

Multi-particle production in small systems from CGC

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7th International Workshop on Multiple Partonic Interactions at the LHC

The Abdus Salam International Centre for Theoretical Physics (ICTP), Trieste, Italy

outline

- Introduction
- Framework of color glass condensate
- Phenomenology of high multiplicity events
- Combining CGC with PYTHIA

Based on the work done in collaboration with :

K. Dusling, L. McLerran, B. Schenke, S. Schlichting & R. Venugopalan

Multi-particle production at high energies

Goal : Study correlated production of particles

$$\left\langle \frac{dN}{dy_1 d^2\mathbf{p}_{\perp 1} \dots dy_q d^2\mathbf{p}_{\perp q}} \right\rangle \Longleftrightarrow \left\langle \frac{dN}{dy_1 d^2\mathbf{p}_{\perp 1}} \right\rangle \dots \left\langle \frac{dN}{dy_q d^2\mathbf{p}_{\perp q}} \right\rangle$$

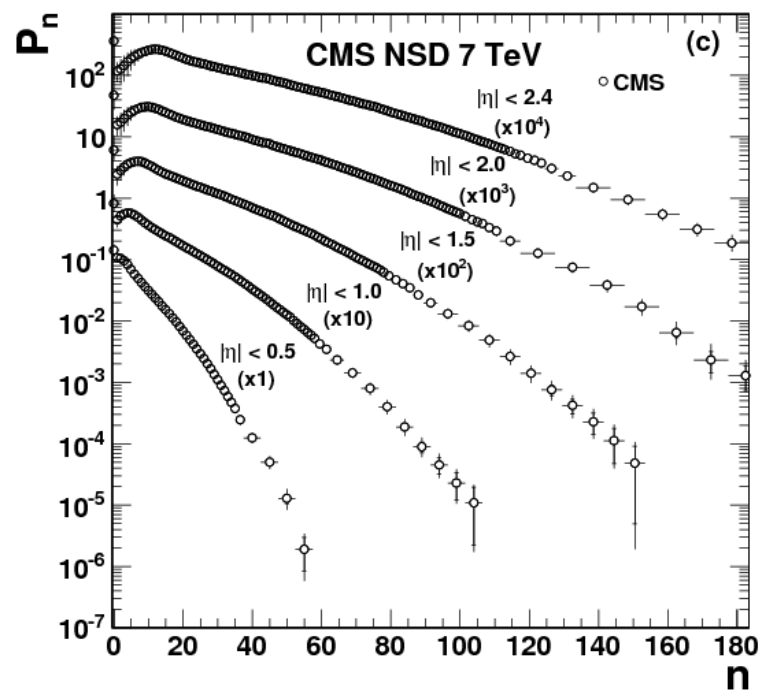
Focus : Collisions of small systems p+p and p+Pb are interesting as final state effects are minimal

We need :

- An *ab-initio* framework of particle production
- Full treatment of different sources of fluctuations
- State-of-the art treatment of fragmentation

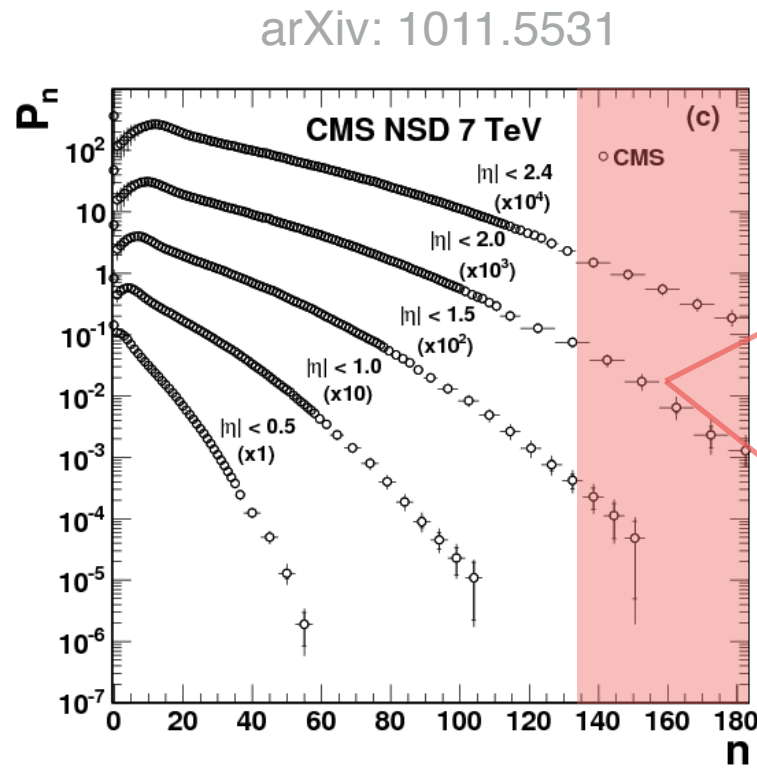
Phenomena we want to describe

arXiv: 1011.5531



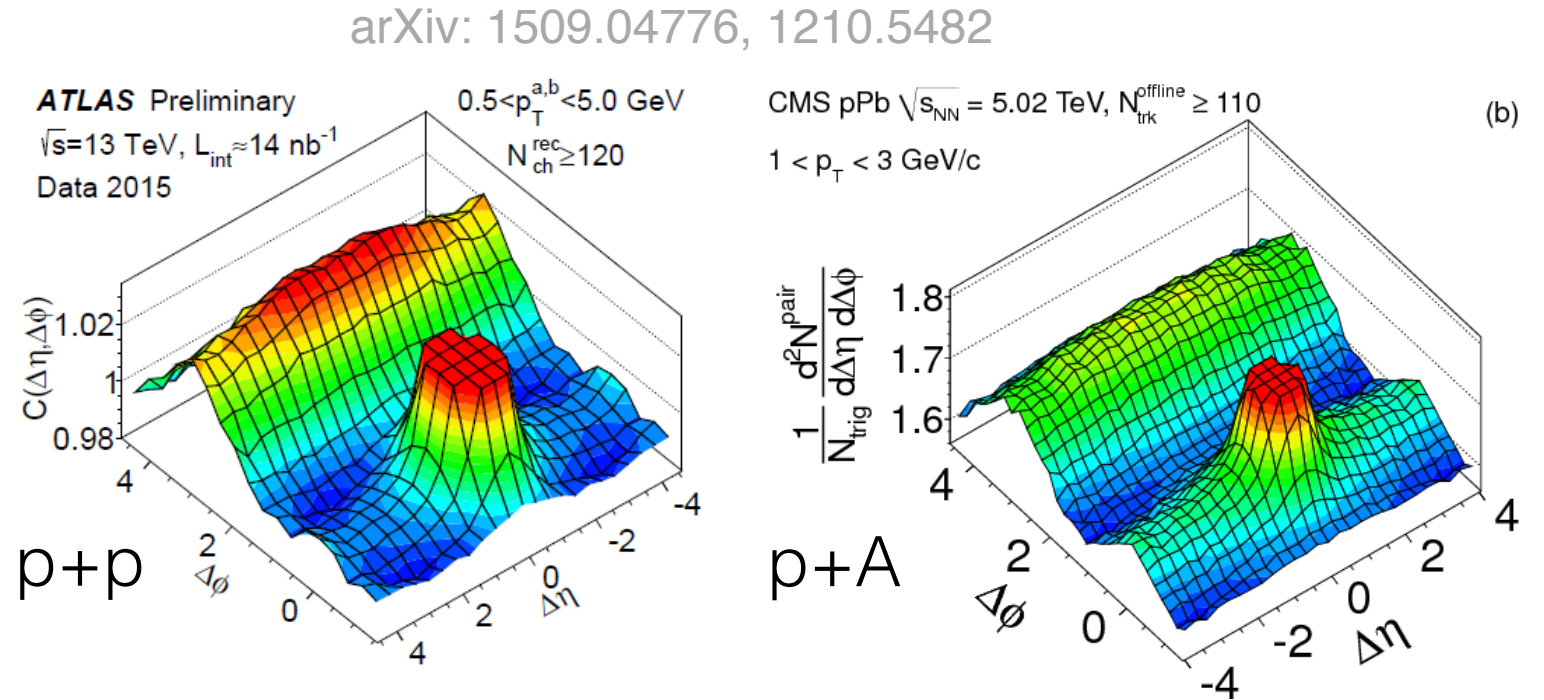
Origin of high multiplicity events

Phenomena we want to describe

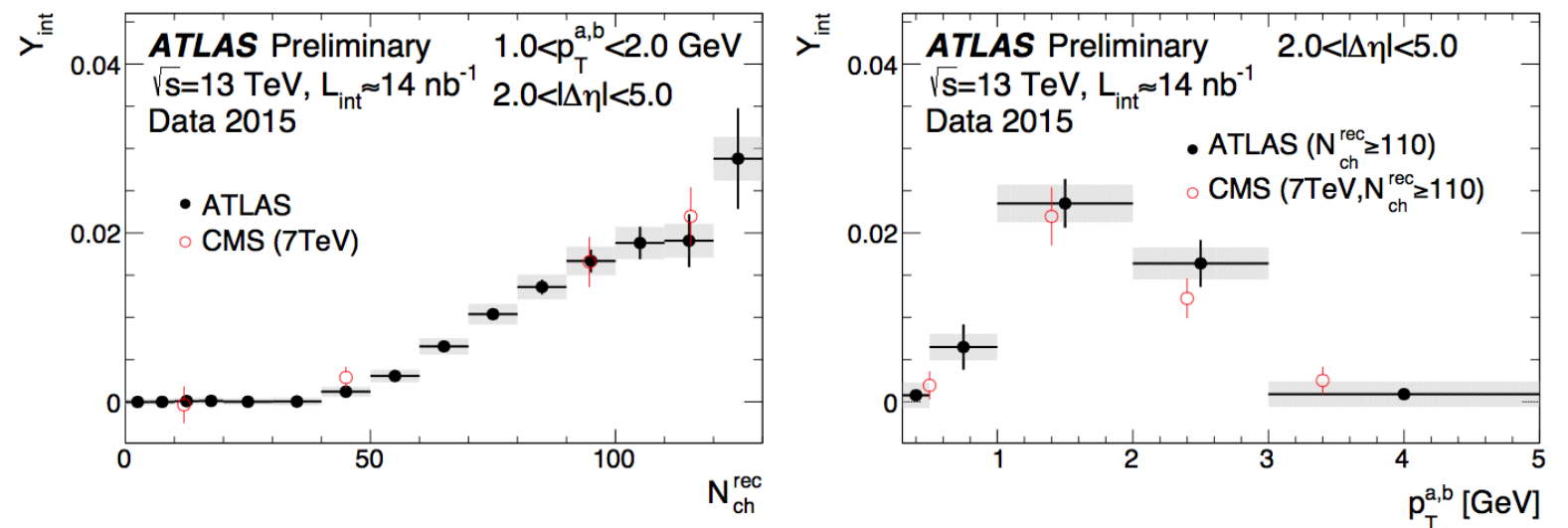


Origin of high multiplicity events

Similar underlying dynamics must drive these phenomenon



Systematics of $\Delta\eta$ - $\Delta\phi$ correlations



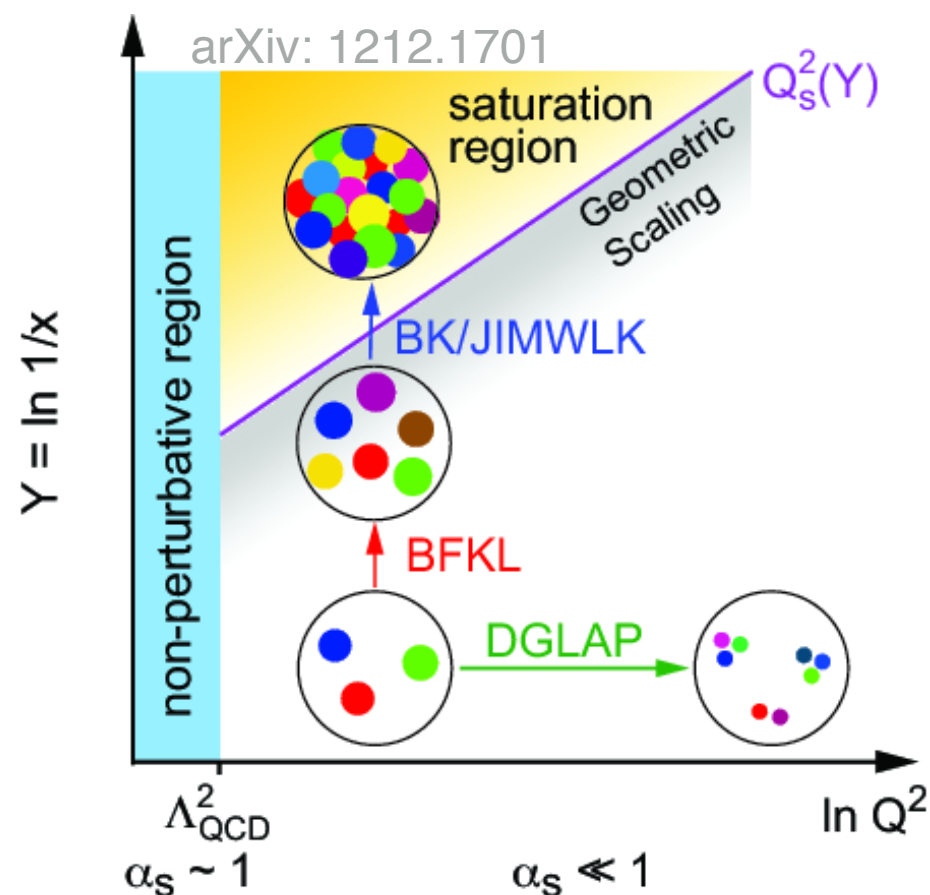
Energy dependence of ridge in p+p

Particle production at high energies

Multi-particle production at high energies in **Regge Gribov limit** ($x \rightarrow 0$)

Colliding hadrons/nuclei :

- Saturation : Non-linear process stops growth of gluons, semi-hard saturation scale $Q_s(x) > \Lambda_{QCD}$

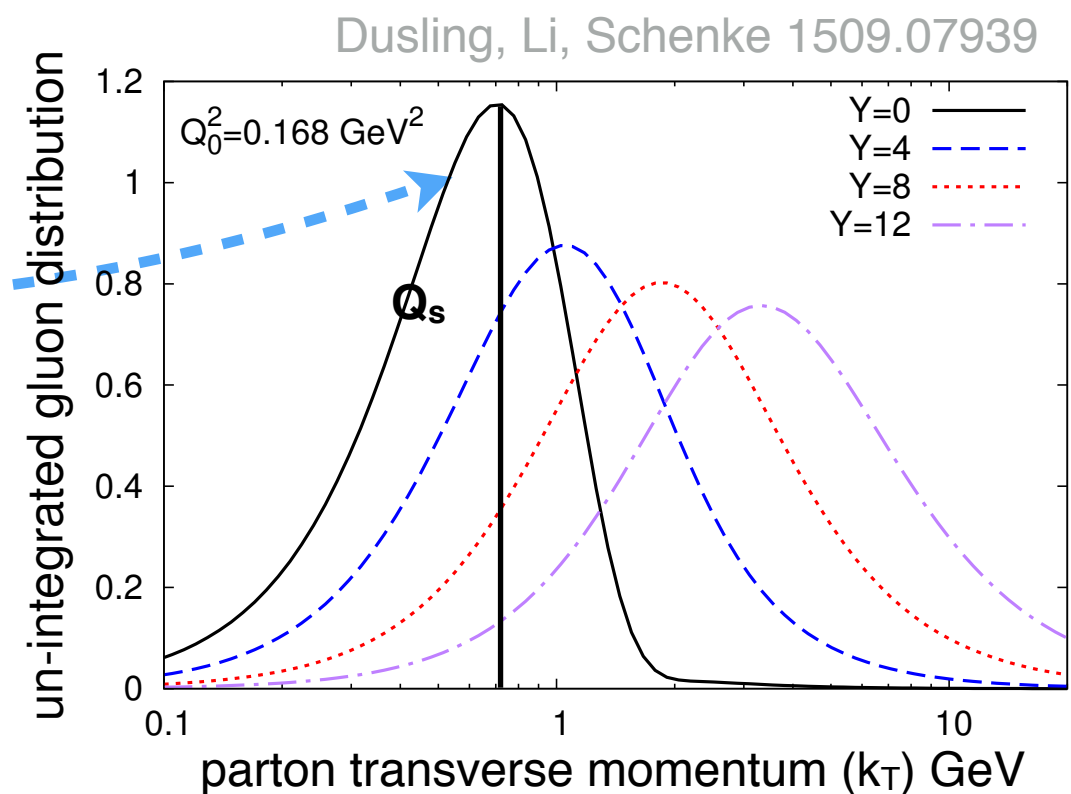
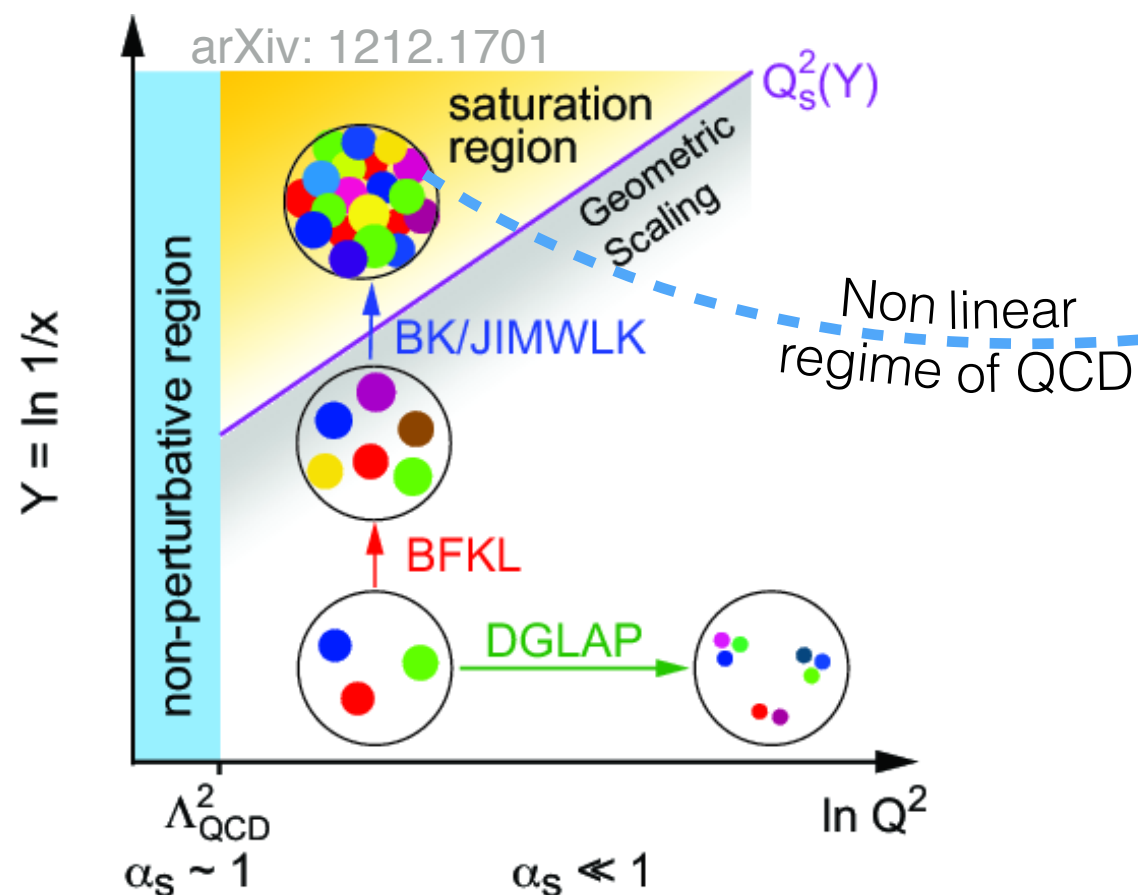


Particle production at high energies

Multi-particle production at high energies in **Regge Gribov limit** ($x \rightarrow 0$)

Colliding hadrons/nuclei :

- Saturation : Non-linear process stops growth of gluons, semi-hard saturation scale $Q_s(x) > \Lambda_{QCD}$
- Gluon dominated wave function, peaked at $Q_s(x \sim x_0 e^{-Y})$

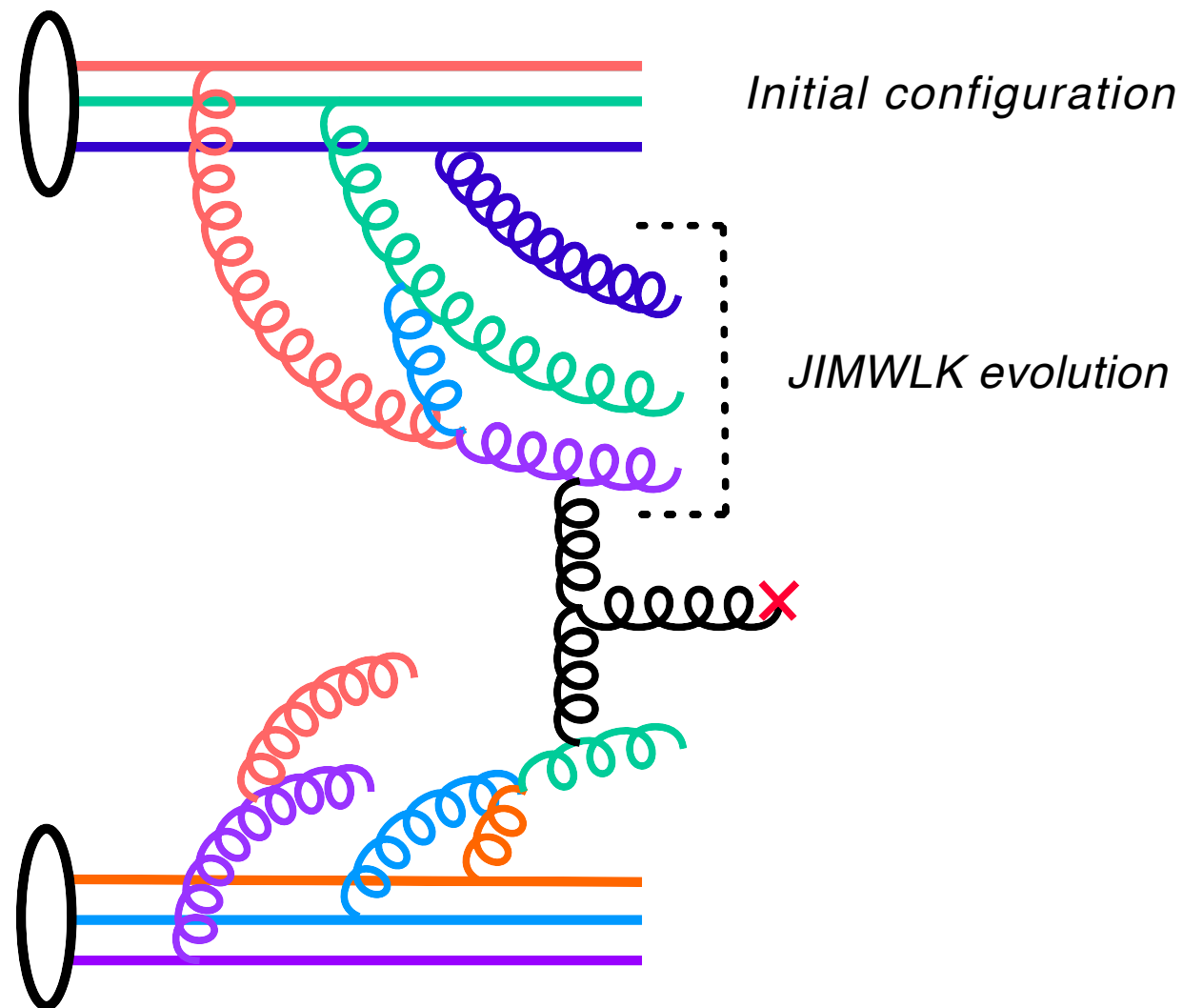


Particle production at high energies

Multi-particle production at high energies in Regge Gribov limit ($x \rightarrow 0$)

Particle production :

- t-channel exchange of ladder like emissions of gluons,



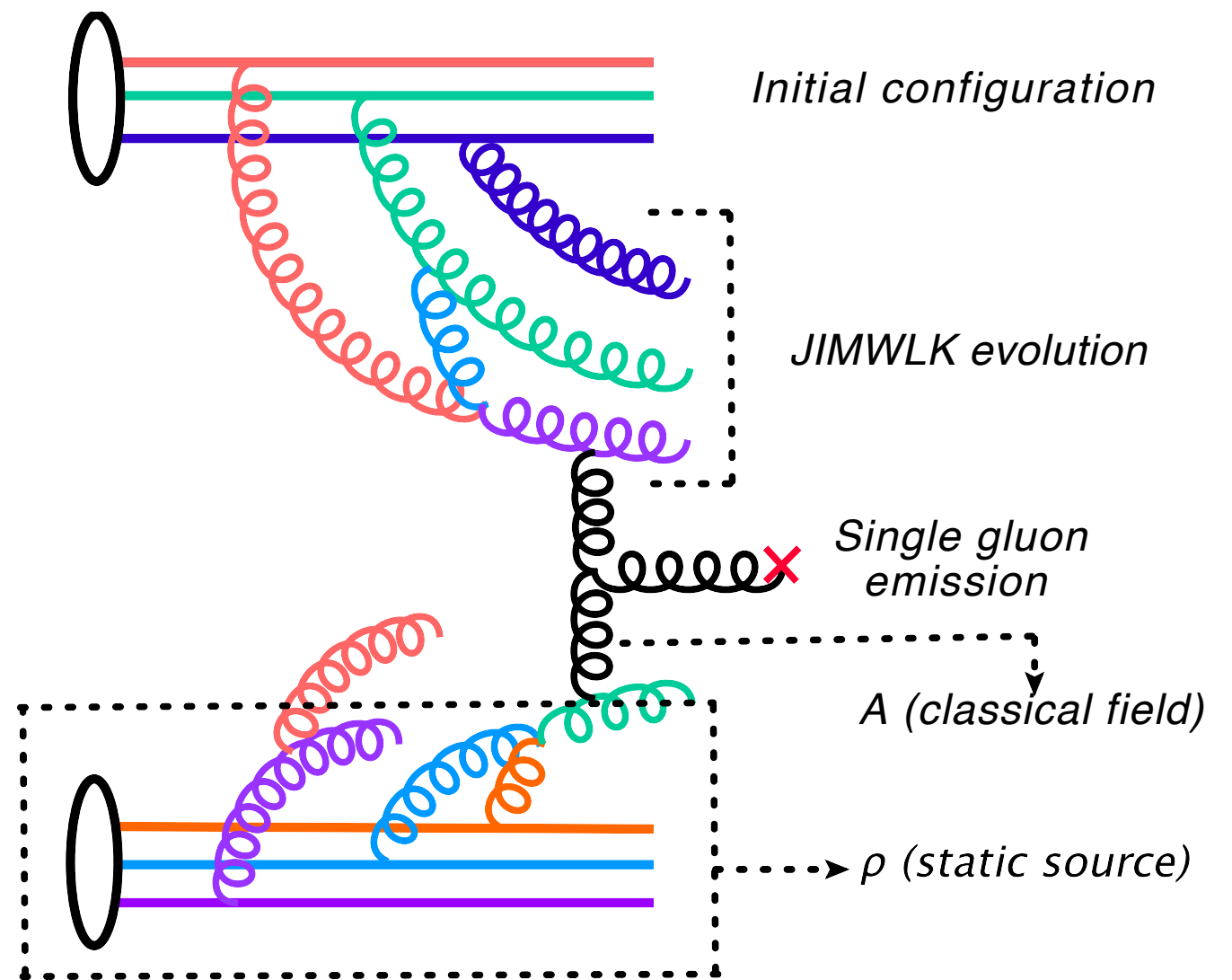
Particle production at high energies

Multi-particle production at high energies in Regge Gribov limit ($x \rightarrow 0$)

Particle production :

- t-channel exchange of ladder like emissions of gluons,
- Strong color fields, weak coupling, *high occupancy of gluonic states* $\sim 1/g^2$

(classical approximation)



McLerran, Venugopalan hep-ph/9309289

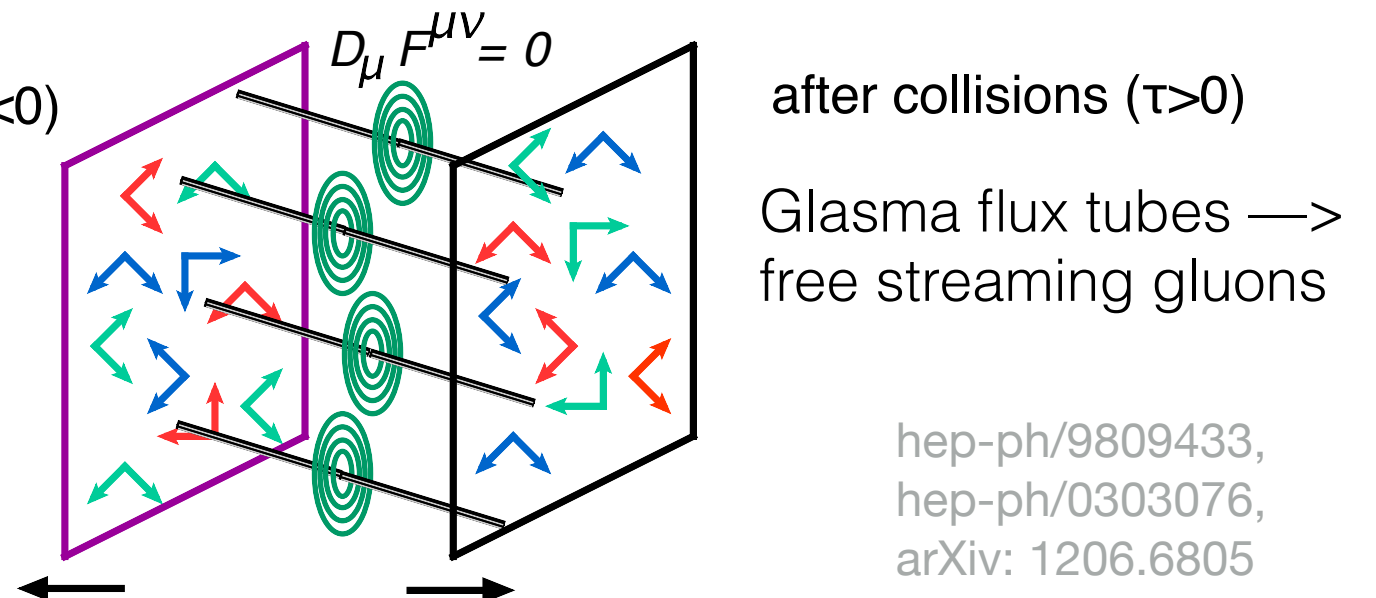
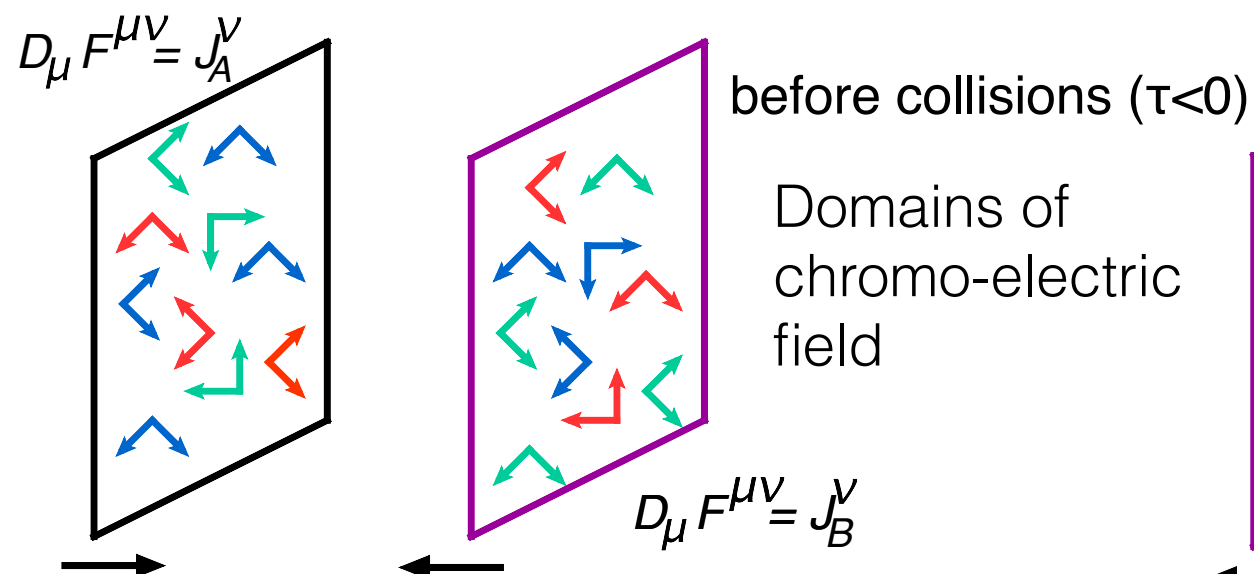
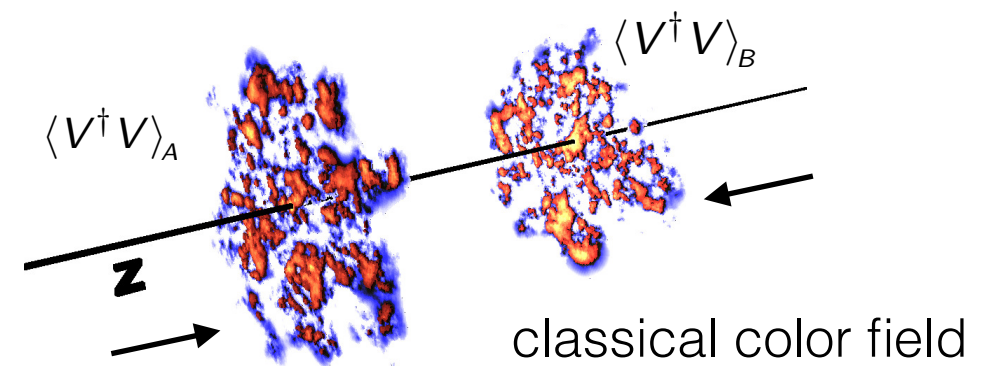
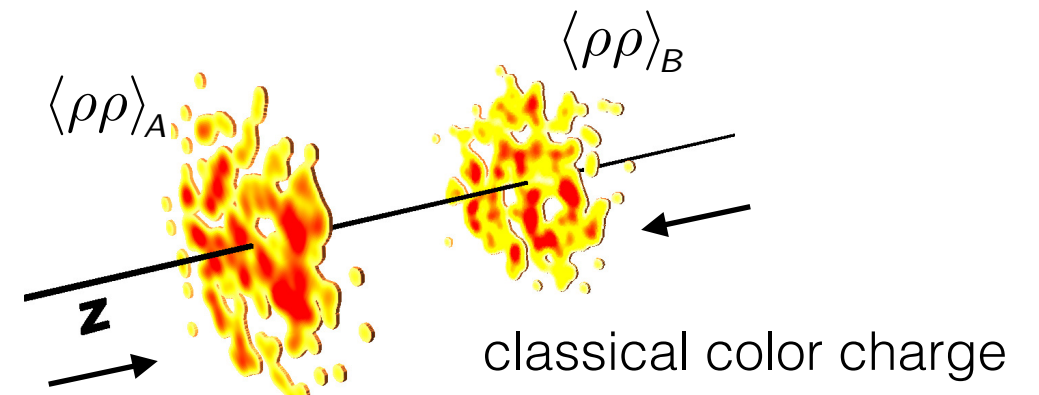
Color Glass condensate effective field theory \rightarrow *ab-initio* framework to this problem

Details of CGC the framework

- Fundamental objects are Color Charge density matrices $\rho^a(\mathbf{x}_\perp, Y)$
Local Gaussian distribution $W[\rho]$
(MV-Model)

$$\langle \rho^a(\mathbf{x}_\perp) \rho^b(\mathbf{y}_\perp) \rangle = \delta^{ab} \delta^2(\mathbf{x}_\perp - \mathbf{y}_\perp) g^2 \mu^2(\mathbf{x}_\perp)$$

- Color field before collisions : solving Yang Mills equations $[D_\mu, F_{\mu\nu}] = J_\nu$ for each configuration of source $\rho(\mathbf{x}_\perp)$



hep-ph/9809433,
hep-ph/0303076,
arXiv: 1206.6805
arXiv: 1202.6646

Details of the CGC framework

Input is constrained by dipole-cross sections in e+p/A collisions

Perturbative approach

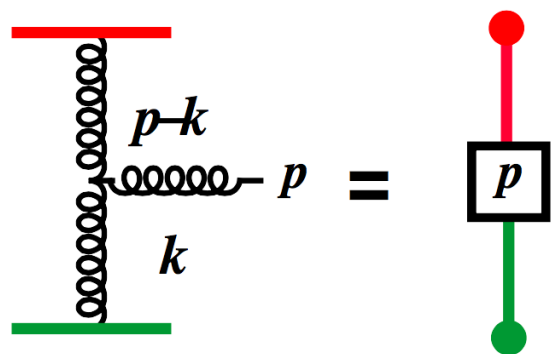
- Employ k_T -factorization ($p_T > Q_s$), dilute-dilute/dense systems

Non-perturbative approach

- Full solutions of CYM on 2+1D lattice : **IP-Glasma**
Monte-Carlo model of initial conditions

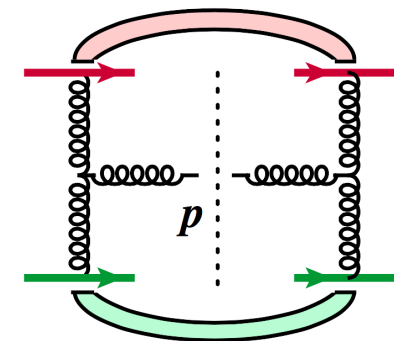
Multi-particle productions

Single-Inclusive



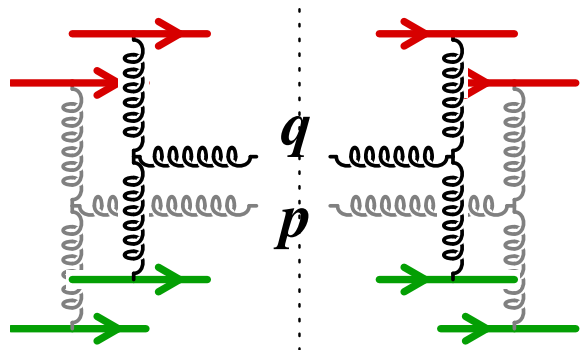
Color Averaging

$$\mathcal{M} \sim \frac{\rho_1(\mathbf{k}_\perp)}{\mathbf{k}_\perp^2} \frac{\rho_2(\mathbf{p}_\perp - \mathbf{k}_\perp)}{(\mathbf{p}_\perp - \mathbf{k}_\perp)^2} L^\gamma(\mathbf{p}, \mathbf{k}_\perp)$$



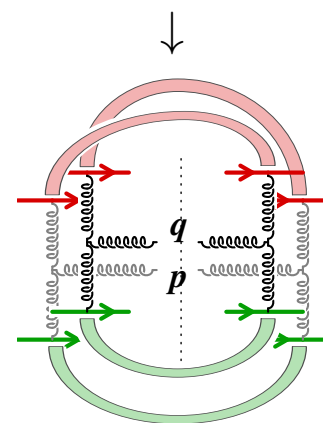
$$\left\langle \frac{dN}{dy_p d^2\mathbf{p}_\perp} \right\rangle \sim \langle |\mathcal{M}|^2 \rangle \sim \langle \rho_1^* \rho_1 \rho_2^* \rho_2 \rangle$$

Double-Inclusive

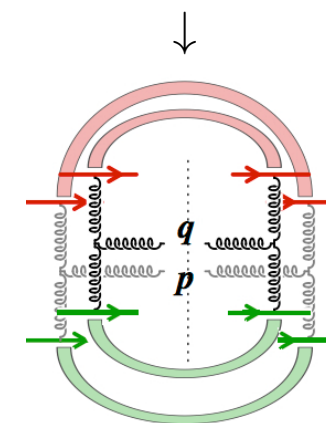


$$\langle |\mathcal{M}|^2 \rangle \rightarrow \langle \rho_1^* \rho_1^* \rho_1 \rho_1 \rho_2^* \rho_2^* \rho_2 \rho_2 \rangle$$

$$C_2(\mathbf{p}, \mathbf{q}) \equiv \left\langle \frac{dN_2}{dy_p d^2\mathbf{p}_\perp dy_q d^2\mathbf{q}_\perp} \right\rangle - \left\langle \frac{dN}{dy_p d^2\mathbf{p}_\perp} \right\rangle \left\langle \frac{dN}{dy_q d^2\mathbf{q}_\perp} \right\rangle$$



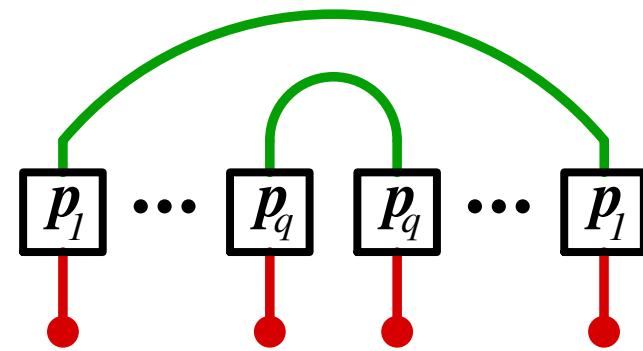
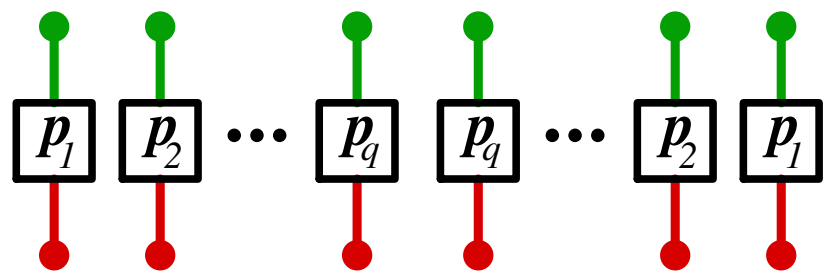
8 topologies



1 topology

n-particle correlations

CGC framework is extendable to n-particle correlations



$2^n(n-1)!$ topologies

Naturally generates Negative Binomial distribution probability distribution

$$P_n^{\text{NB}} = \frac{\Gamma(k+n)}{\Gamma(k)\Gamma(n+1)} \frac{\bar{n}^n k^k}{(\bar{n}+k)^{n+k}}$$

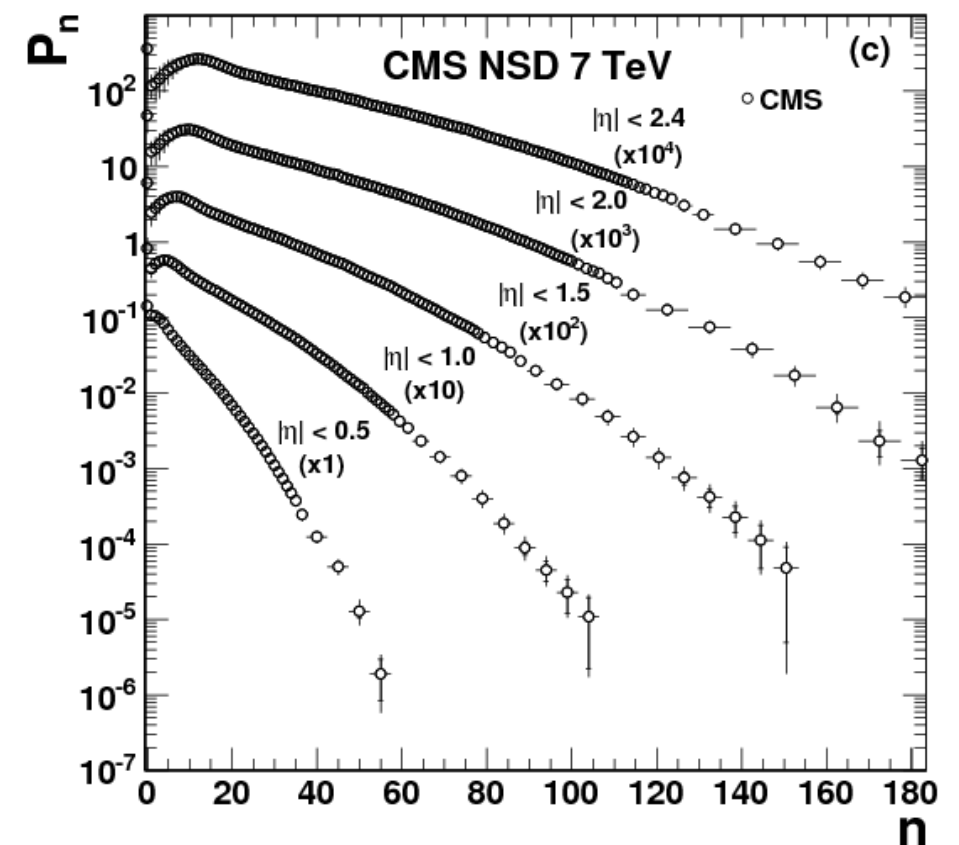
$$k = \kappa \frac{(N_c^2 - 1) Q_s^2 S_\perp}{2\pi}$$

High-multiplicity events \longrightarrow originate from correlated production of n-particles
 \longrightarrow Highly non-perturbative

Description of Multiplicity distribution/ high multiplicity events

IP-Glasma model : combines CGC framework & different sources of initial state fluctuations

1. Collision geometry and impact parameter
2. Color charge
3. Rare Fock-Space configurations



(I) Fluctuation of collision geometry

- Collision geometry is not calculable from first principle
- Eikonal model with thickness profile from HERA data

Proton profile

Schenke, Tribedy, Venugopalan 1311.3636

$$T_p(\mathbf{s}_\perp) = \frac{1}{2\pi B_G} \exp\left(\frac{-\mathbf{s}_\perp^2}{2B_G}\right)$$

Impact parameter distribution

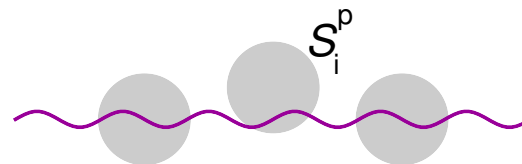
Overlap function

$$T_{pp}(b) = \int d^2\mathbf{s}_\perp T_p^A(\mathbf{s}_\perp) T_p^B(\mathbf{s}_\perp - \mathbf{b}_\perp).$$

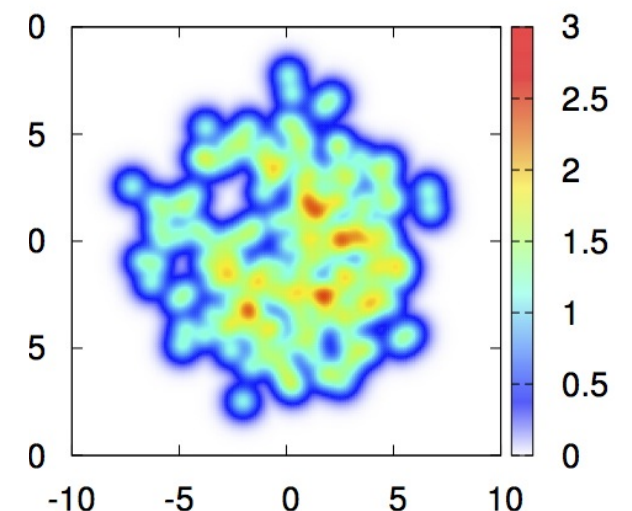
$$\frac{dP}{d^2b}(b) = \frac{1 - e^{-\sigma_{gg} N_g^2 T_{pp}(b)}}{\int d^2b \left(1 - e^{-\sigma_{gg} N_g^2 T_{pp}(b)}\right)}$$

Making Nucleus out of proton scattering

$$S_{\text{dip}}^A(\mathbf{r}_\perp, x, \mathbf{b}_\perp) = \prod_{i=0}^A S_{\text{dip}}^p(\mathbf{r}_\perp, x, \mathbf{b}_\perp)$$



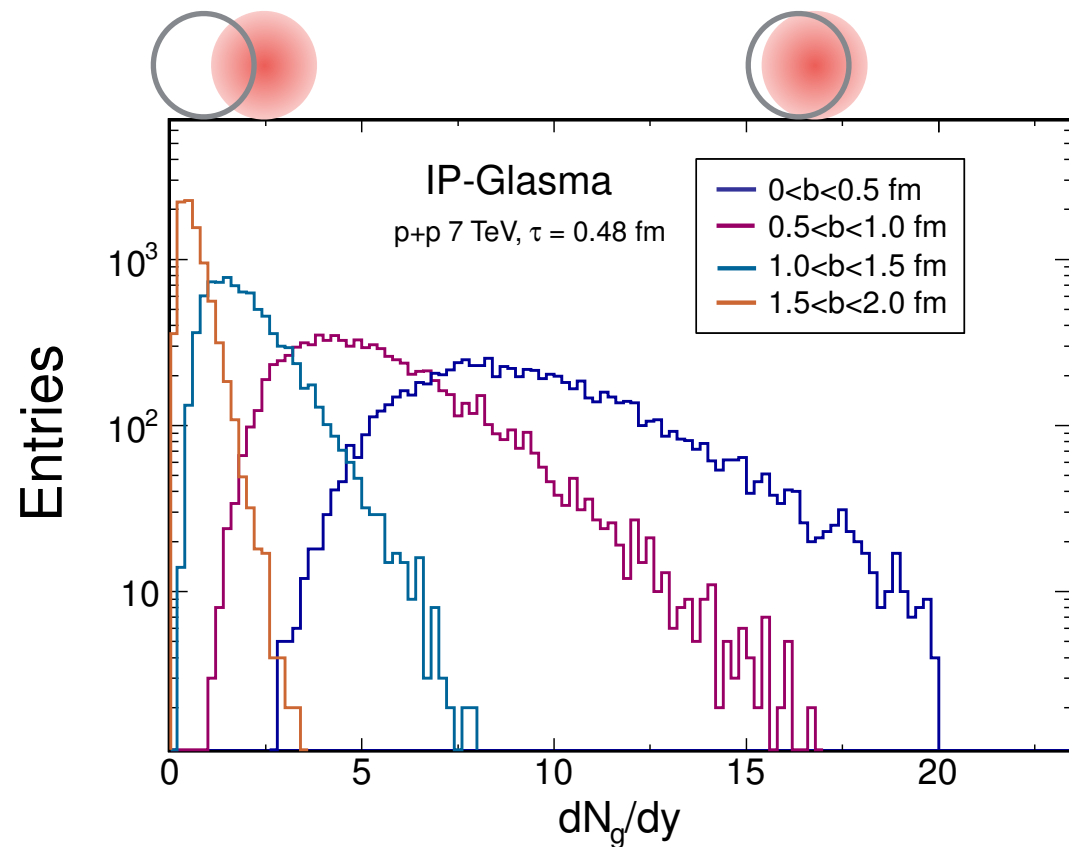
Nuclear saturation scale : $Q_{sA}^2 \sim A^{1/3} Q_{sp}^2$



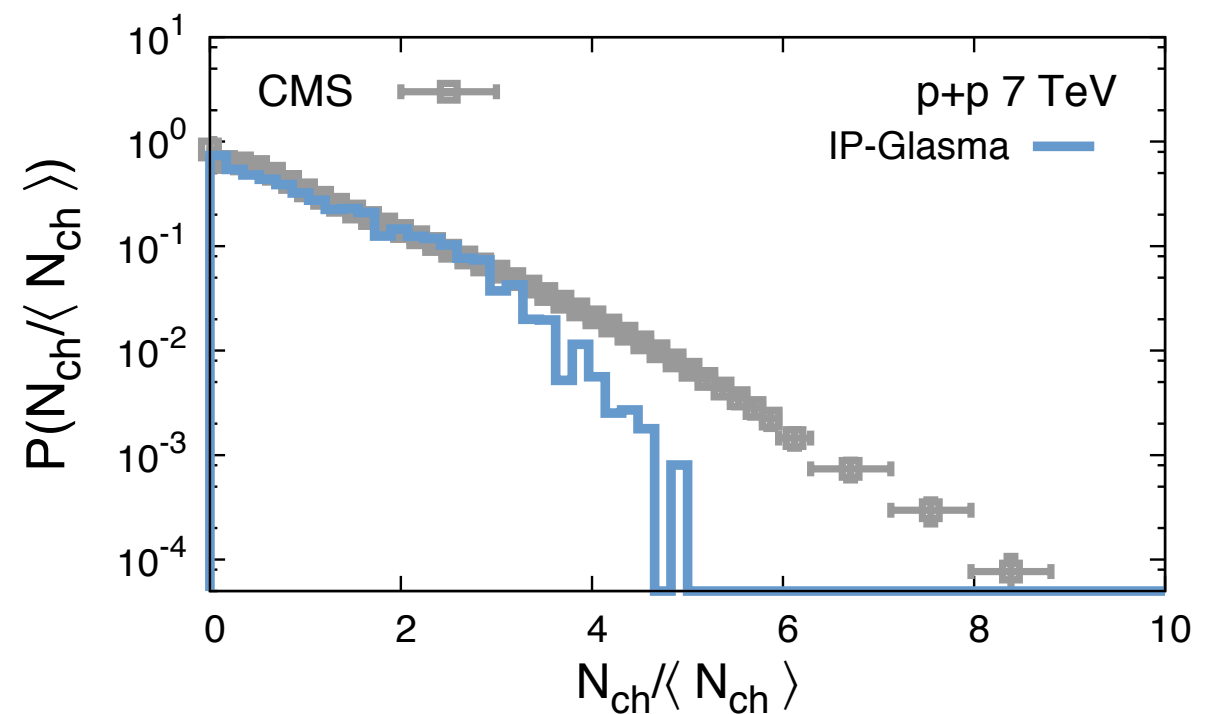
color charge distribution in nucleus

(II) Fluctuation of color charge

For a given geometry fluctuations of color charge
 —> Negative Binomial distribution at each impact parameter



Convolution of many NBDs



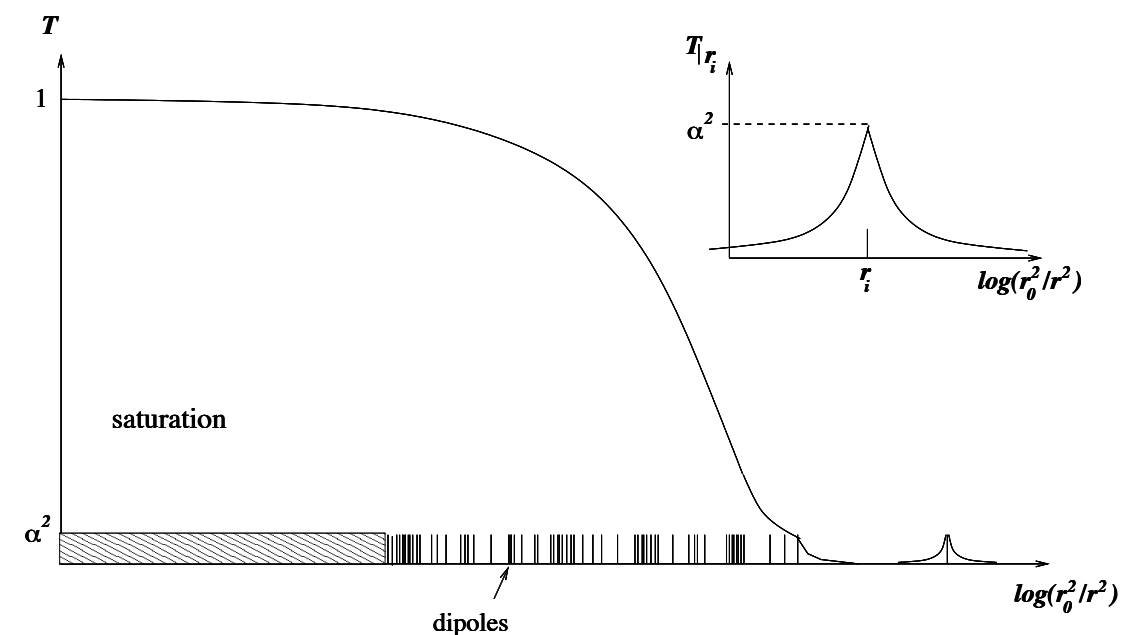
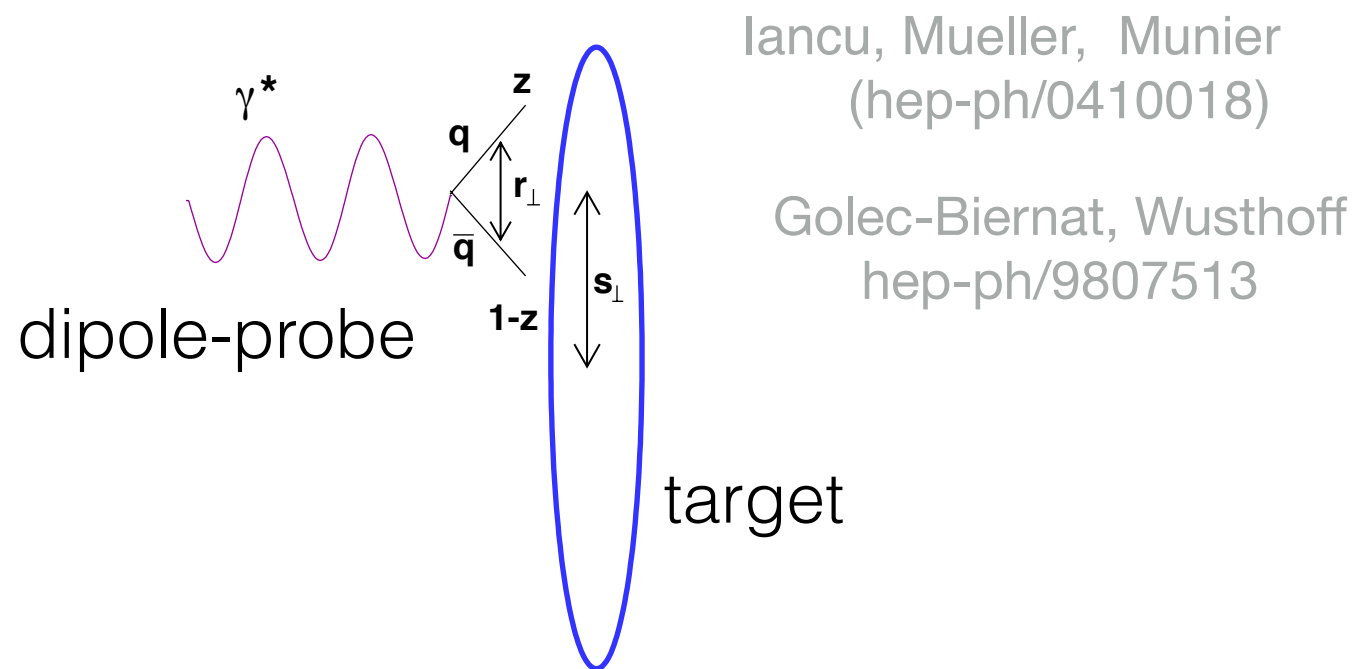
However the distribution is
 not wide enough to describe data

Some sources of fluctuation missing

(III) Intrinsic fluctuations of saturation scale

Input to CGC framework \rightarrow dipole cross section $e+p/A$

Color dipole picture : distribution of partons \rightarrow dist. of color dipoles

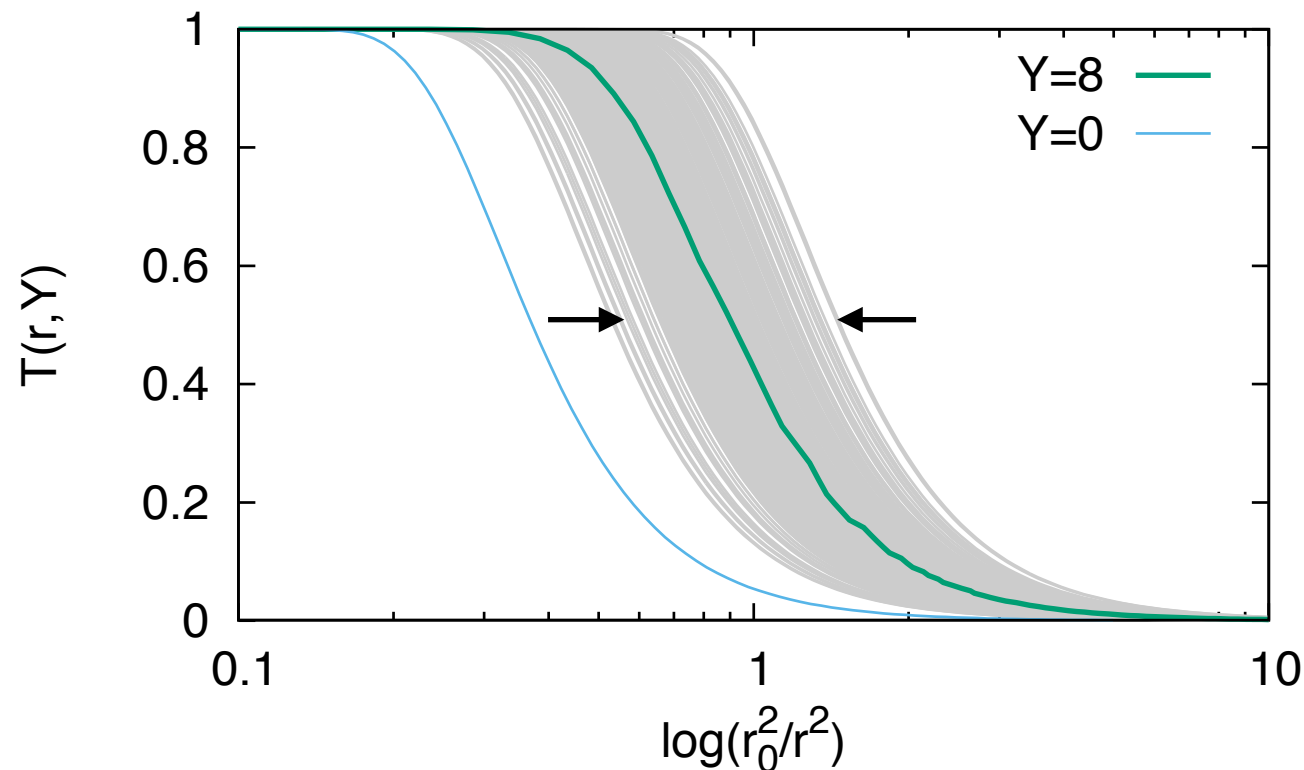


With evolution of rapidity each dipole split with probability $\sim \alpha_s dY$
 \rightarrow dipole splitting is however stochastic

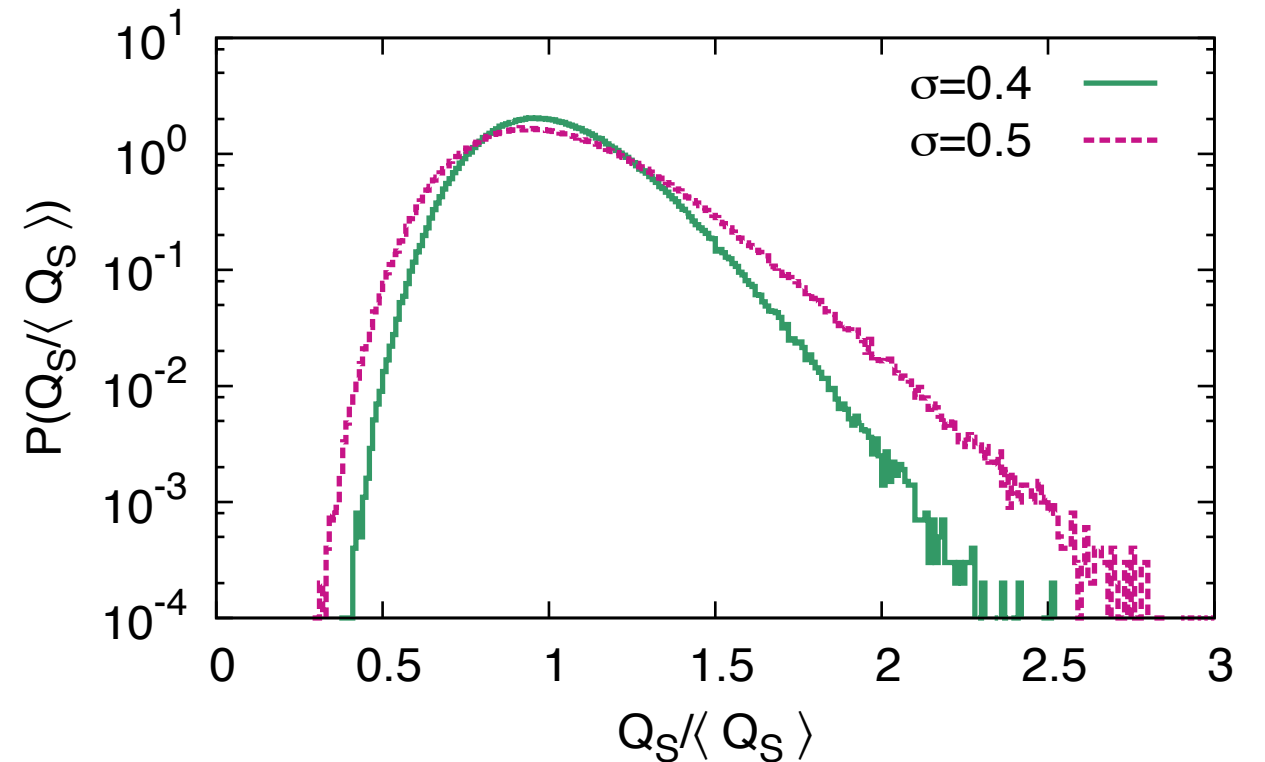
Stochastic dipole splitting \rightarrow not present in BK/JIMWLK \rightarrow beyond CGC

Intrinsic fluctuations of saturation momentum of a proton/nuclei

Dipole amplitude



Saturation scale



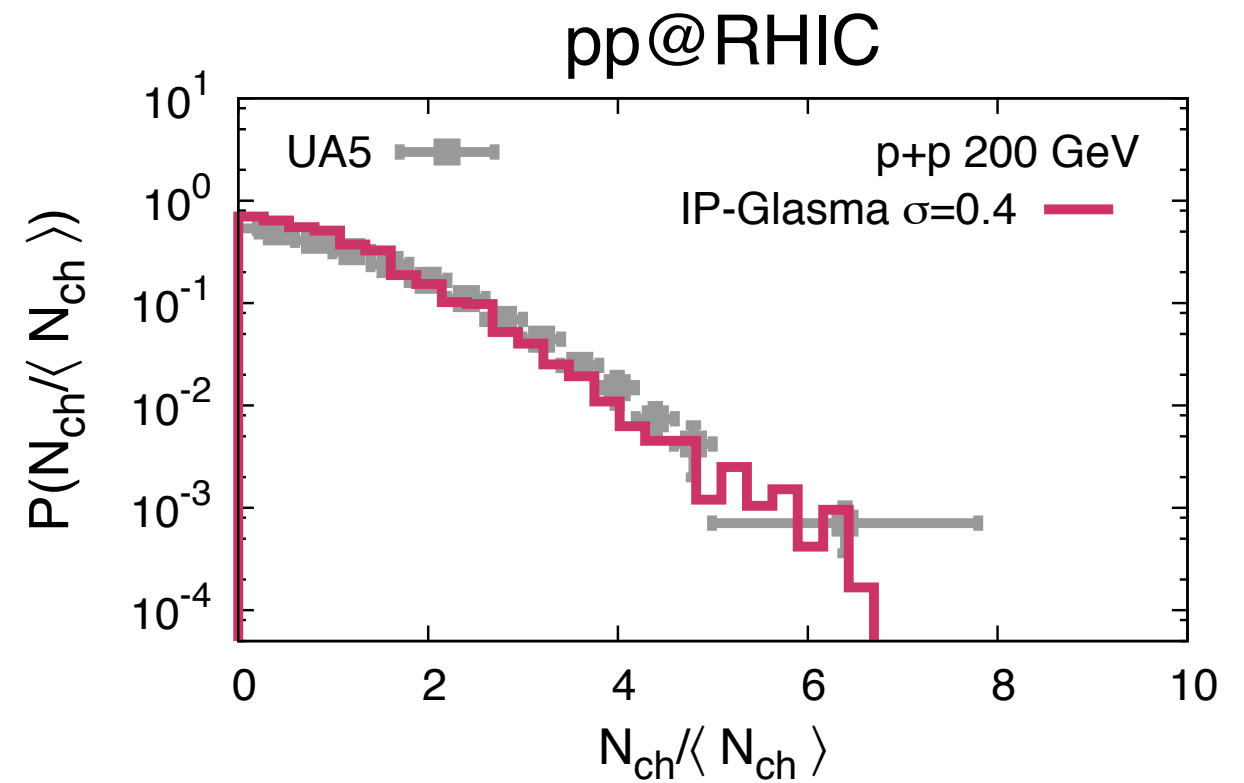
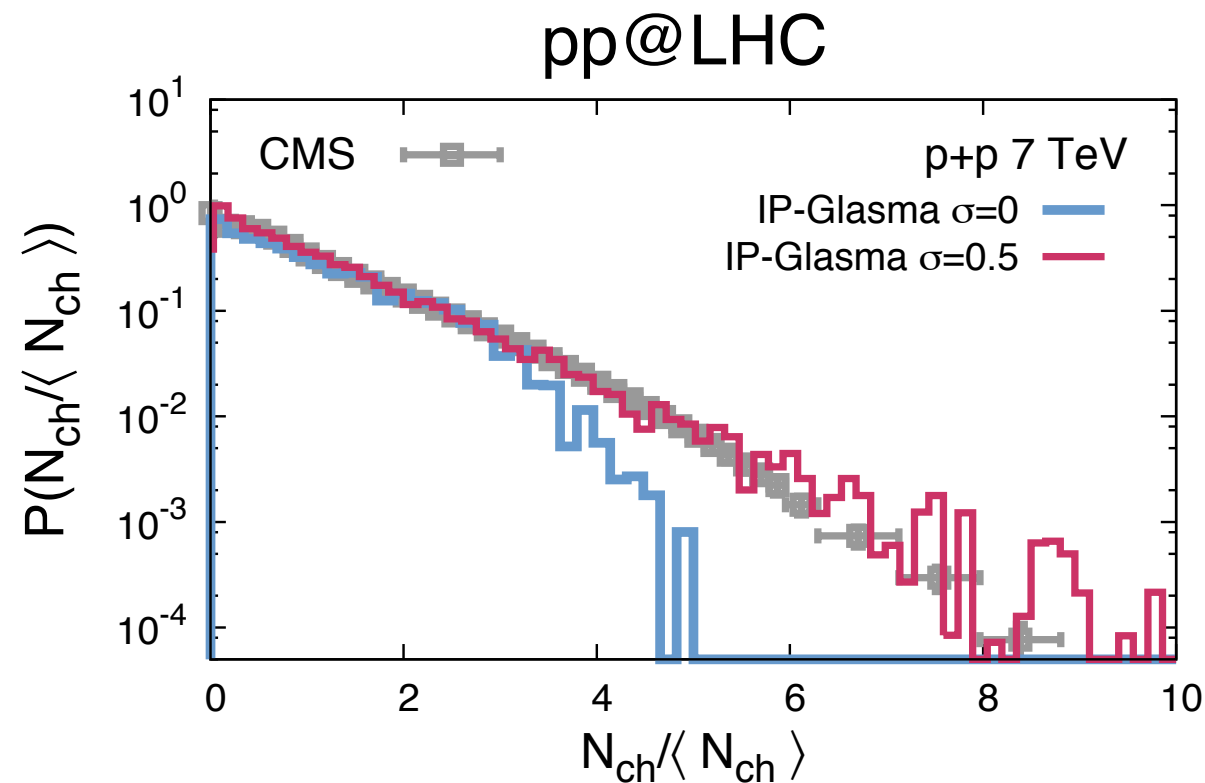
Marquet, Soyez, Xiao hep-ph/0606233

Stochastic splitting of dipole leads to a distribution of Q_s

$$P(\ln(Q_S^2 / \langle Q_S^2 \rangle)) = \frac{1}{\sqrt{2\pi}\sigma} \exp\left(-\frac{\ln^2(Q_S^2(s_\perp) / \langle Q_S^2(s_\perp) \rangle)}{2\sigma^2}\right) \quad \sigma^2(Y) = \sigma_0^2(Y_0) + \sigma_1^2(Y - Y_0)$$

Distribution of multiplicity

McLerran, Tribedy 1508.03292

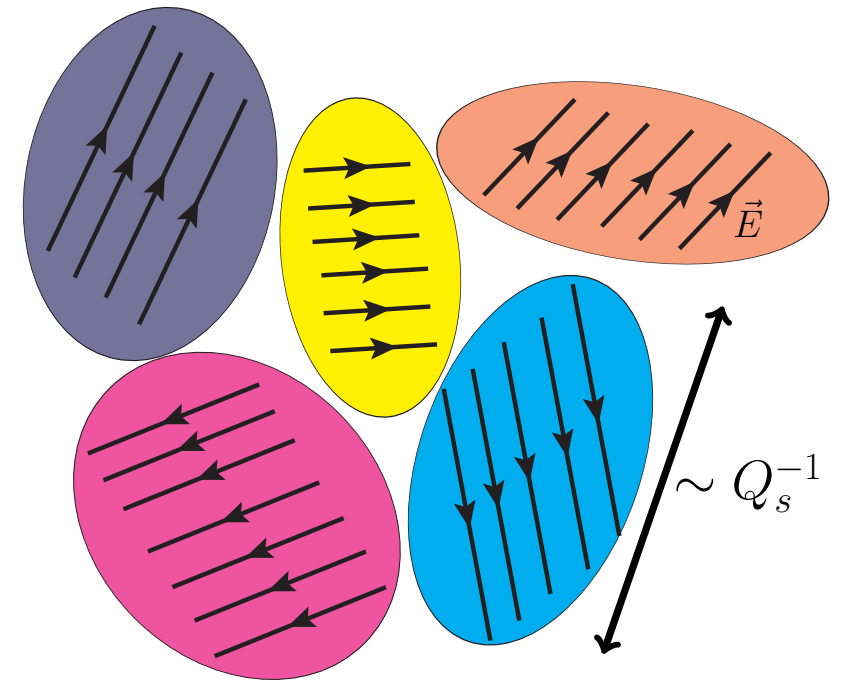


Origin of High multiplicity events (Tail of distributions)

High multiplicity events \rightarrow rare configuration of high color charge density ($1/g^2$)

Azimuthal Correlations in CGC

- Intrinsic momentum space correlation from initial state
- Originate probe scattering off a color domain
- Suppressed by number of color sources/domains

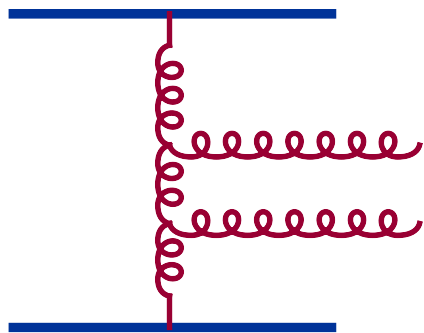


Dumitru, Dusling, Gelis, Jalilian-Marian,
Lappi, Venugopalan 1009.5295
Kovner, Lublinsky 1012.3398
Dusling, Venugopalan 1201.2658
Kovchegov, Wertepny 1212.1195
Dumitru, Giannini 1406.5781
Lappi, Schenke, Schlichting, Venugopalan 1509.03499

Very distinct from Hydrodynamic flow (driven by geometry)

Two particle correlation in CGC

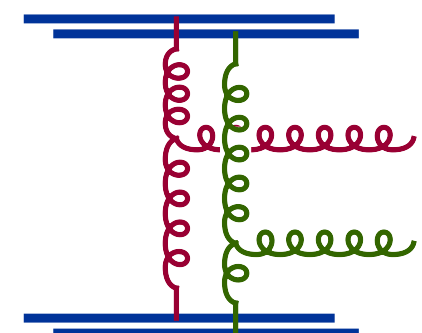
Jet Graph



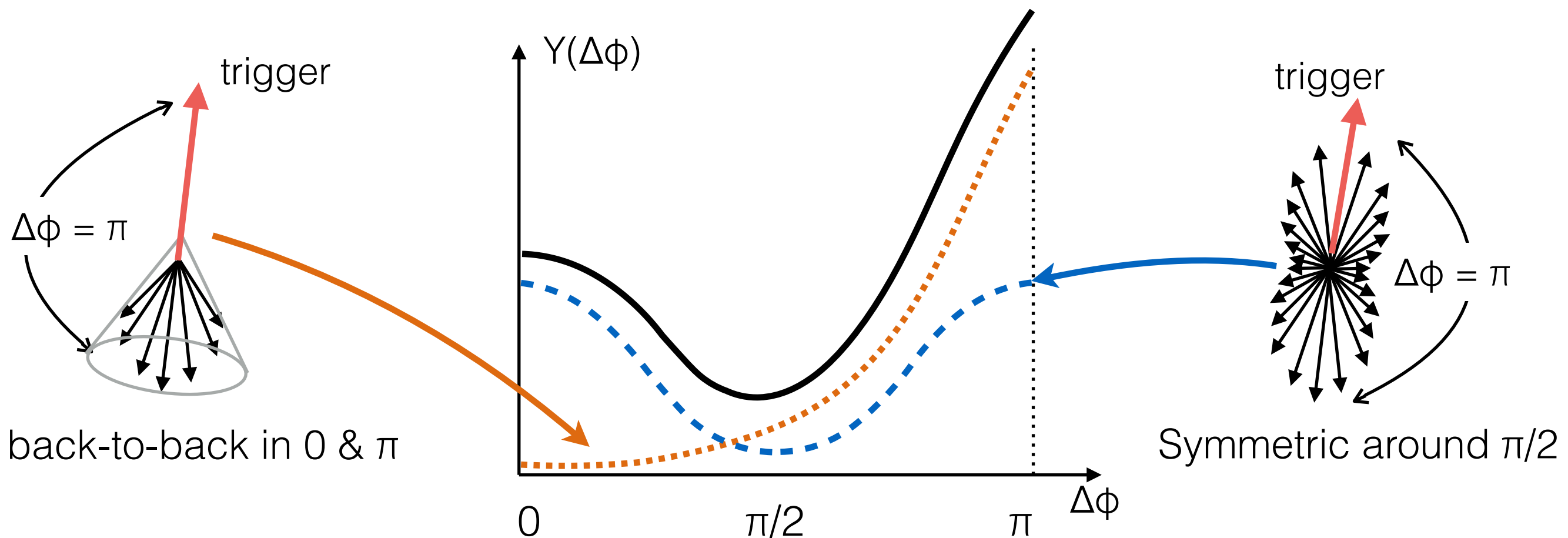
Kinematically constrained (back-to-back)

Dusling, Venugopalan
1201.2658, 1210.3890

Glasma Graph

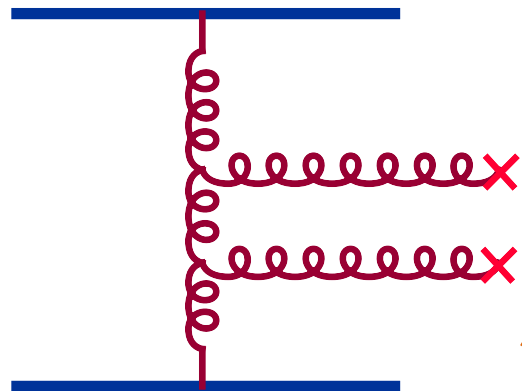


Not kinematically constrained

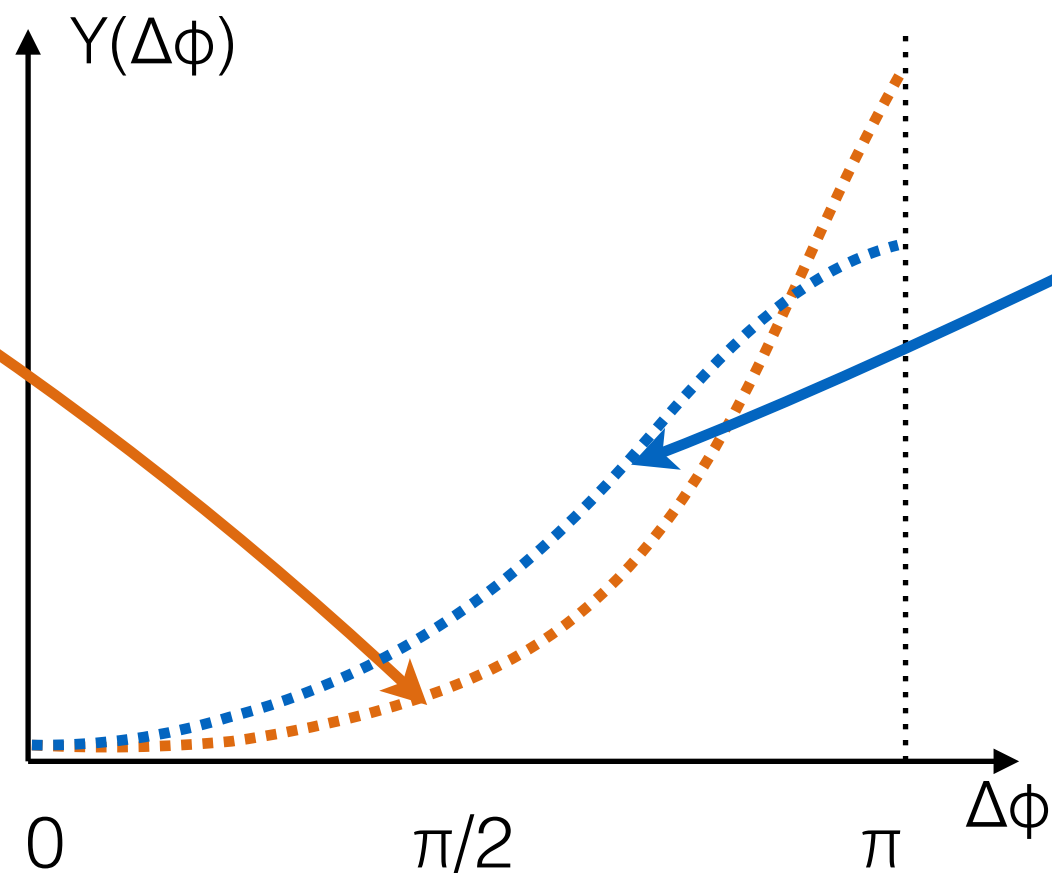
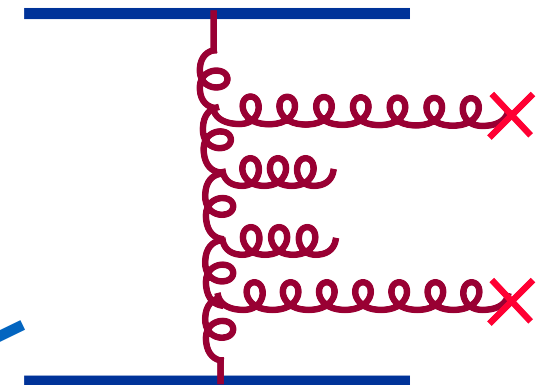


Two particle correlation in CGC

Di-Jet Graph

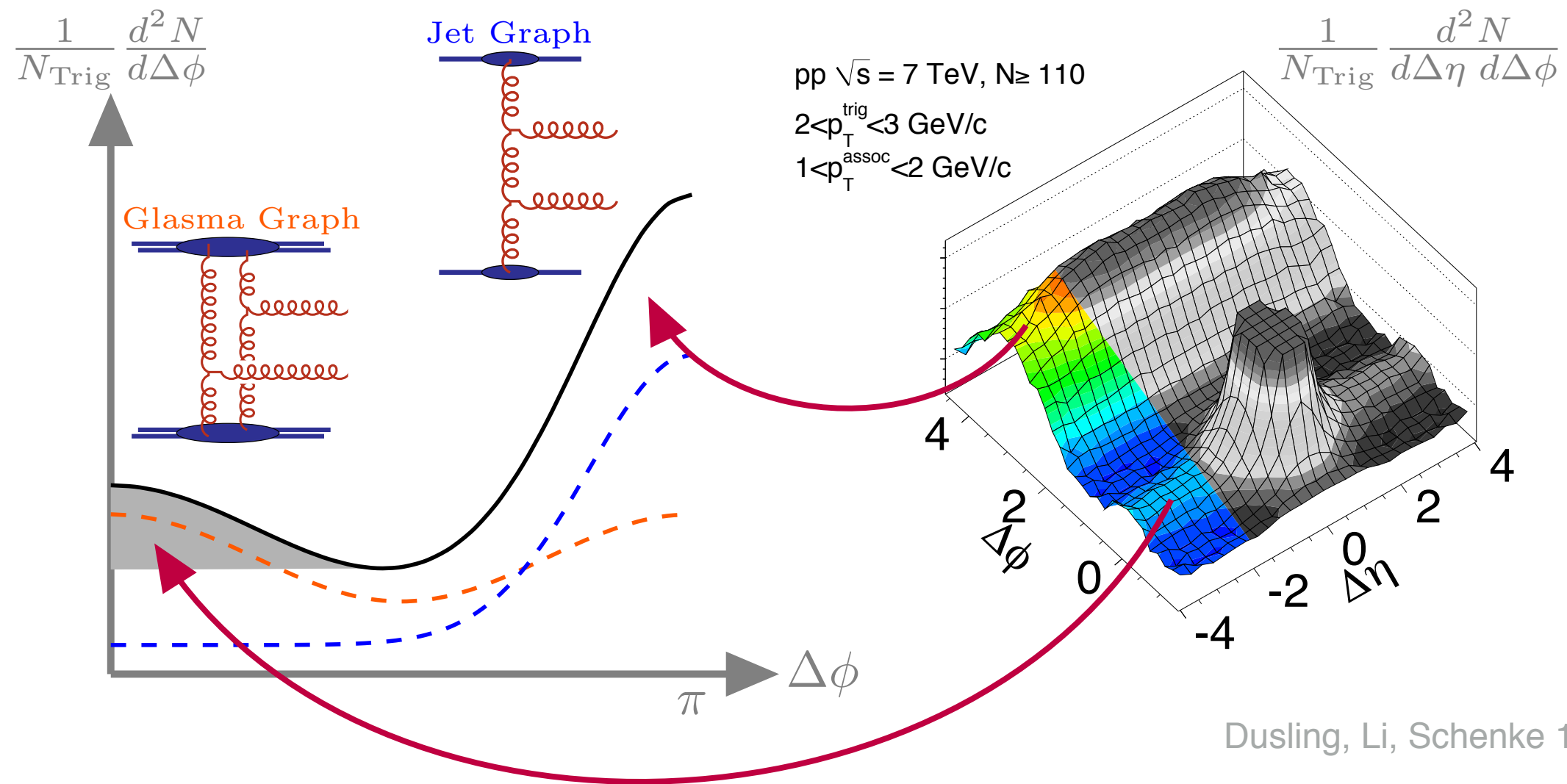


Jet + BFKL emissions



gluon emissions between two triggered hadrons—> broadening of the away side (de-correlation)

Origin of ridge-like correlations

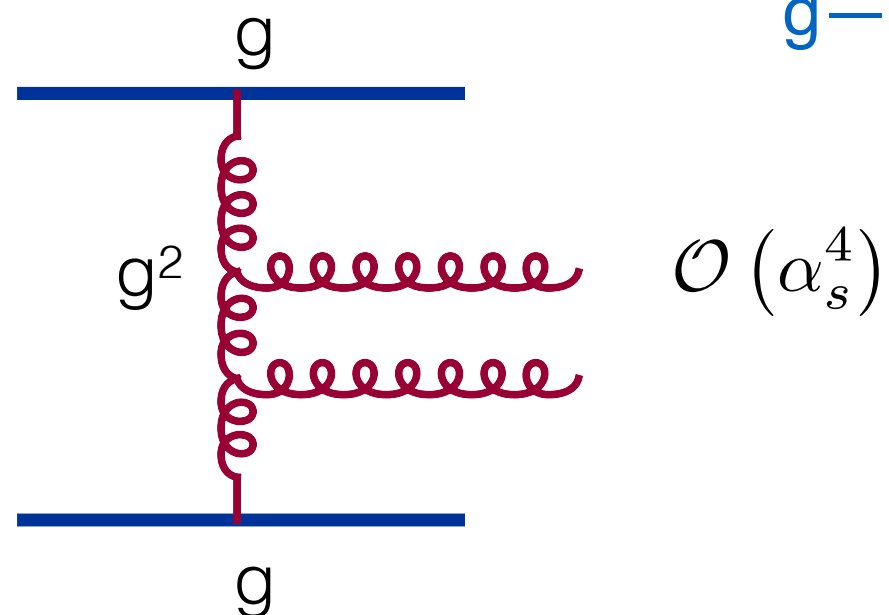
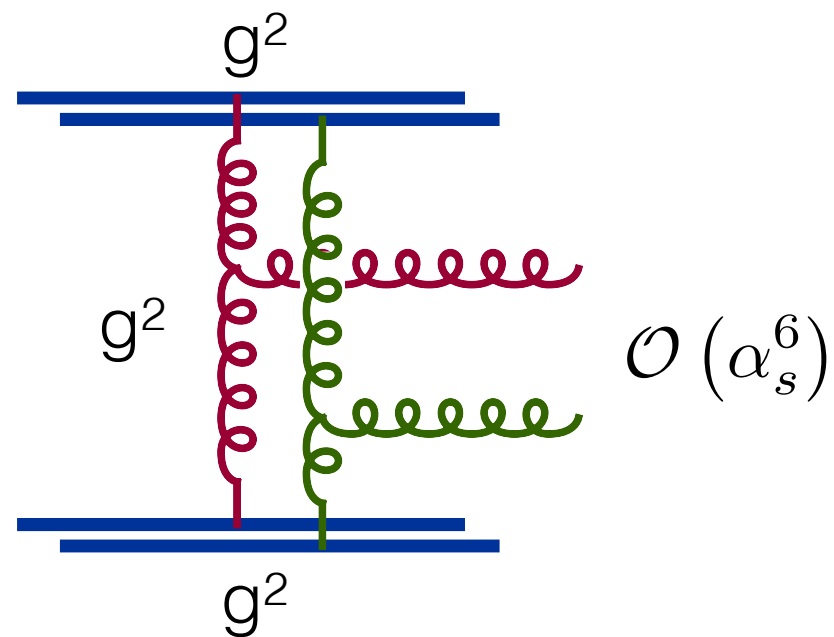


But why ridge appears in high multiplicity events ?

In CGC, high occupancy $\sim 1/g^2 \rightarrow$ effective coupling $1/g^2 \times g = 1/g$

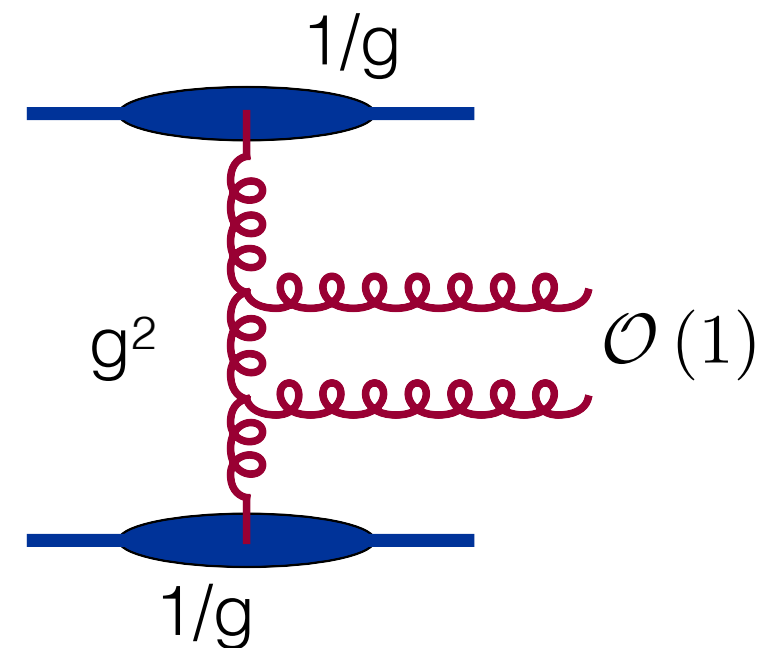
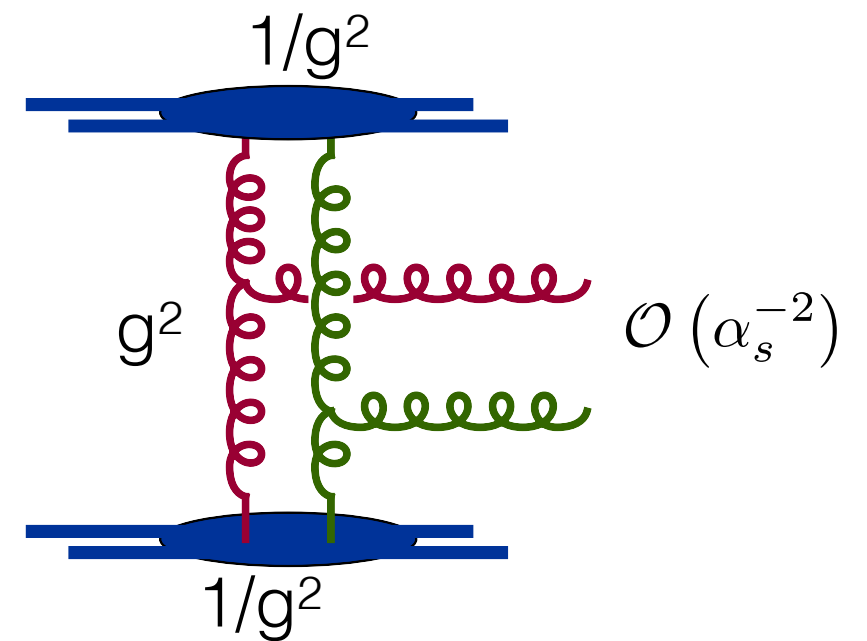
Picture in high multiplicity events

low density (min-bias events)



strong color field in CGC
 $g \rightarrow \rho g \sim 1/g^2$ $g \sim 1/g$

high density (high multiplicity events)



Comparison to data

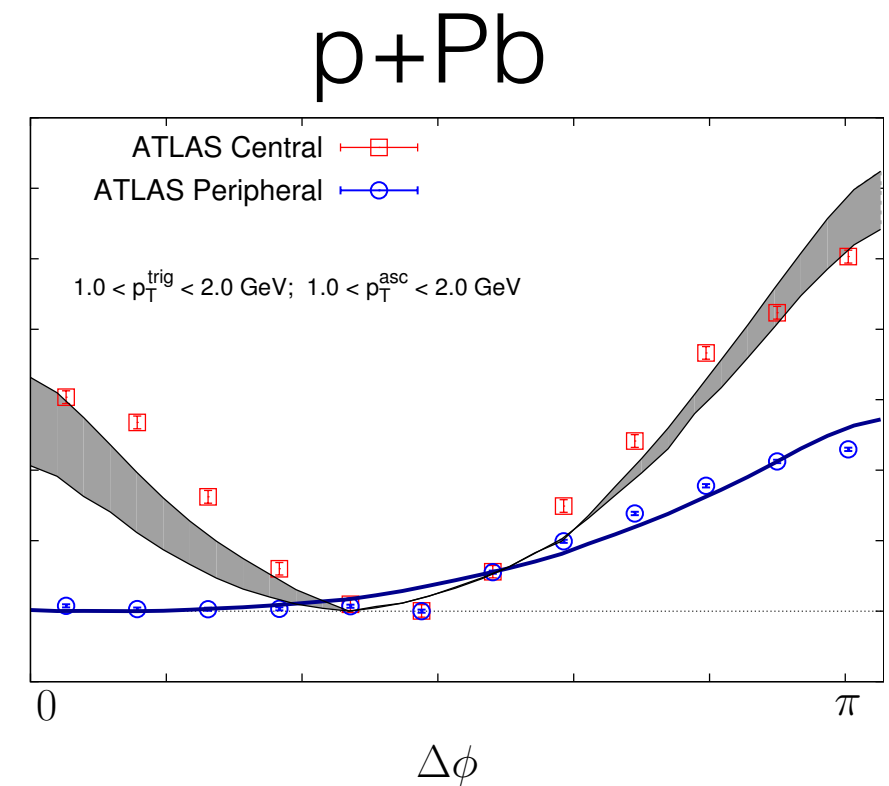
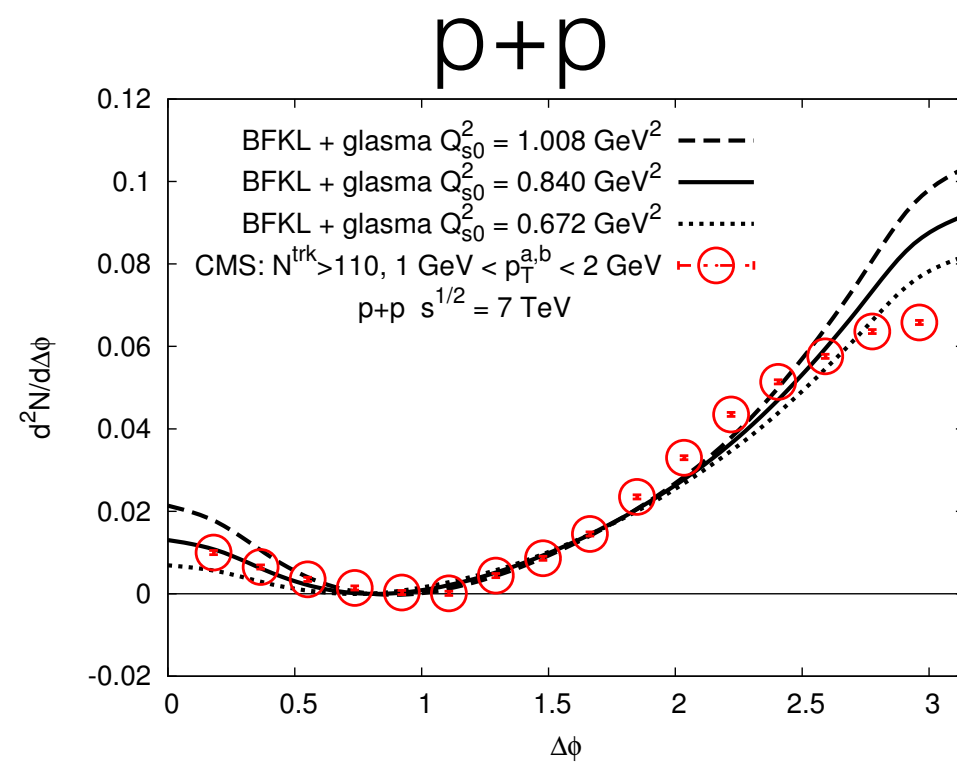


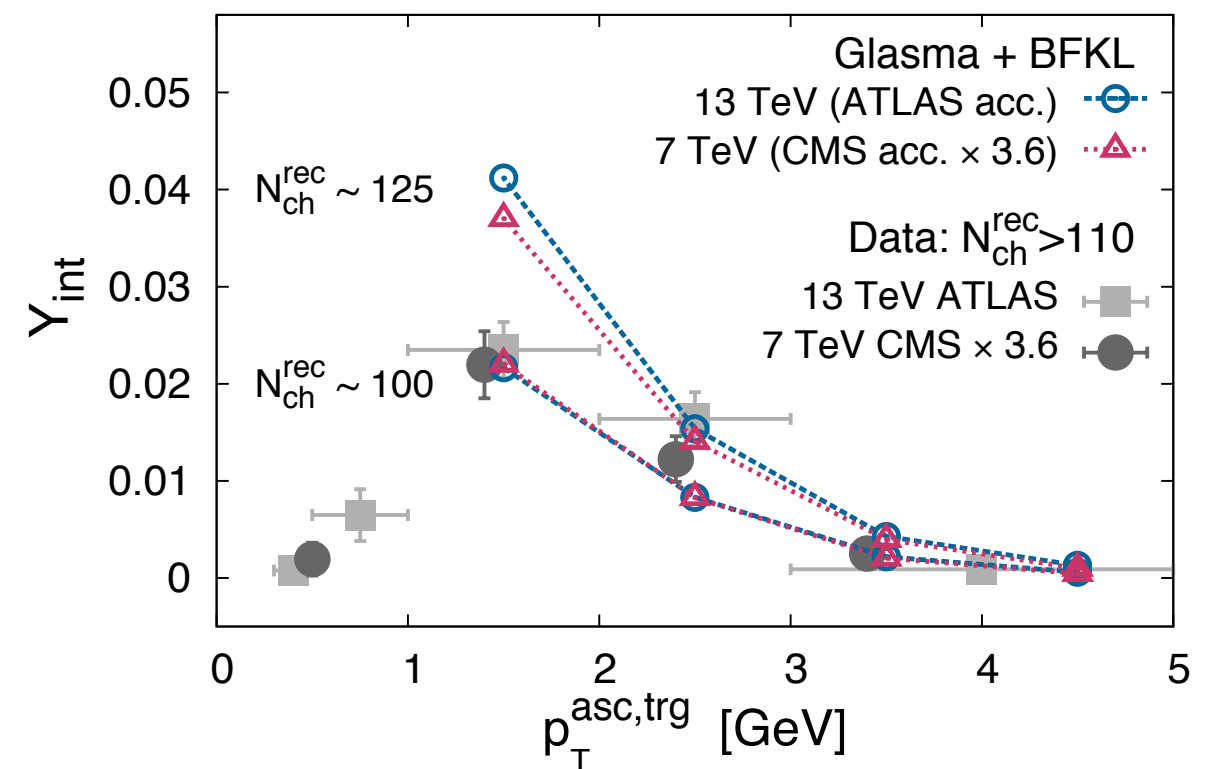
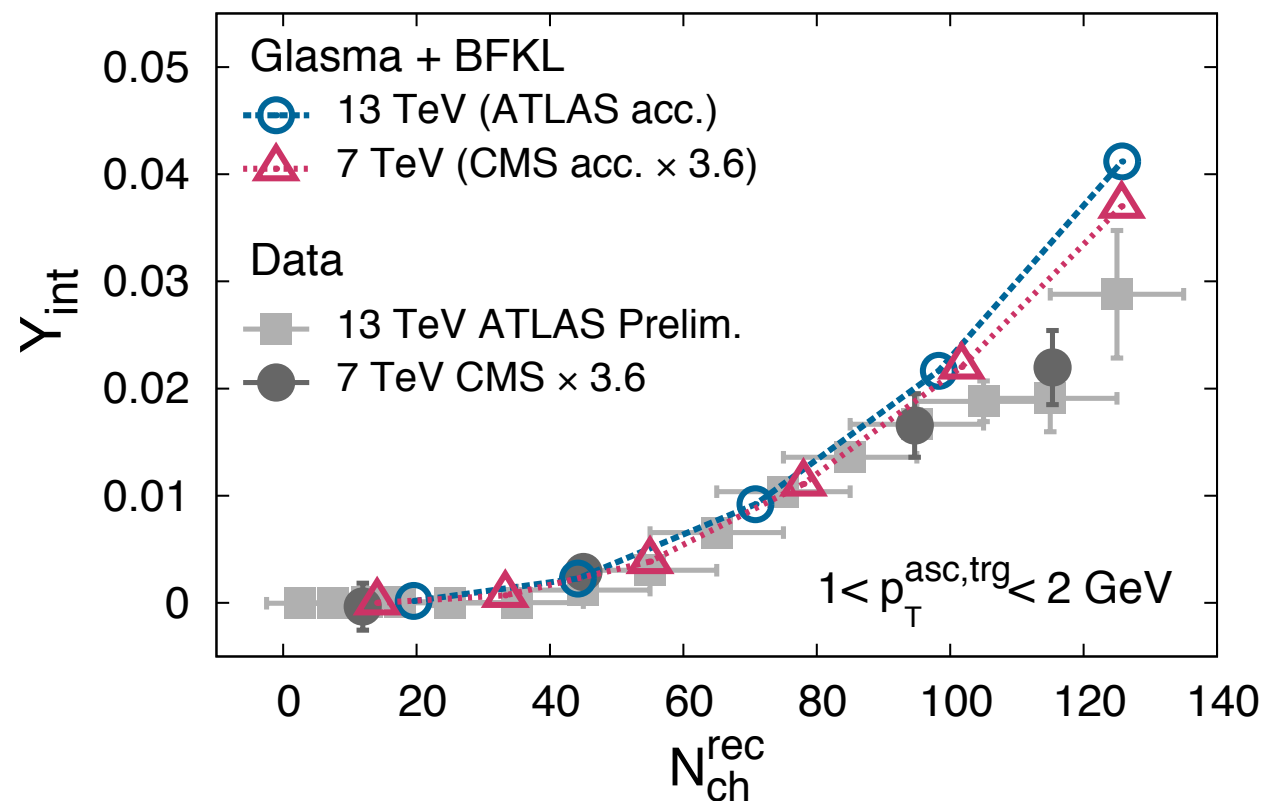
figure: Dusling, Li, Schenke 1509.07939

Dusling, Venugopalan 1201.2658, 1210.3890,
1211.3701, 1302.7018

Consistent explanation in the CGC picture

Energy independence of Ridge

Dusling, Tribedy, Venugopalan 1509.04410

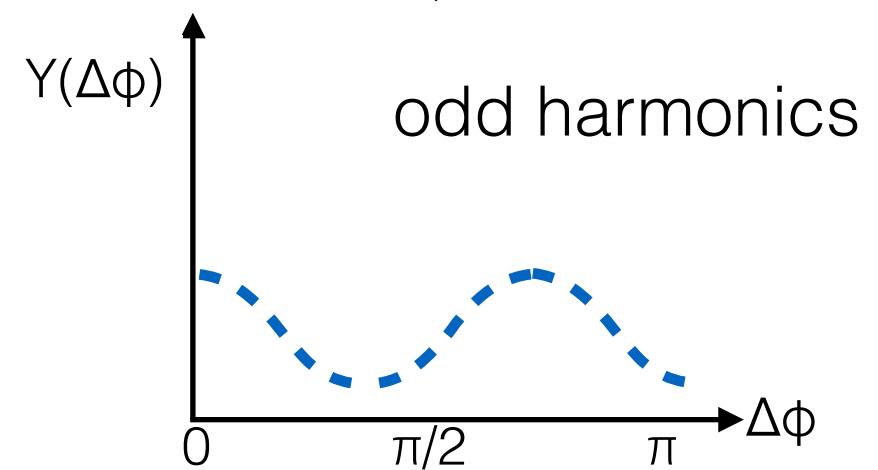
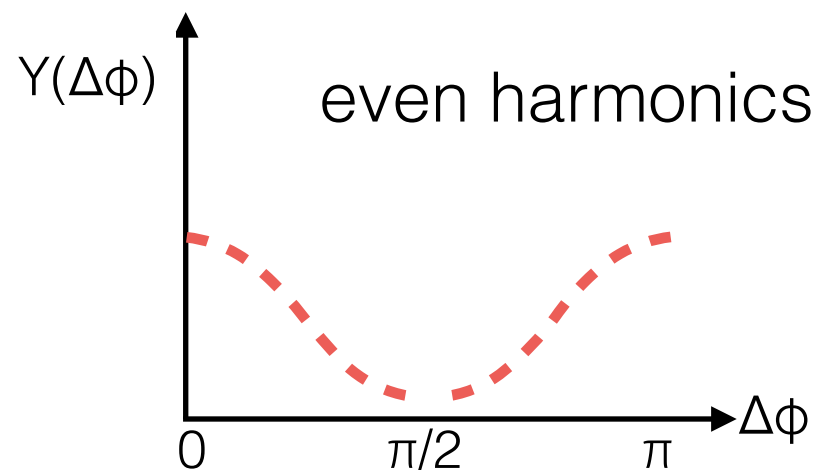
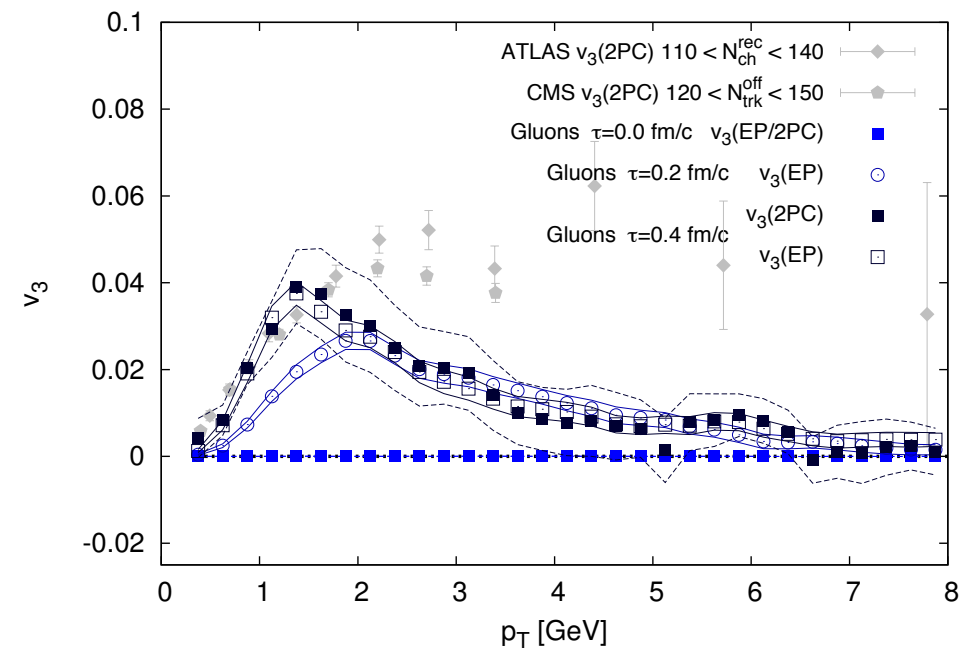
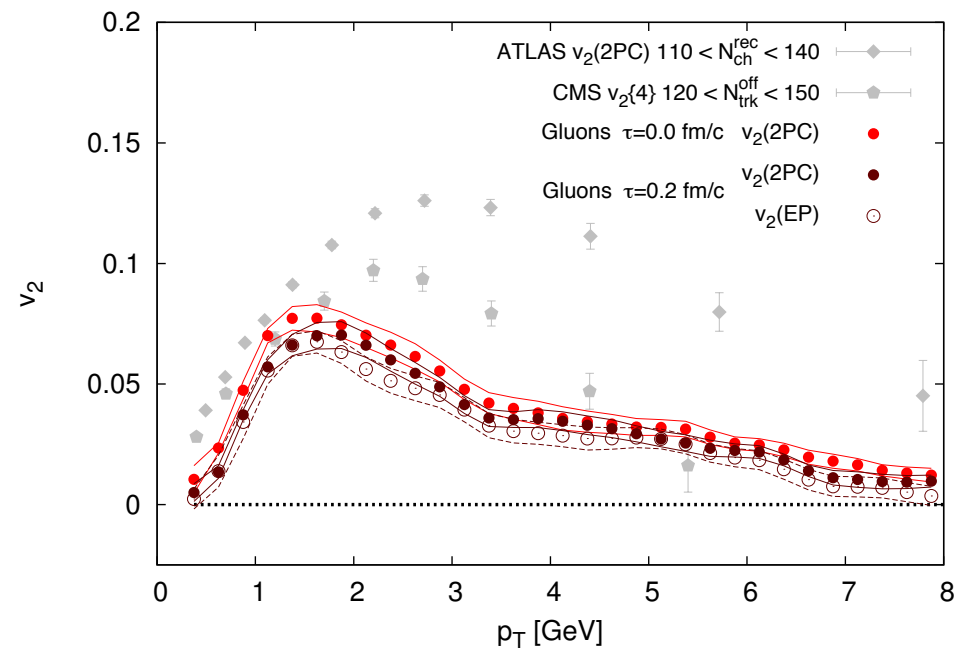
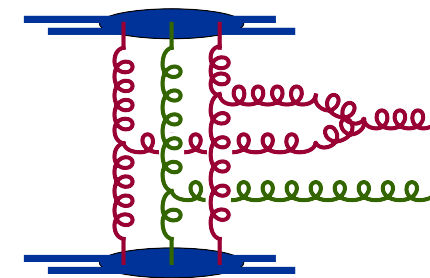
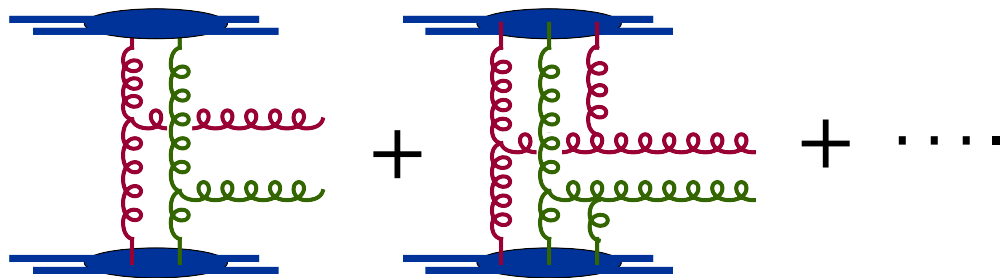


$$N_{\text{ch}}^{\text{rec}}(\sqrt{s}) = N_{\text{ch}}^{\text{rec}}(Q_s^2), \quad Y_{\text{int}}(\sqrt{s}) = Y_{\text{int}}(Q_s^2)$$

Energy dependence enters only through the saturation scale (only scale in problem) : Scaling of near side yield is natural in CGC approach

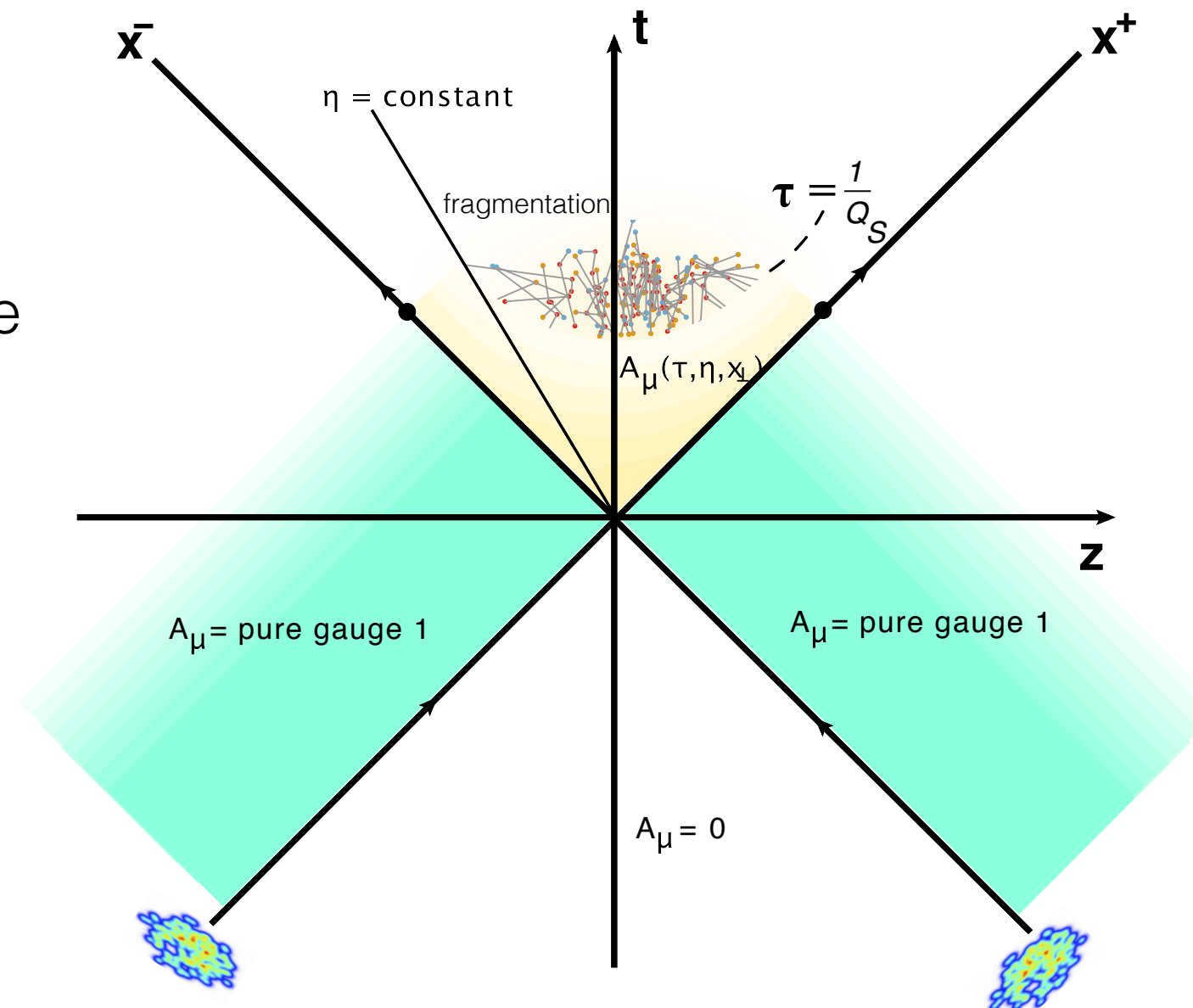
Harmonics of azimuthal correlations

Schenke, Schlichting, Venugopalan 1502.01331

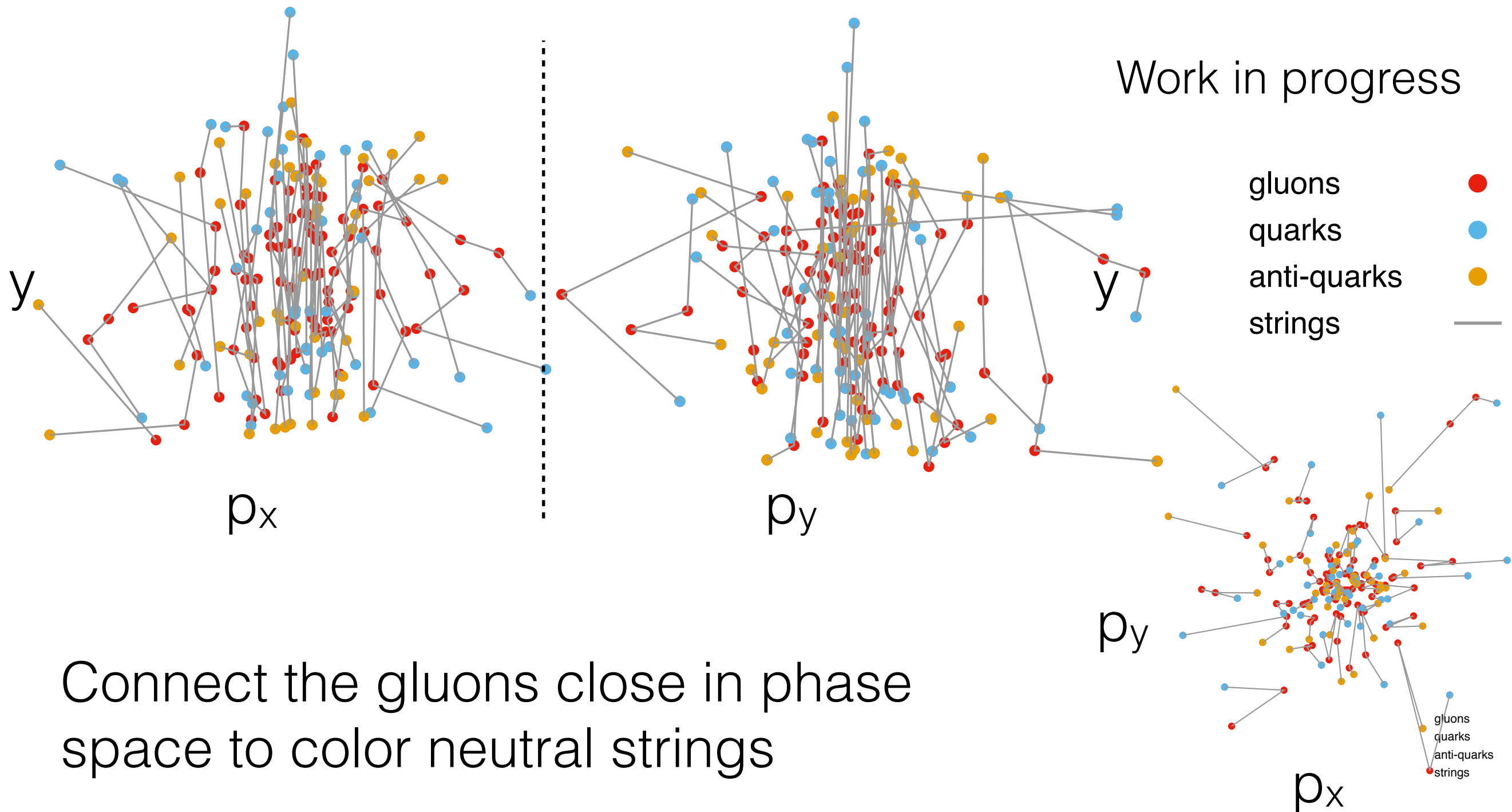


Combining CGC and PYTHIA

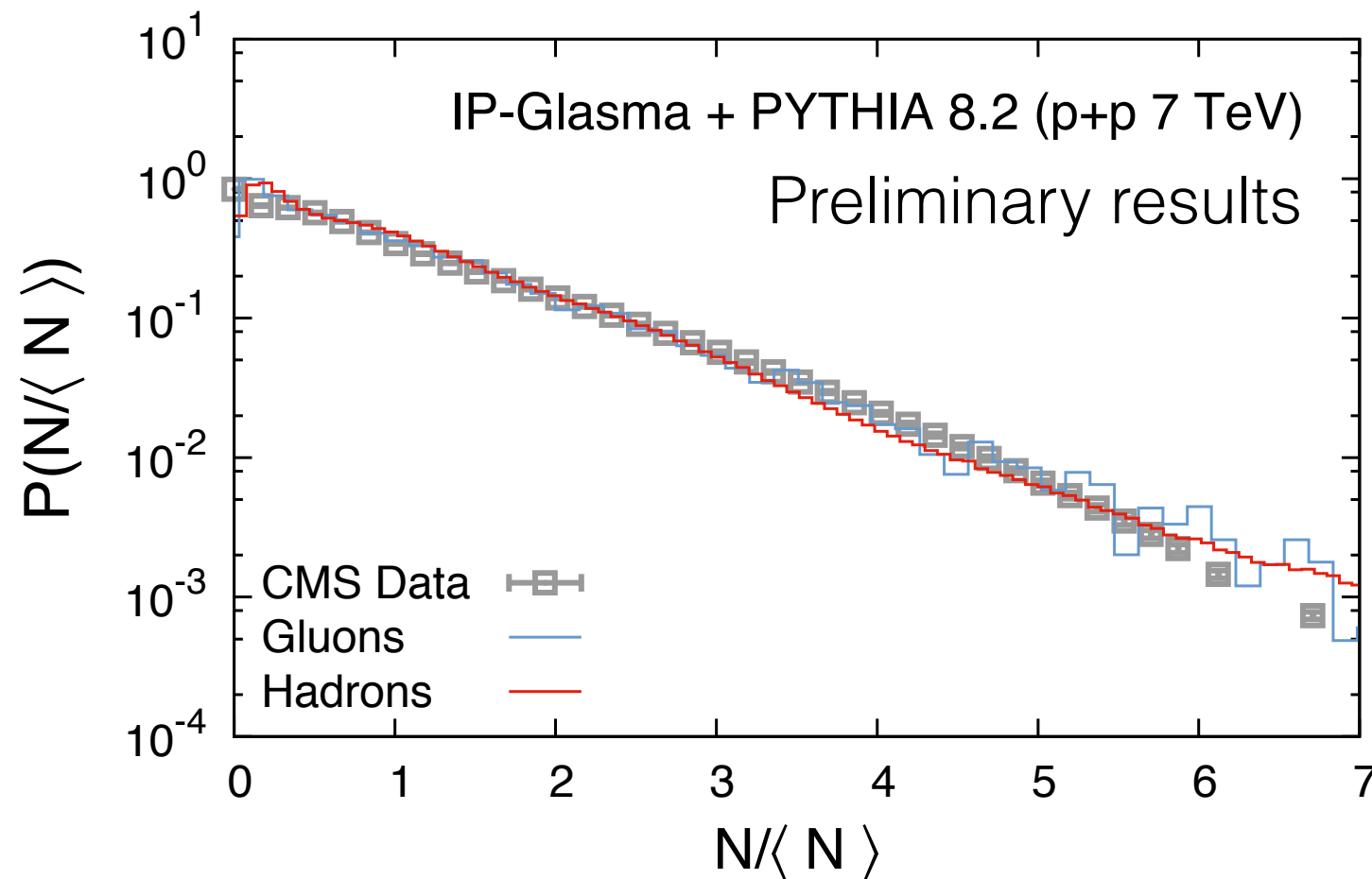
- Output distribution of Gluons from CGC
- Sample gluons in momentum space
- Connect the gluons close in phase space to color neutral strings
- Input to PYTHIA and fragment into final particles



Implementing PYTHIA Strings



Combining CGC and PYTHIA



- Promising results on multiplicity distributions
- Angular correlations, more observables are to be studied

Outlook

- Including diffractive process
- Full 3+1D with JIMWLK rapidity evolutions
- Implement color-reconnection in PYTHIA
- Comparison to more data from LHC

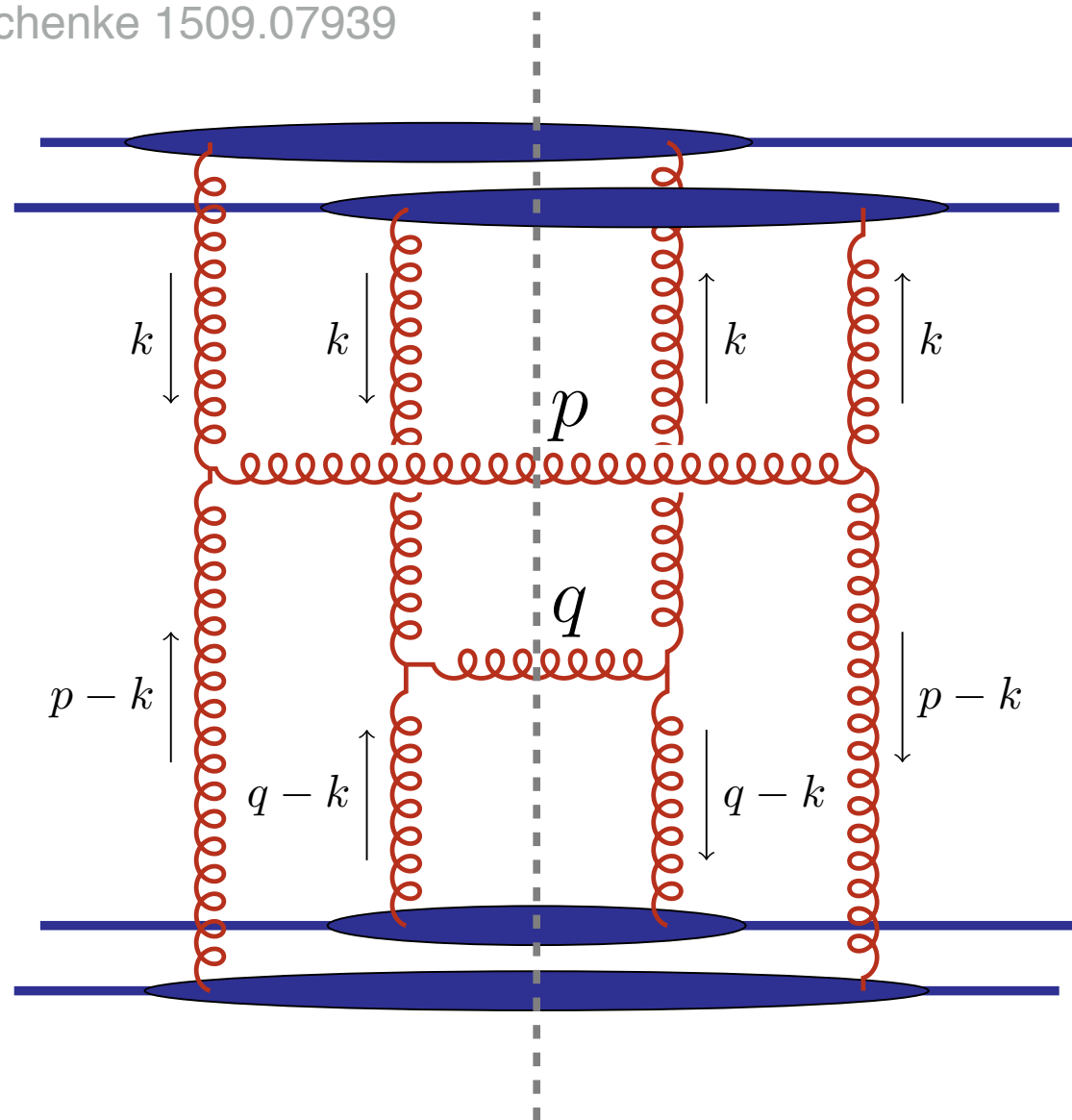
Summary

The *ab-initio* framework of CGC constrained by HERA DIS data provide successful description of the phenomena seen in high multiplicity events at LHC

back-up

Momentum flow in Glasma graph (origin of ridge-like correlation)

Dusling, Li, Schenke 1509.07939



Qualitative picture of correlations in small systems

Soeren Schlichting
Quark Matter 2015

