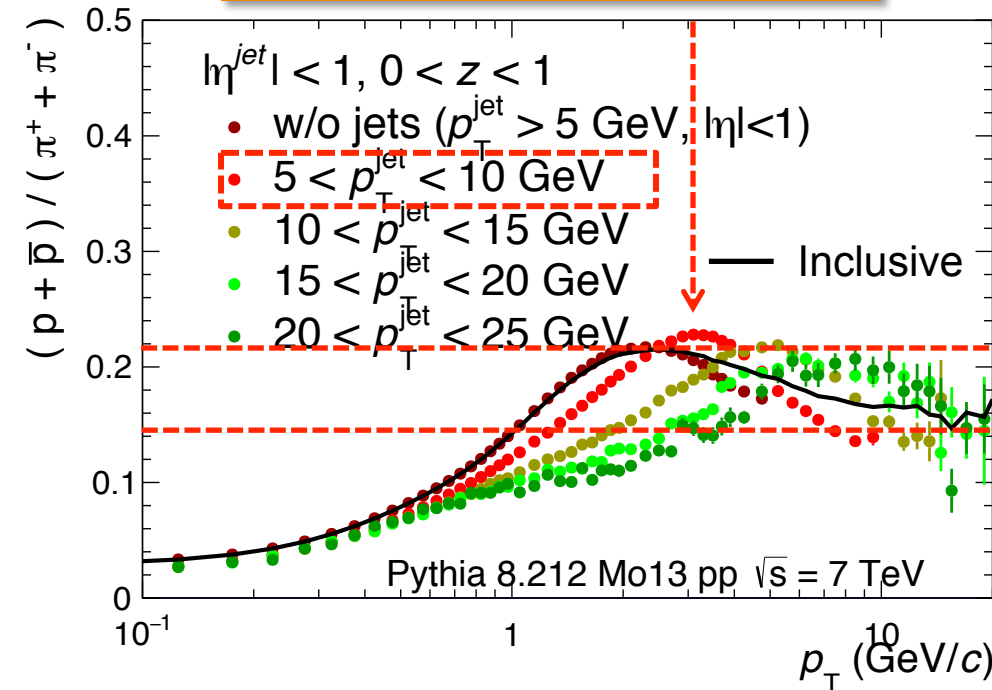


Without Color Reconnection

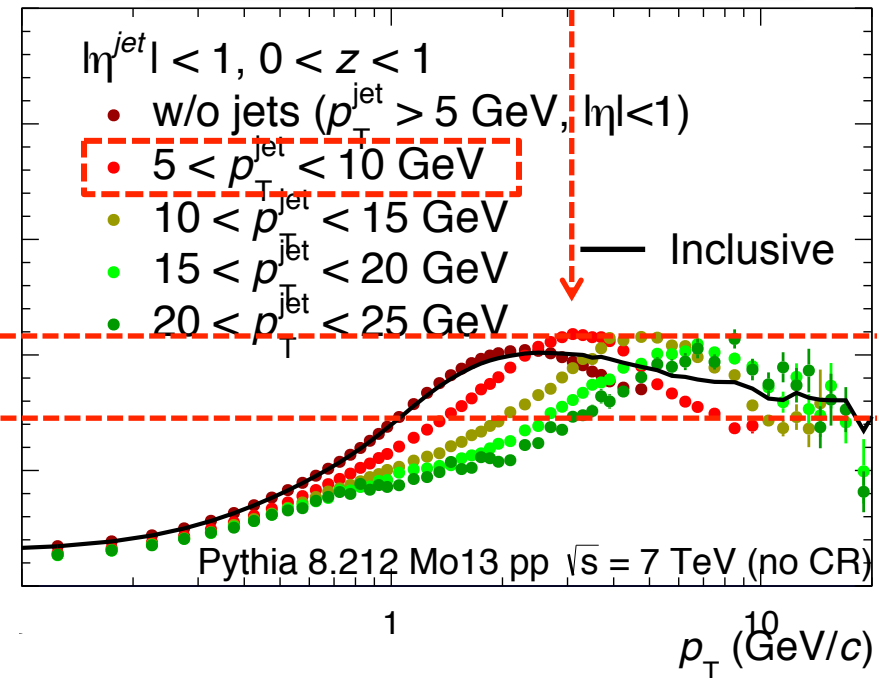


# $p/\pi$ vs. $p_T$ (low multiplicity)

With Color Reconnection



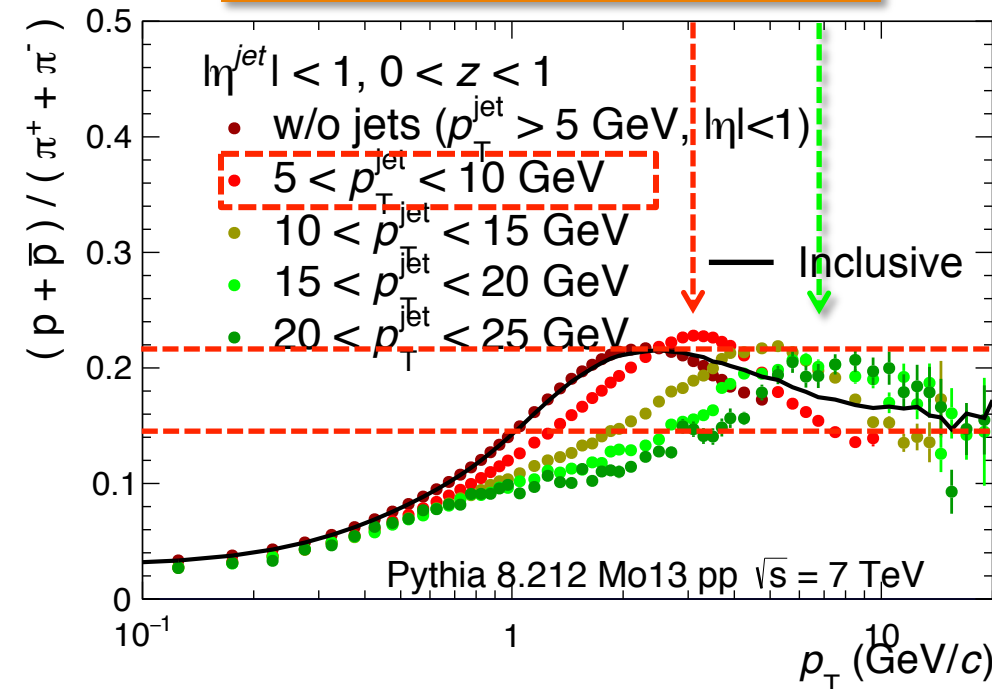
Without Color Reconnection



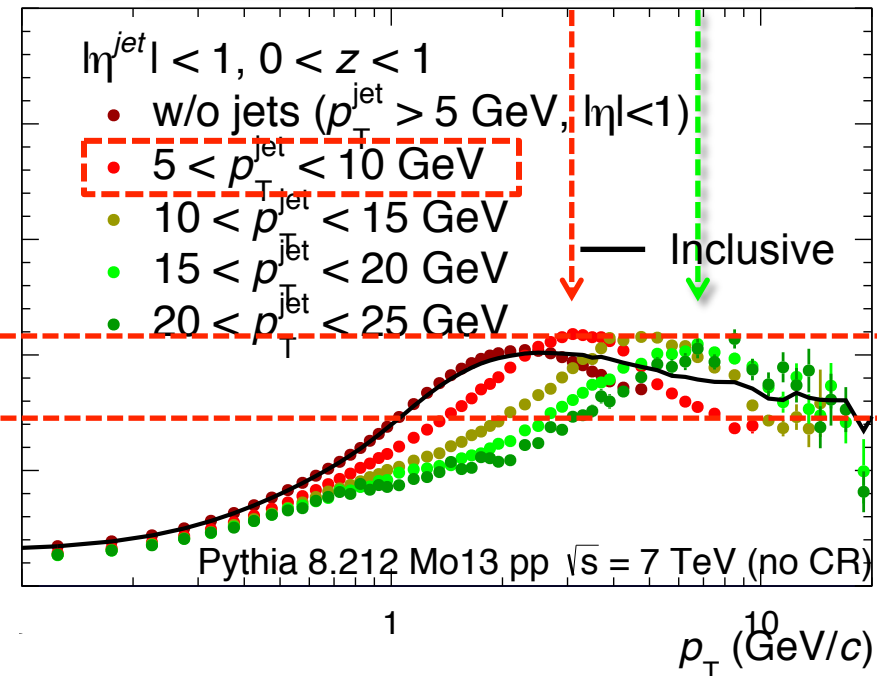
□ CR effects are observer for  $p_T^{jet} < 10 \text{ GeV}$

# $p/\pi$ vs. $p_T$ (low multiplicity)

With Color Reconnection



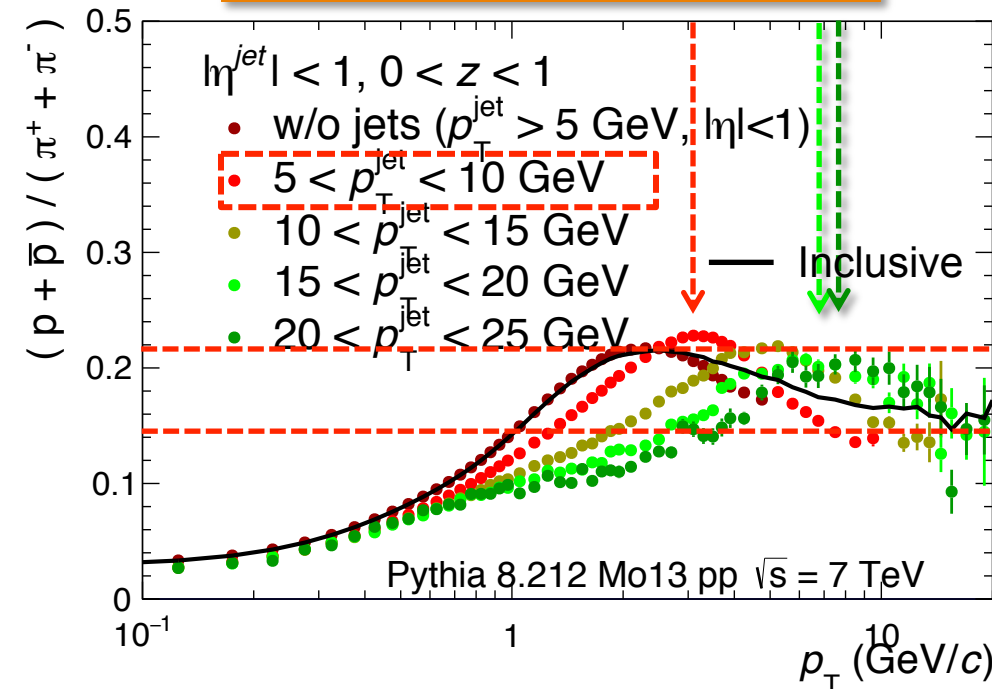
Without Color Reconnection



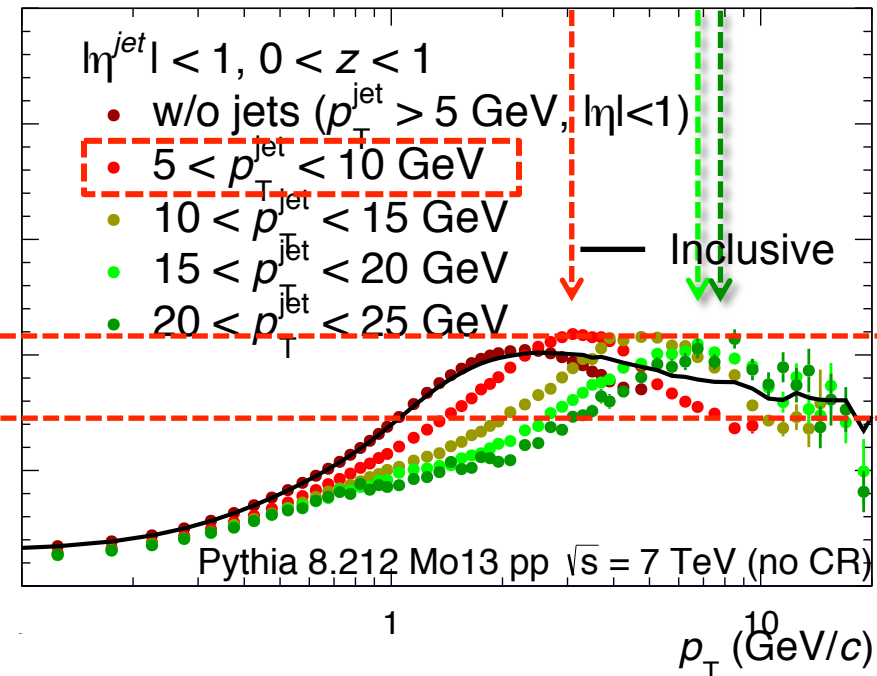
- ❑ CR effects are observer for  $p_T^{\text{jet}} < 10 \text{ GeV}$
- ❑ The position of the peak is shifted to higher  $p_T$  when  $p_T^{\text{jet}}$  increases. The shift is accompanied by an increase of  $\langle \beta_T \rangle$

# $p/\pi$ vs. $p_T$ (low multiplicity)

With Color Reconnection

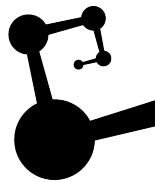


Without Color Reconnection



- ❑ CR effects are observer for  $p_T^{jet} < 10 \text{ GeV}$
- ❑ The position of the peak is shifted to higher  $p_T$  when  $p_T^{jet}$  increases. The shift is accompanied by an increase of  $\langle \beta_T \rangle$  (from Blast-Wave analysis)
- ❑ The effect is very small for  $p_T^{jet} > 15 \text{ GeV}$

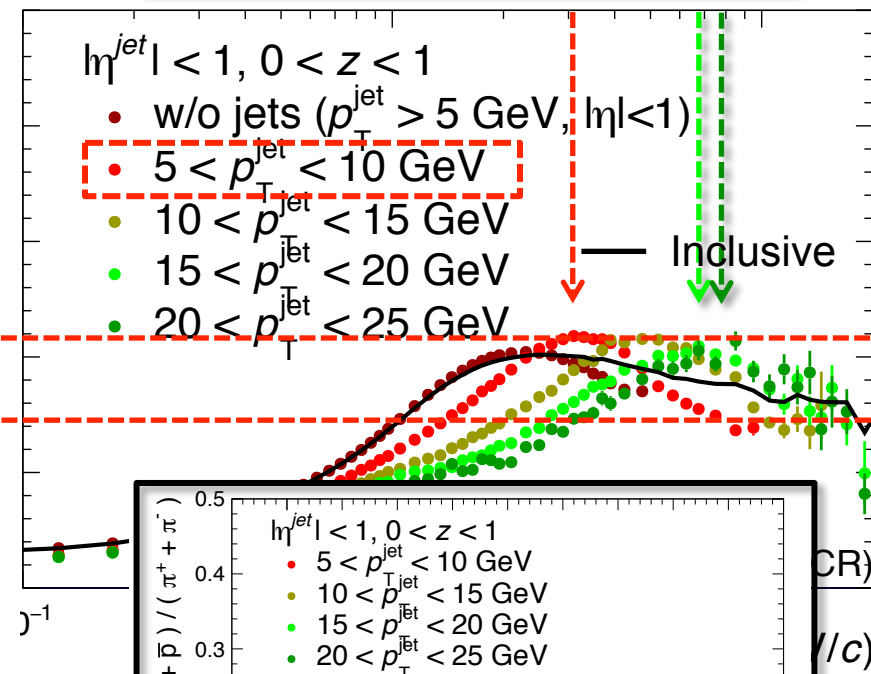
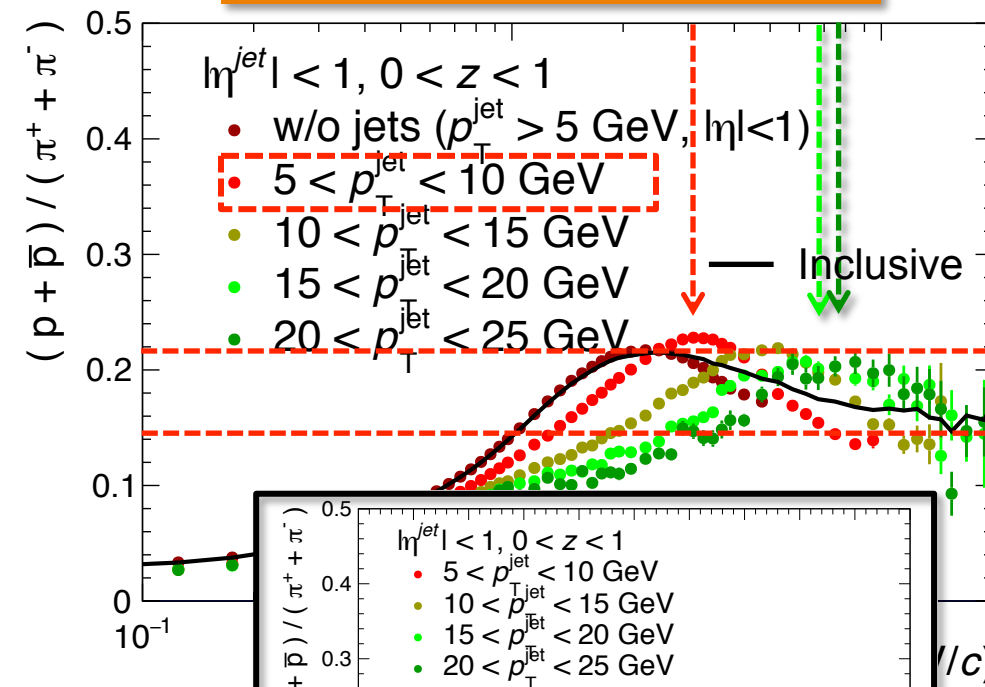
# $p/\pi$ vs. $p_T$ (low multiplicity)



This is a FF effect ( $p/\pi$  vs.  $p_T/p_T^{\text{jet}}$  is  $\approx p_T^{\text{jet}}$  independent)

With Color Reconnection

Without Color Reconnection



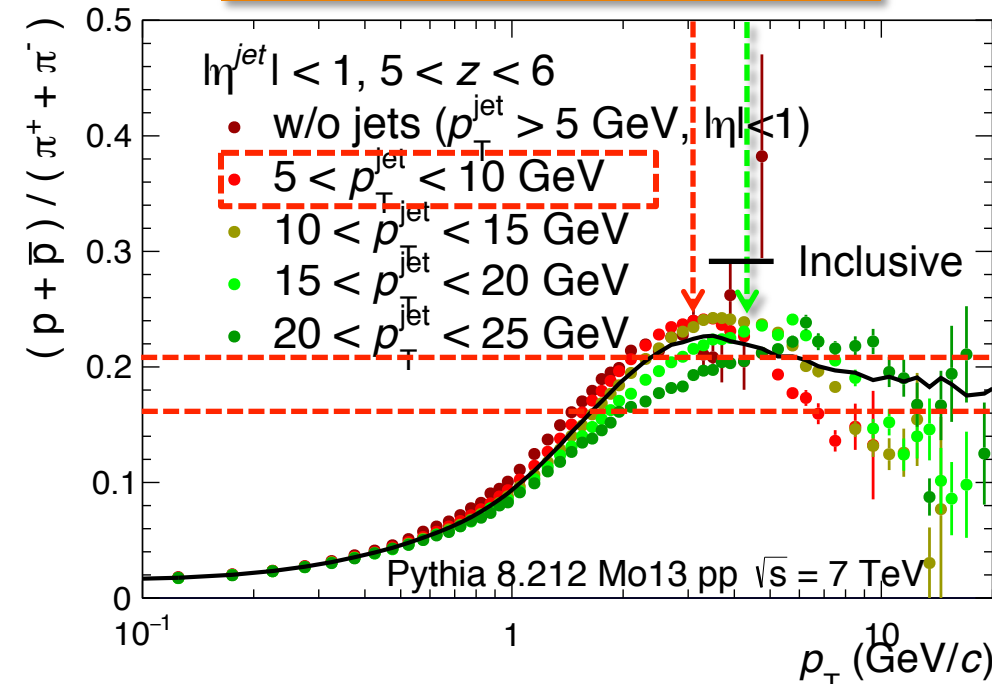
- CR e
- The shift

10 GeV  
to higher  
values of  $\beta_T$   
jet  $> 15 \text{ GeV}$

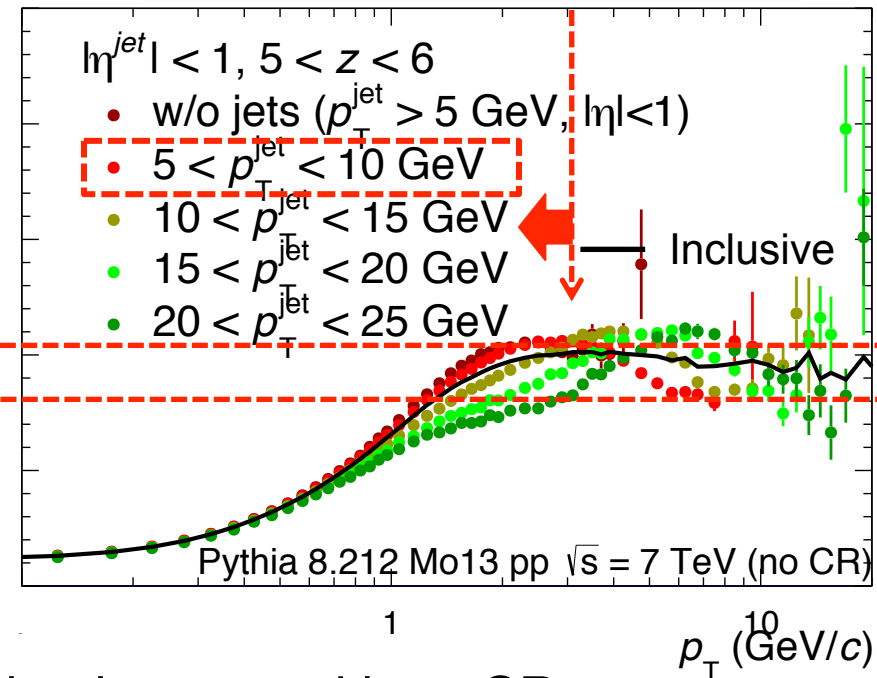


# $\rho/\pi$ vs. $p_T$ (high multiplicity)

With Color Reconnection



Without Color Reconnection



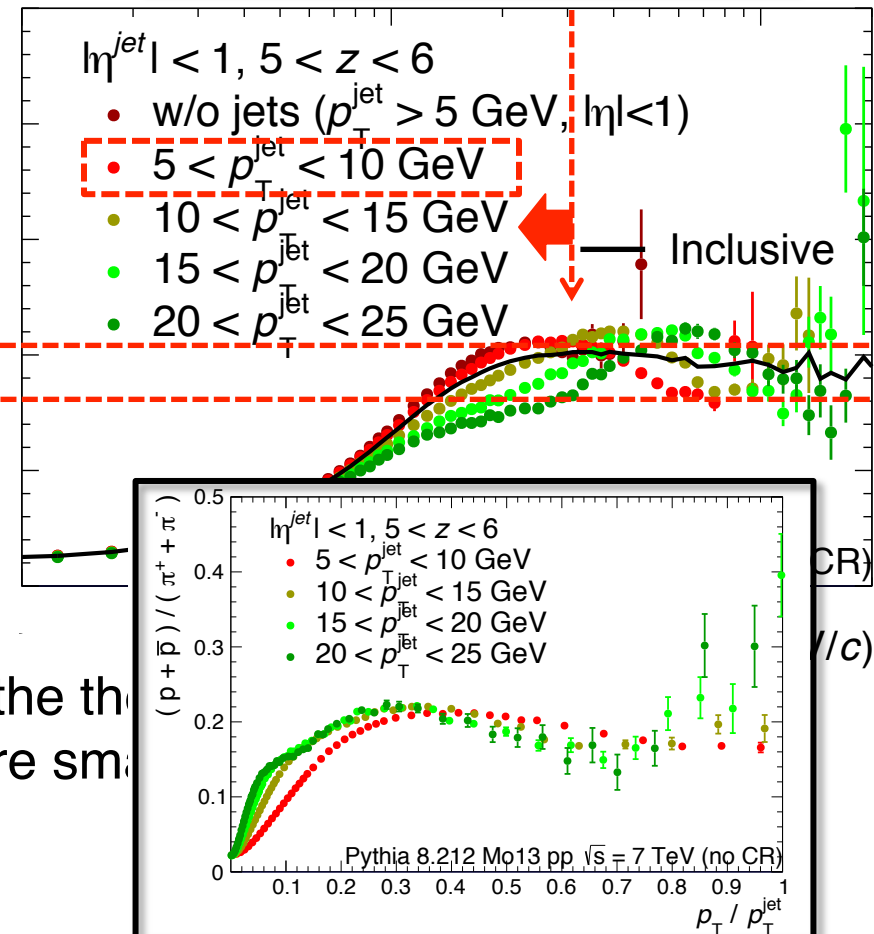
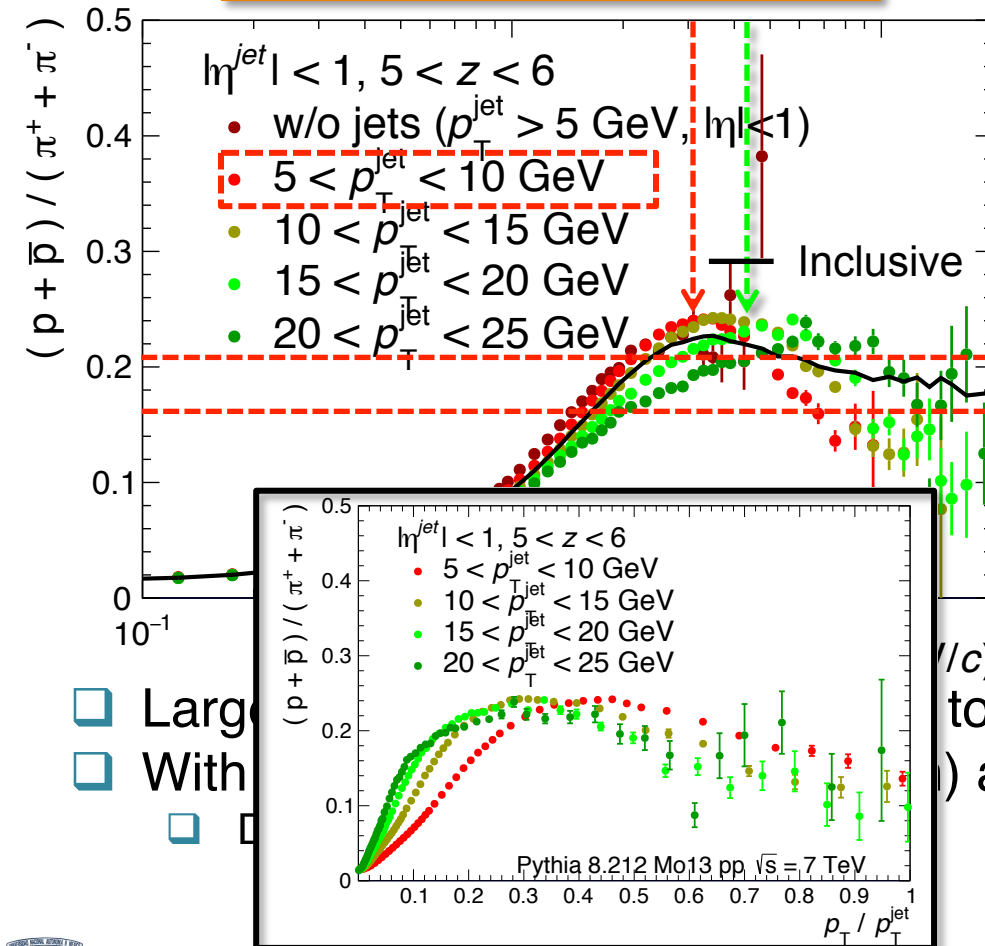
- Larger enhancement with respect to the the case without CR
- With CR, jet effects (peak position) are smaller than in the low  $N_{ch}$  case
  - Dominated by underlying event

# $p/\pi$ vs. $p_T$ (high multiplicity)

Without CR:  $p/\pi$  vs.  $p_T/p_T^{\text{jet}}$  is  $\approx p_T^{\text{jet}}$  independent (FF)

With Color Reconnection

Without Color Reconnection



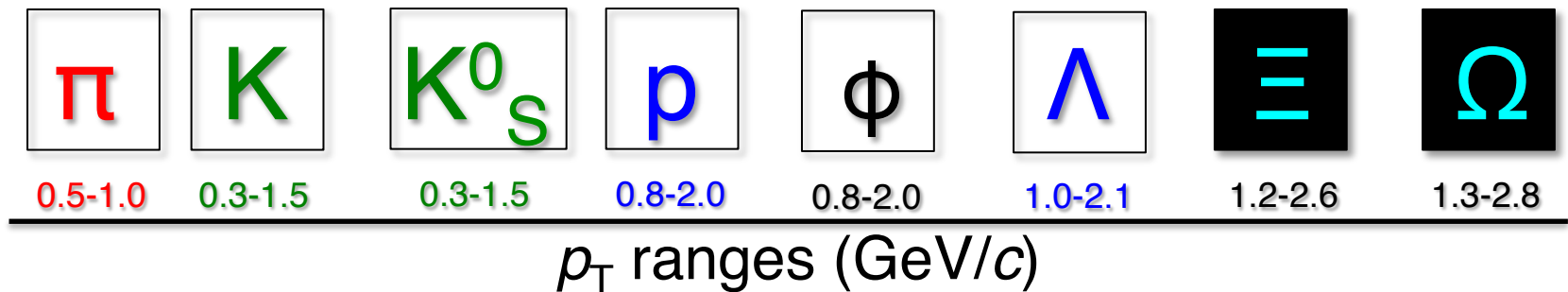
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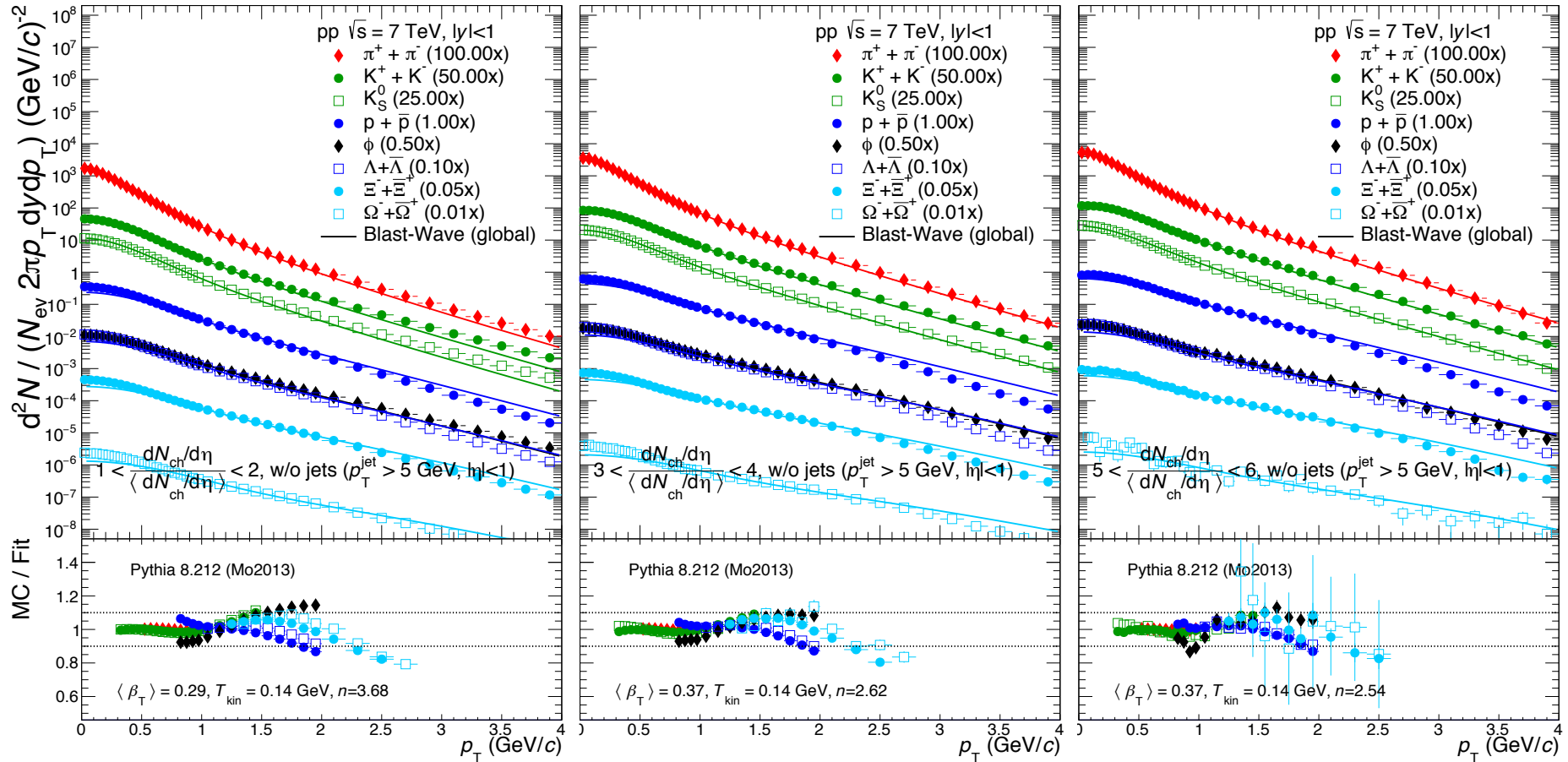
# Study of the inclusive light flavored hadron production

Results from the Blast-Wave analysis are presented, for this a simultaneous fit of the BW function to the the  $p_T$  spectra is performed in order to extract  $\langle\beta_T\rangle$ . The fitting ranges are the following:



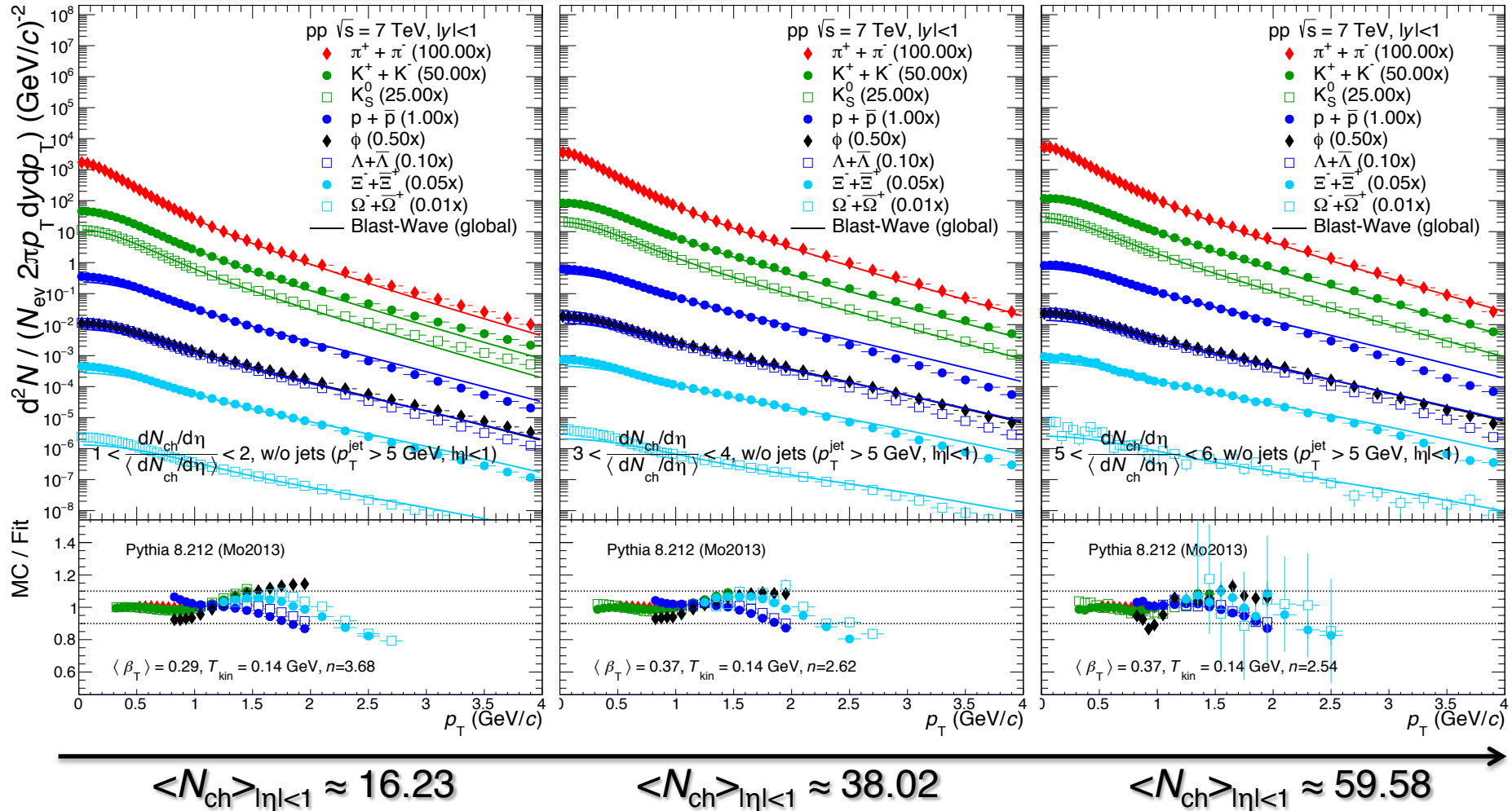
(Same  $p_T$  ranges as in: [G. Paic, E. Cuautle and A. O. NPA 941 \(2015\) 78-86](#), where the  $p_T$  spectra in high multiplicity events were described by BW model within 10%)

# Without Jets



# Without Jets

$$\langle \beta_T \rangle \approx 0.34, \langle T_{\text{kin}} \rangle \approx 0.14, \langle n \rangle \approx 2.94$$



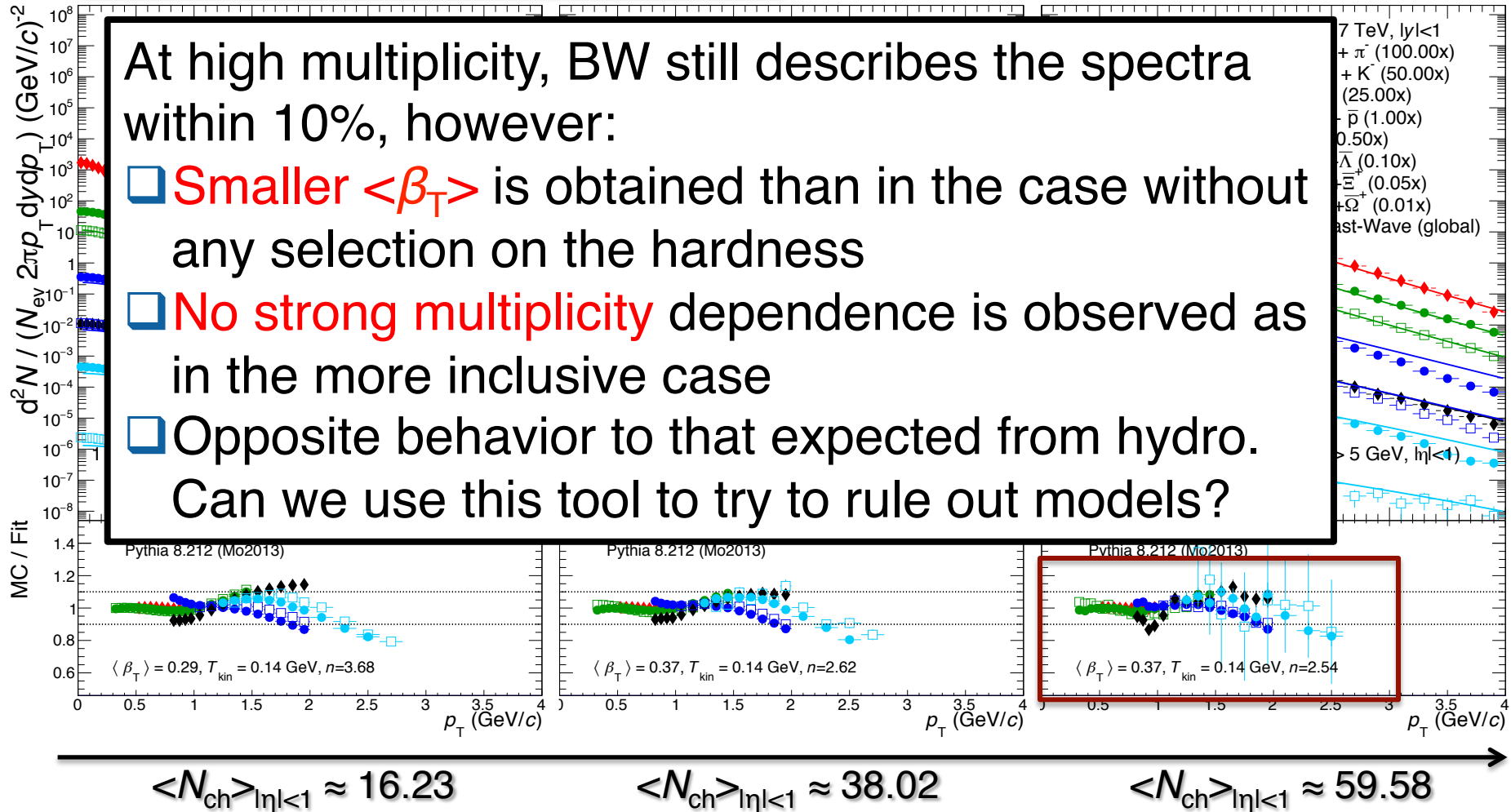


# Without Jets

$$\langle \beta_T \rangle \approx 0.34, \quad \langle T_{\text{kin}} \rangle \approx 0.14, \quad \langle n \rangle \approx 2.94$$

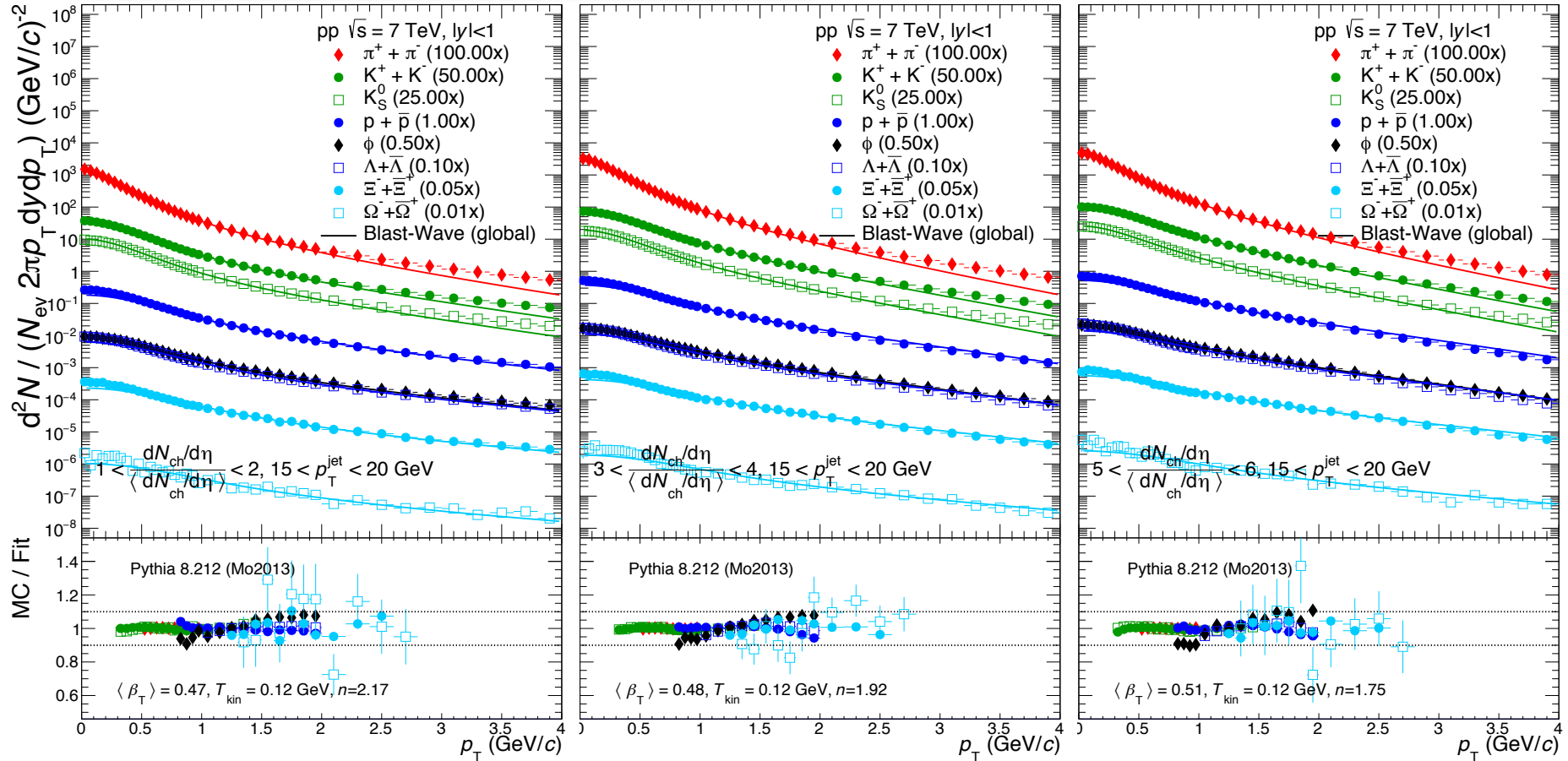
At high multiplicity, BW still describes the spectra within 10%, however:

- ❑ **Smaller  $\langle \beta_T \rangle$**  is obtained than in the case without any selection on the hardness
- ❑ **No strong multiplicity** dependence is observed as in the more inclusive case
- ❑ Opposite behavior to that expected from hydro.  
Can we use this tool to try to rule out models?



# $15 < p_T^{\text{Jet}} < 20 \text{ GeV}$

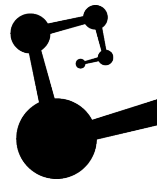
$$\langle \beta_T \rangle \approx 0.48, \langle T_{\text{kin}} \rangle \approx 0.12, \langle n \rangle \approx 1.94$$



$$\langle N_{\text{ch}} \rangle_{|y| < 1} \approx 16.23$$

$$\langle N_{\text{ch}} \rangle_{|y| < 1} \approx 38.02$$

$$\langle N_{\text{ch}} \rangle_{|y| < 1} \approx 59.58$$

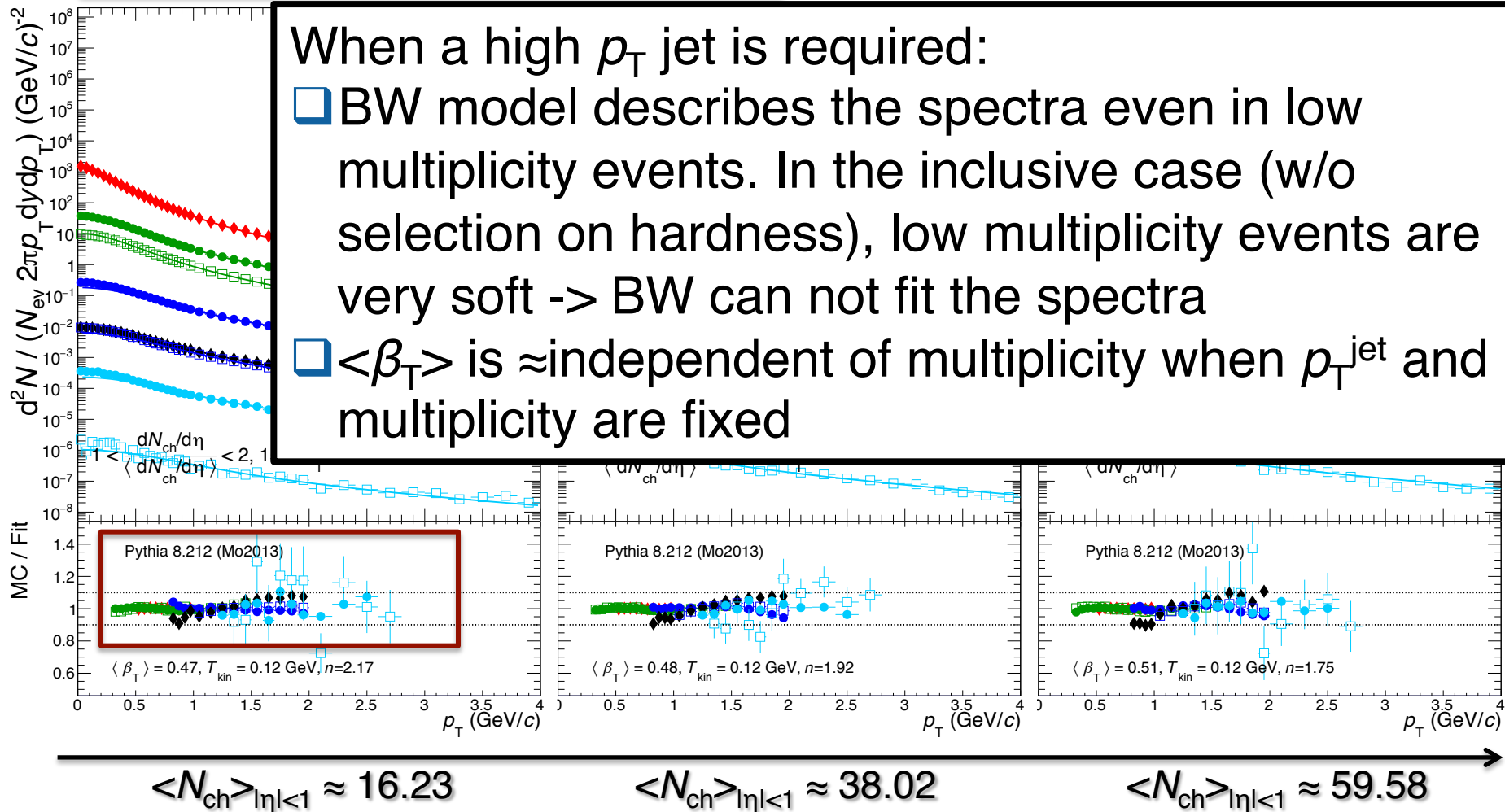


$$15 < p_T^{\text{Jet}} < 20 \text{ GeV}$$

$$\langle \beta_T \rangle \approx 0.48, \langle T_{\text{kin}} \rangle \approx 0.12, \langle n \rangle \approx 1.94$$

When a high  $p_T$  jet is required:

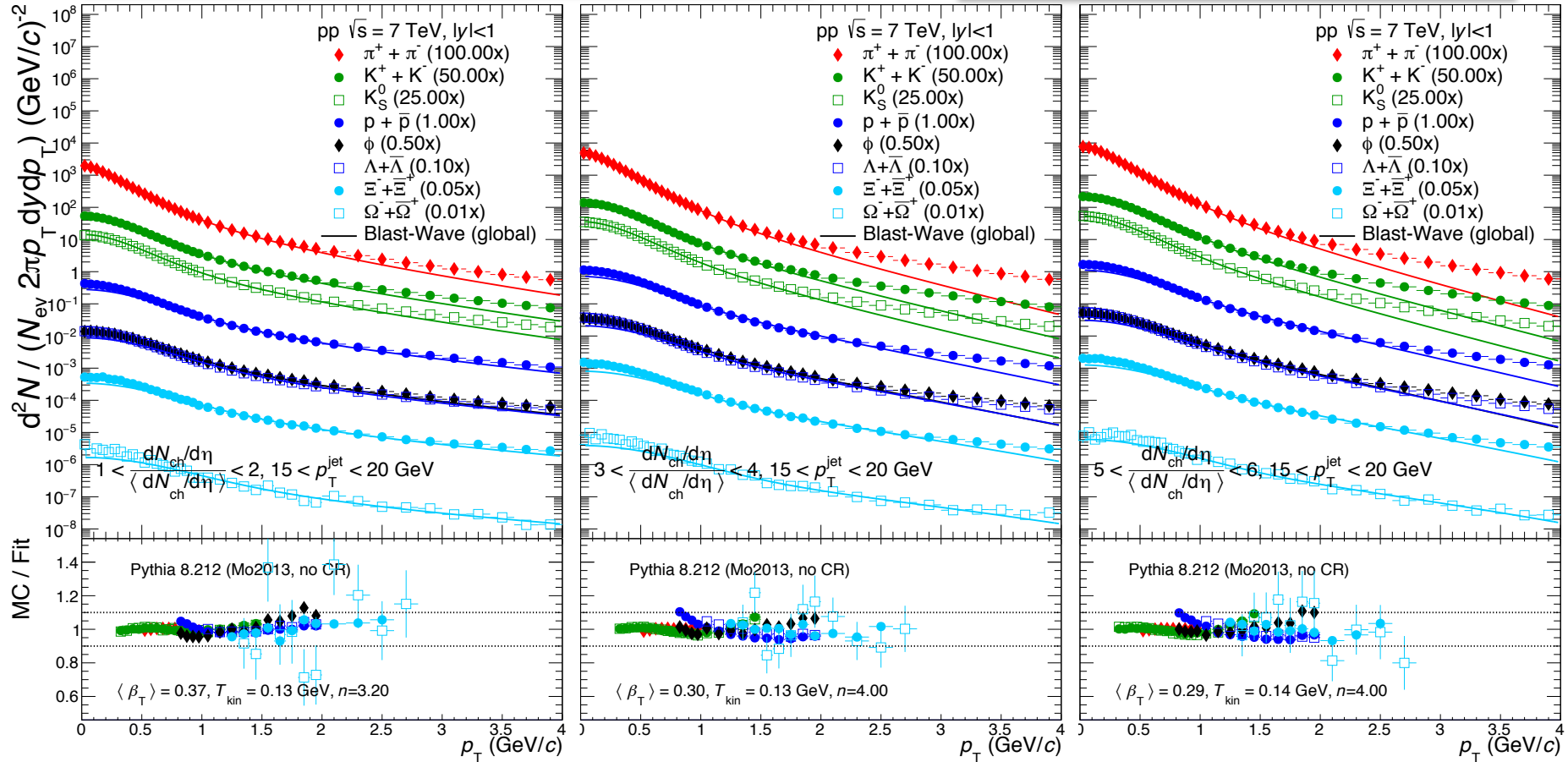
- ❑ BW model describes the spectra even in low multiplicity events. In the inclusive case (w/o selection on hardness), low multiplicity events are very soft  $\rightarrow$  BW can not fit the spectra
- ❑  $\langle \beta_T \rangle$  is  $\approx$  independent of multiplicity when  $p_T^{\text{jet}}$  and multiplicity are fixed



$$15 < p_T^{\text{Jet}} < 20 \text{ GeV}$$

Smaller  $\langle \beta_T \rangle \approx 0.32$  & larger  $\langle n \rangle \approx 3.7$

Without Color Reconnection



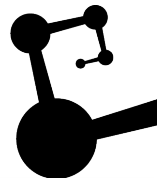
Slight increase of  $\langle \beta_T \rangle$

# ENERGY DEPENDENCE

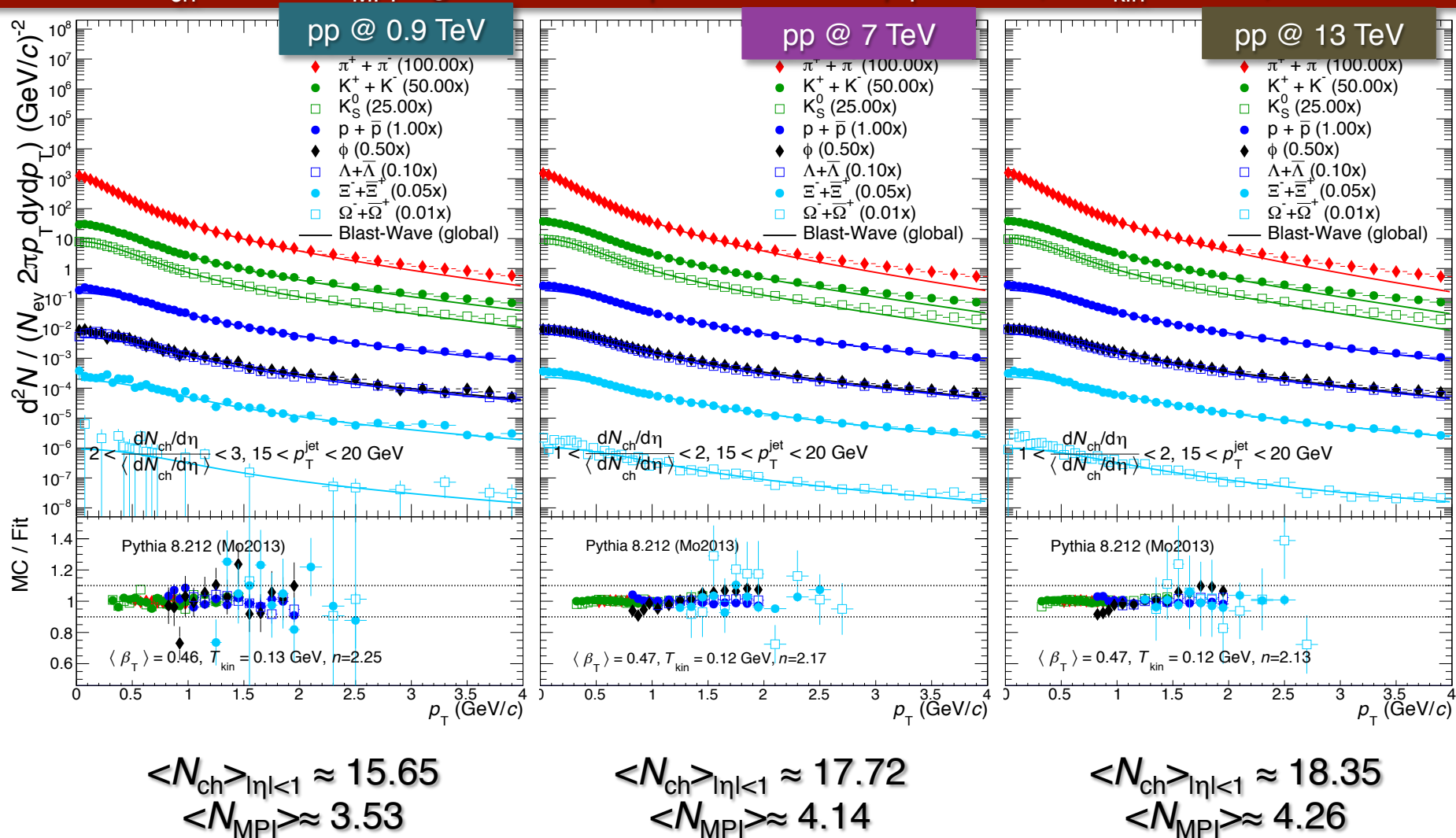




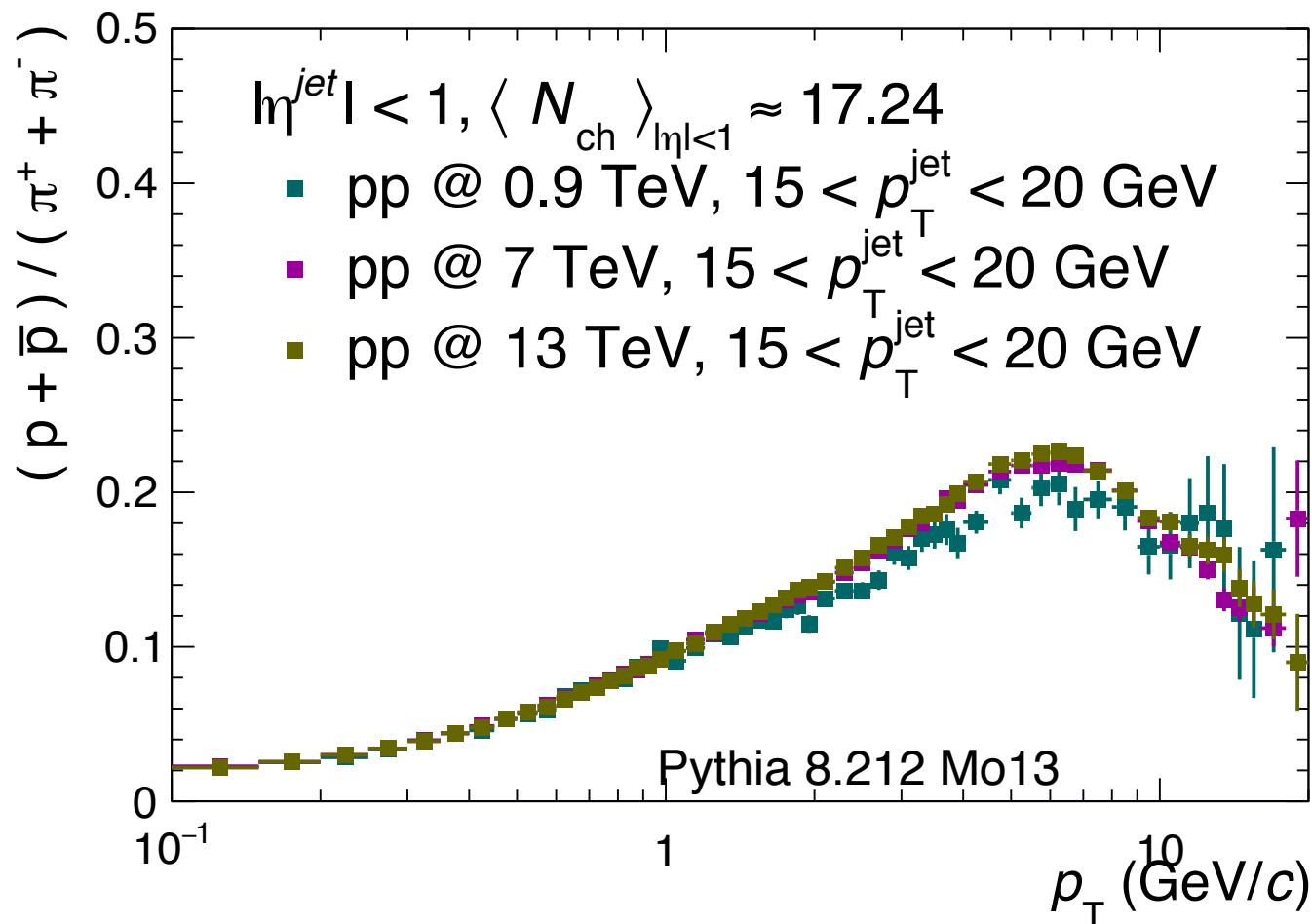
# $15 < p_T^{\text{Jet}} < 20 \text{ GeV}$



Similar  $\langle N_{\text{ch}} \rangle$  and  $\langle N_{\text{MPI}} \rangle$  gives similar parameters:  $\langle \beta_T \rangle \approx 0.47$ ,  $\langle T_{\text{kin}} \rangle \approx 0.12$ ,  $\langle n \rangle \approx 2.18$

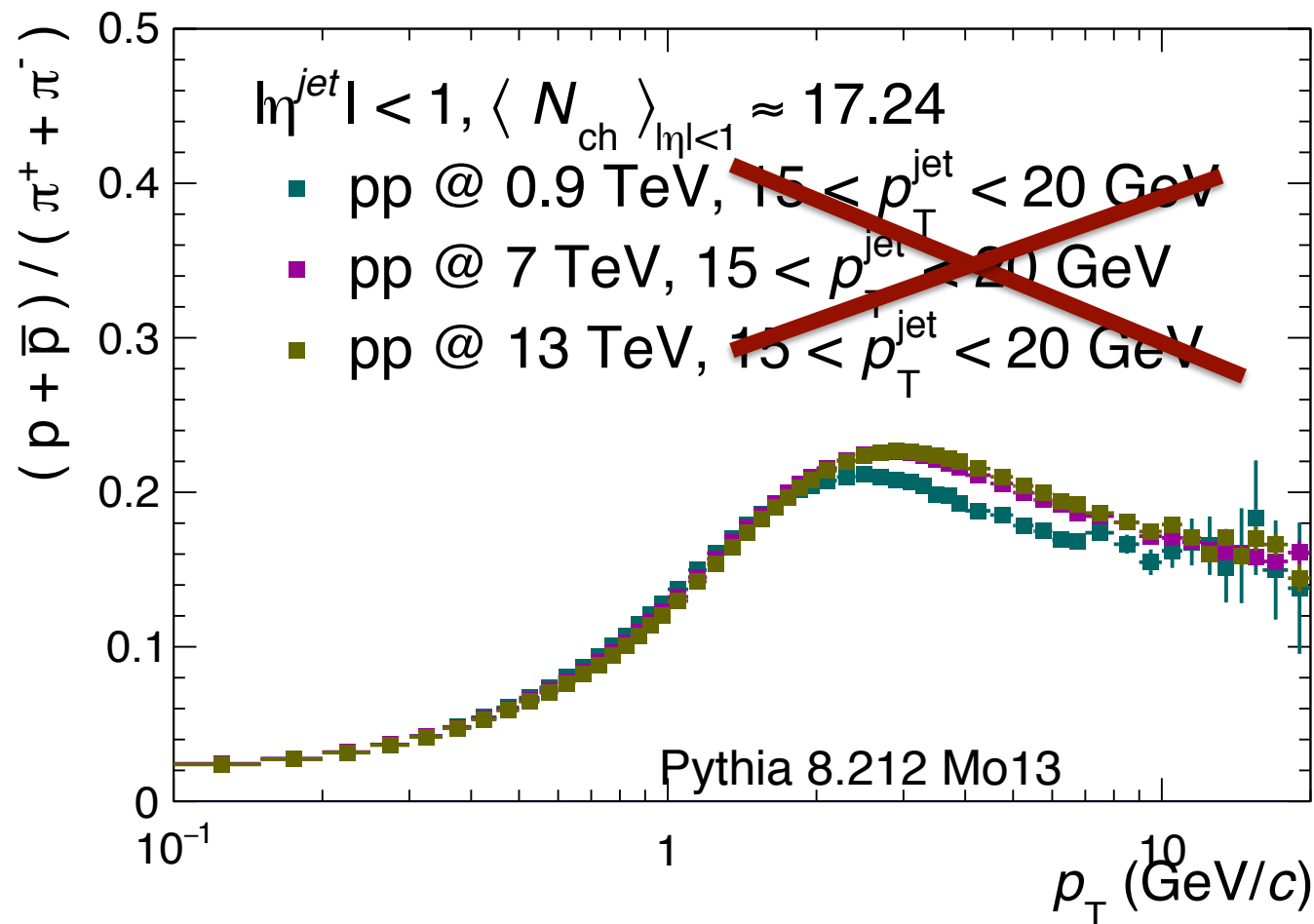


# Spectral shapes



Proton-to-pion ratio show little or no dependence with  $\sqrt{s}$   
 ( $p_T$  position of the peak is the same for the three colliding systems)

# Spectral shapes



Without the jet requirement, the ratios look more different due to the different jet biases

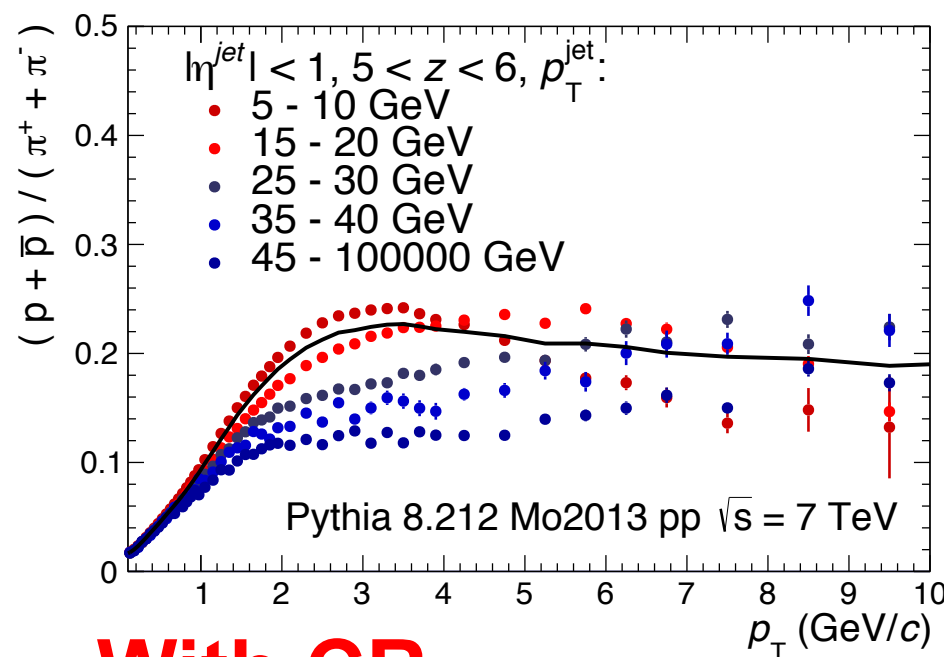
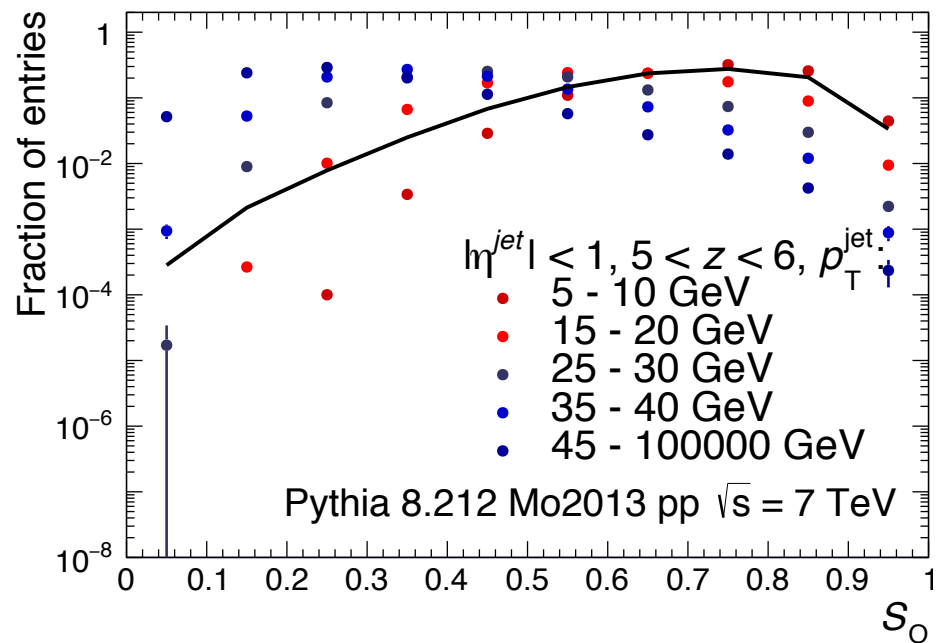
# Summary

- ❑ In Pythia, MPI (semi-hard and **hard partonic scatterings**) and CR produce flow-like effects
- ❑ The result of the interaction between the soft and hard component could be used as a tool to validate or rule out models which produce flow(-like) effects in small systems, e.g. hydro vs. color reconnection (important for HI physics)
- ❑ Same physics is obtained when a selection on multiplicity and hardness is implemented

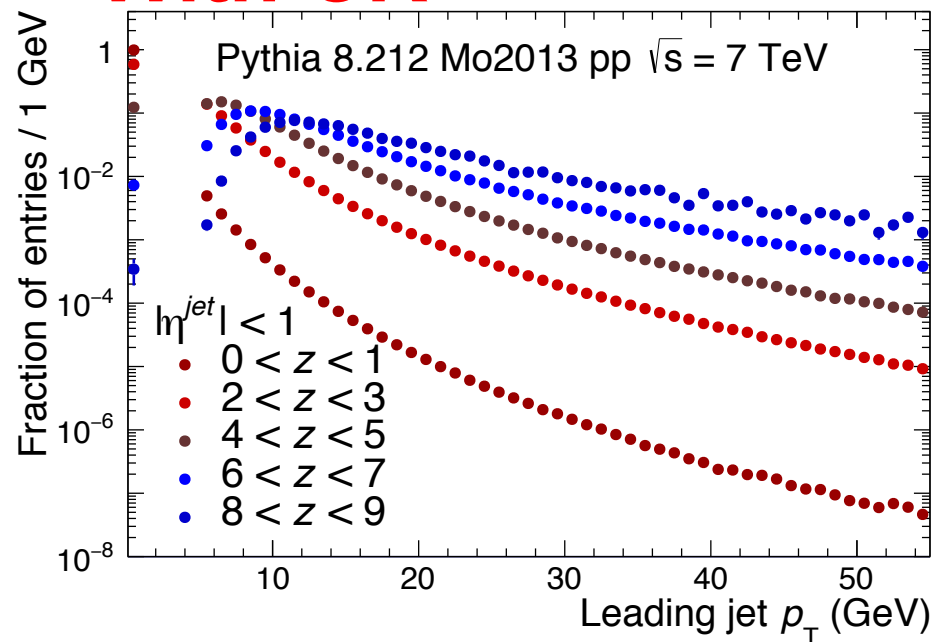
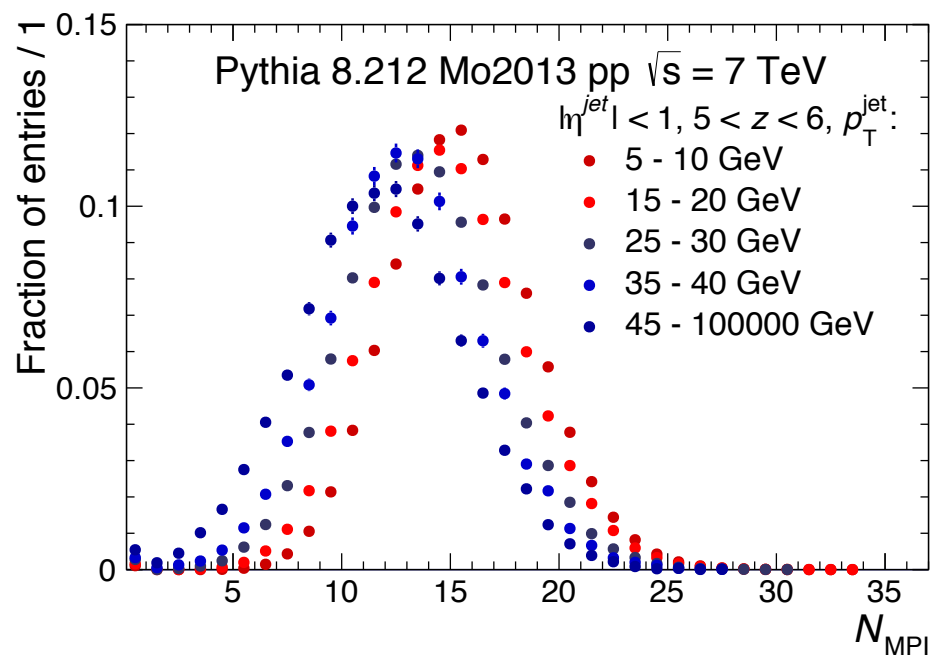
Guy Paić, Peter Christiansen, Andreas Morsch and Eleazar Cuautle are acknowledged for the useful discussions

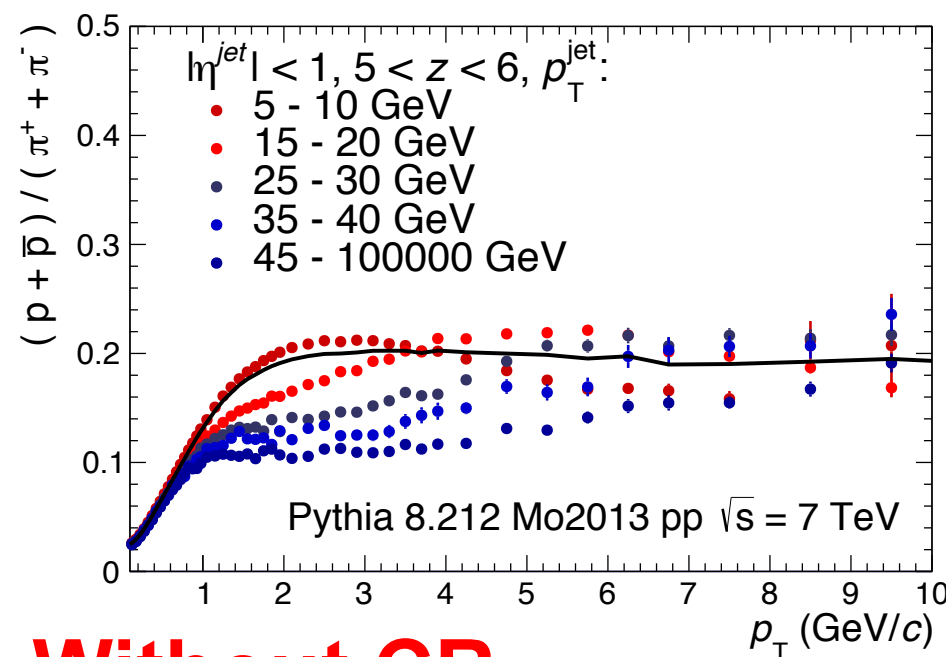
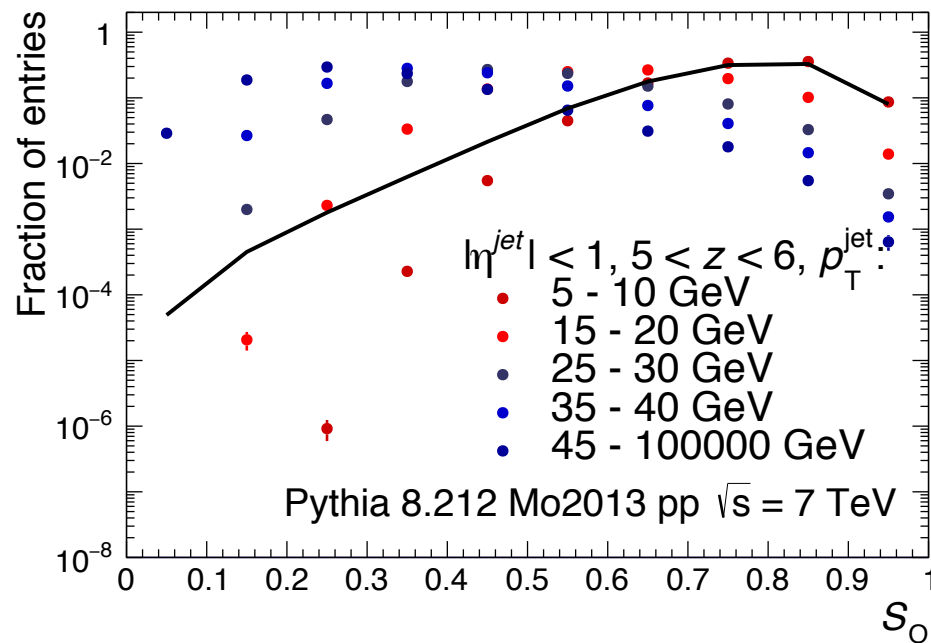
# BACKUP



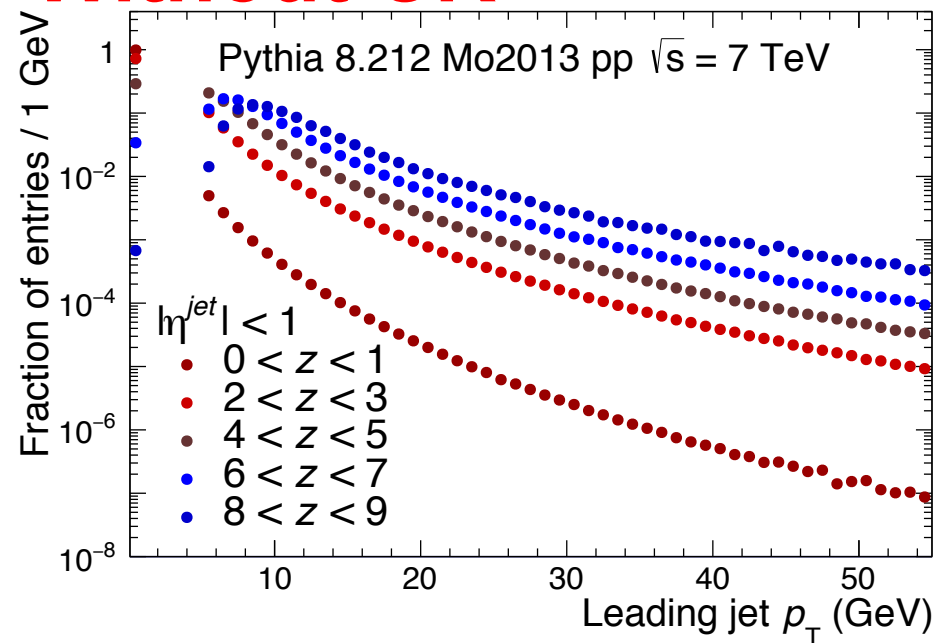
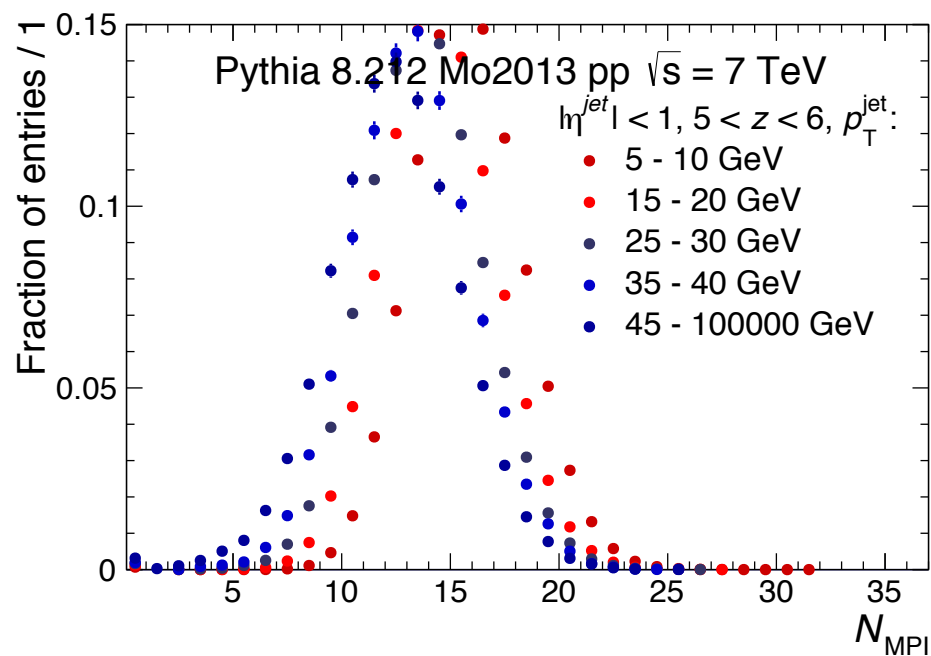


**With CR**





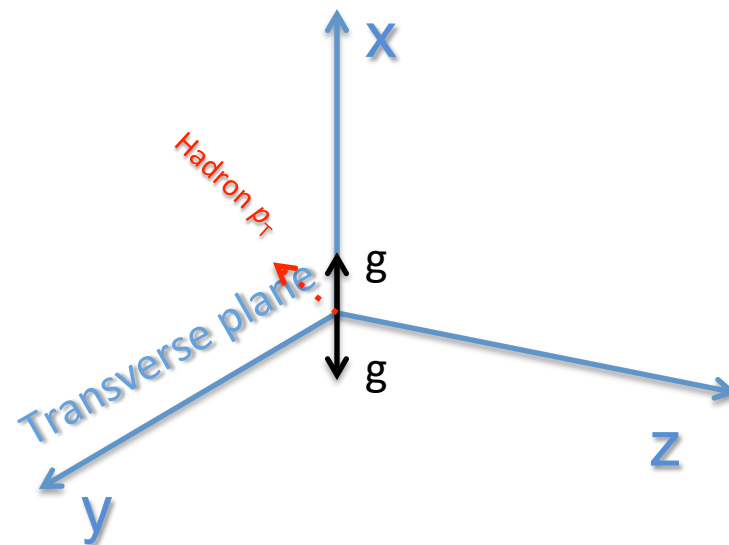
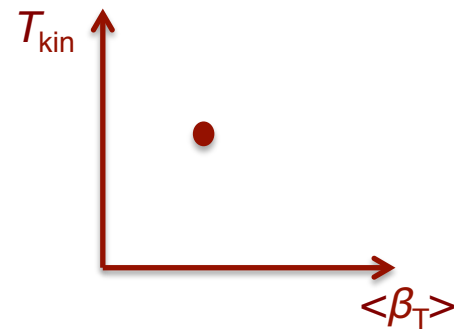
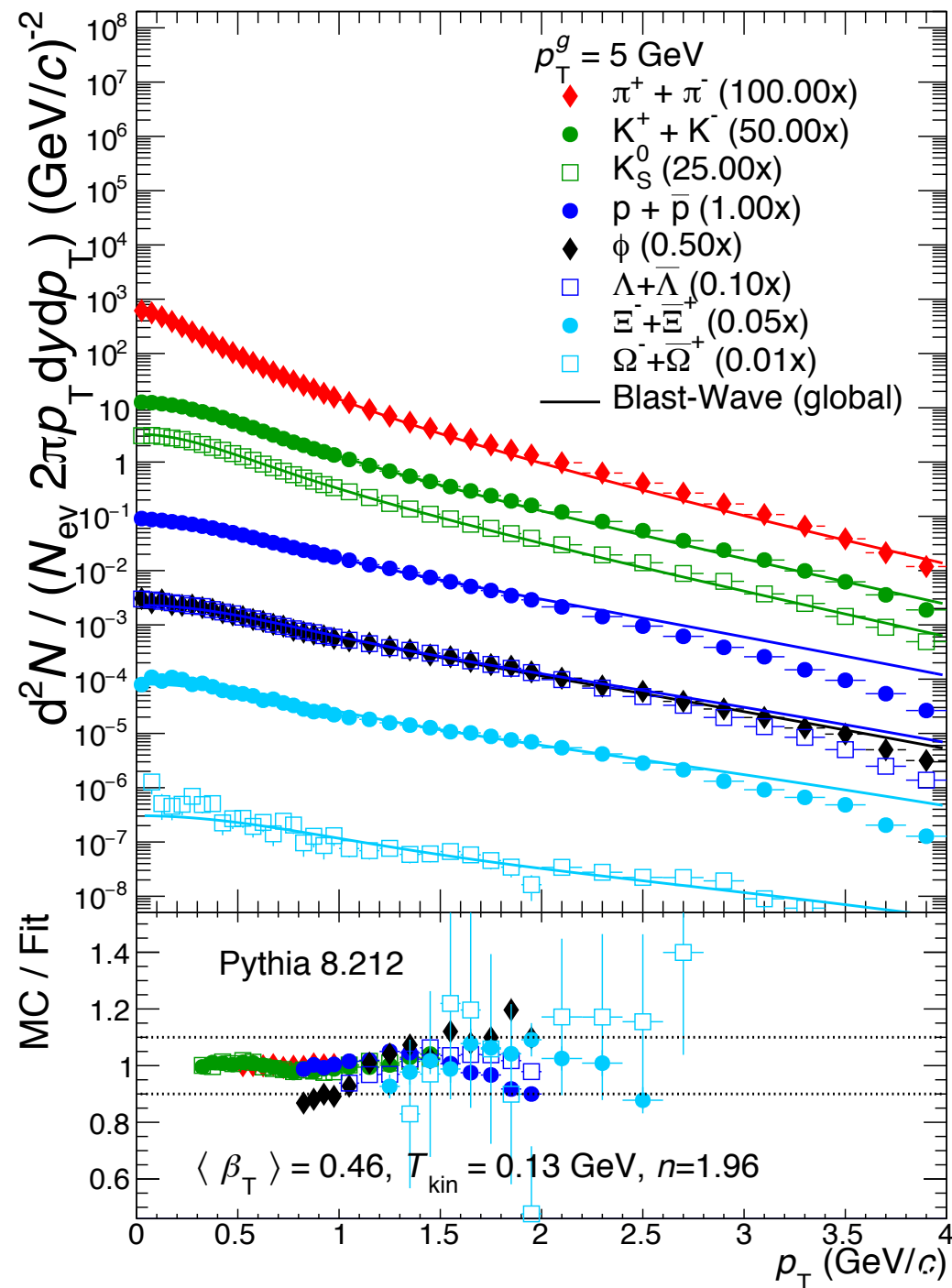
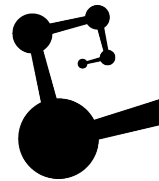
**Without CR**



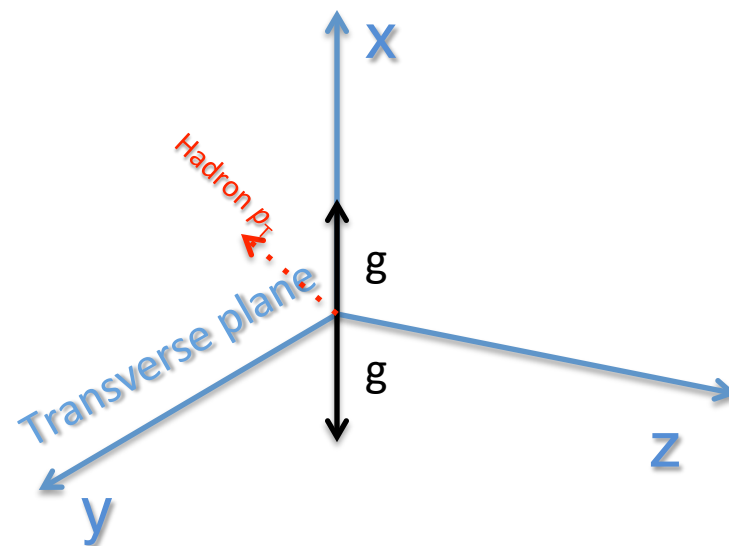
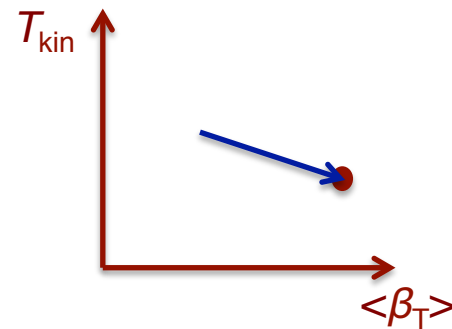
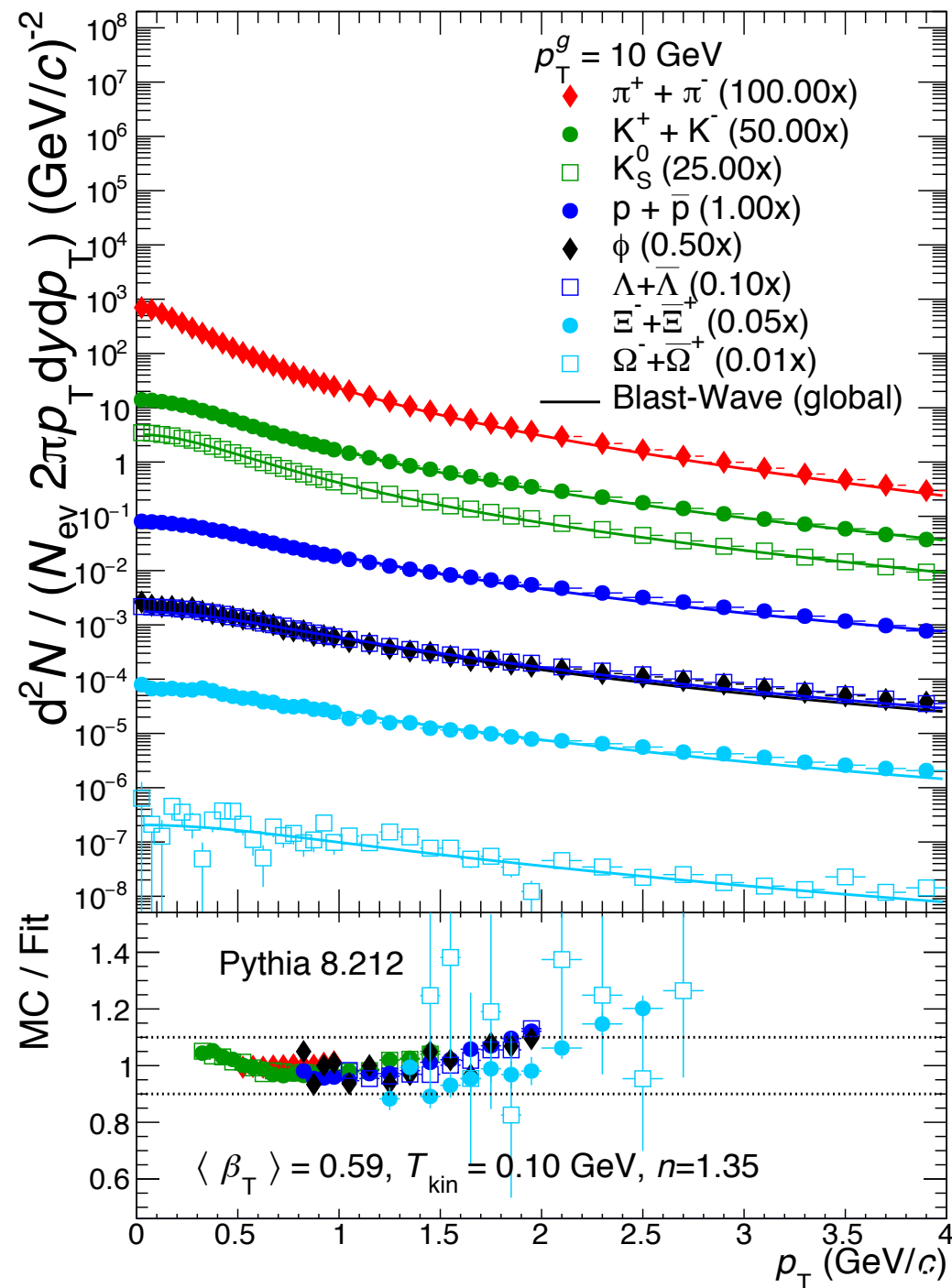
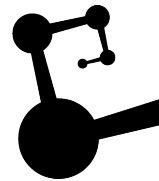


# HADRONIZATION IN A CLEAN PARTONIC CONFIGURATION

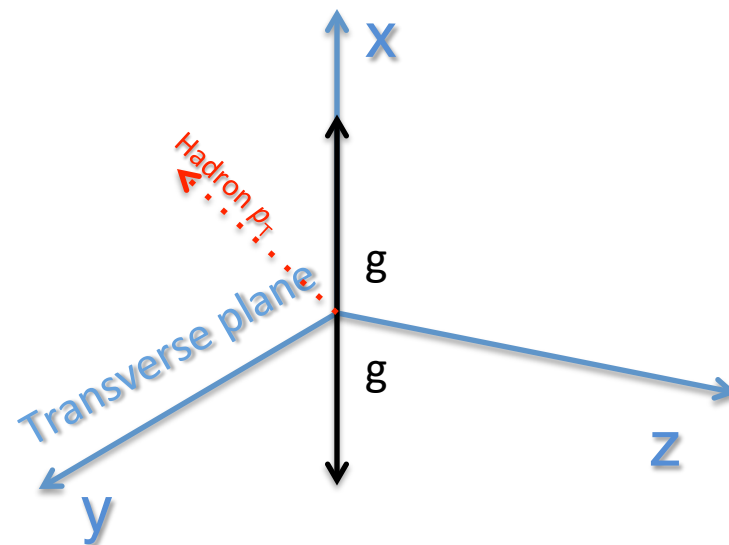
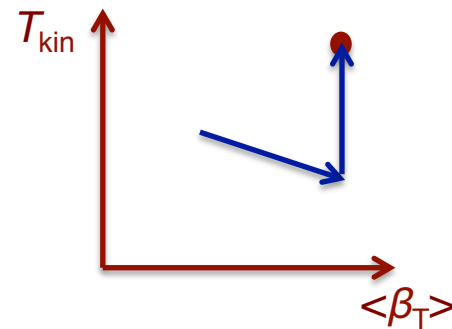
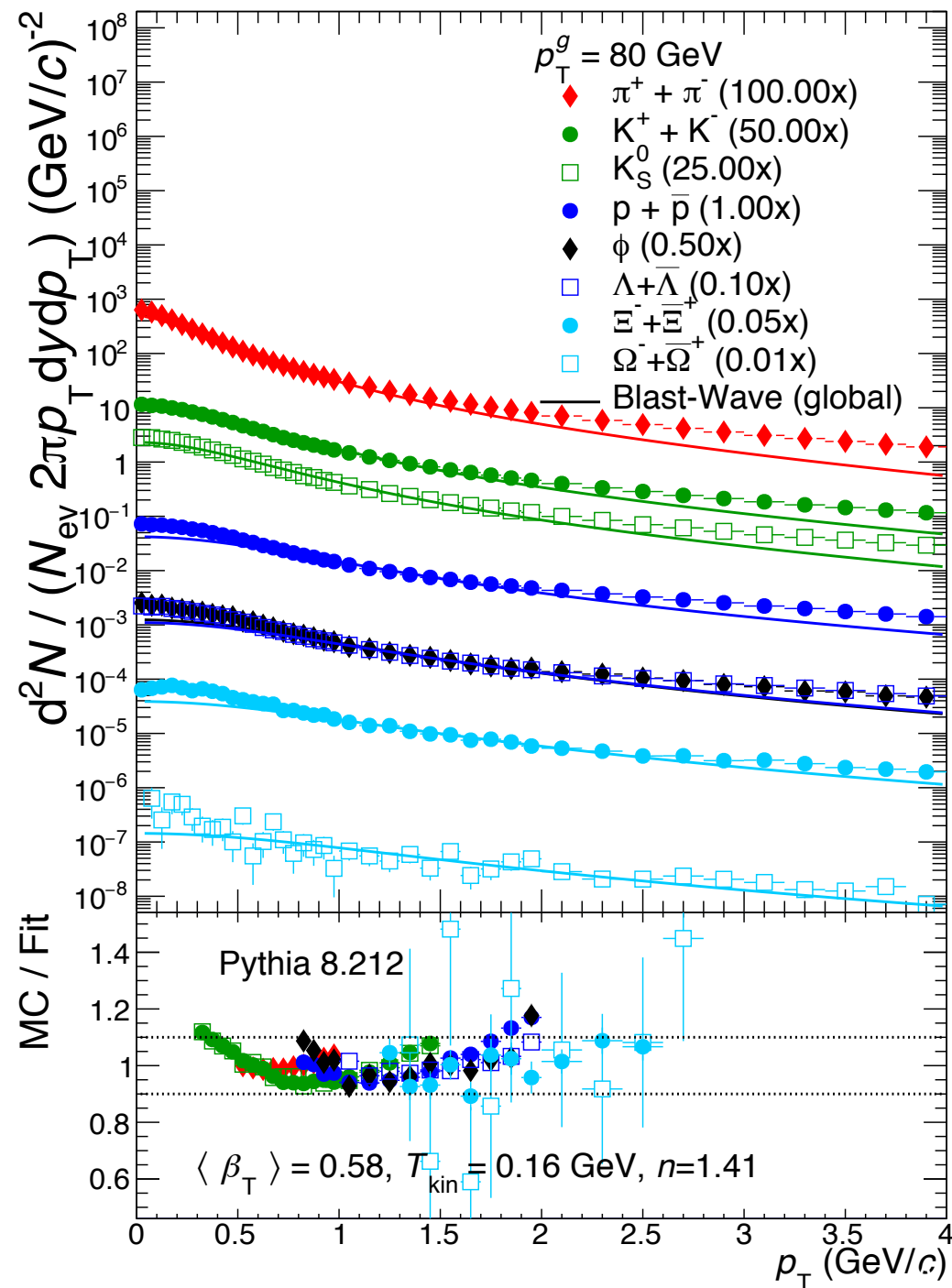
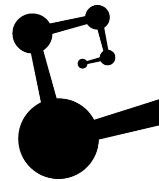




Parton-level configurations as  
direct input for hadronization



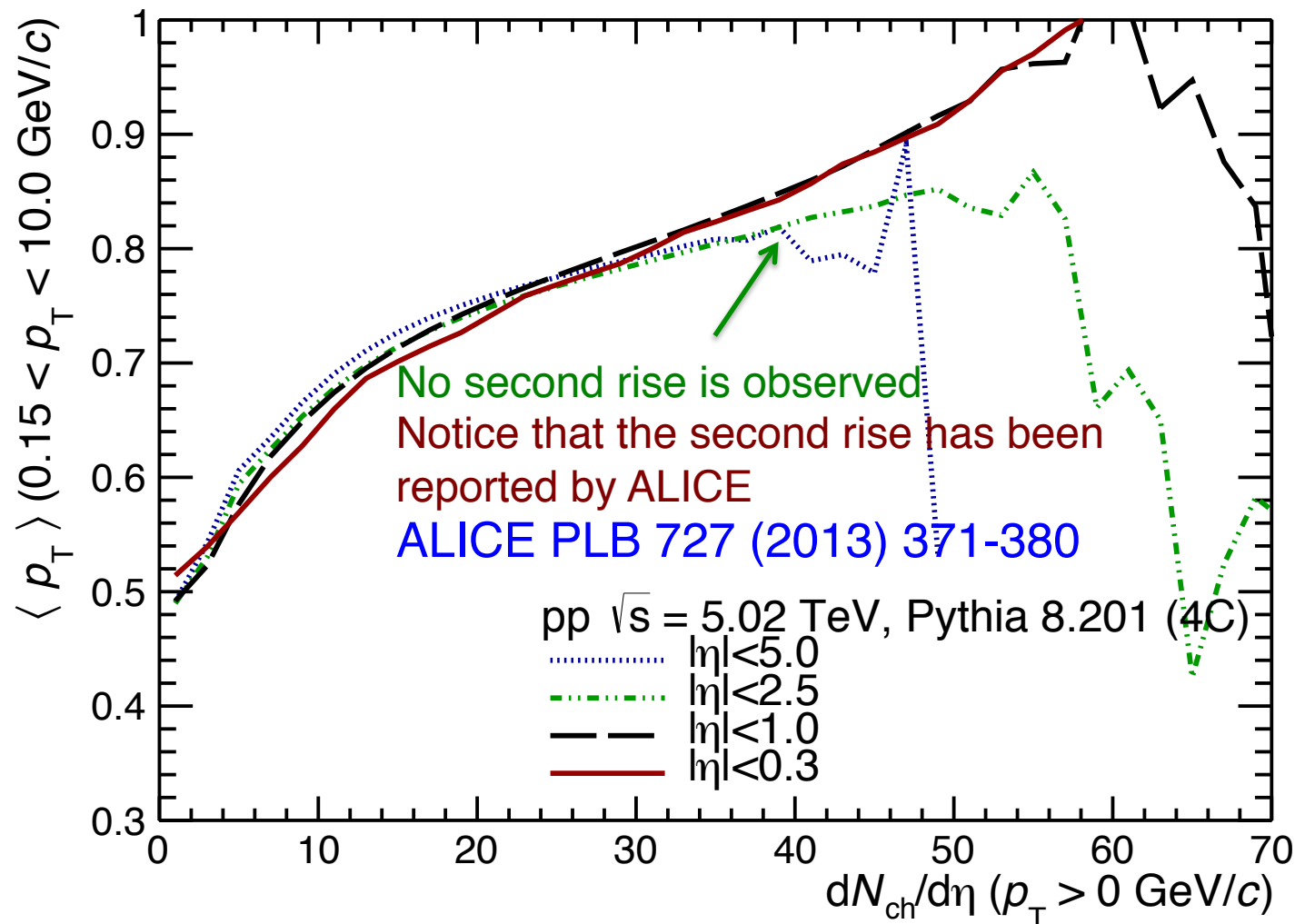
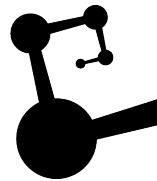
Parton-level configurations as  
direct input for hadronization



Parton-level configurations as  
direct input for hadronization

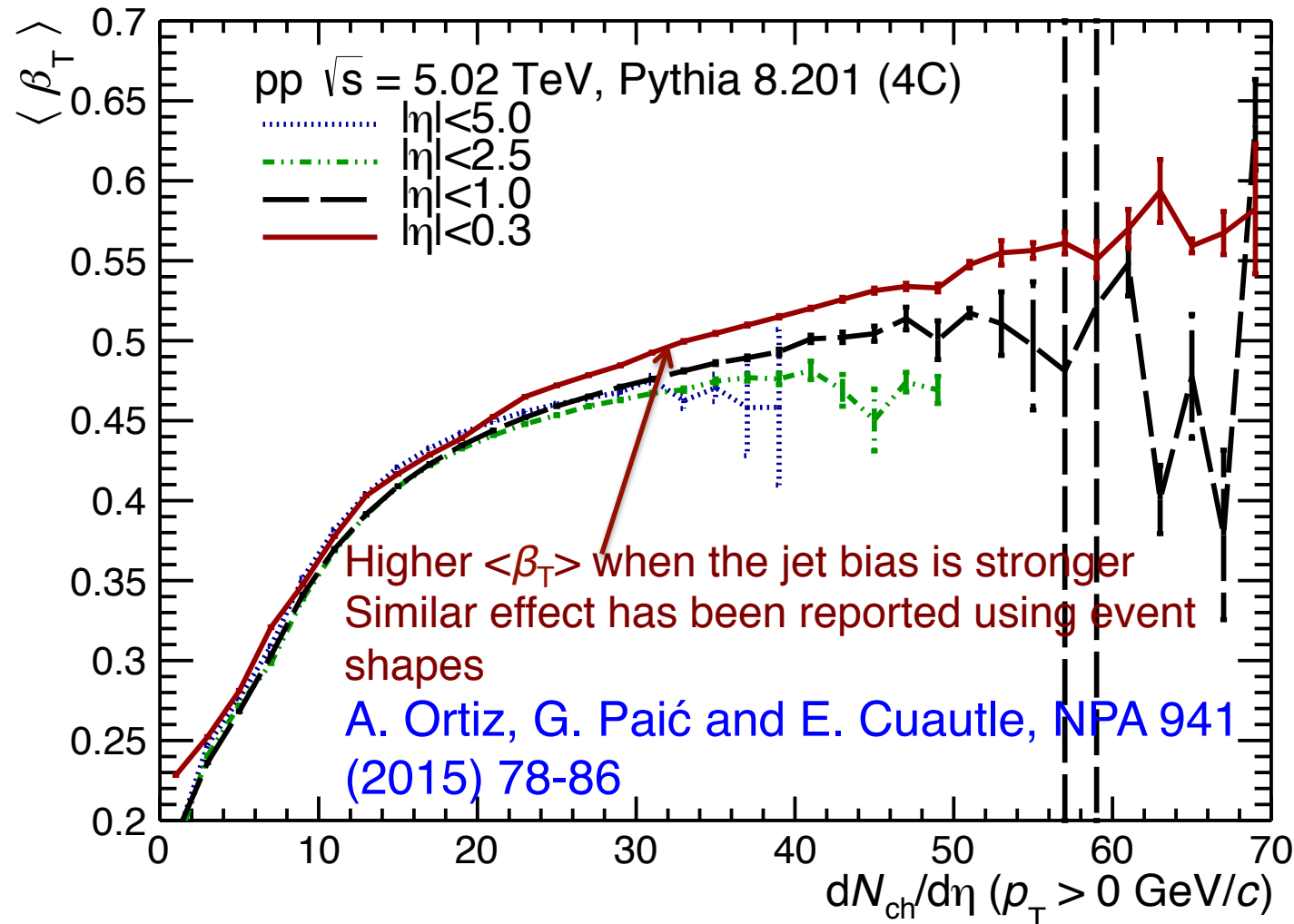
# OTHER APPROACHES

# Jet effects can be also seen in a more inclusive analysis



- $\langle p_T \rangle$  was computed with charged particles within  $|\eta| < 0.3$
- $dN_{ch}/d\eta$  was computed using different  $\eta$  windows

# Jet effects can be also seen in a more inclusive analysis



- $\square$   $\langle \beta_T \rangle$  was computed with charged particles within  $|\eta| < 0.3$
- $\square$   $dN_{ch}/d\eta$  was computed using different  $\eta$  windows



# PID in charged jets

Xianguo Lu, for the ALICE Collaboration, NPA 00 (2014), 1-4

