

- Light Waves and Matter Waves - Atom Interferometers and Optical Clocks

Guglielmo M. Tino

Dipartimento di Fisica & Astronomia e LENS – Università di Firenze

Istituto Nazionale di Fisica Nucleare, Sezione di Firenze

<http://coldatoms.lens.unifi.it/>

International Centre for Theoretical Physics

Winter College on Optics:

Light: a bridge between Earth and Space

Trieste, 9 - 20 February 2015

Outline of the Lectures

- *Lecture I: Atom Interferometry*
Light waves and matter waves, atom interferometry, methods, experiments on Earth and in space.
- *Lecture II: Optical Atomic Clocks*
Atomic clocks, basics, methods, optical clocks, experiments on Earth and in space.

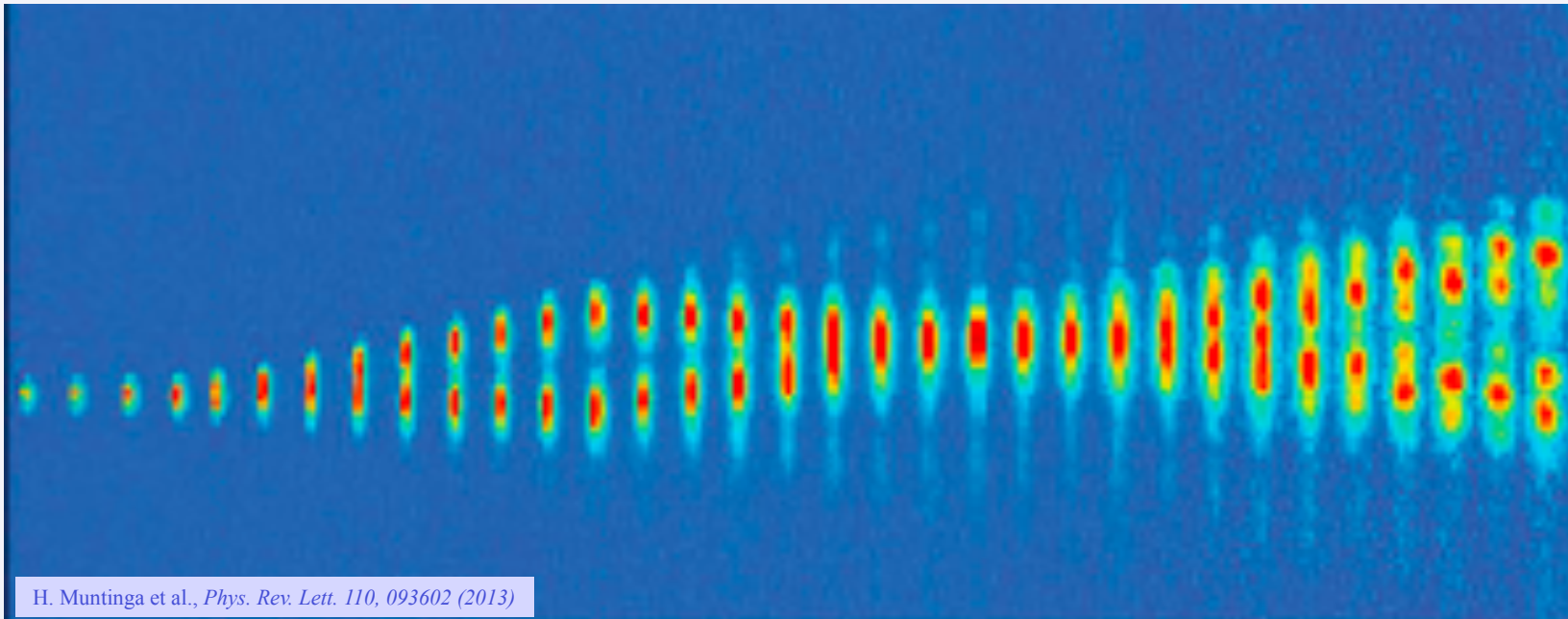
Lecture I: Atom Interferometry

- Introduction
- Basic concepts
- Experimental methods
- Atom interferometers
- Experiments on Earth and in space

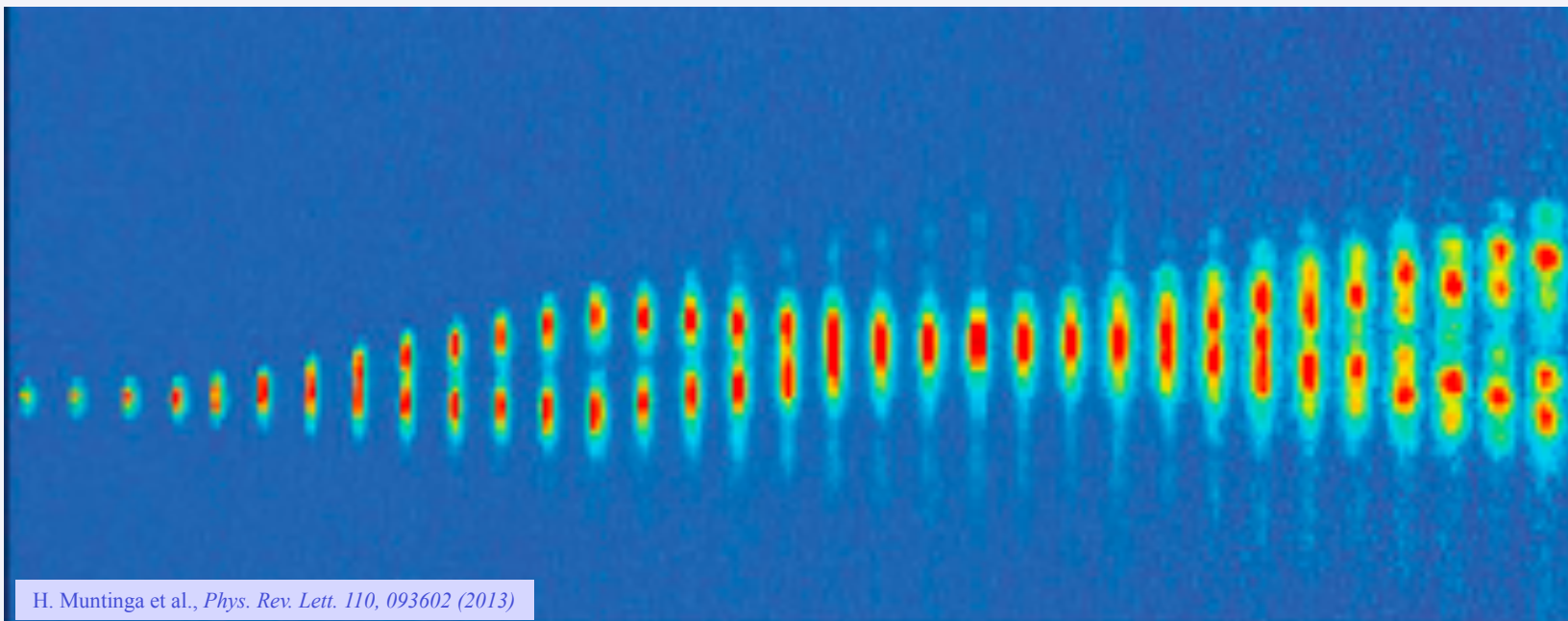
Main references

- A. D. Cronin, J. Schmiedmayer, D. E. Pritchard, *Optics and interferometry with atoms and molecules*, Rev. Mod. Phys. 81, 1051 (2009).
- J. Schmiedmayer, *Interferometry with atoms*, Lectures at the E. Fermi School on *Atom Interferometry*, Varenna (2013).
- G. M. Tino, M. A. Kasevich (eds). *Atom Interferometry*. Proc. International School of Physics ‘Enrico Fermi’, Course CLXXXVIII, Varenna 2013, SIF and IOS (2014).

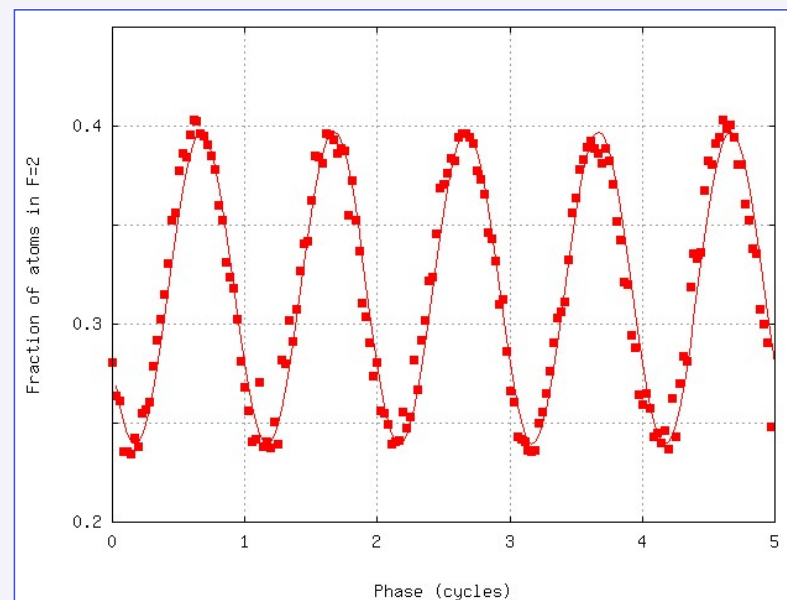
Atom Interferometry



Atom Interferometry

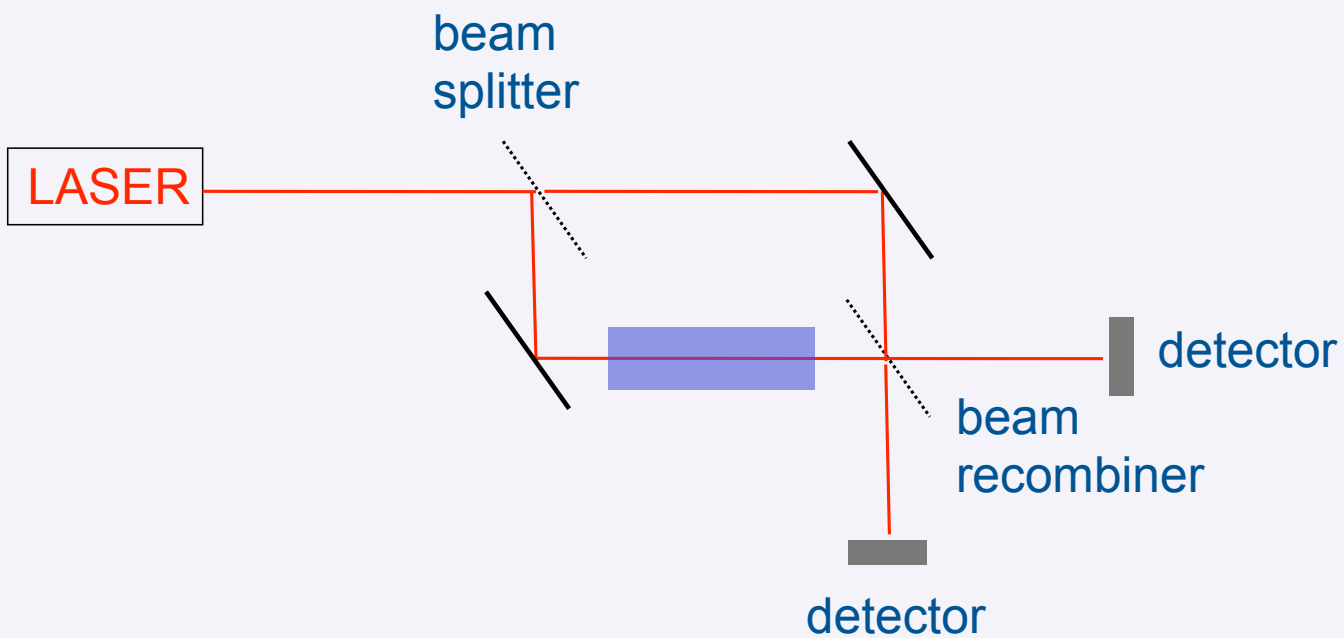


H. Muntinga et al., *Phys. Rev. Lett.* 110, 093602 (2013)

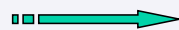


Interference fringes – Firenze 2006

Optical interferometry

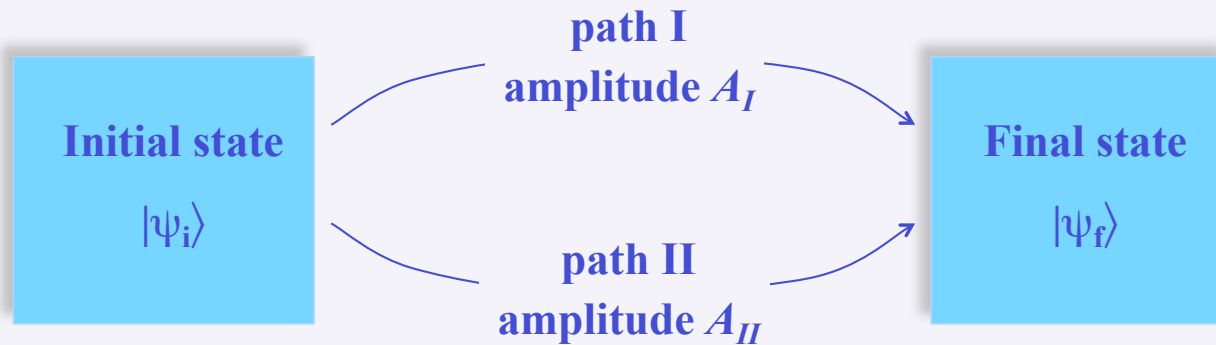


Change of optical path, length,
pressure, temperature,...



Change of phase
of interference pattern

Quantum interference



Interference of transition amplitudes

$$P(|\psi_i\rangle \Rightarrow |\psi_f\rangle) = |A_I + A_{II}|^2 = |A_I|^2 + |A_{II}|^2 + 2 \operatorname{Re}(A_I A_{II}^*)$$

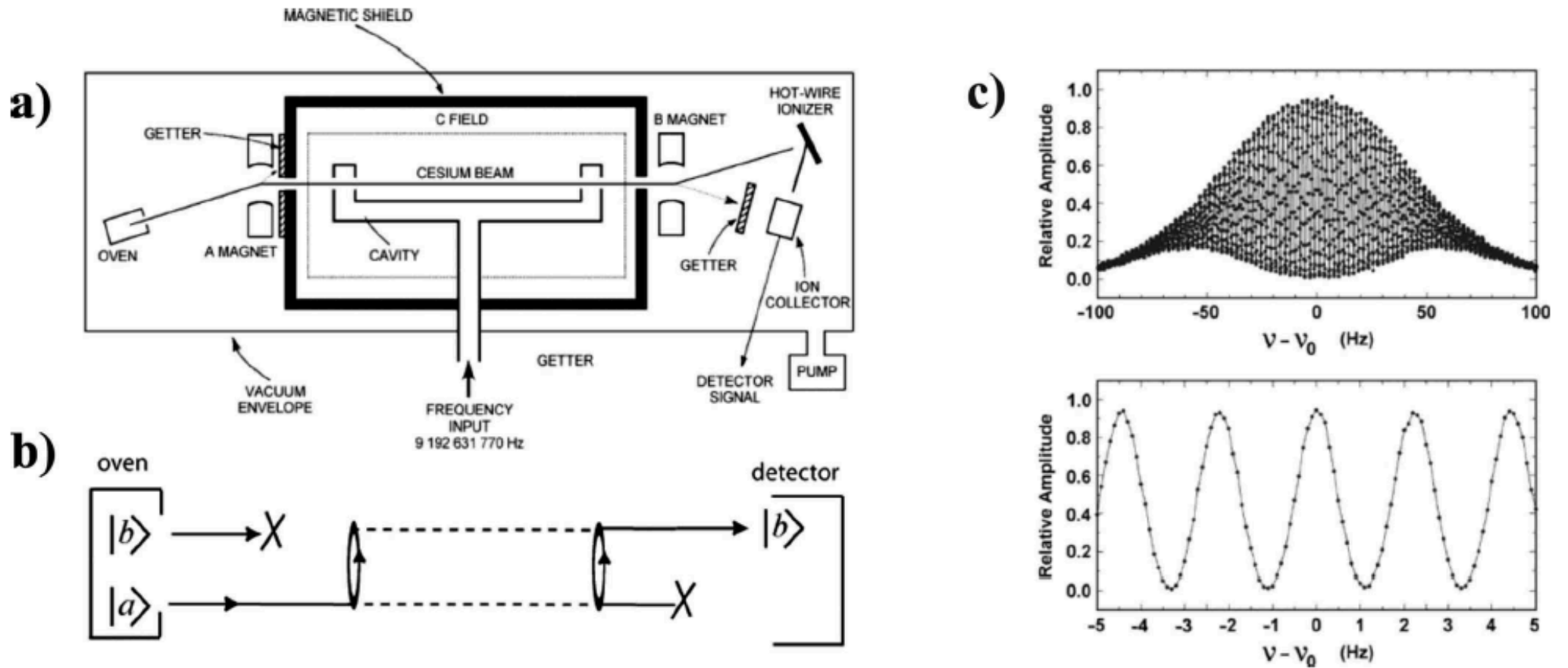


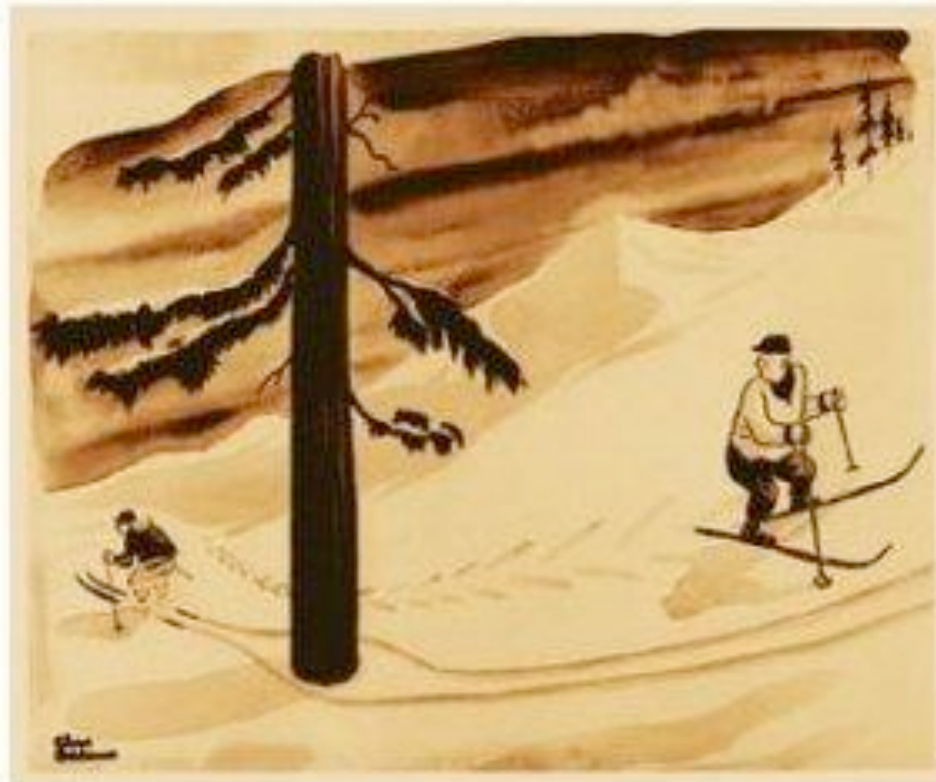
FIG. 1. Interferometry with internal quantum states of atoms. (a) Ramsey's separated oscillatory fields experiment. (b) The same experiment depicted as an interferometer for internal states. (c) The detected atom count rate exhibits interference fringes as a function of the applied rf frequency. These interference fringes, from the NIST-F1 fountain clock (Sullivan *et al.*, 2001), demonstrate the precision obtained with interference techniques. From Sullivan *et al.*, 2001.

Yakir Aharonov
Daniel Rohrlich

WILEY-VCH

Quantum Paradoxes

Quantum Theory for the Perplexed



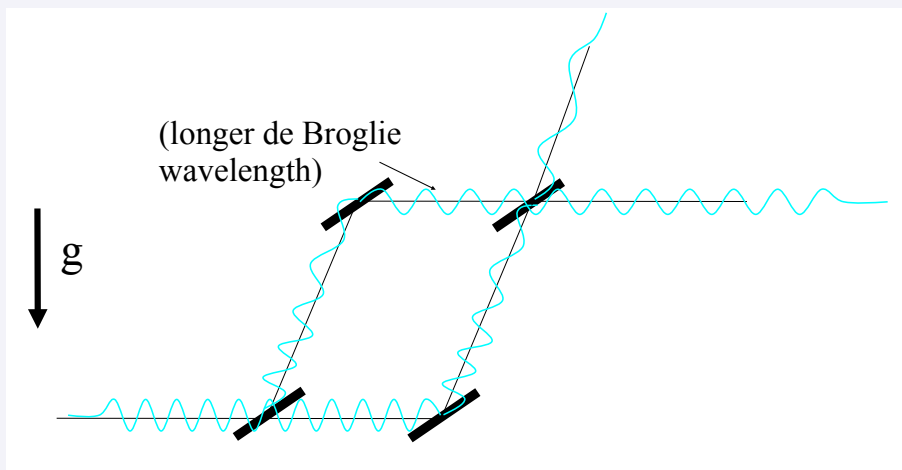
Atom interferometer force sensors

The quantum mechanical wave-like properties of atoms can be used to sense inertial forces.

$$\lambda_{DB} = \frac{h}{Mv}$$

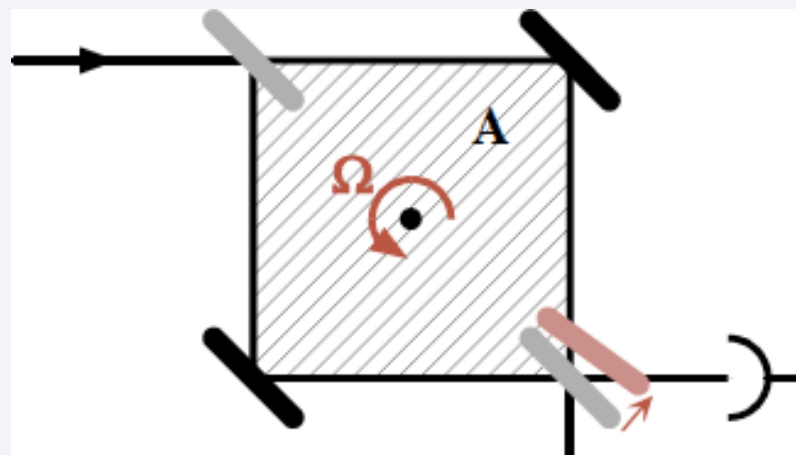
Gravity/Accelerations

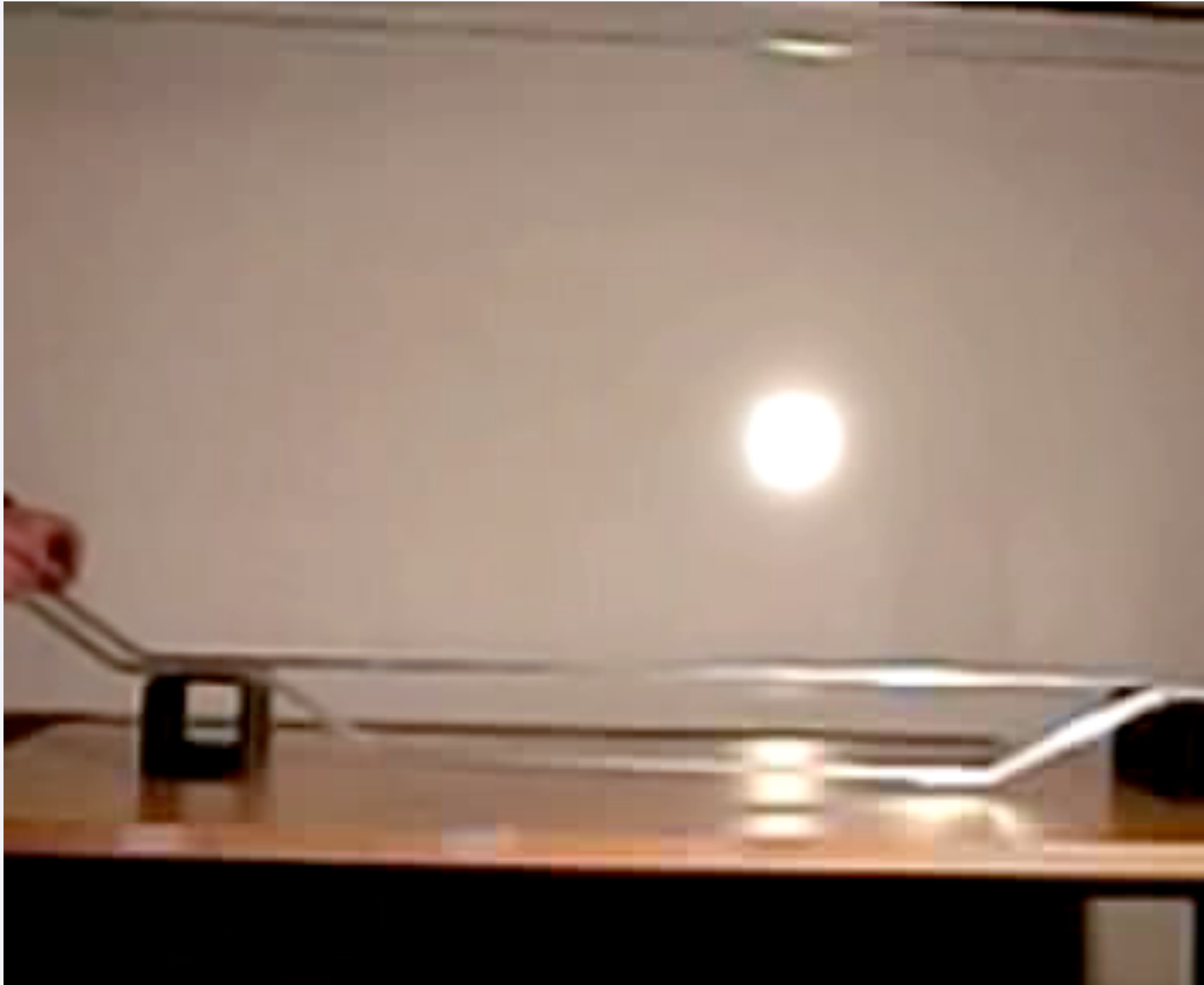
As atom climbs gravitational potential, velocity decreases and wavelength increases



Rotations

Rotations induce path length differences by shifting the positions of beam splitting optics

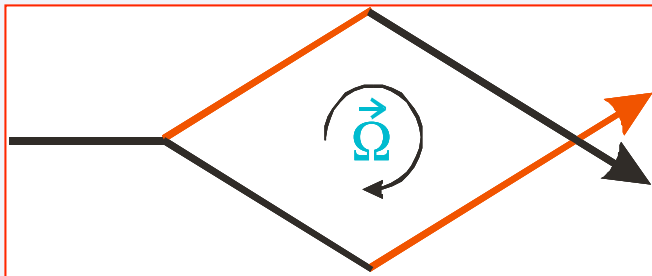




Courtesy of M. Gisselbrecht

Matter wave sensors

rotations:



$$\Delta\Phi_{\text{rot}} = 2\pi \frac{2 m_{\text{at}}}{h} \mathbf{A} \times \boldsymbol{\Omega}$$

$$\frac{\Delta\varphi_{\text{mat}}}{\Delta\varphi_{\text{ph}}} \sim \frac{m_{\text{at}} \times \lambda \times c}{h} \approx 5 \times 10^{10}$$

$\Delta\varphi$ effects

- Accelerations
- Rotations
- Laser frequency detuning
- Laser phase
- Photon recoil
- Electric/magnetic fields
- Interactions with atoms and molecules

Atomic vs optical interferometers

- The atoms move much more slowly in the interferometer compared to the photons. The interaction times are therefore much longer and the sensitivity can be much higher, for example to inertial effects (acceleration, rotation, gravitation).
- The actual size of optical interferometers can be much larger compared to atom interferometers. Examples are LIGO/Virgo interferometers and optical fiber gyroscopes.
- For the atoms, possibility to control the internal state and detection sensitive to internal state.
- For the atoms a larger variety of internal states compared to the photons which have only the two polarizations.
- Higher fragility of the atoms. In an atom interferometer it is important to avoid all processes of spontaneous emission or collisions that can destroy the atomic coherence.
- The flux of photons in a laser beam is typically much larger than the flux of atoms in a beam.

Matter Waves

Luis de-Broglie 1924:

Particles with rest-mass $m_0 > 0$ (electrons, neutrons, atoms, molecules) are elementary quanta of a wave field $\Psi(x,t)$

relativistic:

$$\hbar\omega = E = mc^2 = \gamma m_0 c^2$$

$$\hbar|\vec{k}| = \frac{h}{\lambda} = |\vec{p}| = mv = \gamma m_0 v$$

non - relativistic:

$$\lambda_{dB} = \frac{h}{mv} \quad \hbar k = mv$$

Optics with
Matter Waves:
Some numbers

Particle	Energy	Velocity	Wave length
Neutron	0.025 eV	2200 m/s	2.2 Å
Electron	100 eV	$6 \cdot 10^6$ m/s	1.2 Å
Na (atomic beam)	0.11 eV	1000 m/s	0.17 Å
Cs (laser cooled)	$7 \cdot 10^{-11}$ eV	1 cm/s	3000 Å

Wave Optics: comparison Light – Matter Waves

Light:
Maxwell equations

$$\left[\nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right] \vec{A}(r, t) = 0$$

Matter waves:
Schrödinger equation

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(r, t) \right] \Psi(r, t) = i\hbar \frac{\partial \Psi(r, t)}{\partial t}$$

Wave equation in
time independent formulation:

$$\left[\nabla^2 + k^2(r) \right] \Psi(r) = 0$$

Wave vector for matter waves:

$$k(r) = \frac{1}{\hbar} \sqrt{2m(E - V(r))}$$

$$k = \frac{2\pi}{\lambda}$$

Coherence length and thermal de Broglie wavelength

Coherence length $\xi \approx \frac{\hbar}{\Delta p}$

Thermal de Broglie wavelength $\lambda_T = \frac{h}{\sqrt{2\pi m k_B T}}$

Experimental Methods

Atomic Beams

Atom Sources

- Atoms enter the atomic beam machine from a (gaseous) reservoir through a small opening
- Velocity (wave length) of the atoms is given by the thermal distribution of the atoms passing through the opening.
 - Effusive sources
wide thermal velocity distribution
 - Supersonic expansion
cooling in the expansion leads to a narrow velocity distribution

$$\frac{\Delta v}{v} = \frac{\Delta \lambda}{\lambda} \approx 100\%$$

$$\frac{\Delta v}{v} = \frac{\Delta \lambda}{\lambda} \leq 10\%$$

Atom Beam

- For free propagation of atoms we need at least HV conditions: $p < 10^{-7}$ torr
- In the gravitational field of the earth propagation is in free fall
- Good collimation can be achieved with narrow slits (10 μ m slits separated 1m apart $\sim 10^{-5}$ rad)

Experimental Methods

Cold Atoms

Laser Cooling

Neutral atoms can be cooled by interacting with monochromatic light (~thermal equilibrium with the light)

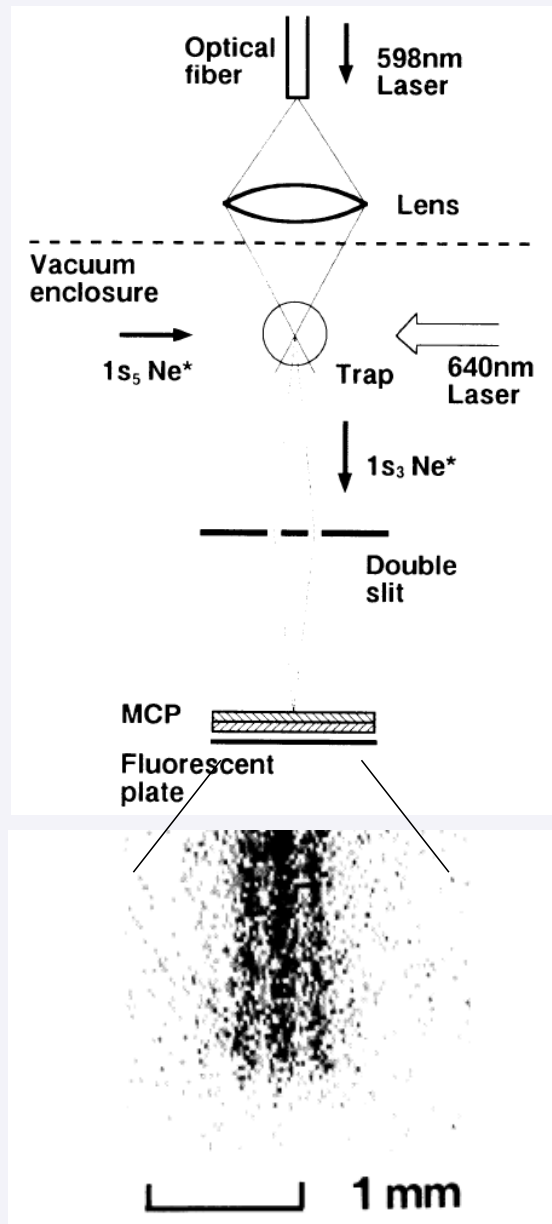
- Temperature $1\text{mK} \Leftrightarrow 1\mu\text{K}$
- Velocity $0.5\text{m/s} \Leftrightarrow 1\text{mm/s}$
- deBroglie wavelength $10\text{nm} \Leftrightarrow 500\text{nm}$
- Typical samples 10^8 atoms @ 10^{11} atoms/cm³

BEC

Cooling in a magnetic trap by removing the hottest atoms and thermal equilibration (evaporative cooling)

- Typical samples $>10^5$ atoms @ 10^{14} atoms/cm³
- Temperature $<1\mu\text{K}$
- deBroglie wavelength $>1\mu\text{m}$

Young's double slit interference with cold metastable atoms



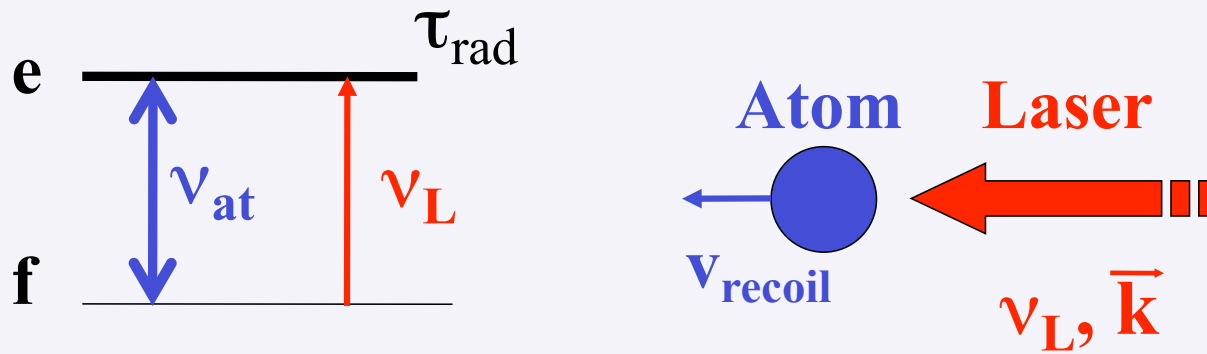
F. Shimizu et al.,
Phys. Rev. A 46, R17 (1992)

Cold atom source
 → large λ_{dB}
 → large fringe spacing

Metastable atoms
 → single atom detection
 → wave/particle duality for atoms

Laser cooling and manipulation of atoms

Radiation pressure



Momentum conservation $\rightarrow \vec{v}_{\text{recoil}} = \hbar \vec{k} / M$

Isotropic emission $\rightarrow \langle \vec{v}_{\text{em}} \rangle = 0$

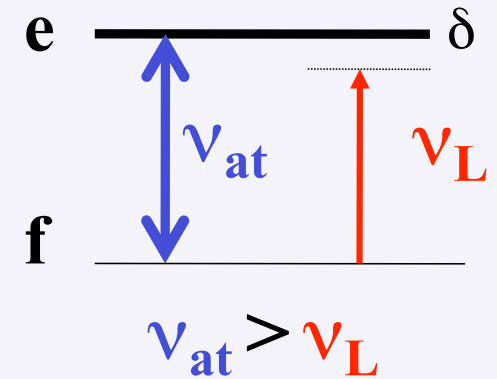
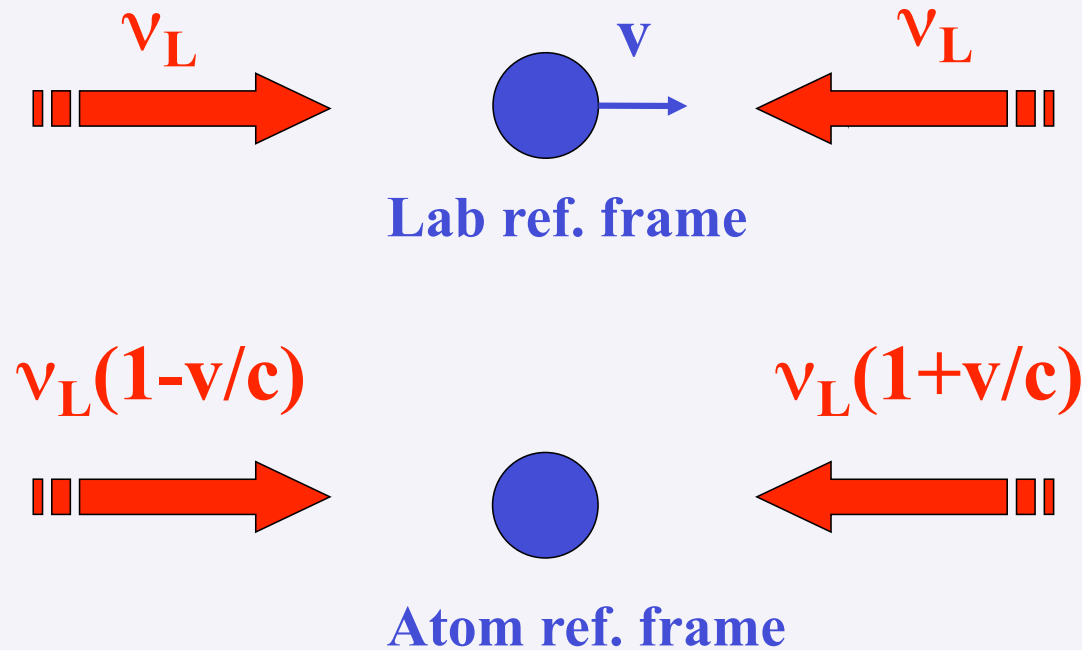
\rightarrow Radiation pressure force: $\vec{F} = \hbar \vec{k} / 2 \tau_{\text{rad}} = M \vec{a} \rightarrow a = \hbar k / 2 M \tau_{\text{rad}}$

Example, Na atom : $\lambda \approx 589 \text{ nm}$, $M = 23 \text{ a.m.u.}$, $\tau_{\text{rad}} \approx 16 \text{ ns}$

$\rightarrow v_{\text{recoil}} \approx 3 \text{ cm/s}$, $a \approx 10^6 \text{ m/s}^2 \approx 10^5 g$

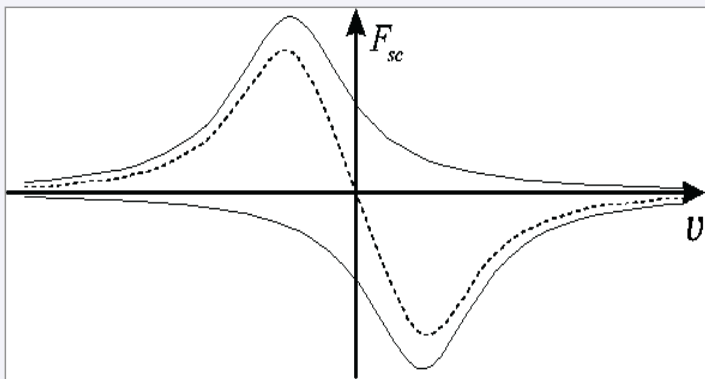
$t_{\text{stop}} = v_{\text{in}} / a \approx 1 \text{ ms}$, $L_{\text{stop}} = v_{\text{in}}^2 / 2a \approx 0.5 \text{ m}$

Optical molasses



$$(I/I_0 \ll 1)$$

$$F(\nu) \approx \frac{h\nu_L}{c} \times \frac{1}{2\pi} \times \left[\frac{I/I_0}{1 + I/I_0 + \frac{4}{\Gamma^2} (\delta - \frac{\nu_L}{c} \nu)^2} - \frac{I/I_0}{1 + I/I_0 + \frac{4}{\Gamma^2} (\delta + \frac{\nu_L}{c} \nu)^2} \right]$$



$$F(\nu) \approx \frac{h}{4\pi^2} \frac{\omega_L^2 8\delta}{c^2 \Gamma} \frac{I/I_0}{[1 + (\frac{2\delta}{\Gamma})^2]^2} \nu = -\alpha \nu$$

Idea: T.W. Hänsch, A. Schawlow, 1975

Exp. demonstration: S. Chu et al., 1985

Laser cooling: atomic temperatures

Atomic Temperature : $k_B T = M v_{\text{rms}}^2$

Minimum temperature for Doppler cooling:

$$k_B T_D = \frac{h\Gamma}{2}$$

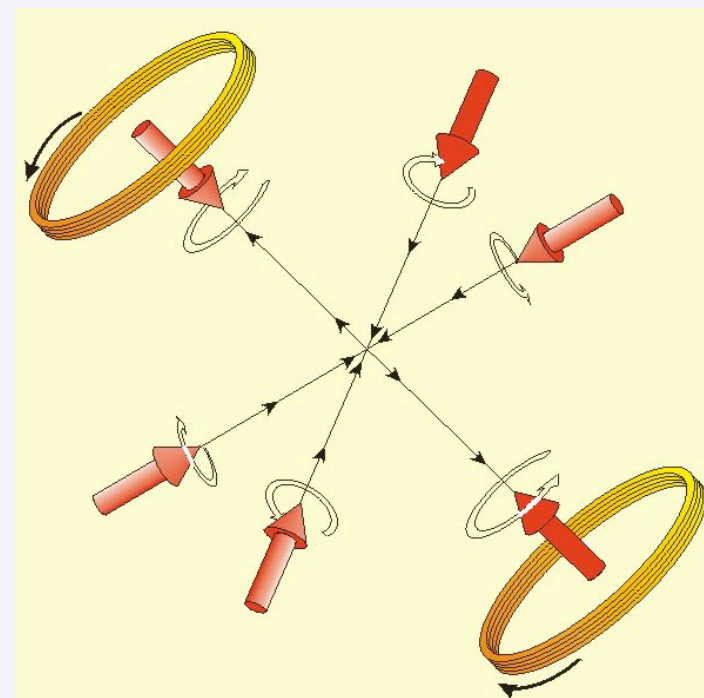
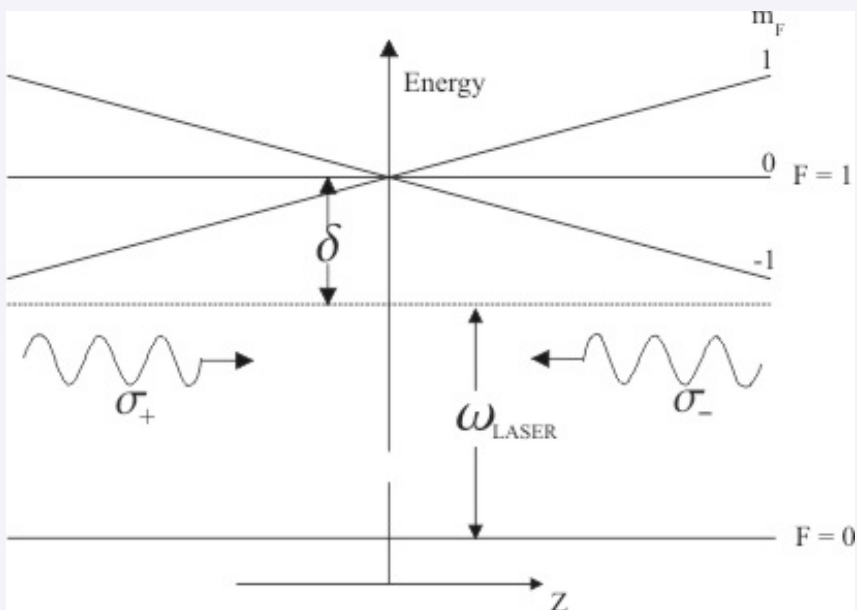
Single photon recoil temperature:

$$k_B T_r = \frac{1}{M} \left(\frac{h\nu_L}{c} \right)^2$$

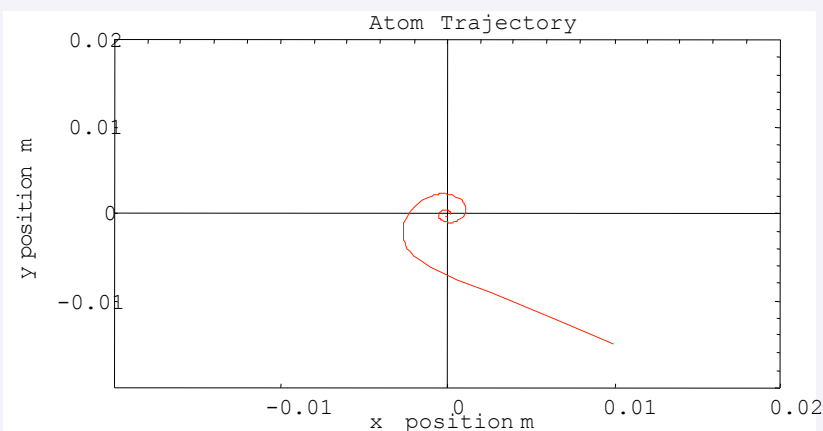
Examples:

	T_D	T_r
Na	240 μK	2.4 μK
Rb	120 μK	360 nK
Cs	120 μK	200 nK
Sr (intercombination transition)	180 nK	460 nK

Magneto-Optical Trap (MOT)



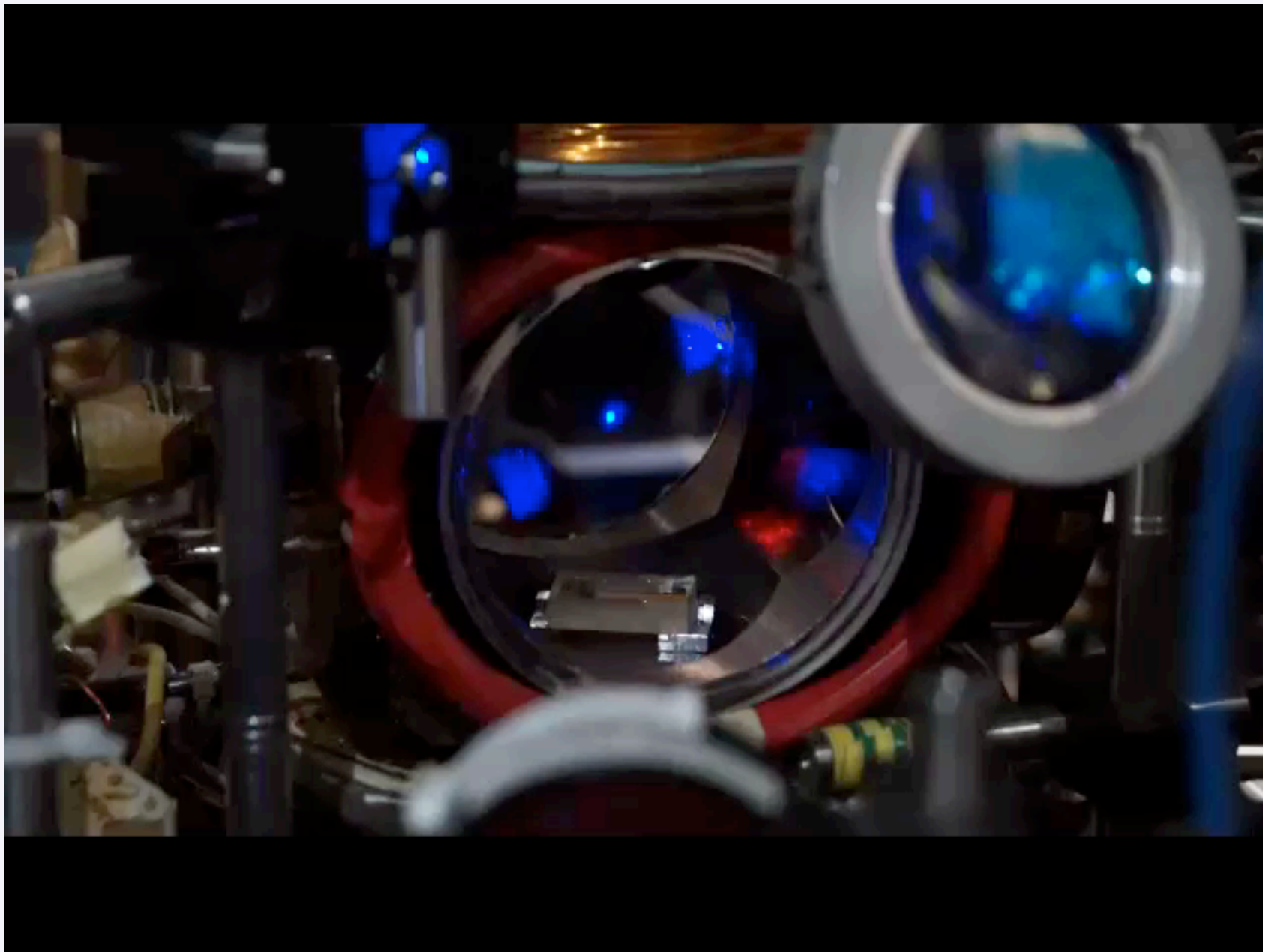
$$F(z, v) \approx \frac{4hk I \delta}{\pi I_0 \Gamma} \frac{kv + \beta z}{[1 + (\frac{2\delta}{\Gamma})^2]^2}$$



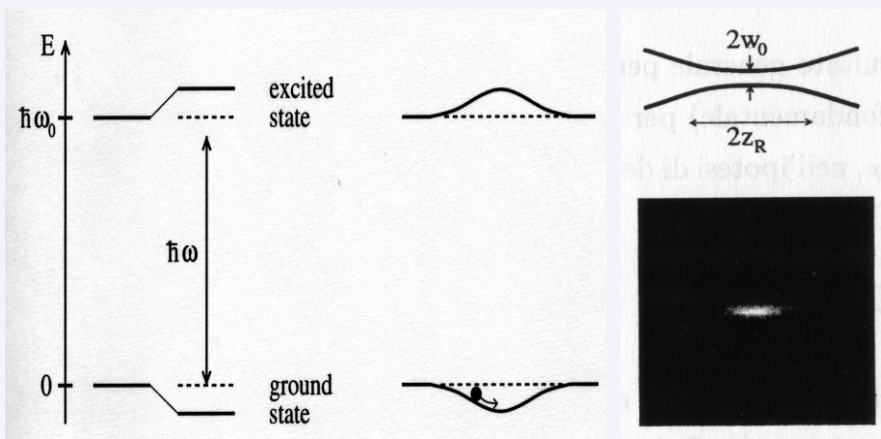
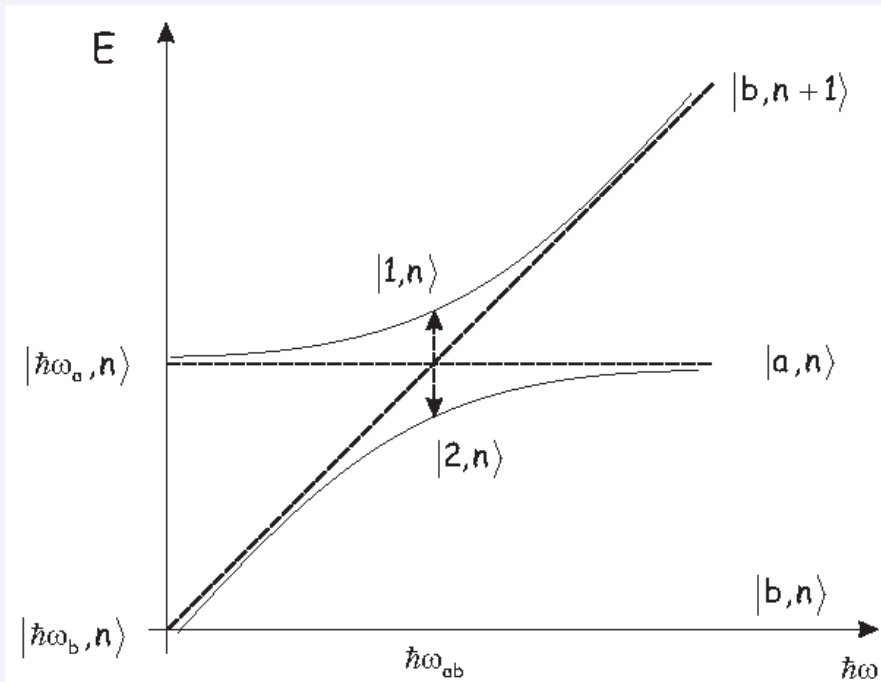
density $n \approx 10^{11} \text{ cm}^{-3}$
 temperature $T \approx 100 \mu\text{K}$
 size $\Delta x \approx 1 \text{ mm}$

E. Raab *et al.*, Phys. Rev. Lett. **59**, 2631 (1987)

Sr Magneto-Optical Trap (MOT) LENS - Firenze

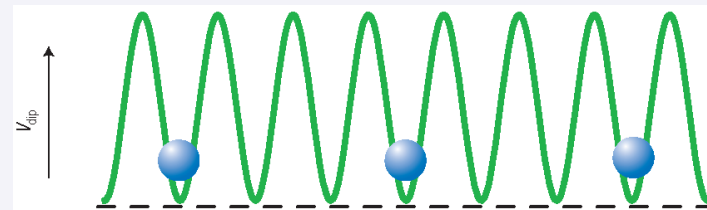


Light shifts and optical traps

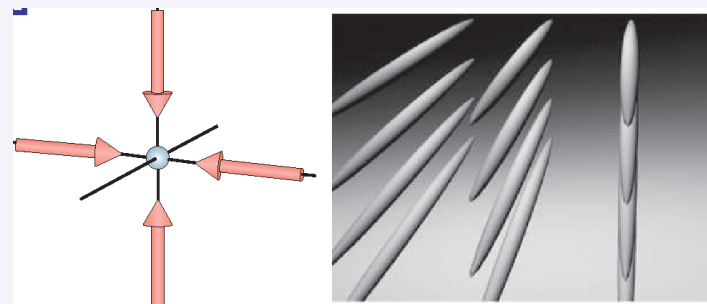


$$V_{\text{dip}}(\mathbf{r}) = -\mathbf{d} \cdot \mathbf{E}(\mathbf{r}) \propto \alpha(\omega_L) |\mathbf{E}(\mathbf{r})|^2$$

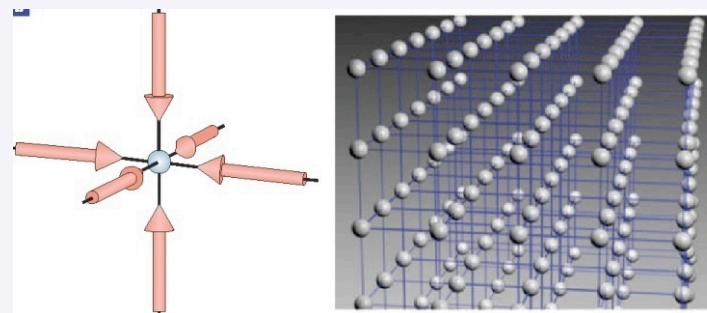
optical lattices



1D optical lattice \Rightarrow array of 2D disk-like trapping potentials



2 D optical lattice \Rightarrow array of 1D potential tubes



3 D optical lattice \Rightarrow 3D simple cubic array of h.o. potentials

First exp. demonstration: S. Chu et al., 1986

Review: I. Bloch, 2005

Cooling and trapping atoms with laser light

The Nobel Prize in Physics 1997



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Web-Adapted Version of the Nobel Poster from the Royal Swedish Academy of Sciences

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The Nobel Prize in Physics 1997

▶▶

The Royal Swedish Academy of Sciences has awarded the 1997 Nobel Prize in Physics jointly to

Steven Chu, Claude Cohen-Tannoudji and William D. Phillips

for their developments of methods to cool and trap atoms with laser light.



Steven Chu
Stanford University, Stanford, California, USA



Claude Cohen-Tannoudji
Collège de France and École Normale Supérieure, Paris, France

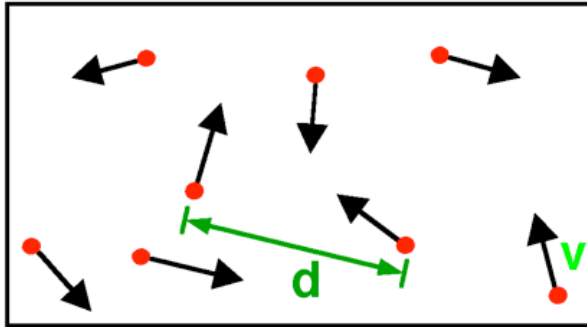


William D. Phillips
National Institute of Standards and Technology, Gaithersburg, Maryland, USA

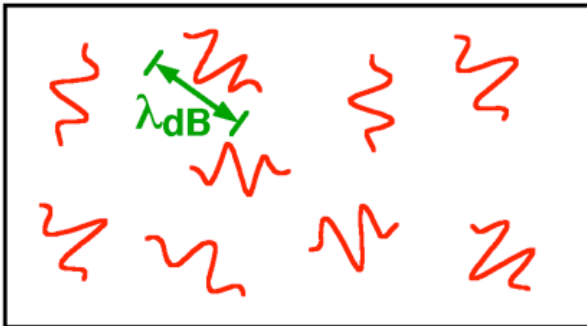


This year's Nobel laureates in physics have developed methods of cooling and trapping atoms by using laser light. Their research is helping us to study fundamental phenomena and measure important physical quantities with unprecedented precision.

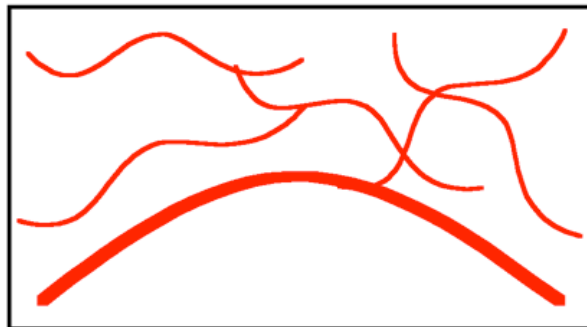
What is Bose-Einstein condensation (BEC)?



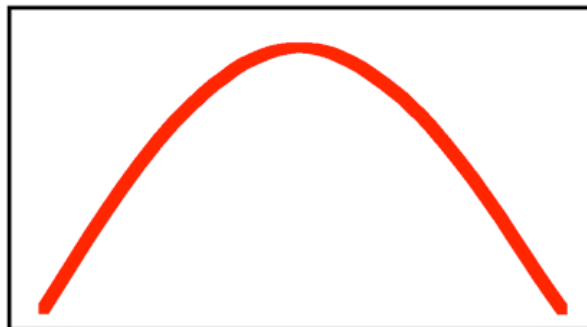
**High
Temperature T:**
thermal velocity v
density d^{-3}
"Billiard balls"



**Low
Temperature T:**
De Broglie wavelength
 $\lambda_{dB} = h/mv \propto T^{-1/2}$
"Wave packets"



$T = T_{crit}$:
**Bose-Einstein
Condensation**
 $\lambda_{dB} \approx d$
"Matter wave overlap"



$T = 0$:
**Pure Bose
condensate**
"Giant matter wave"

from W. Ketterle

Bose-Einstein condensation

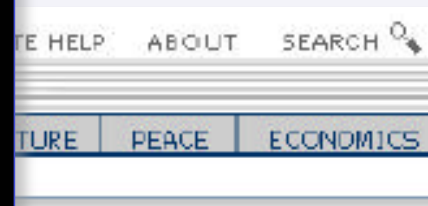
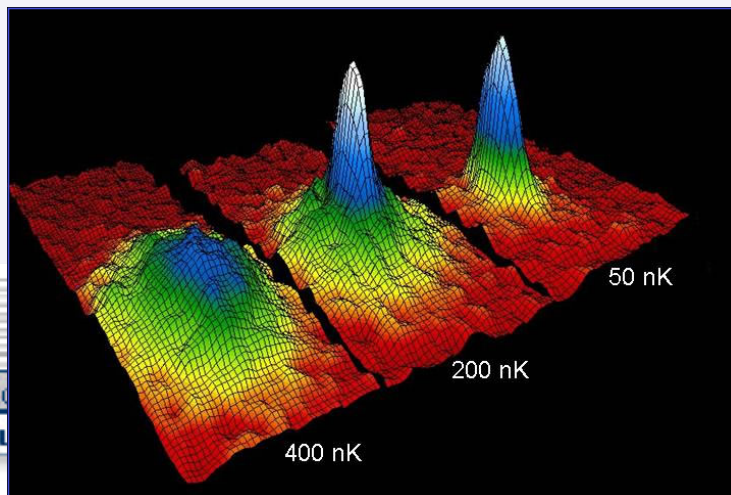
The atoms with an even number of electrons + protons + neutrons at very low temperatures occupy all the ground state of the system.

This new state of matter is called *Bose-Einstein condensate*.

The atoms are called *bosons*.

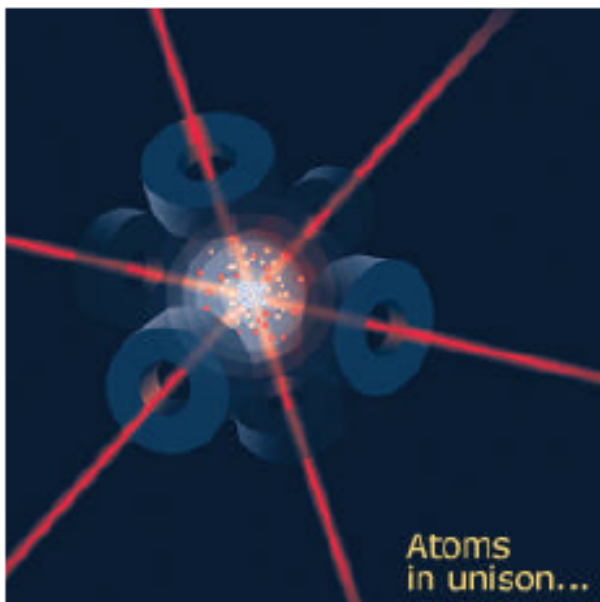
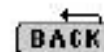


A. Einstein and S.N. Bose (1925)



The Nobel Prize in Physics 2001

The Royal Swedish Academy of Sciences has awarded the Nobel Prize in Physics for 2001 jointly to Eric A. Cornell, Wolfgang Ketterle and Carl E. Wieman "for the achievement of Bose-Einstein condensation in dilute gases of alkali atoms, and for early fundamental studies of the properties of the condensates".



Eric A. Cornell
JILA and National Institute of Standards and Technology (NIST), Boulder, Colorado, USA.

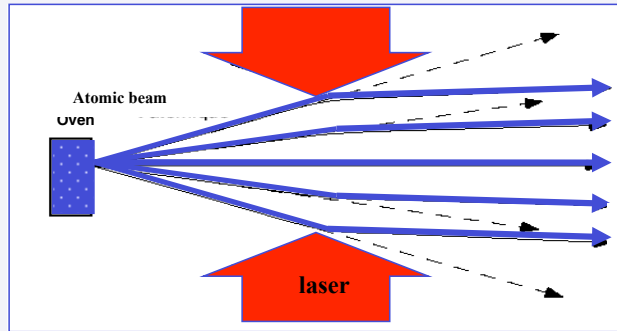
Carl E. Wieman
JILA and University of Colorado, Boulder, Colorado, USA.

Wolfgang Ketterle
Massachusetts Institute of Technology (MIT), Cambridge, Massachusetts, USA.

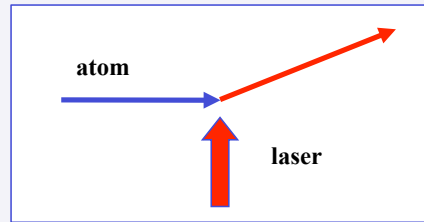
Contents:

Atom optics

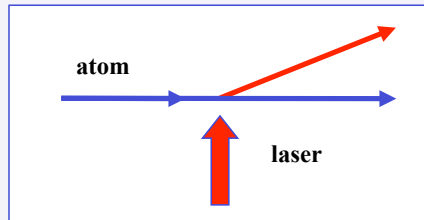
lenses



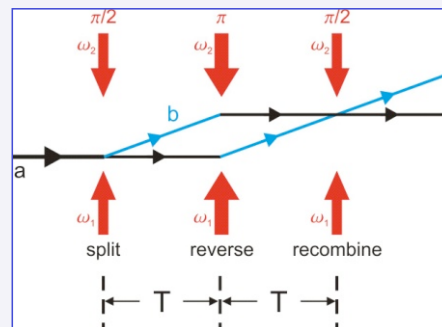
mirrors



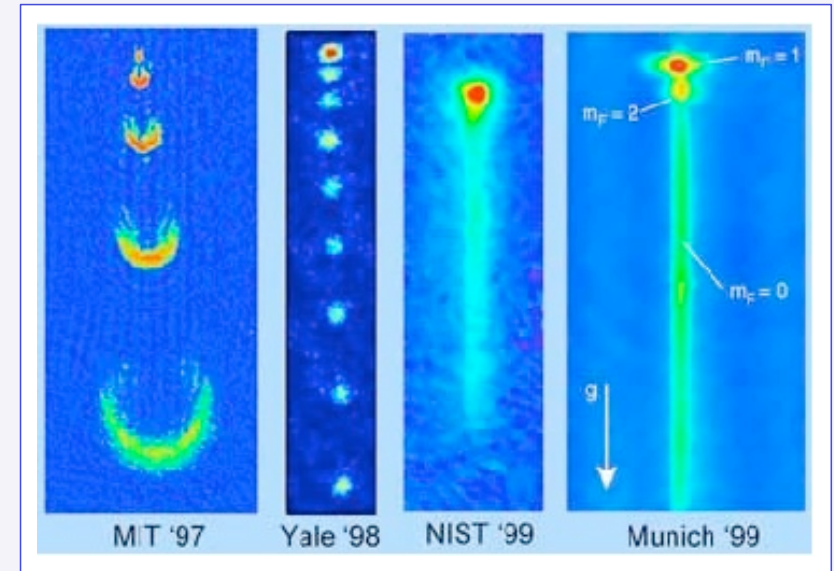
beam-splitters



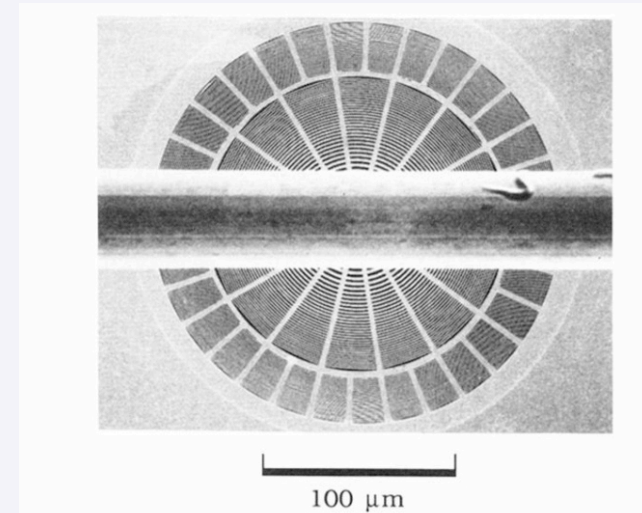
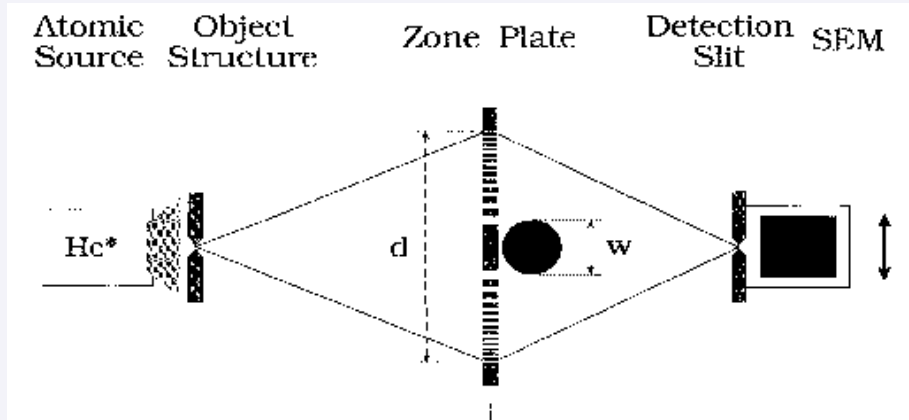
interferometers



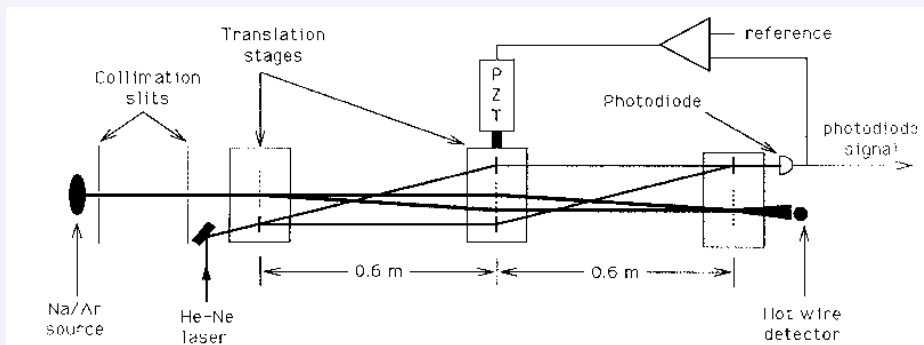
atom laser



Microfabricated atom optics



O. Carnal, M. Sigel, T. Sleator, H. Takuma, J. Mlynek, *Phys. Rev. Lett.* 67, 3231 (1991)



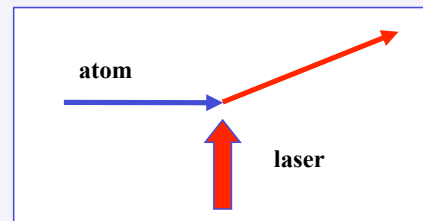
D. W. Keith, C. R. Ekstrom, Q. A. Turchette, D. E. Pritchard, *Phys. Rev. Lett.* 66, 2693 (1991)

Physics and applications of ultracold atoms

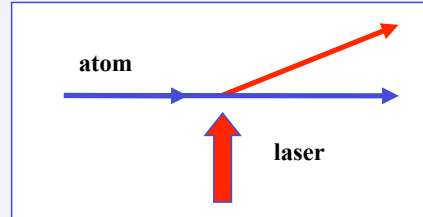
- BEC, degenerate Fermi gases, collective quantum effects
- Ultracold interactions and collision dynamics
- Ultracold molecules
- Surface physics and quantum reflection
- Entanglement and quantum information
- Precision spectroscopy
- Ultrasensitive isotope trace analysis
- Atom optics and atom laser
- Atom lithography
- **Atomic clocks**
- **Atom interferometers**

Optical pulse atom optics

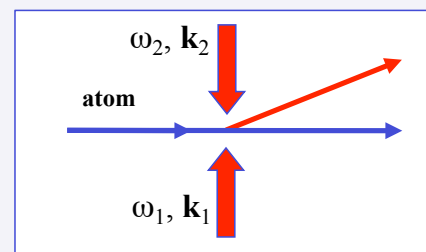
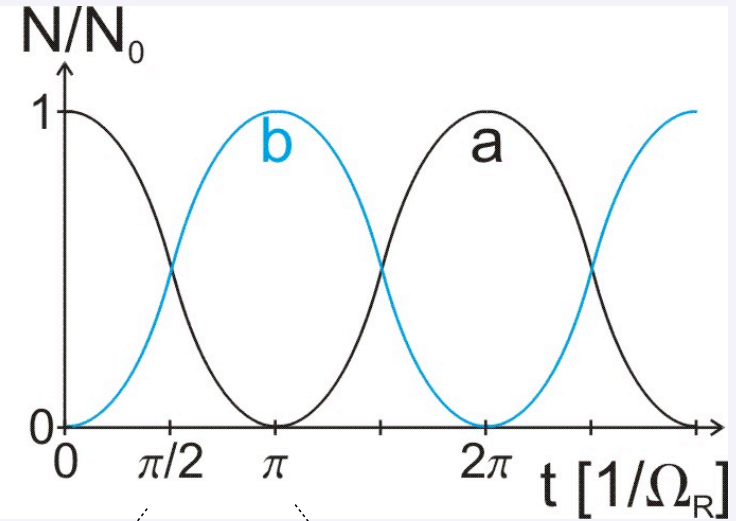
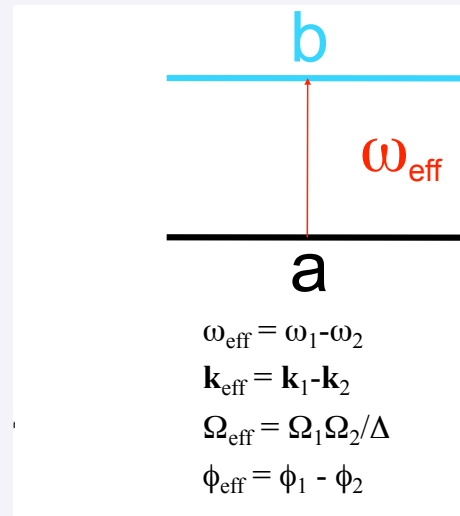
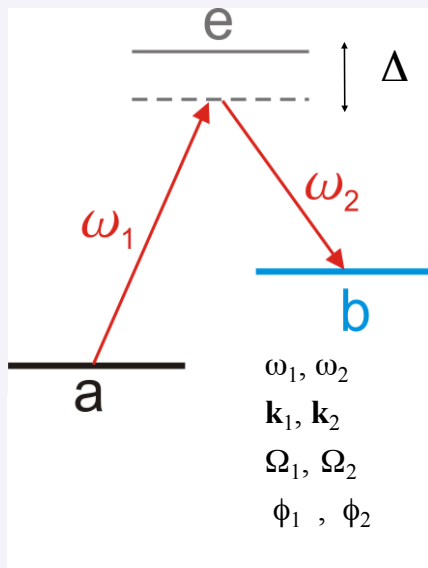
mirror



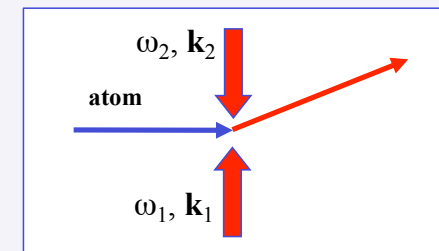
beam-splitter



Raman pulses

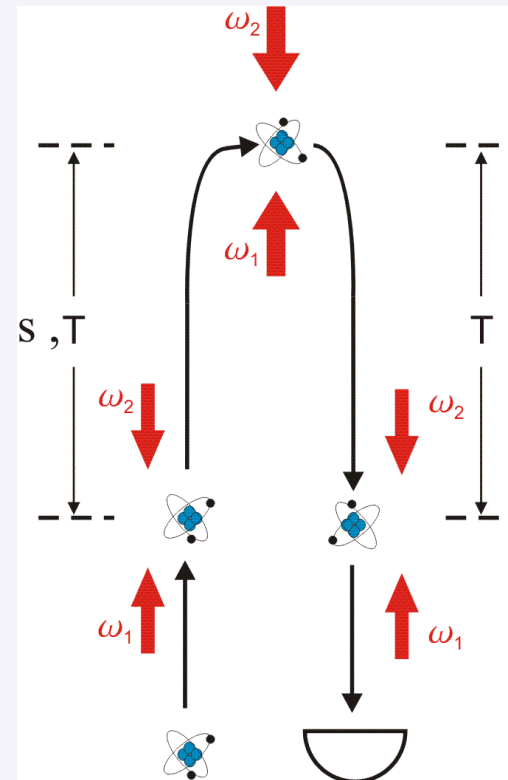
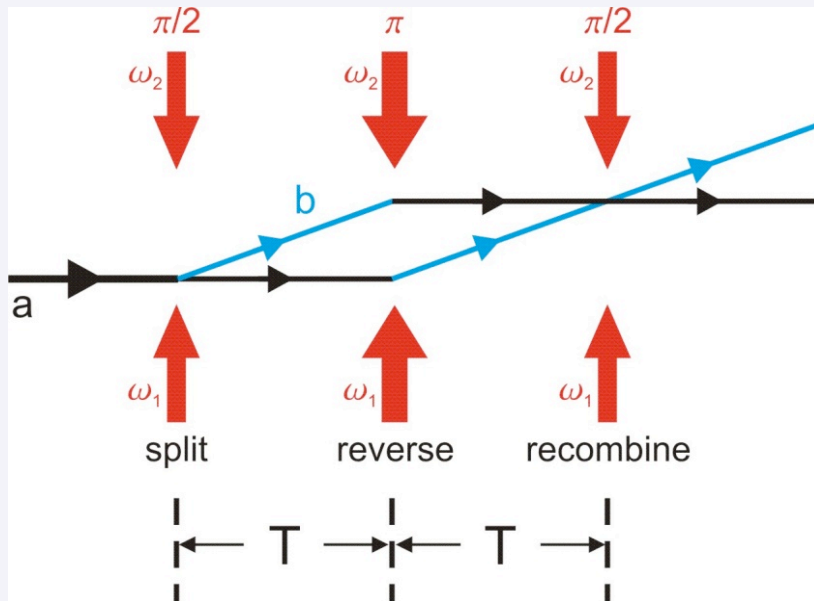
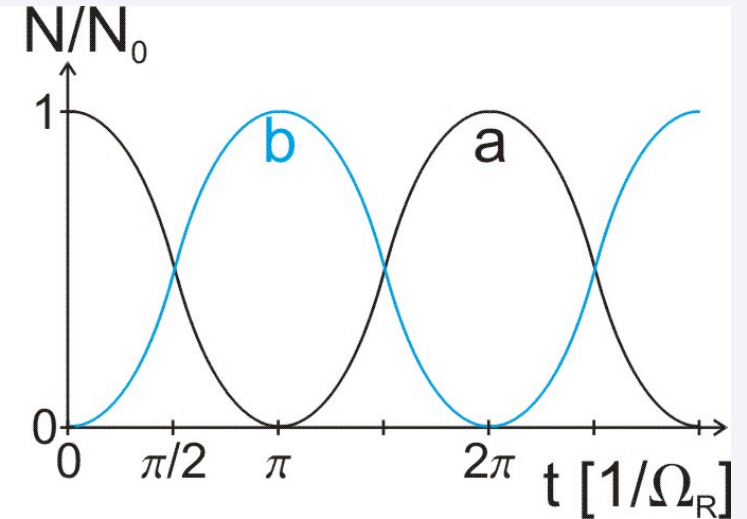
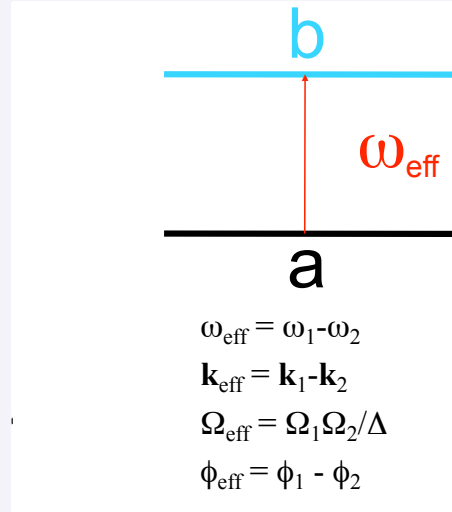
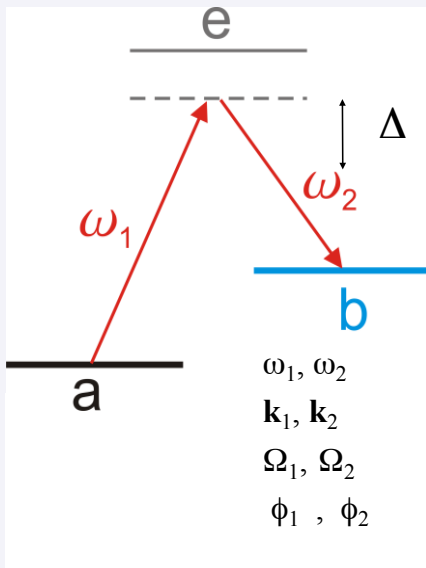


beam-splitter



mirror

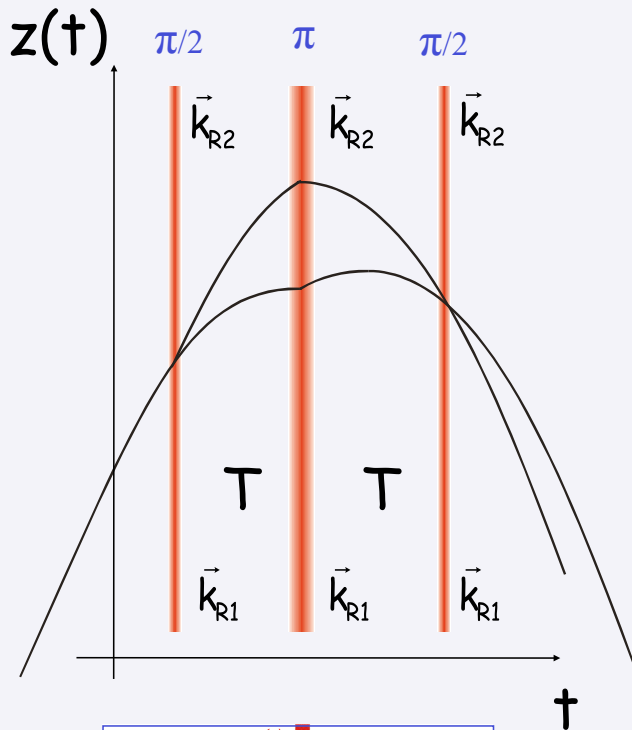
Raman pulse interferometer



M. Kasevich, S. Chu, *Appl. Phys. B* **54**, 321 (1992)

A. Peters, K.Y. Chung and S. Chu, *Nature* **400**, 849 (1999)

Raman interferometry in an atomic fountain



Phase difference between the paths:

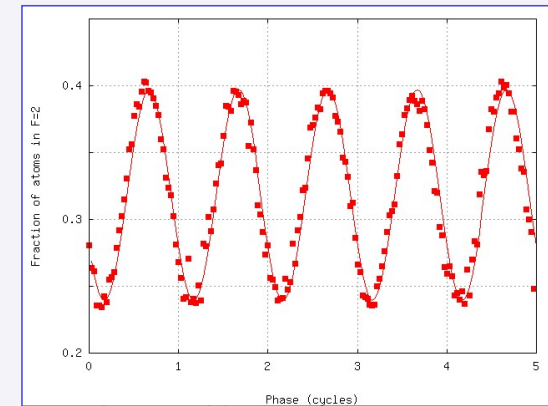
$$\Delta\Phi = k_e[z(0) - 2z(T) + z(2T)] + \Phi_e \quad k_e = k_1 - k_2, \quad \omega_e = c k_e$$

with $z(t) = -g t^2/2 + v_0 t + z_0$ & $\Phi_e = 0 \Rightarrow \Delta\Phi = k_e g T^2$

$$g = \Delta\Phi / k_e T^2$$

Final population:

$$N_a = N/2 (1 + \cos[\Delta\Phi])$$

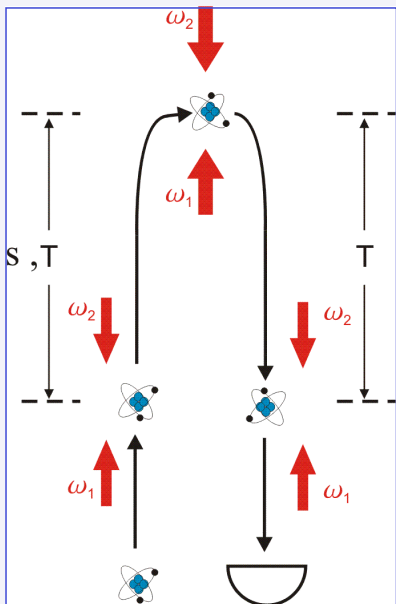


Interference fringes – Firenze 2006

$$T = 150 \text{ ms} \Rightarrow 2\pi = 10^{-6} g$$

$$S/N = 1000$$

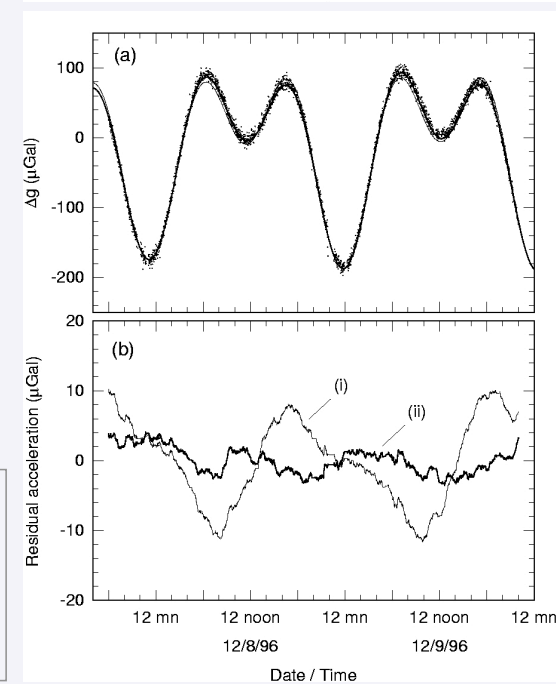
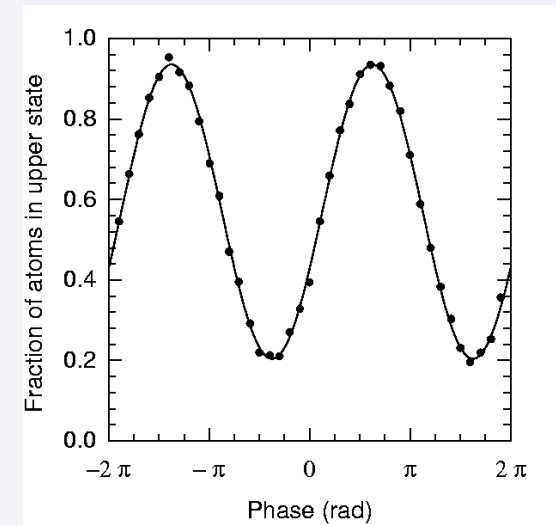
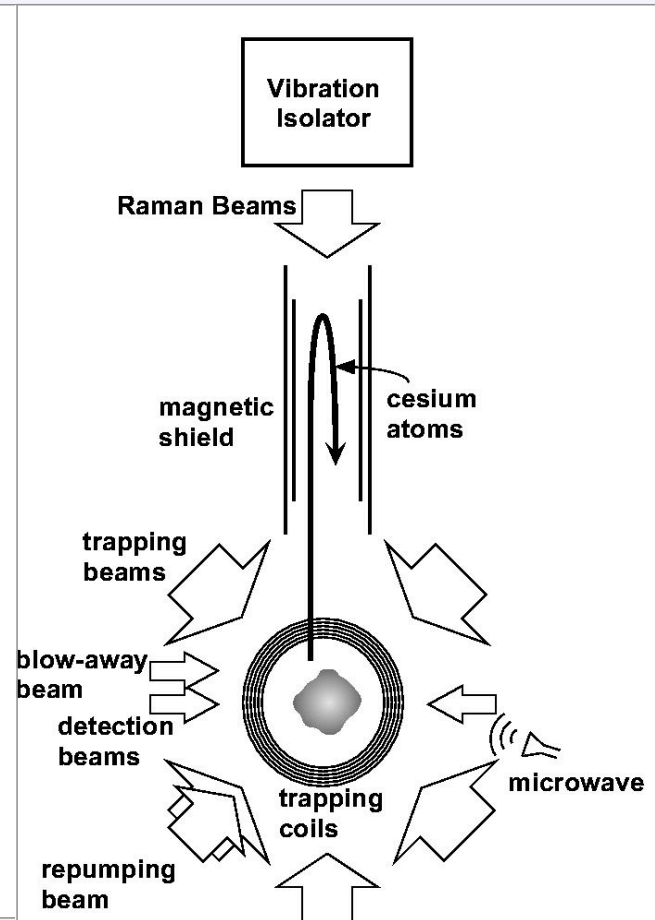
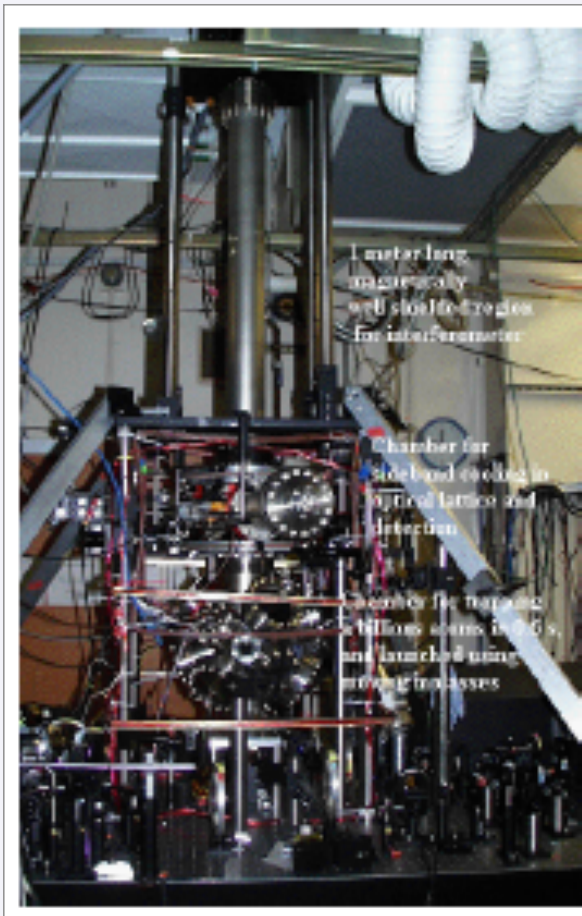
$$\Rightarrow \text{Sensitivity } 10^{-9} \text{ g/shot}$$



M. Kasevich, S. Chu, *Appl. Phys. B* **54**, 321 (1992)

A. Peters, K.Y. Chung and S. Chu, *Nature* **400**, 849 (1999)

Stanford atom gravimeter



Resolution: 3×10^{-9} g after 1 minute

Absolute accuracy: $\Delta g/g < 3 \times 10^{-9}$

A. Peters, K.Y. Chung and S. Chu, *Nature* **400**, 849 (1999)

Atom interferometry – gravity measurements

● spatial gravity variations

- gravity gradient $\sim 3 \cdot 10^{-7} \text{ g / m}$
- global scale $\sim 10^{-3} \text{ g}$
- regional scale $\sim 10^{-6} \text{ g}$

– navigation
– finding oil, water, minerals,
archeological sites, ...

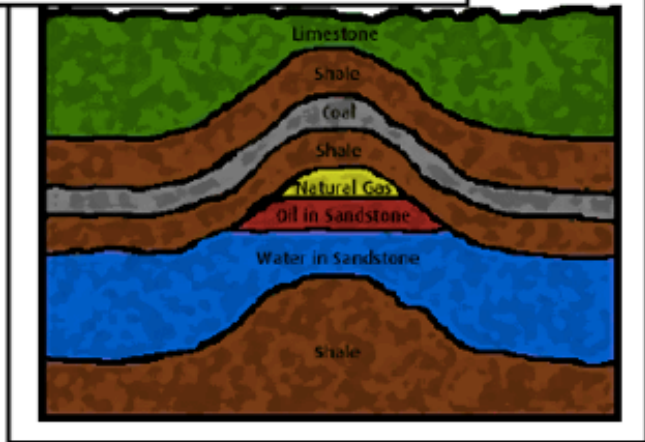
● temporal gravity variations

- tides $\sim 10^{-7} \text{ g}$
- man-made changes $\sim 10^{-9} \text{ g}$
- atmospheric pressure $\sim 10^{-10} \text{ g / mbar}$
- local water table $\sim 10^{-8} \text{ g}$
- ... $\sim 10^{-9} \text{ g}$

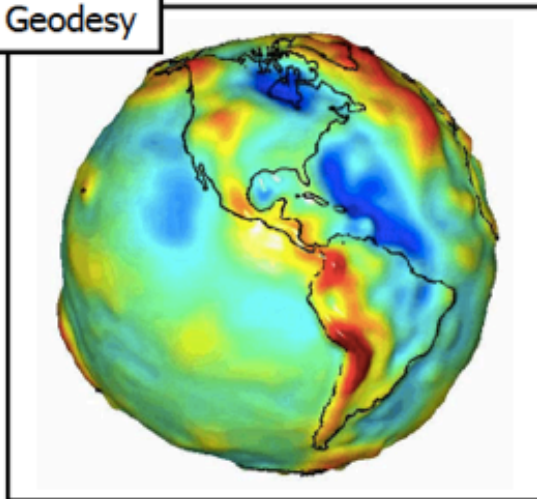
(1 Gal = 1 cm/s² → 1 μGal ≈ 10⁻⁹ g)

Gravimetry

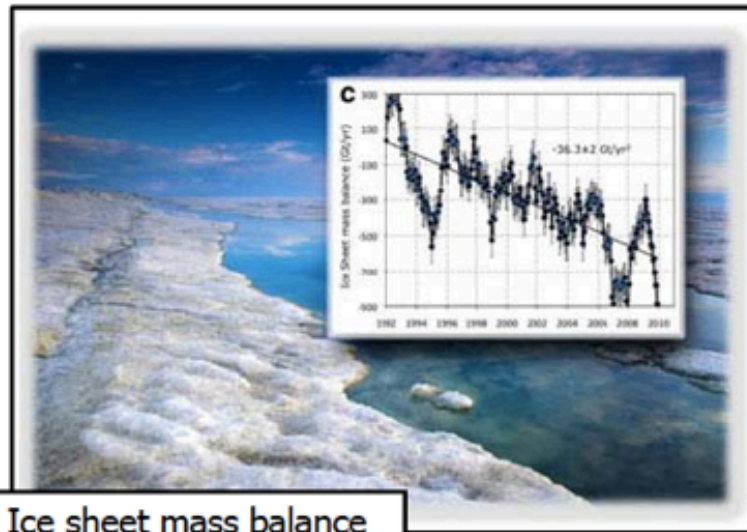
Mapping of natural resources



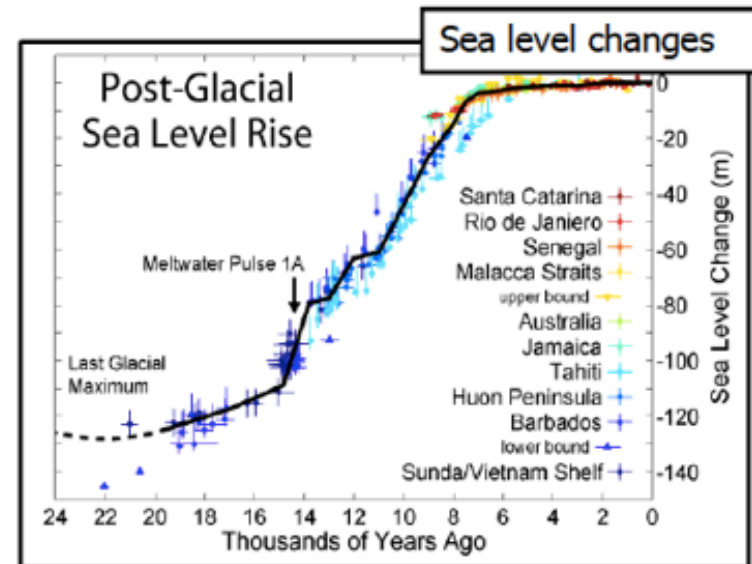
Geodesy



Navigation



Ice sheet mass balance



Images: www.esa.int, Google Images

TOPICAL REVIEW

Precision gravimetry with atomic sensors

M de Angelis^{1,2}, A Bertoldi³, L Cacciapuoti⁴, A Giorgini^{2,5},
 G Lamporesi⁶, M Prevedelli⁷, G Saccorotti⁸, F Sorrentino²
 and G M Tino²

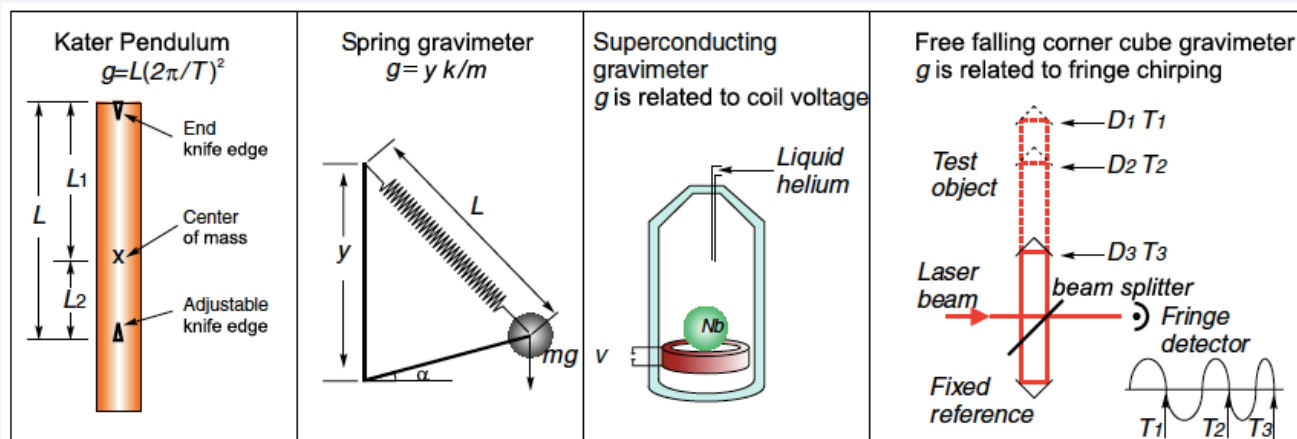
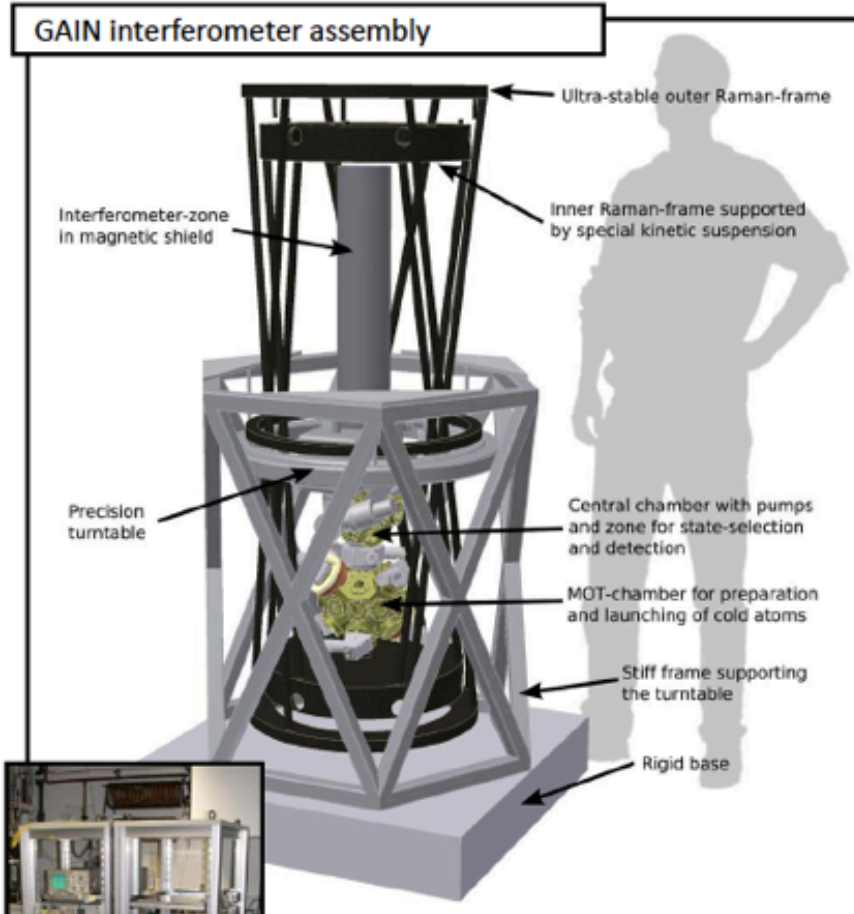


Table 1. Summary of error sources level and technical budgets for most used commercial gravimeters.

	Spring [94]	Superconducting [68, 95]	Free falling [69, 72]
Noise $(\Delta g/g)/\sqrt{\text{Hz}}$	5×10^{-9}	1×10^{-12}	5×10^{-8}
Drift $(\Delta g/g)$	1.5×10^{-6} per month	1×10^{-9} per year	–
Accuracy $\Delta g/g$	–	–	4×10^{-9}
Measurement	Relative	Relative	Absolute
Size (m ³)	0.04	~1.5	1.5
Weight (kg)	14	321	127
Power (W)	24	400	350
Error sources	Temperature and random seasonal drift. Calibration varies in time and position	No field operation. Magnetic and electrostatic effects	Thermal drift. Magnetic and electrostatic effects

Portable atomic quantum gravimeter GAIN



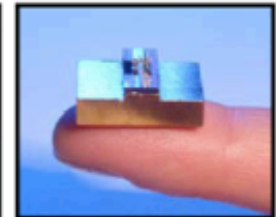
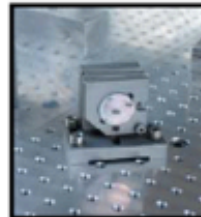
- **Compact:** three $\sim 1 \text{ m}^3$ Modules (interferometers assembly + two 19" racks for laser system and electronics)
- **Robust:** critical components based on technology developed for the high g-loads in drop tower experiments
- **Mobile:** designed to be „truckable“ and for use at a variety of interesting locations

Targeted sensitivity:

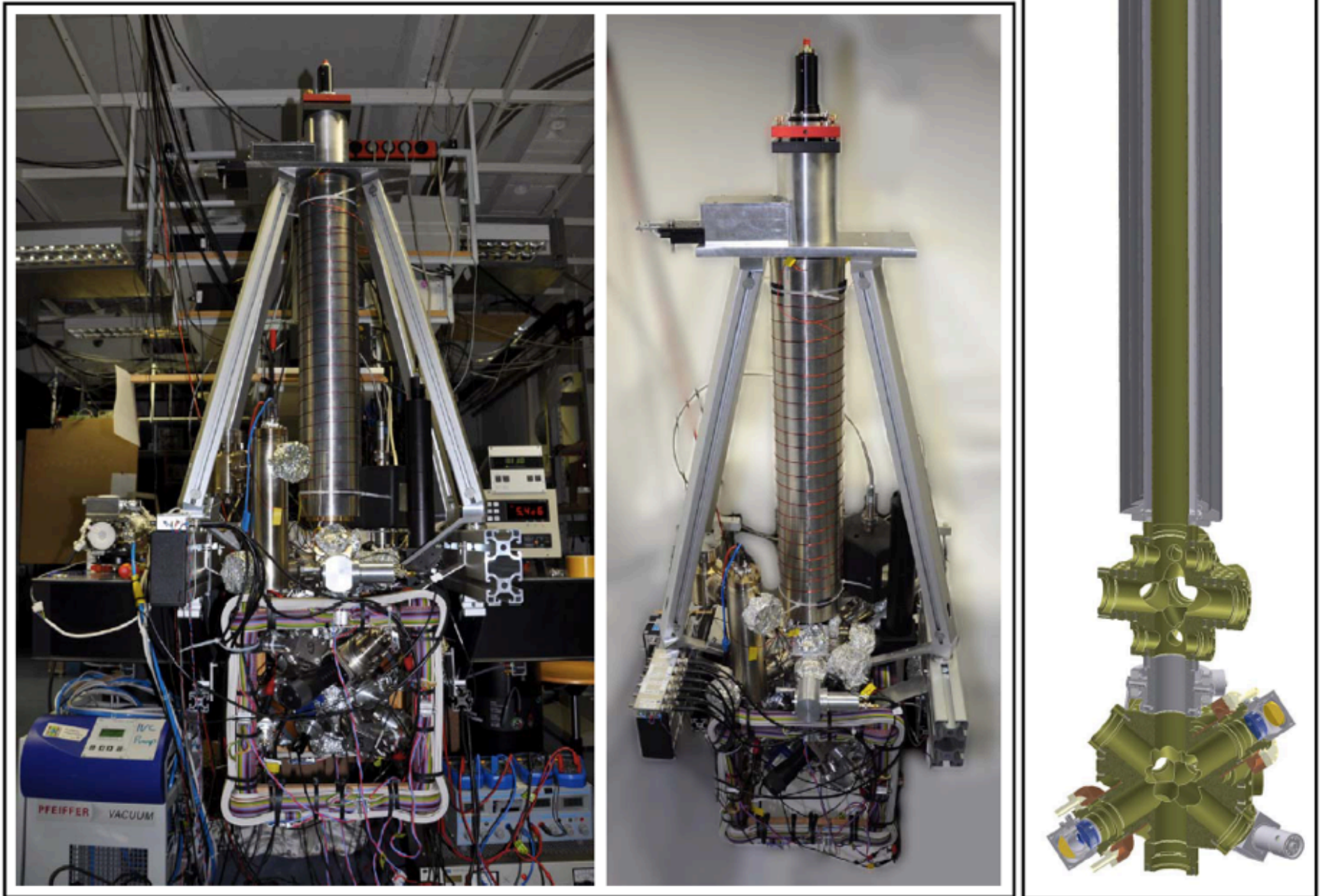
$1 \cdot 10^{-9} \text{ g} / \text{sqrt}(\text{Hz})$ at a SNR of 300:1
(intrinsic noise only)

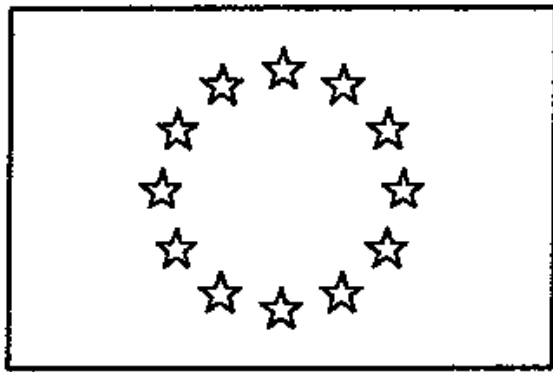
$1 \cdot 10^{-8} \text{ g} / \text{sqrt}(\text{Hz})$ at a SNR of 30:1
(under realistic vibration conditions)

Targeted absolute accuracy: $5 \cdot 10^{-10} \text{ g}$



GAIN – physics package





EUROPEAN COMMISSION

Information Society and Media Directorate-General

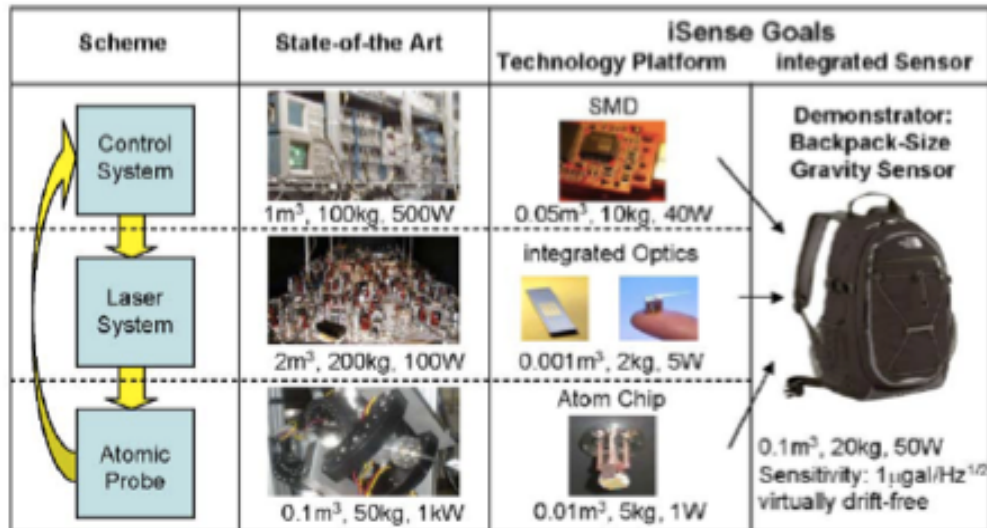
Emerging Technologies and Infrastructures

Future and Emerging Technologies (FET) - Open

iSense – Integrated Quantum Sensors

7th Framework Programme - Theme 3 "Information and Communication Technologies"

Call identifier: FP7-ICT-2009- C FET-Open

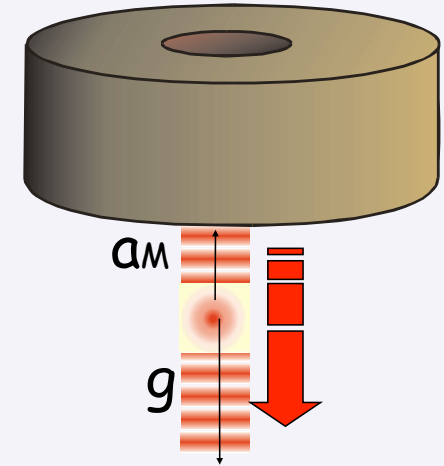


Participant no. *	Participant organisation name	Part. short name	Country
1 (Coordinator)	The University of Birmingham	Bham	UK
2	QinetiQ	QinetiQ	UK
3	University of Hamburg	UHH	D
4	Centre National de la Recherche Scientifique ¹	CNRS	F
5	University of Florence	UNIFI	I
6	Leibniz University Hannover	LUH	D
7	Institute for quantum optics and quantum information - Austrian Academy of Sciences	IQOQI-OEAW	A
8	Ferdinand-Braun-Institut für Höchstfrequenztechnik im Forschungsverbund Berlin e.V.	FBH	D
9	University of Nottingham	Nham	UK

MAGIA

(MISURA ACCURATA di G MEDIANTE INTERFEROMETRIA ATOMICA)

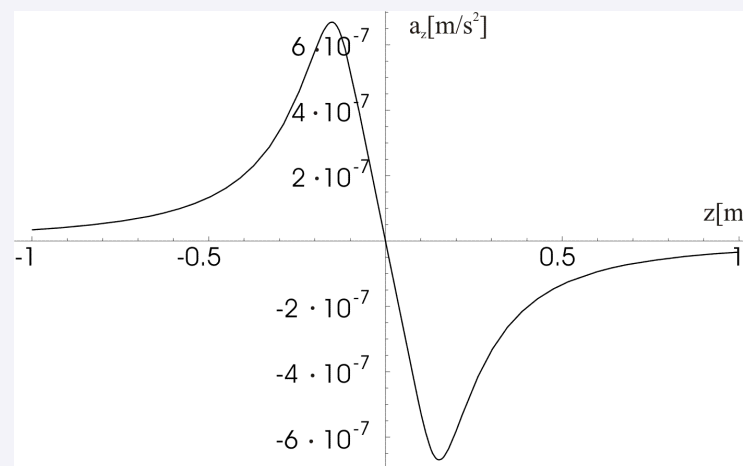
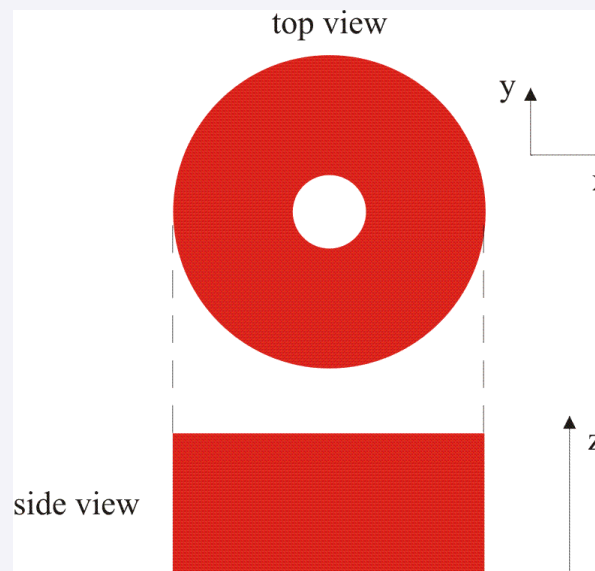
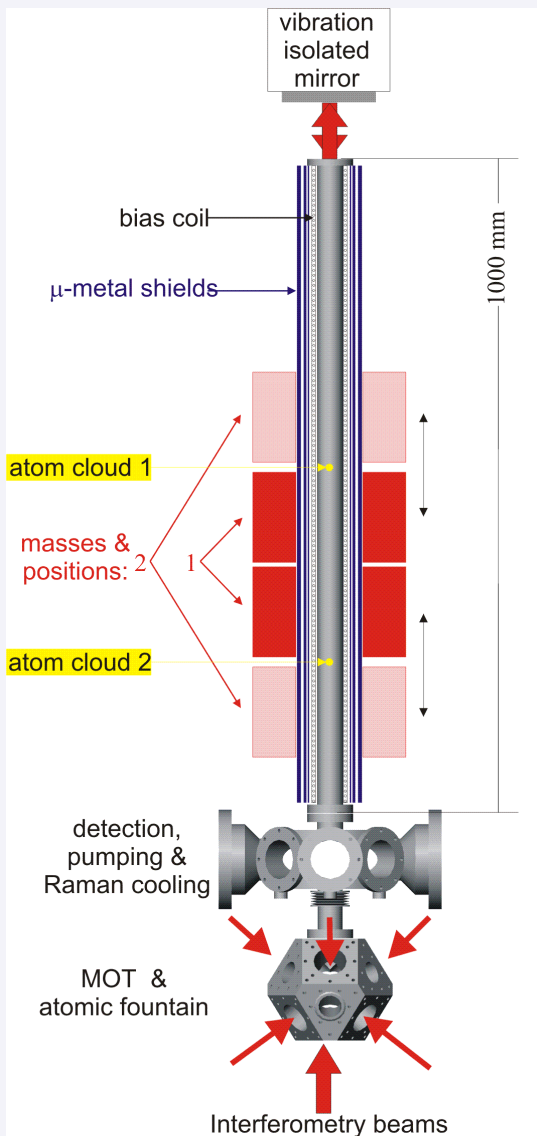
- Measure g by atom interferometry
- Add source mass
- Measure change of g



➤ *Precision measurement of G*

$$F(r) = G \frac{M_1 M_2}{r^2}$$

MAGIA: atom gravimeter + source mass



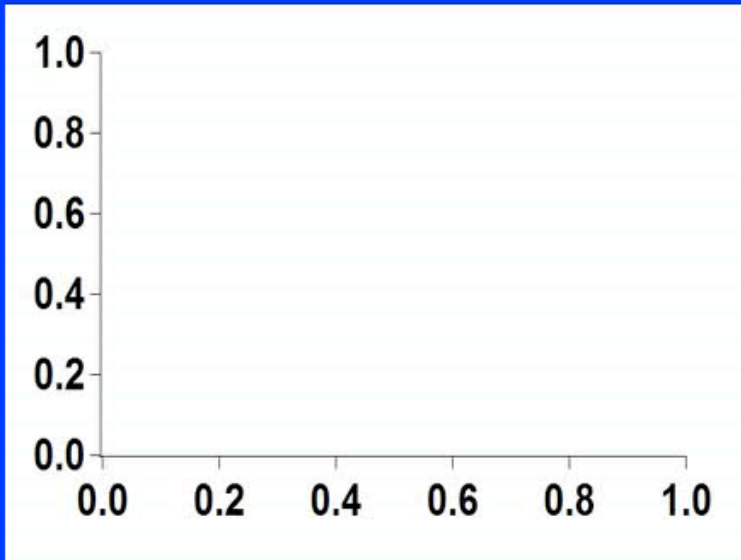
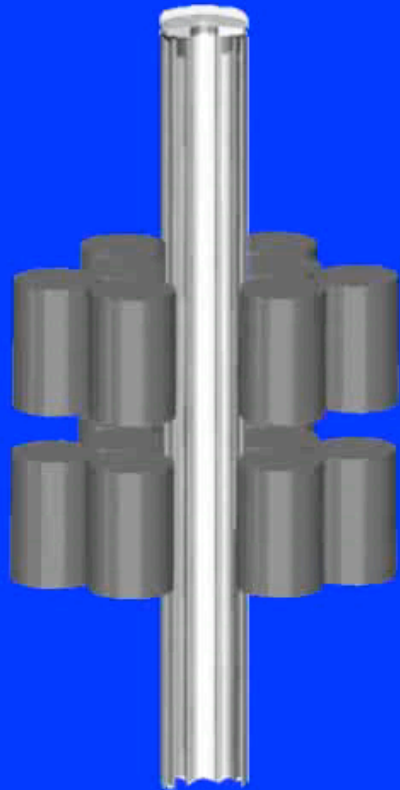
500 kg tungsten mass

Sensitivity 10^{-9} g/shot

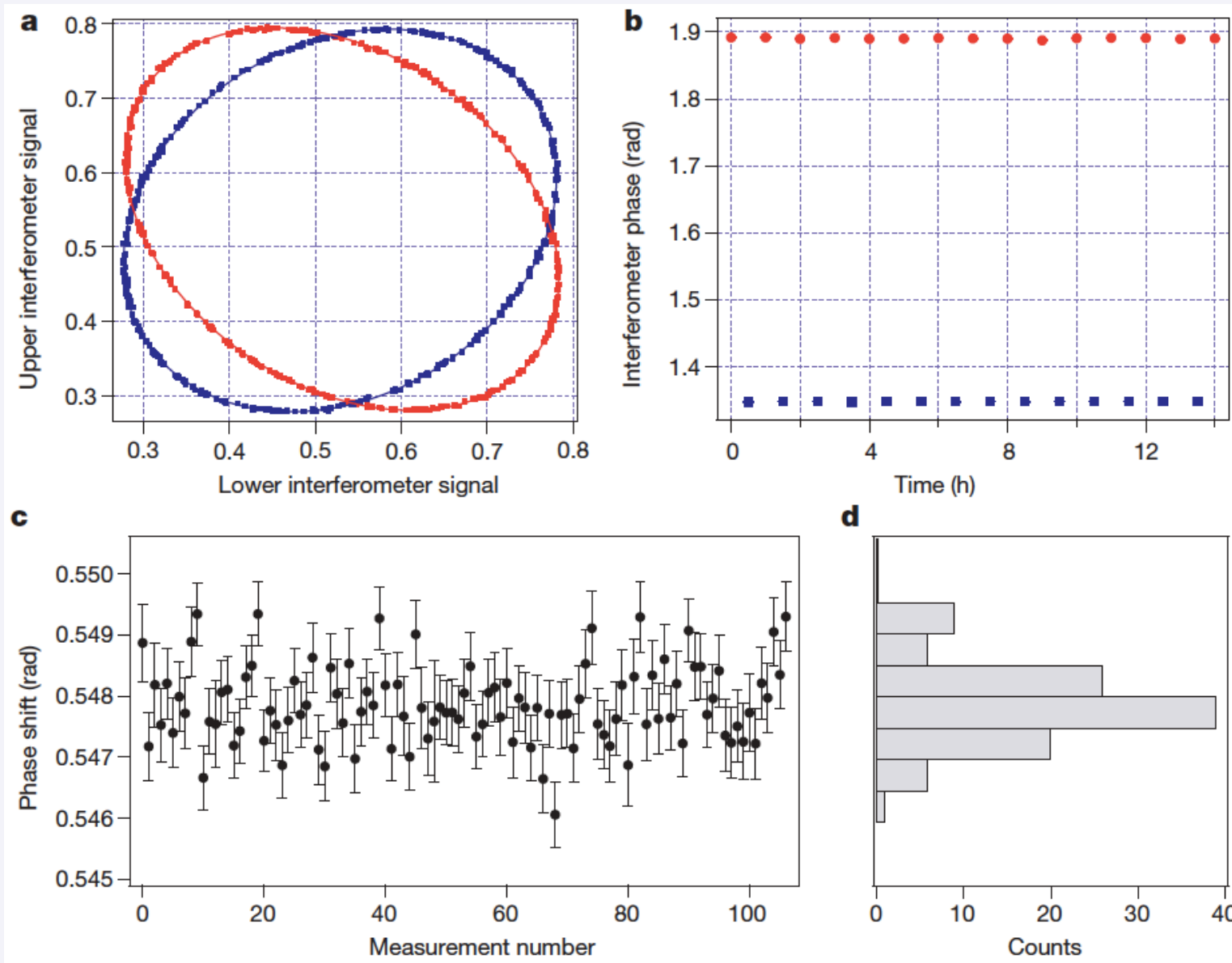
Peak mass acceleration $a_G \approx 10^{-7}$ g

one shot $\Rightarrow \Delta G/G \approx 10^{-2}$

10000 shots $\Rightarrow \Delta G/G \approx 10^{-4}$



Measurement of G



(July 2013)

Relative uncertainty ~ 116 ppm (statistical)

LETTER

doi:10.1038/nature13433

Precision measurement of the Newtonian gravitational constant using cold atoms

G. Rosi¹, F. Sorrentino¹, L. Cacciapuoti², M. Prevedelli³ & G. M. Tino¹

About 300 experiments have tried to determine the value of the Newtonian gravitational constant, G , so far, but large discrepancies in the results have made it impossible to know its value precisely¹. The weakness of the gravitational interaction and the impossibility of shielding the effects of gravity make it very difficult to measure G while keeping systematic effects under control. Most previous experiments performed were based on the torsion pendulum or torsion balance scheme as in the experiment by Cavendish² in 1798, and in all cases macroscopic masses were used. Here we report the precise determination of G using laser-cooled atoms and quantum interferometry. We obtain the value $G = 6.67191(99) \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$ with a relative uncertainty of 150 parts per million (the combined standard

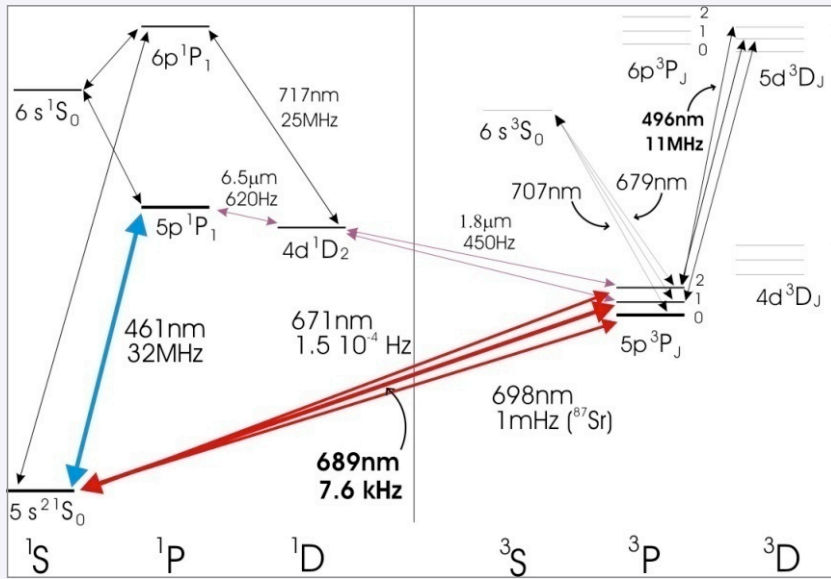
the relevant gravitational signal. An additional cancellation of common-mode spurious effects was obtained by reversing the direction of the two-photon recoil used to split and recombine the wave packets in the interferometer¹⁸. Efforts were devoted to the control of systematics related to atomic trajectories, the positioning of the atoms and effects due to stray fields. The high density of tungsten was instrumental in maximizing the signal and in compensating for the Earth's gravitational gradient in the region containing the atom interferometers, thus reducing the sensitivity of the experiment to the vertical position and size of the atomic probes.

The atom interferometer is realized using light pulses to stimulate ⁸⁷Rb atoms at the two-photon Raman transition between the hyperfine levels $F=1$ and $F=2$ of the ground state¹⁹. The light field is generated

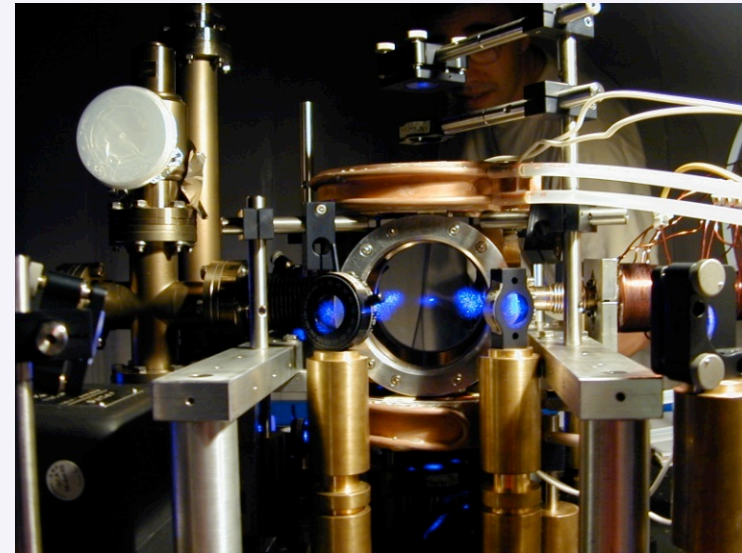
$$G = 6.67191(99) \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$$

Relative uncertainty: 150 ppm

Ultracold Sr – The experiment in Firenze



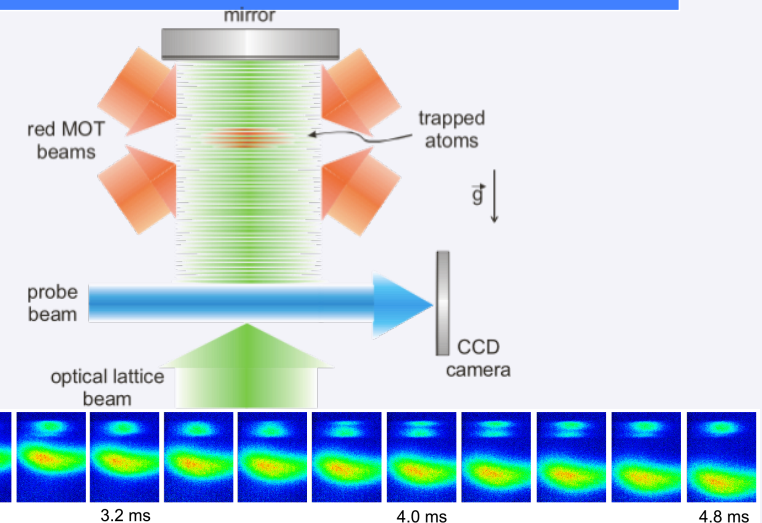
- Optical clocks using visible intercombination lines



- New atomic sensors for fundamental physics tests

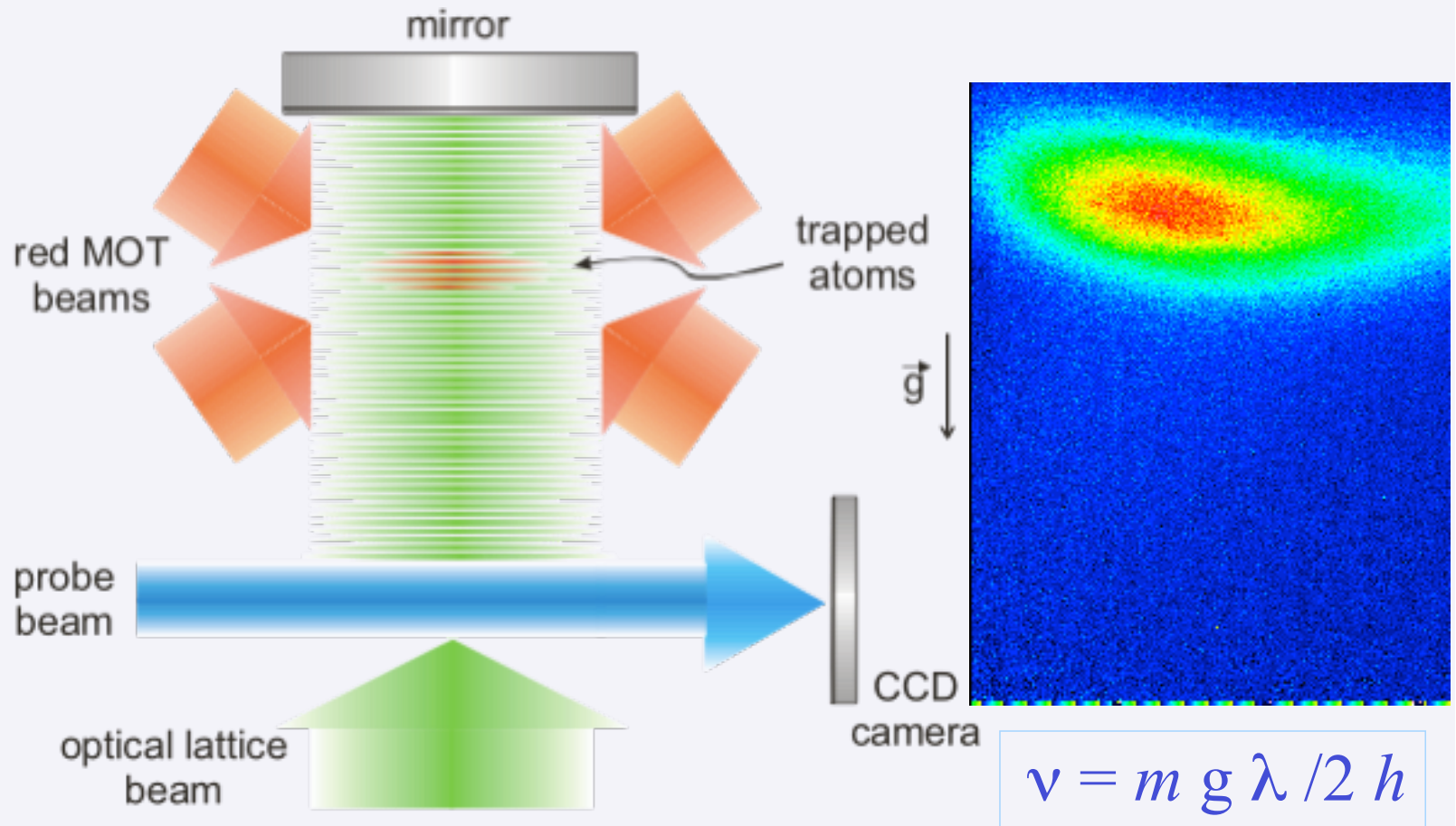


G. Ferrari, P. Cancio, R. Drullinger, G. Giusfredi, N. Poli, M. Prevedelli, C. Toninelli, G.M. Tino, *Precision Frequency Measurement of Visible Intercombination Lines of Strontium*, Phys. Rev. Lett. 91, 243002 (2003)



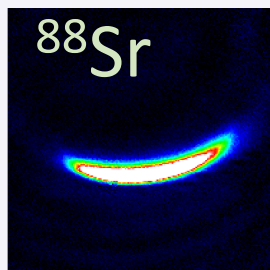
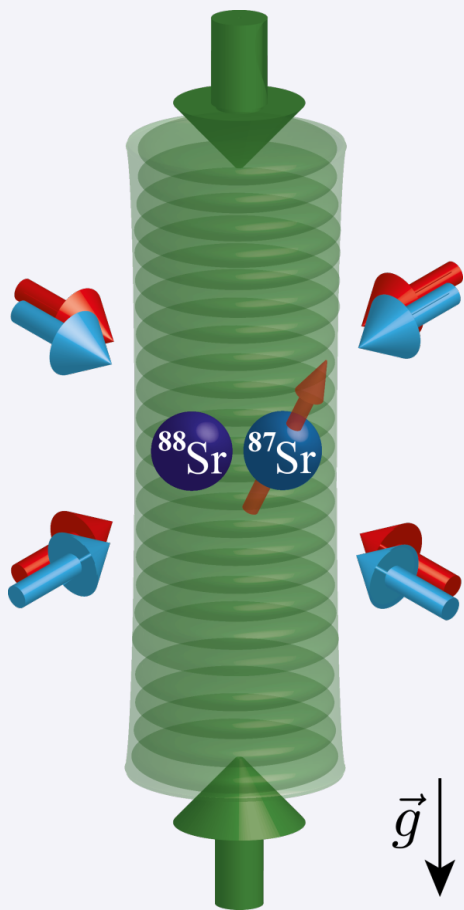
G. Ferrari, N. Poli, F. Sorrentino, and G. M. Tino, *Long-lived Bloch oscillations with bosonic Sr atoms and application to gravity measurement at micrometer scale*, Phys. Rev. Lett. 97, 060402 (2006)

Precision gravity measurement at μm scale with Bloch oscillations of Sr atoms in an optical lattice

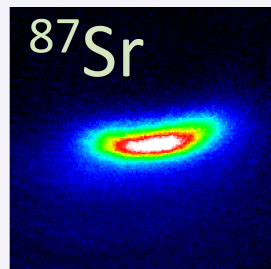


G. Ferrari, N. Poli, F. Sorrentino, G. M. Tino, *Long-Lived Bloch Oscillations with Bosonic Sr Atoms and Application to Gravity Measurement at the Micrometer Scale*, *Phys. Rev. Lett.* **97**, 060402 (2006)

Test of Einstein equivalence principle for 0-spin and half-integer-spin Sr atoms: Search for spin-gravity coupling effects



8×10^6 atoms
T: 1 μK



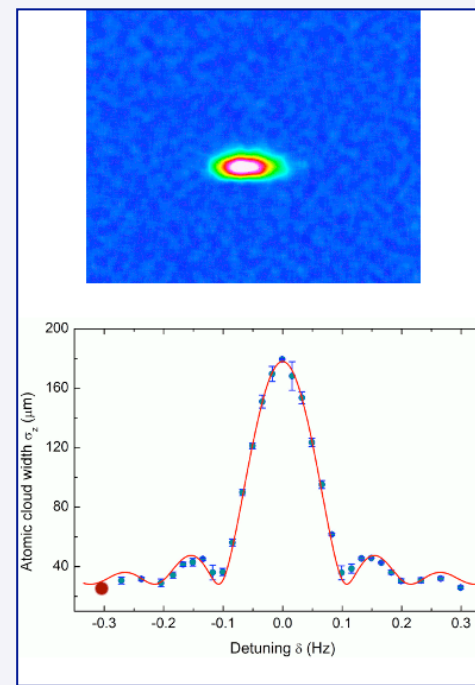
1×10^6 atoms
T: 1.4 μK



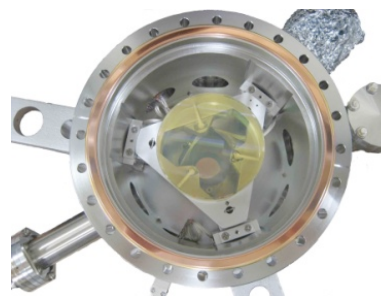
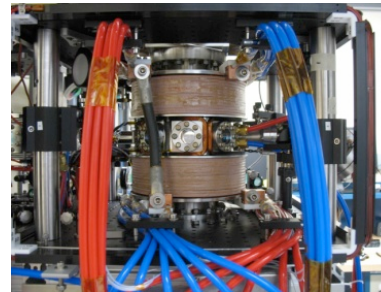
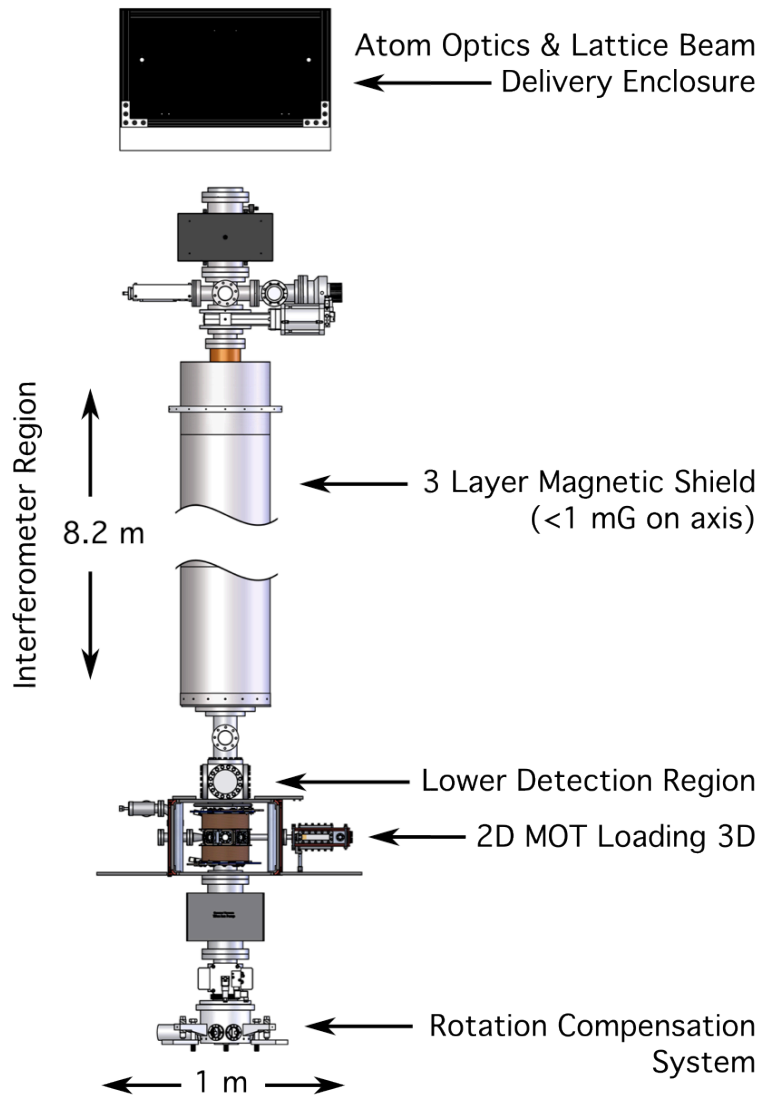
\vec{g}

Final result:
 $\eta = (0.2 \pm 1.6) \times 10^{-7}$

$\Rightarrow k = (0.5 \pm 1.1) \times 10^{-7}$



Apparatus



Ultracold atom source

$>10^6$ atoms at 50 nK
 $3e5$ atoms at 1.6 nK

Optical Lattice Launch

13.1 m/s with 2372 photon recoils to 9 m

Atom Interferometry

2 cm $1/e^2$ radial waist
6 W total power

Dynamic nrad control of laser angle with precision piezo-actuated stage

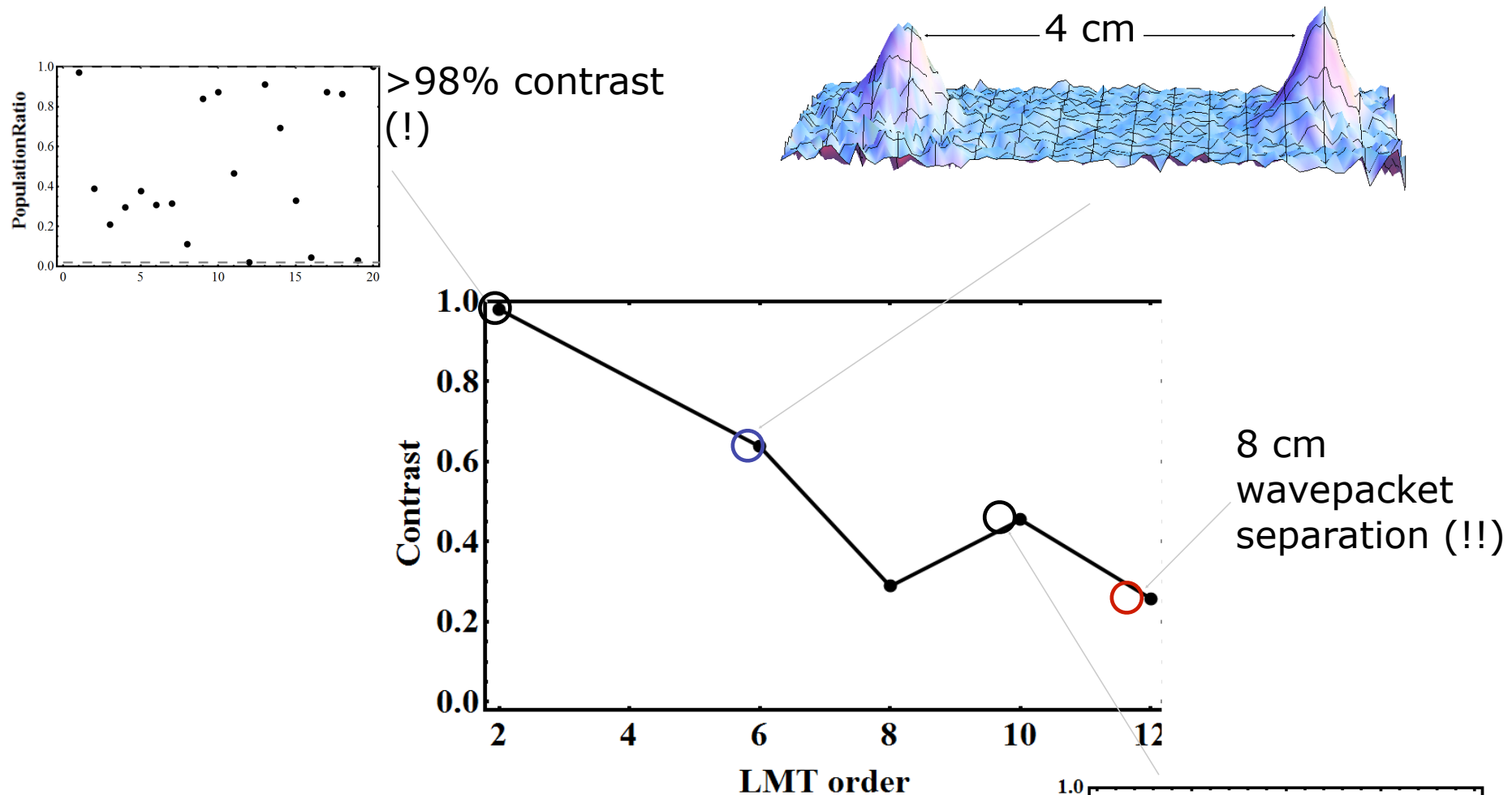
Detection

Spatially-resolved fluorescence imaging

Two CCD cameras on perpendicular lines of sight

Current demonstrated statistical resolution, $\sim 5e-13$ g in 1 hr (87Rb)

Contrast vs. momentum recoil at $2T = 2.3$ s



Large momentum transfer demonstration at $2T = 2.3$ s (unpublished).

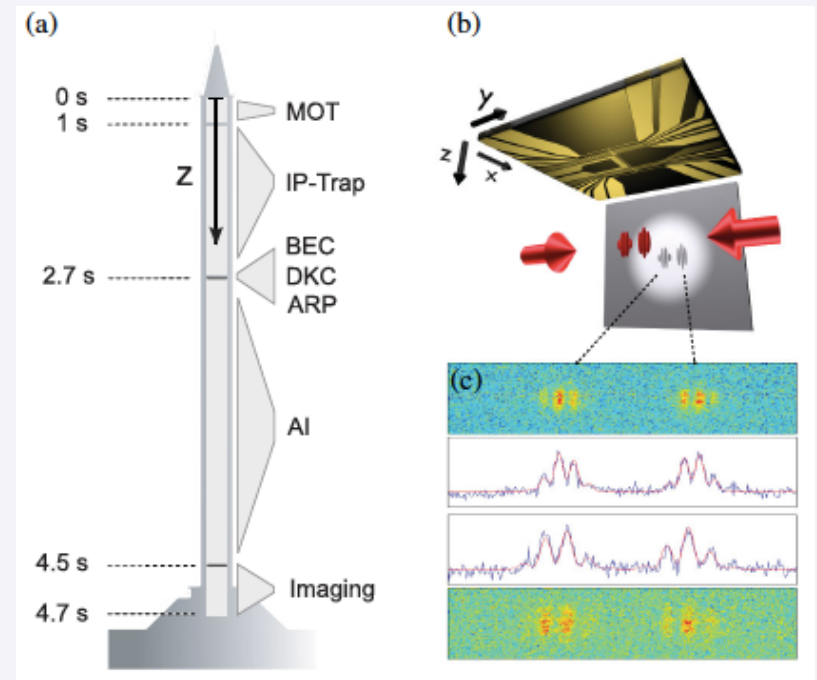
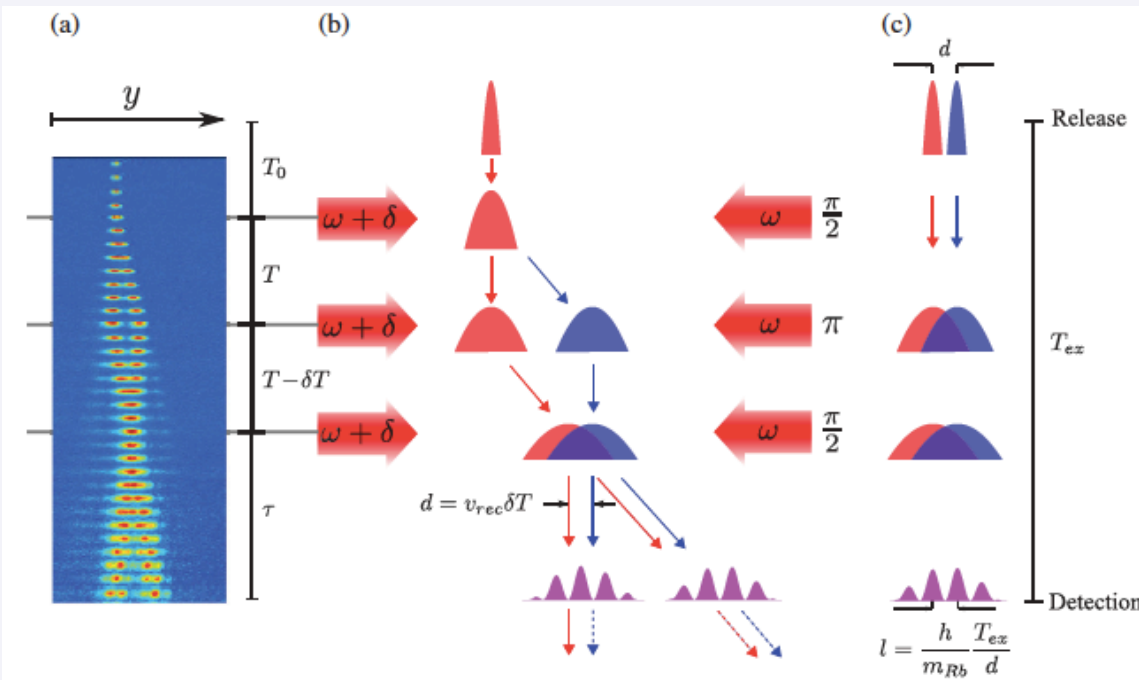




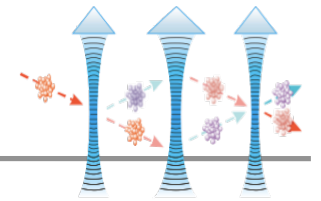
Interferometry with Bose-Einstein Condensates in Microgravity

H. Müntinga,¹ H. Ahlers,² M. Krutzik,³ A. Wenzlawski,⁴ S. Arnold,⁵ D. Becker,² K. Bongs,⁶ H. Dittus,⁷ H. Duncker,⁴ N. Gaaloul,² C. Gherasim,⁸ E. Giese,⁵ C. Grzeschik,³ T. W. Hänsch,⁹ O. Hellmig,⁴ W. Herr,² S. Herrmann,¹ E. Kajari,^{5,10} S. Kleinert,⁵ C. Lämmerzahl,¹ W. Lewoczko-Adamczyk,³ J. Malcolm,⁶ N. Meyer,⁶ R. Nolte,⁸ A. Peters,^{3,11} M. Popp,² J. Reichel,¹² A. Roura,⁵ J. Rudolph,² M. Schiemangk,^{3,11} M. Schneider,⁸ S. T. Seidel,² K. Sengstock,⁴ V. Tamma,⁵ T. Valenzuela,⁶ A. Vogel,⁴ R. Walser,⁸ T. Wendrich,² P. Windpassinger,⁴ W. Zeller,⁵ T. van Zoest,⁷ W. Ertmer,² W. P. Schleich,⁵ and E. M. Rasel^{2,*}

Atom interferometers covering macroscopic domains of space-time are a spectacular manifestation of the wave nature of matter. Because of their unique coherence properties, Bose-Einstein condensates are ideal sources for an atom interferometer in extended free fall. In this Letter we report on the realization of an asymmetric Mach-Zehnder interferometer operated with a Bose-Einstein condensate in microgravity. The resulting interference pattern is similar to the one in the far field of a double slit and shows a linear scaling with the time the wave packets expand. We employ delta-kick cooling in order to enhance the signal and extend our atom interferometer. Our experiments demonstrate the high potential of interferometers operated with quantum gases for probing the fundamental concepts of quantum mechanics and general relativity.



Towards an Atom Interferometer in Space



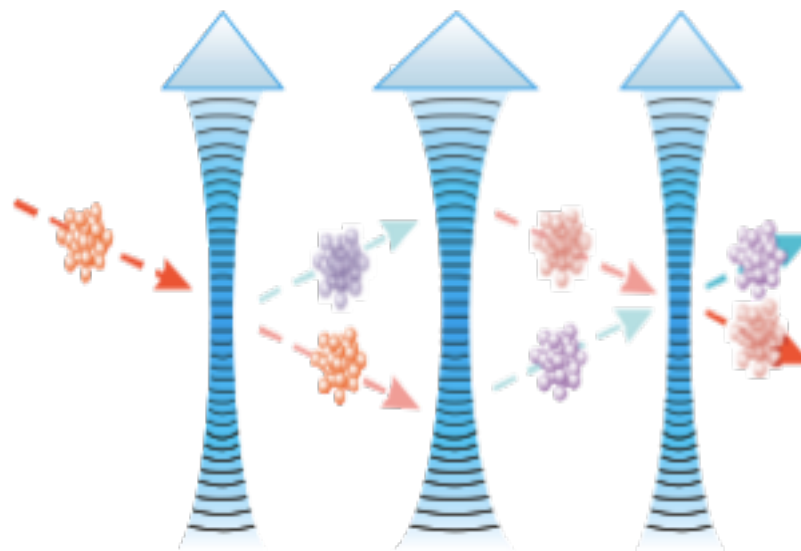
Space Atom Interferometer - SAI

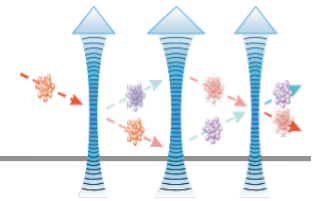
Space Atom Interferometer: Pre-phase A study of a space instrument based on matter-wave interferometry for inertial sensing in space

Team: Firenze Univ. (I), IOTA (F), IQ (D), Hamburg Univ. (D), HU Berlin (D), SYRTE (F), LENS (I), Ulm Univ. (D), ZARM (D)

Objective: Ground based prototype of an atom interferometer for precision measurements

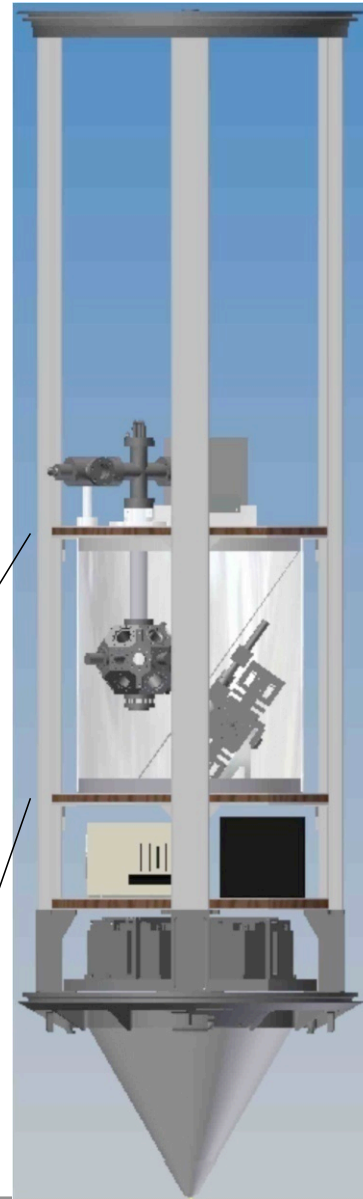
Duration: 3 years, funded within the ELIPS-2 Programme



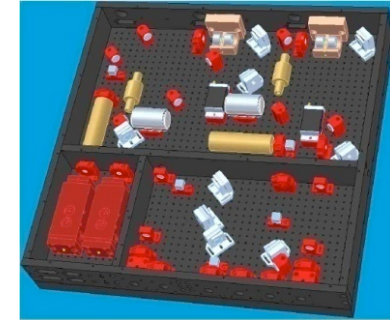


SAI - Space Atom Interferometer

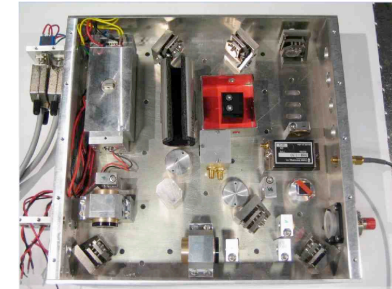
- Single-axis accelerometer
- Sensitivity target $10^{-7} \text{ m/s}^2 @ 1\text{s}$
- Repetition rate $\approx 2 \text{ Hz}$
- Modular laser system + optical fibers
- MOT + atomic fountain
- Same chamber for MOT and detection
- Load from 2D-MOT
- **Compatible with drop-tower capsule**



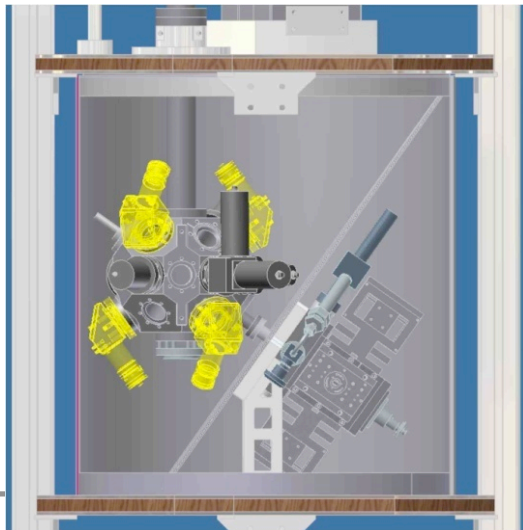
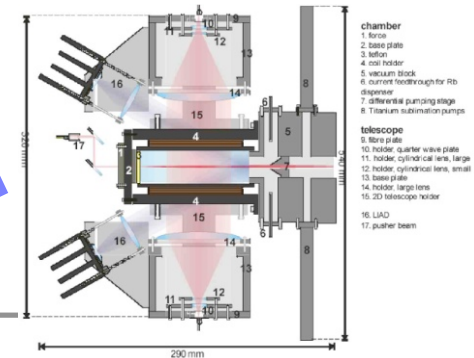
Raman laser system

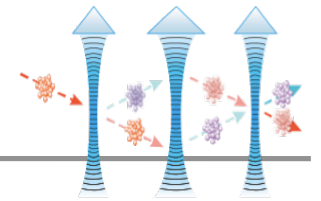


Modular laser system



2D-MOT

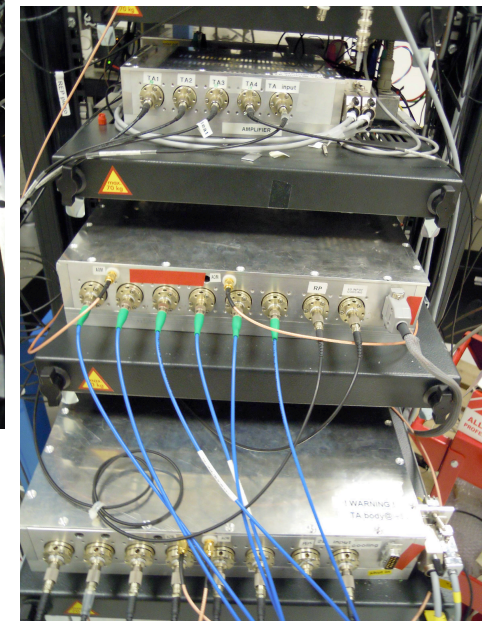
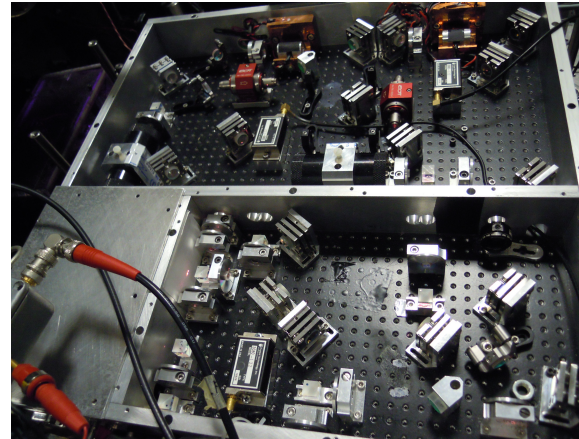
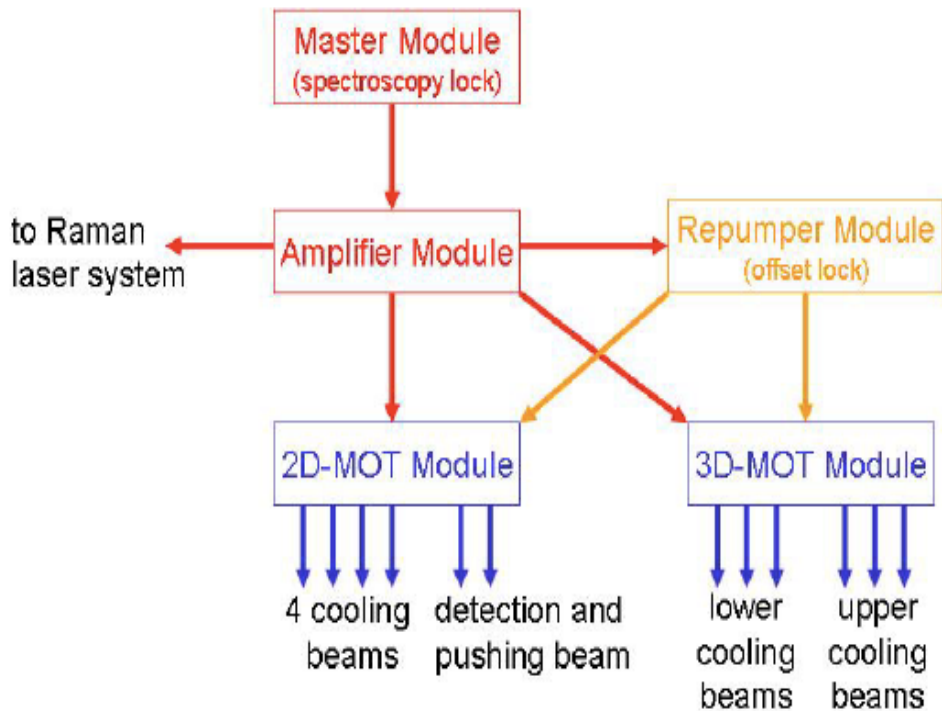
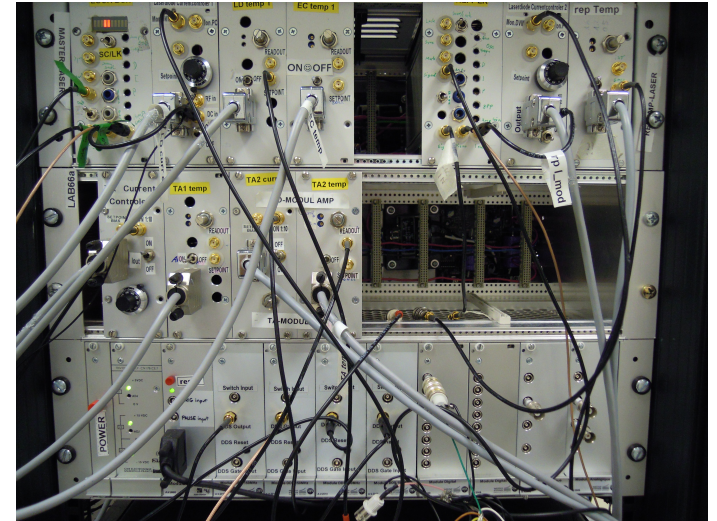
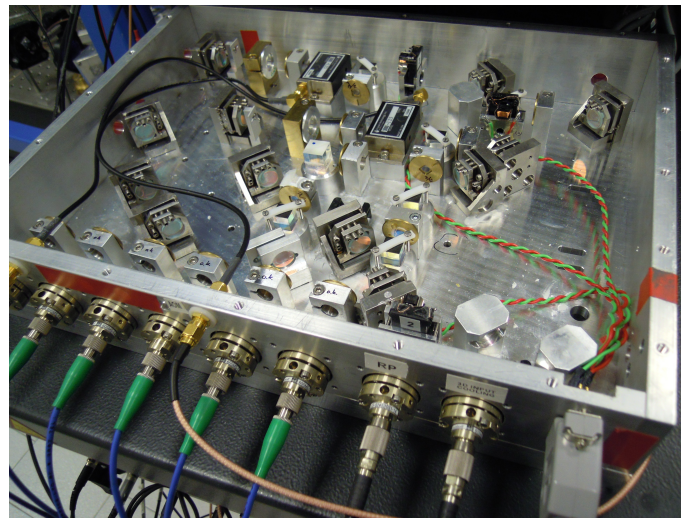
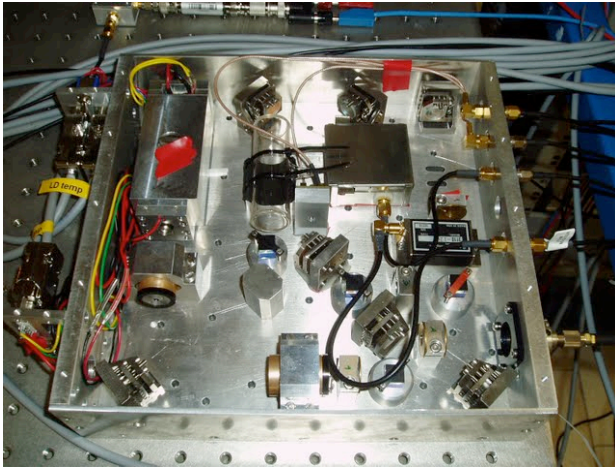
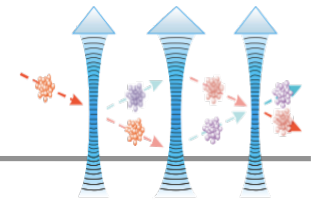


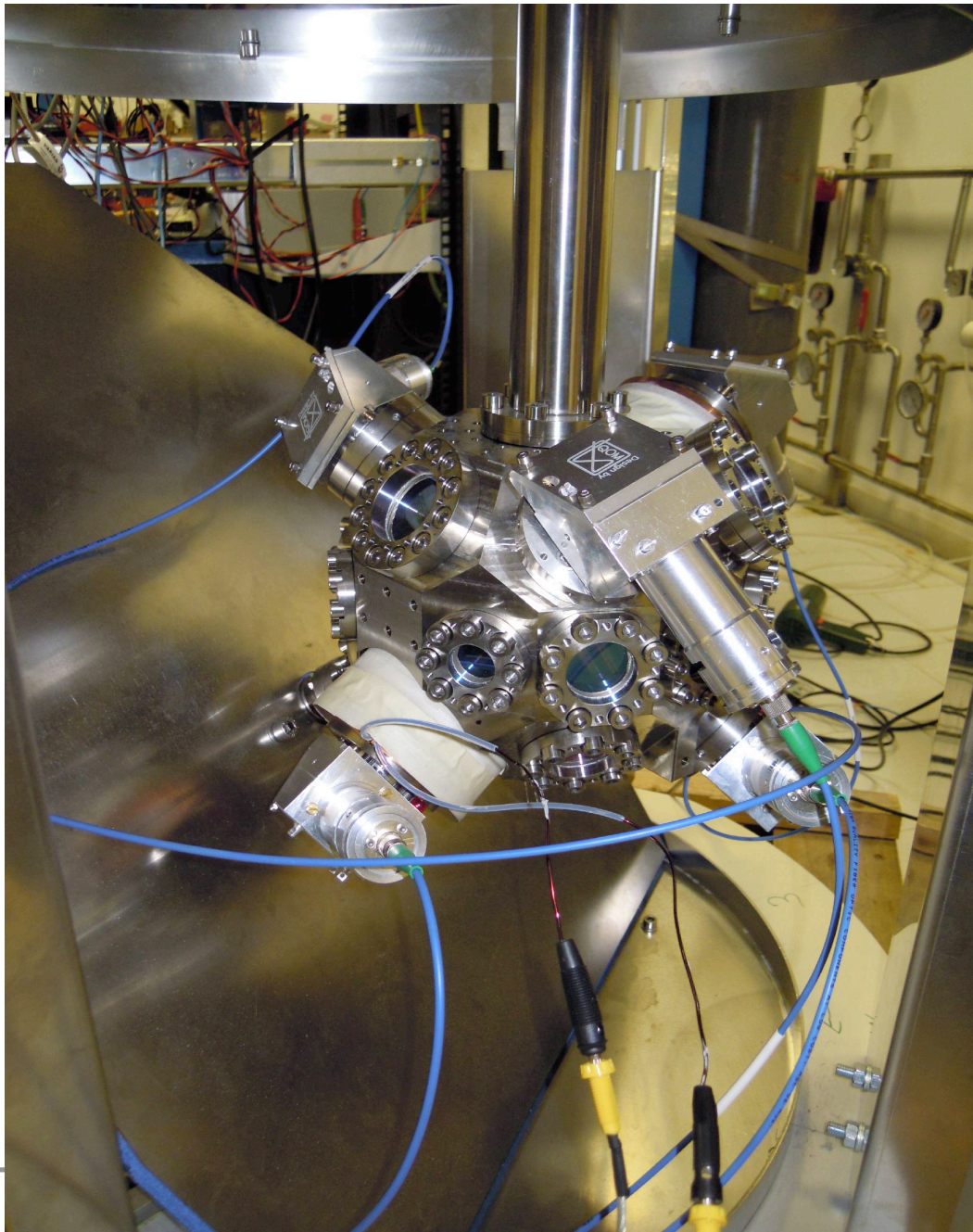
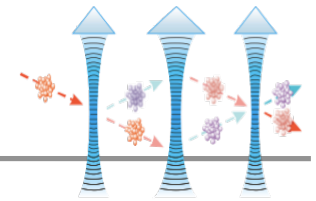


- Sensor hardware assembled
- Main subcomponents
 - Raman laser module
 - Modular laser system
 - 6.8 GHz source
 - High-flux atomic source (2D-MOT)
 - Vacuum system
 - Mechanical structure & magnetic shield
 - Electronic control system
- Theoretical investigations



F. Sorrentino et al., *A compact atom interferometer for future space missions*,
 Microgravity Sci. Technol. **22**, 551 (2010)





- SAI design compatible with all-optical evaporation to BEC
- Two viewports dedicated to 1560 nm FORT with two crossed beams of waist 150 μm and 30 μm



DRL-funded program until 2014

Experiments at the drop tower and catapult

QUANTUS I:

Drop experiments on chip-based interferometry

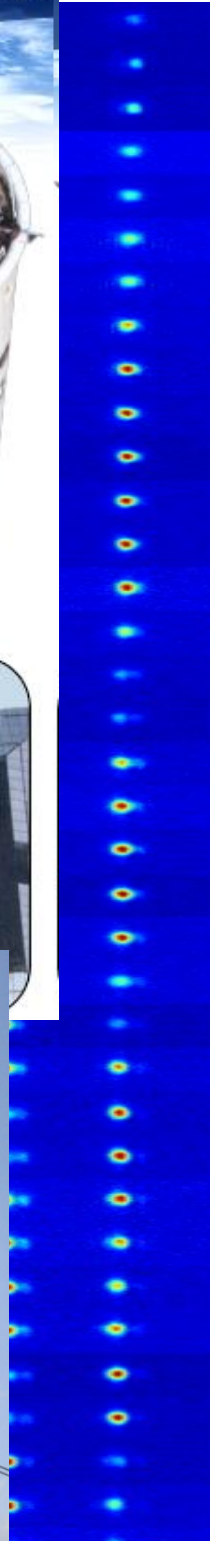
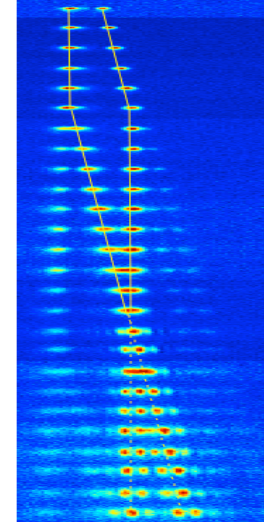
Physics with ultra cold gases

QUANTUS II

Drop & catapult experiments on dual species sources and interferometry

MAIUS

Atom interferometer hosted on a rocket

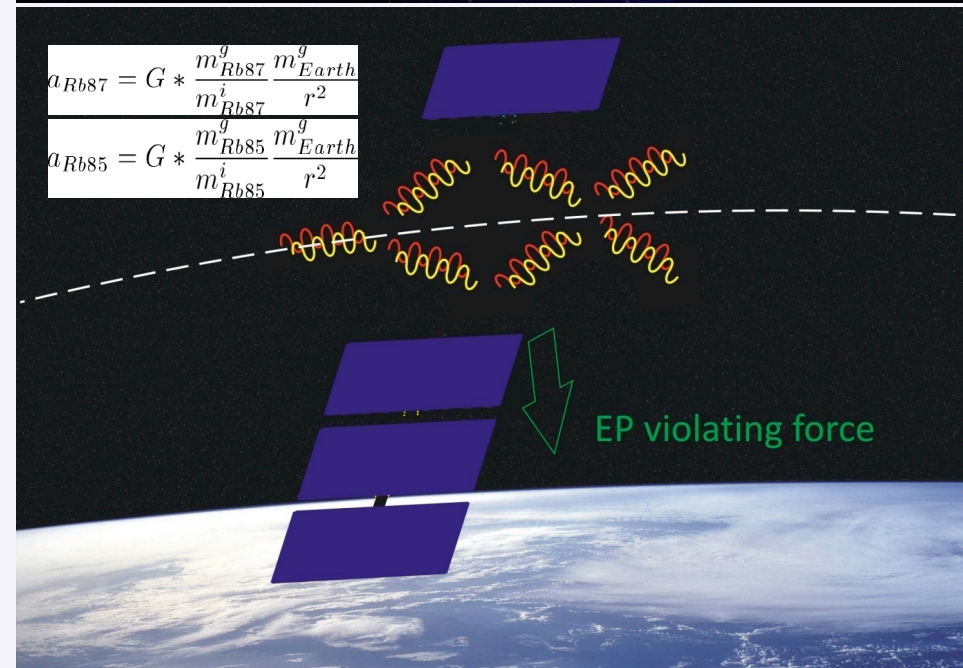
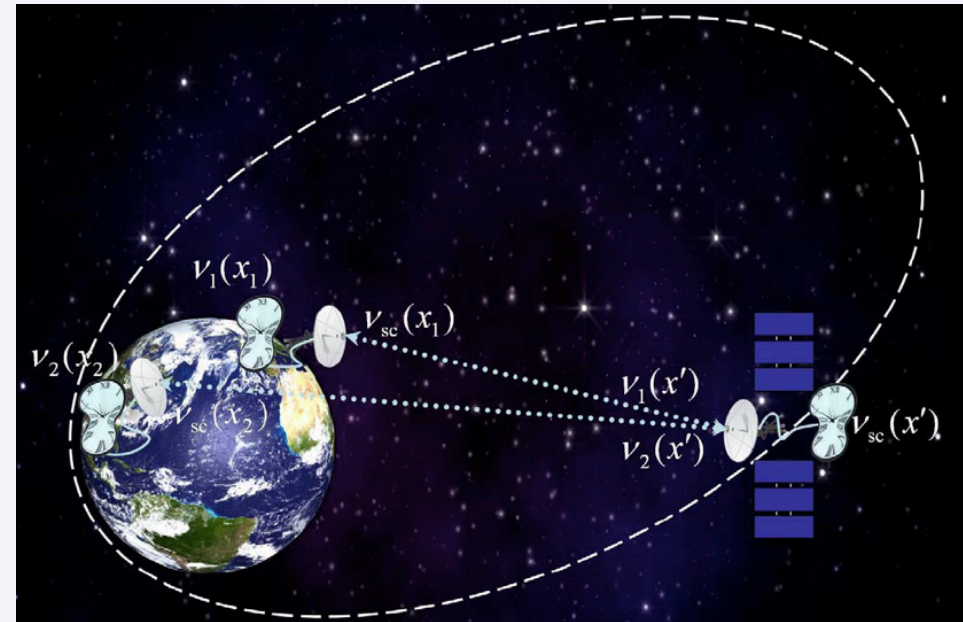
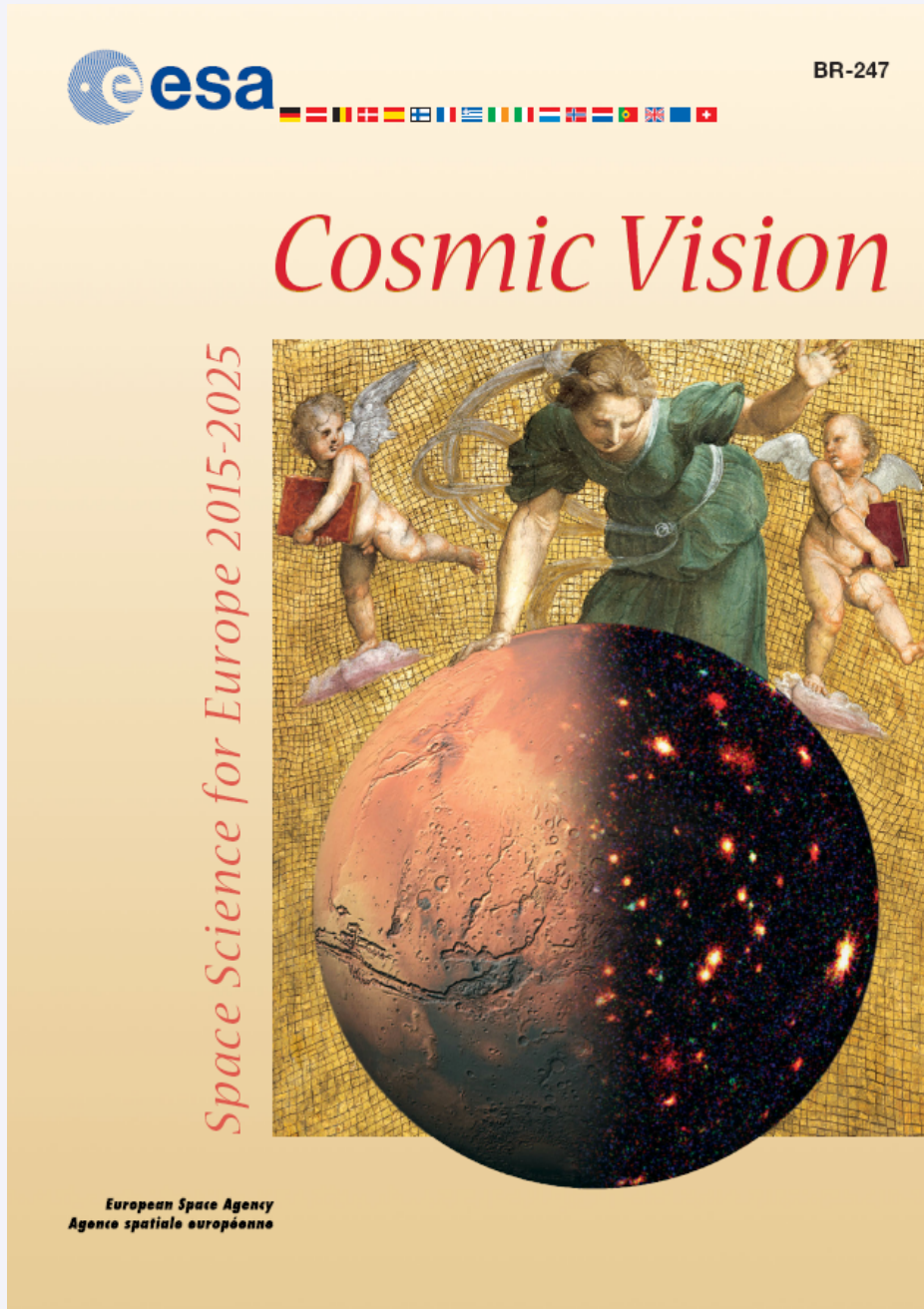


Q-WEP: towards a mission on the ISS

- Two studies in competition (1 year, funded to Astrium & Thales Alenia Space) for designing a mission architecture and addressing the design drivers imposed by the ISS-specific environment
- June 2012: joint ESR document delivered from scientific teams with definition of mission objectives with associated performance figures and derivation of the driving scientific requirements
 - **Primary objective:** WEP test at 10^{-14} on the Eötvös ratio in less than 4 months integration time
 - **Secondary objectives:** demonstration of atom interferometry differential acceleration measurement with sensitivity better than $2 \cdot 10^{-10} \text{ m/s}^2/\sqrt{\text{Hz}}$; demonstration of long interaction time ($1 \text{ s} < T < 5 \text{ s}$) with $> 50\%$ contrast.
 - **Baseline design:** double diffraction Raman differential interferometer on ^{85}Rb - ^{87}Rb with $T=1 \text{ s}$; CMRR for vibration noise up to 10^8 with scale factor compensation using equal k -vectors; tip-tilt mirror to compensate for ISS rotation @ 0.1% .
 - **Driving instrument requirements:** average relative displacement of atomic clouds $< 4 \text{ nm}$; average relative velocity $< 0.5 \text{ nm/s}$; atomic temperature $\sim 1 \text{ nK}$ with 10^6 atoms; retro-reflection mirror rotation $< 10^5 \text{ rad/s}$; radius of curvature of Raman laser beams $\sim 250 \text{ km}$
- ...for a mission opportunity on the ISS in the 2018-2020 time frame

- STE-QUEST Mission -

Test of Gravitational Red Shift and Equivalence Principle



Detection of Gravitational Waves by Atom Interferometry

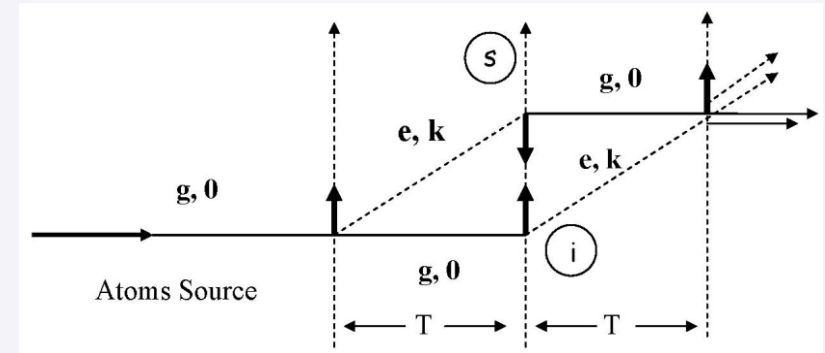
Main ideas

- Detection of GWs by matter waves
- Drastic reduction of critical noise sources
- Addressing new interesting frequency ranges

Two possible schemes

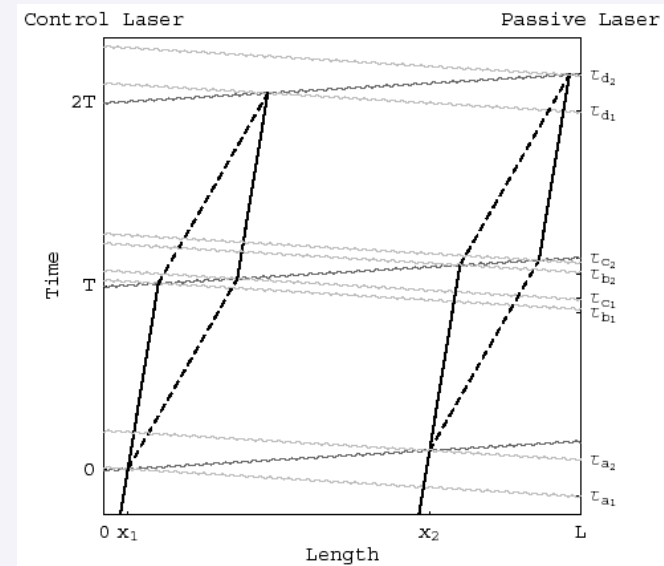
- *Single atom interferometer*

G.M. Tino and F. Vetrano, *Is it possible to detect gravitational waves with atom interferometers?* *Class. Quantum Grav.* 24, 2167 (2007)

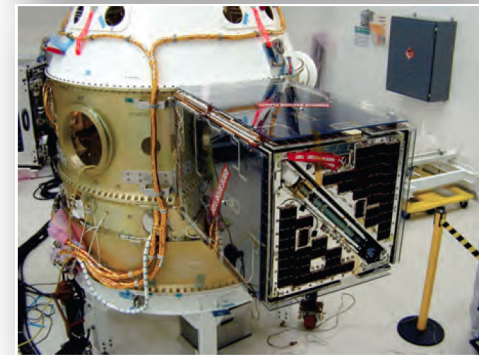
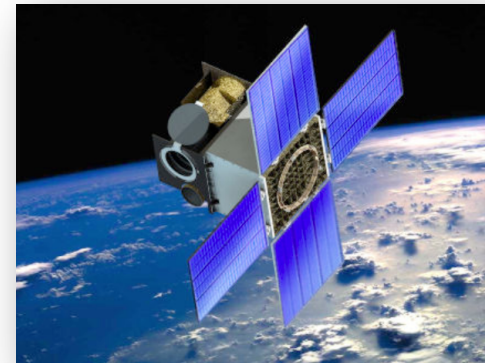
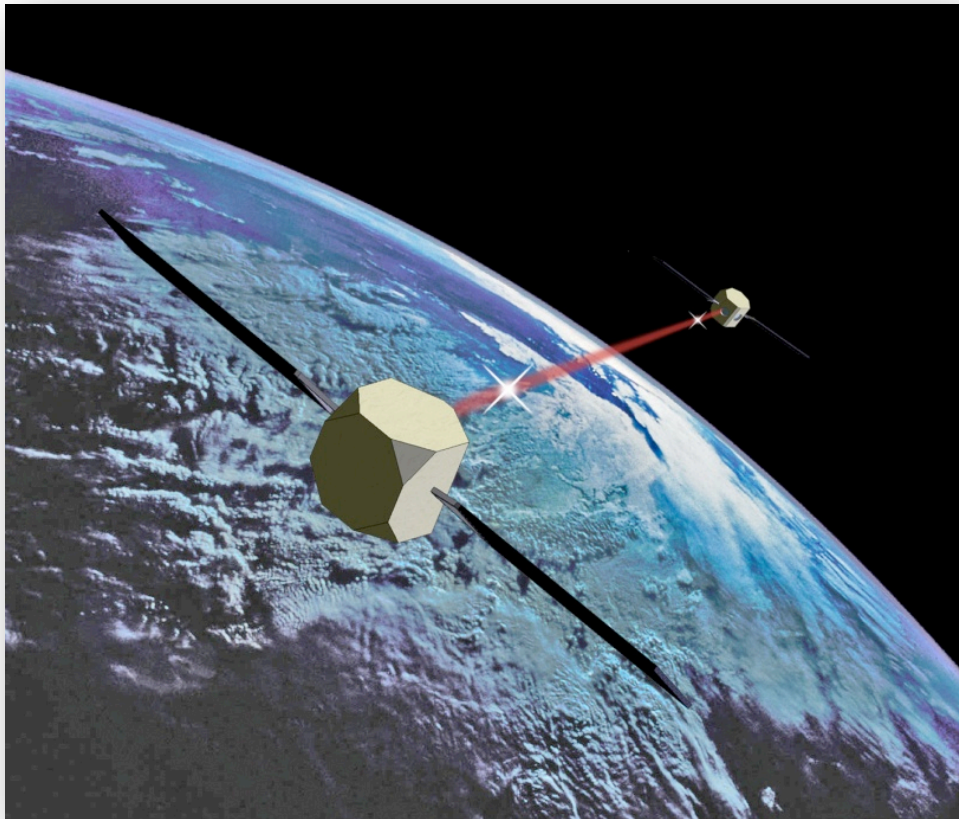
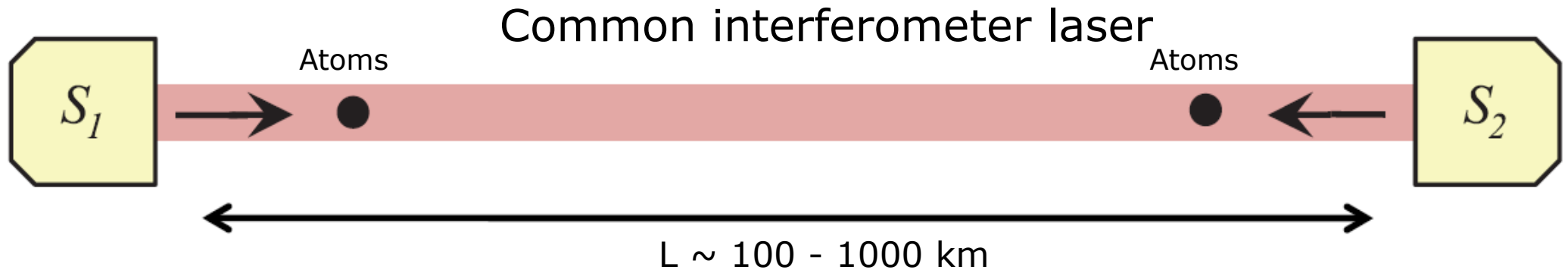


- *Differential scheme*

S. Dimopoulos, P. W. Graham, J. M. Hogan, M. A. Kasevich, S. Rajendran, *Atomic gravitational wave interferometric sensor*, *Phys. Rev. D* 78, 122002 (2008)



Satellite GW Antenna



JMAPS bus/ESPA
deployed



October 14, 2008

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Gravitational Waves Detection with Atom Interferometry

Conference

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[Schedule](#)

Organizers:

Guglielmo M. Tino, University of Firenze, Italy Flavio Vetrano, University of Urbino, Italy

Period: from 23-02-2009 to 24-02-2009

Deadline: 15-01-2009










Note: The number of participants is limited to 50 The participation fee for the Workshop is 150 Euros including registration, coffee-breaks, lunches and the social dinner. The fee should be paid cash on arrival at the registration desk

Abstract

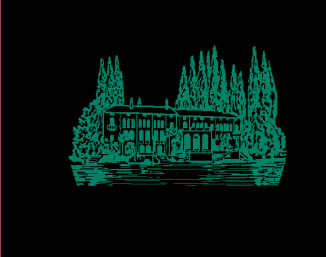
The possibility of using atom interferometers to detect gravitational waves is attracting increasing interest as an alternative to other detectors. Several papers were published discussing theoretical and experimental aspects. Although the results show that dedicated technological developments are still needed to achieve the required sensitivity values which are beyond those presently available, newschemes for atom interferometers, beam splitters and high flux coherent atomic sources could lead to an increase in sensitivity and make atom interferometers competitive with other gravitational wave detectors. The Workshop on "Gravitational Waves Detection with Atom Interferometry" will bring together scientists interested in theoretical and experimental aspects, to discuss different points of view and possible experimental implementations in Earth Laboratories.

Special issue on
Gravitational Waves Detection with Atom Interferometry
G.M. Tino, F. Vetrano, C. Laemmerzahl Editors,
General Relativity and Gravitation **43**, 1901 (2011)

Applications of new quantum sensors based on atom interferometry

- Measurement of fundamental constants $\begin{matrix} \longrightarrow & G \\ \longrightarrow & \alpha \end{matrix}$ 
- New definition of kg 
- Test of equivalence principle 
- Measurement of the gravitational redshift 
- Tests of quantum gravity 
- Short-distances forces measurement 
- Search for electron-proton charge inequality 
- New detectors for gravitational waves ? 
- Development of transportable atom interferometers $\begin{matrix} \longrightarrow & \text{geophysics} \\ \longrightarrow & \text{space} \end{matrix}$ 

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Atom Interferometry

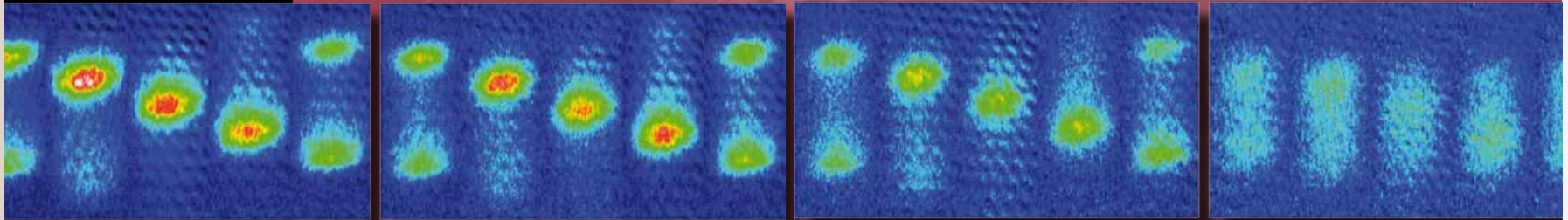
edited by G. M. Tino and M. A. Kasevich

15-20 July 2013

Villa Monastero
Varenna, Lake Como

edited by G. M. Tino and M. A. Kasevich

Atom Interferometry



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