

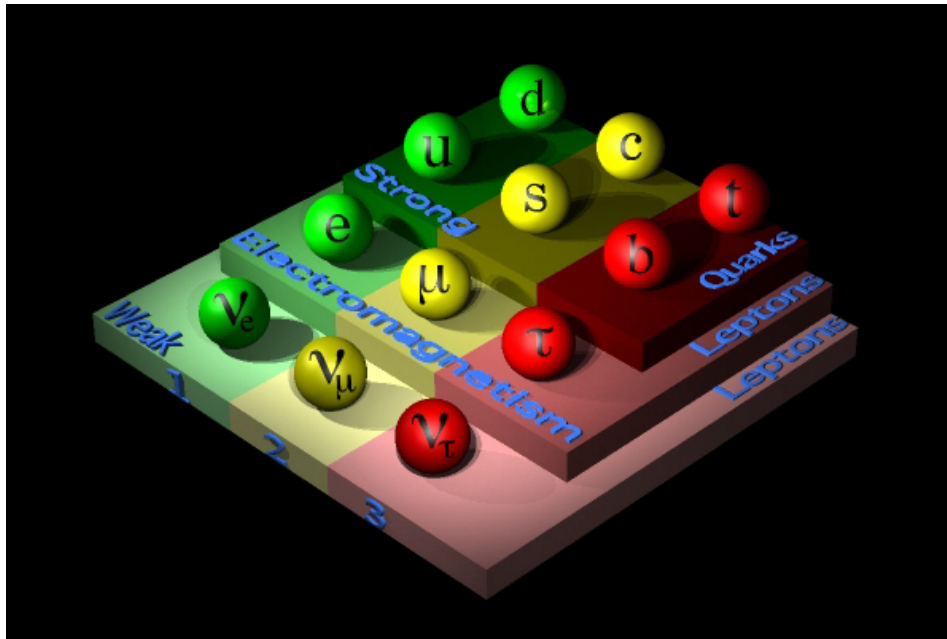
Use of neutrinos while still learning about them



David Lhuillier
CEA-Saclay



Fundamental properties of ν 's



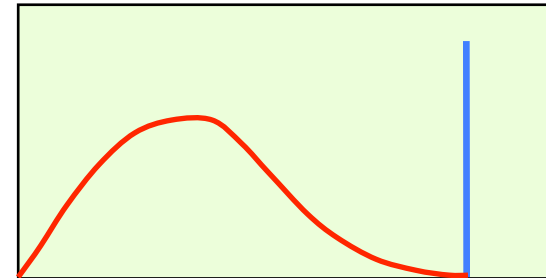
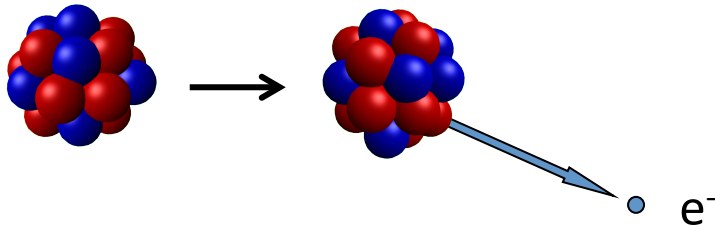
- 3 flavors of neutrinos
- Interact only via the weak force and gravitation
- $\sigma_{\text{int}} \sim 10^{-43} \text{ cm}^2$ @ 1 MeV

→ Very penetrating particles

- Challenging detection technics
- Remote observation of intense sources (possibly inaccessible otherwise)

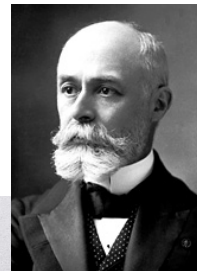
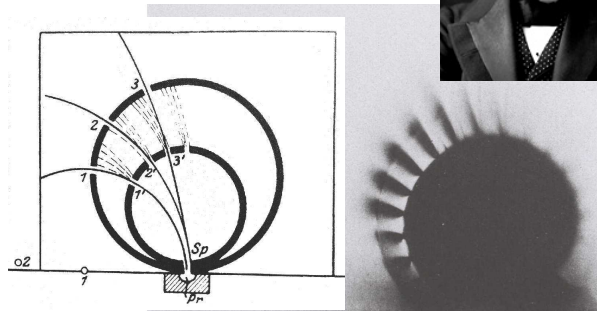
Neutrino footprints

Missing energy in β decays

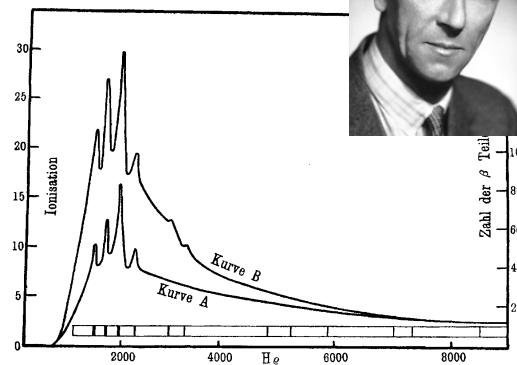


Measurements:

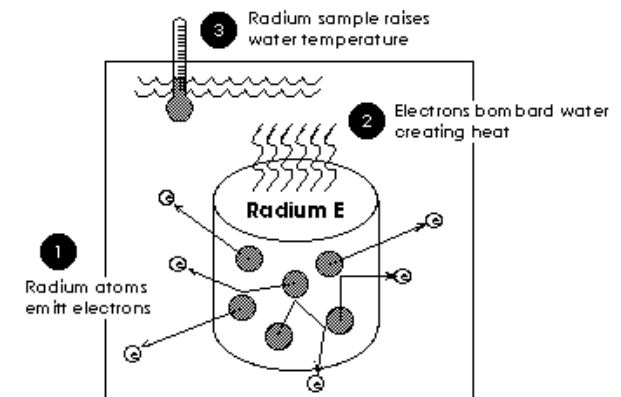
Becquerel, 1901



Chadwick, 1914



Ellis & Wooster, 1927



Neutrino footprints

Discovery at nuclear reactors

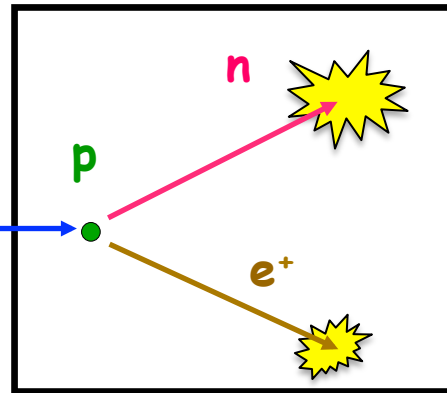
Savannah River
Nuclear Reactor



1 GW, $\sim 10^{20}$ ν /s

← 25 m →

- Discovery in 1956
- Nobel prize 1995



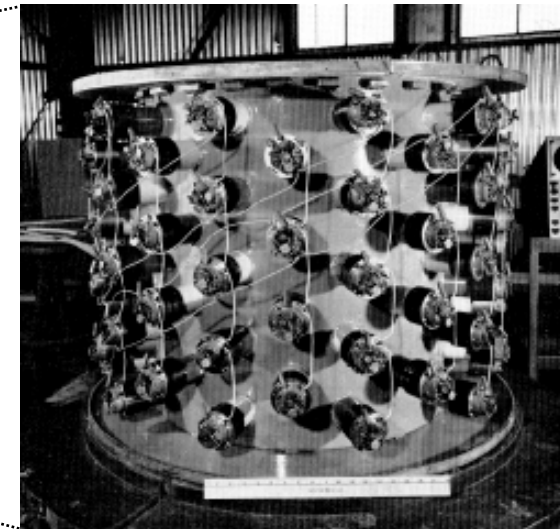
Inverse β decay process



Cowan



Reines

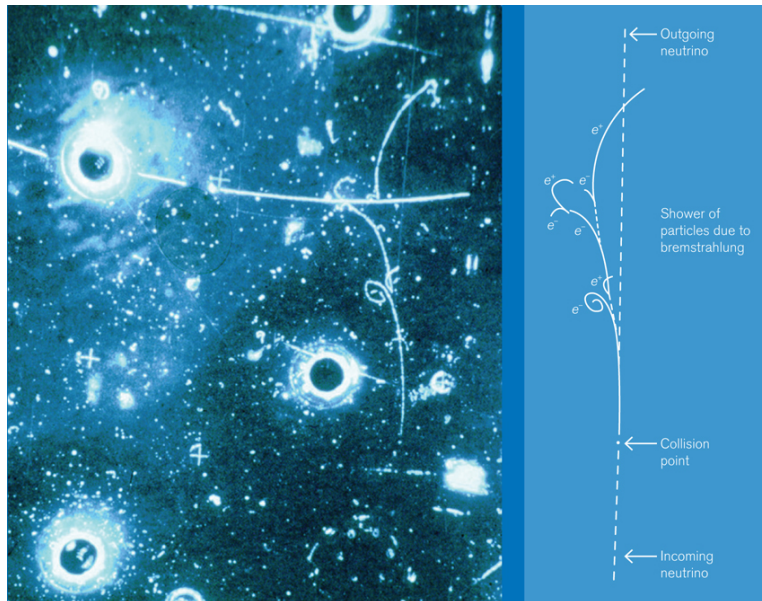
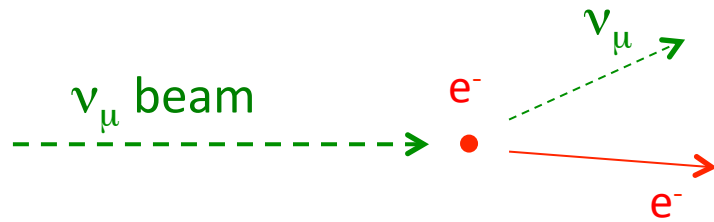


1 m³ liquid scintillator

~ 3 ν /hour correlated to reactor operation

Neutrino footprints

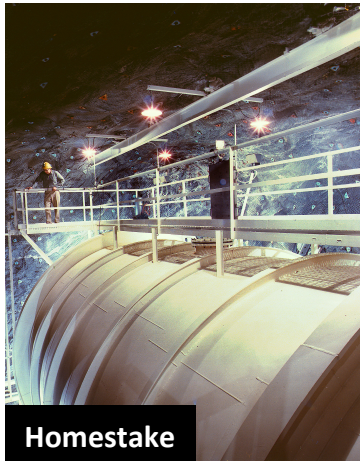
Observation of a neutrino weak neutral current



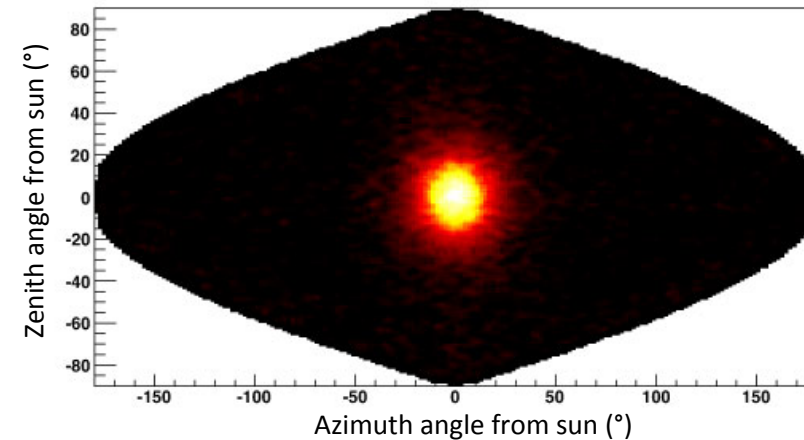
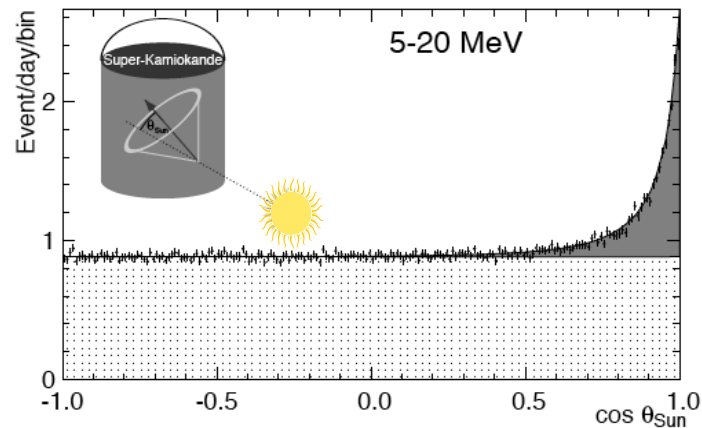
- Evidence of the weak neutral current in the Gargamelle bubble chamber at CERN – 1972

Neutrino footprints

Solar neutrinos

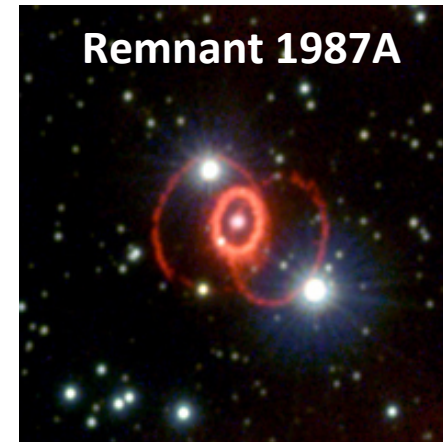
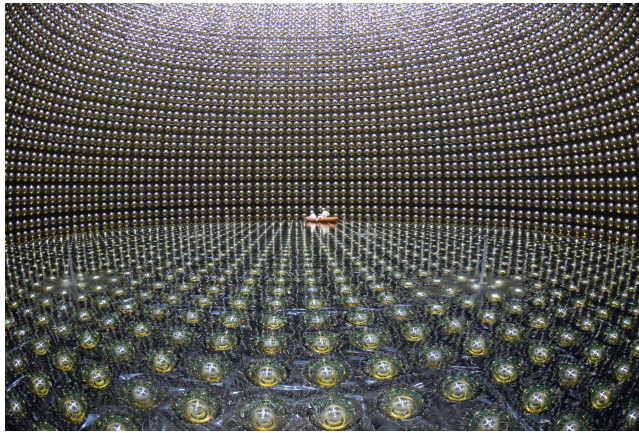


- Kamiokande can detect the Cerenkov light emitted by recoiling electrons struck by the most energetic solar neutrinos (^8B).
- The reconstructed direction of incoming neutrinos builds up the first neutrino photo of the sun, from deep underground (SK data)



Neutrino footprints

Supernova neutrinos



- The neutrino flux of SN1987A was detected by the large neutrino detectors

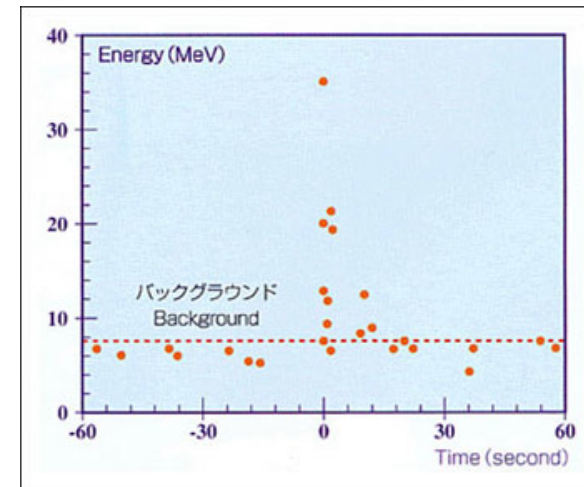


R. Davis

*Nobel 2002
Cosmic ν*

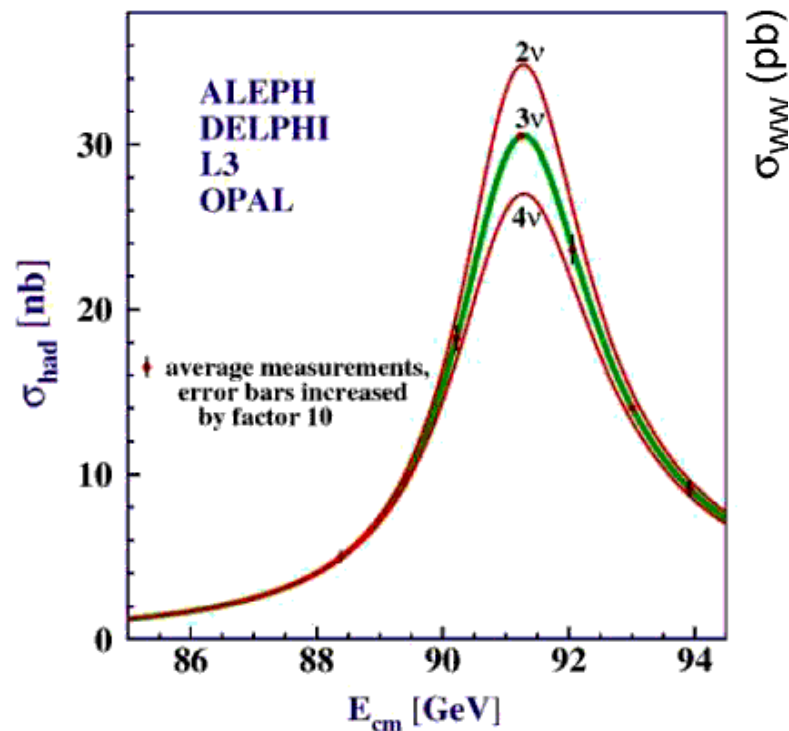


M. Koshihara



Neutrino footprints

Width of the Z boson decay



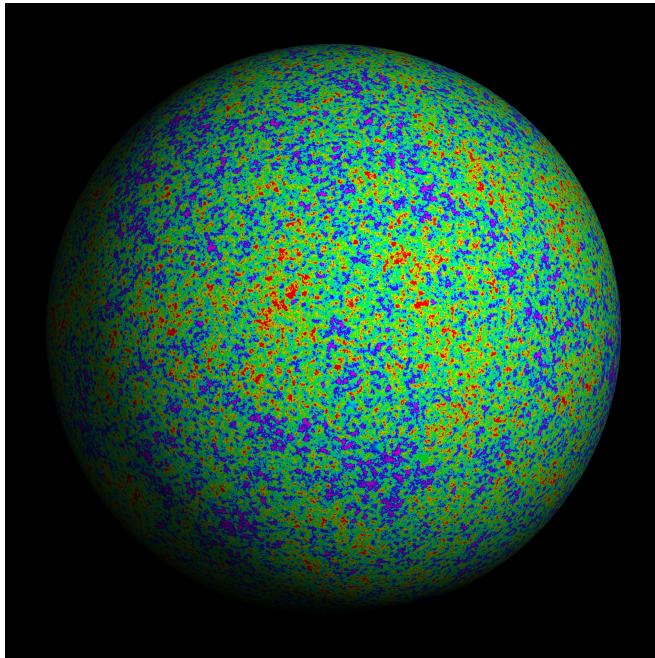
- $Z \rightarrow \nu \bar{\nu}$ decays contribute to the observed total width of the Z boson

- $N_{\nu} = 2.9840 \pm 0.0082$

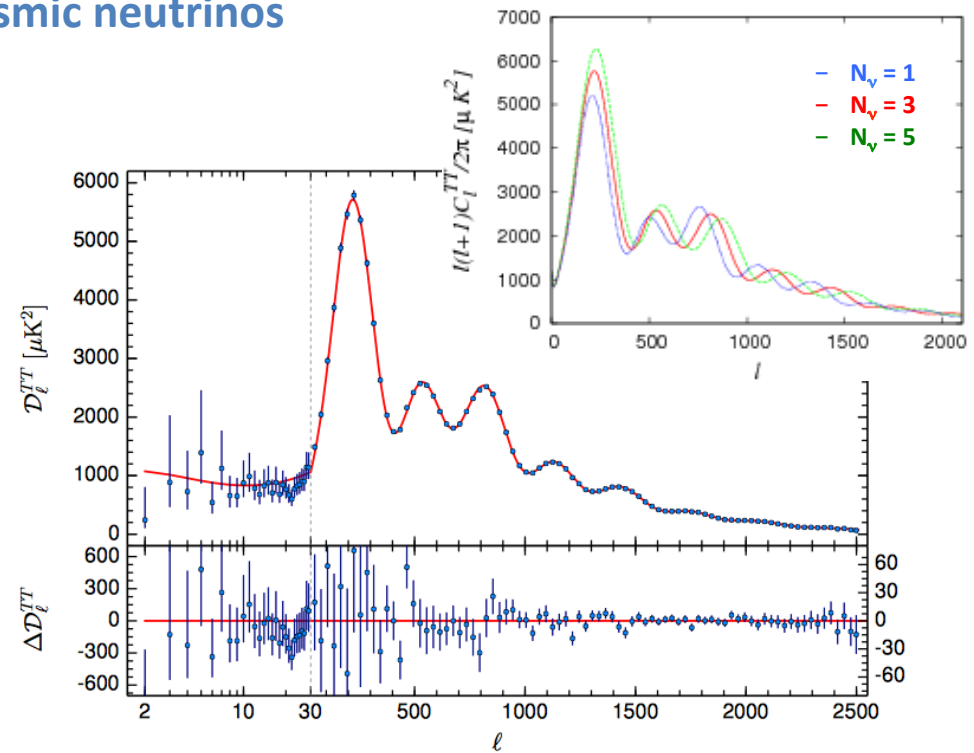
→ 3 neutrino flavors couple to the Z^0 boson

Neutrino footprints

Cosmic neutrinos



Map of the Cosmic Microwave Background
from the Planck satellite



<http://arxiv.org/abs/1502.01589>

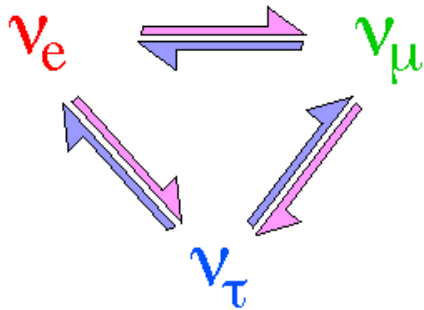
$$N_{\text{eff}} = 3.15 \pm 0.23$$

Neutrinos are the most abundant matter particles in the universe

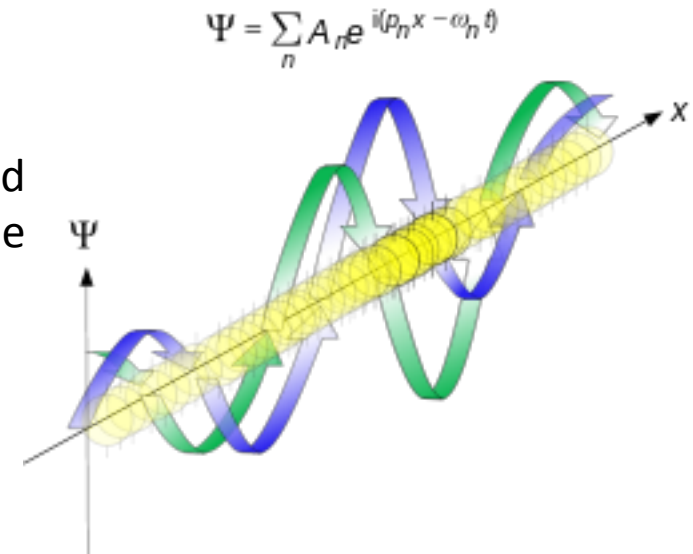
Neutrino oscillations

Propagation of coherent states

- 3 neutrino states carrying the same conserved quantum numbers
→ Coherent superposition of the 3 eigen states could lead to an oscillation phenomenon along the propagation of the neutrinos.

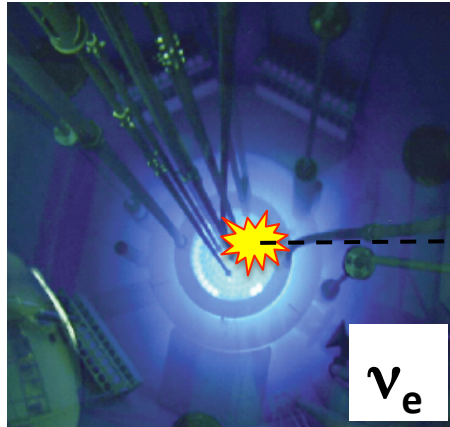


B. Pontecorvo



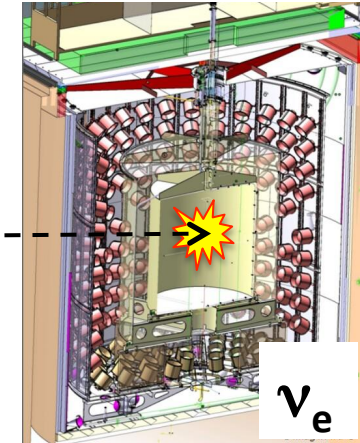
- Requires different masses hence at least two non zero.
- The coherence of the wave packets over macroscopic length scales requires Small mass splitting (prevents the oscillations of charged leptons)

Formalism



Interaction / Flavor

ν_1, ν_2, ν_3



Interaction / Flavor

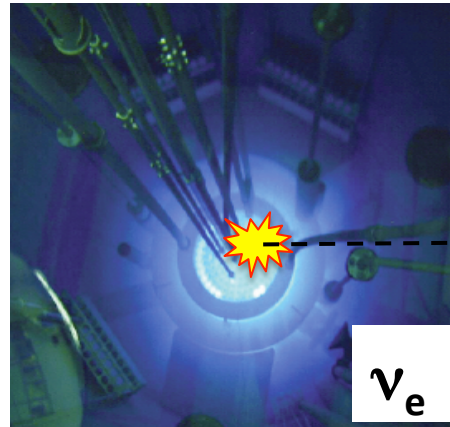
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{bmatrix} \theta \end{bmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \quad \dots \quad \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix} \begin{bmatrix} \theta \end{bmatrix}^{-1} = \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix}$$

Δm^2

3 mixing angles: $\theta_{12}, \theta_{13}, \theta_{23}$

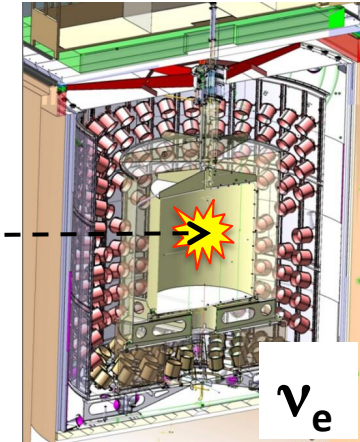
2 non-zero square mass splitting : $\Delta m_{21}^2, \Delta m_{31}^2$

Mixing parameters



Flavor

ν_1, ν_2, ν_3



Mass

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Large
mixing
angles

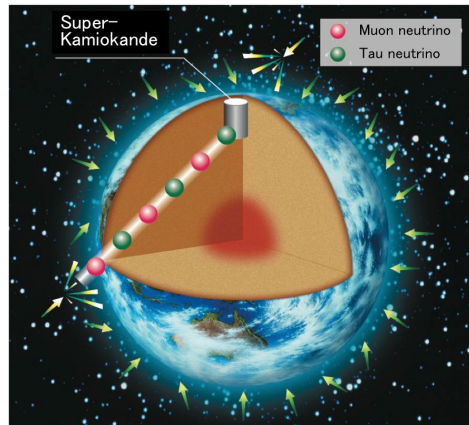
$$\begin{aligned} \theta_{12} &= 33.5 \pm 1.5^\circ \\ \theta_{23} &= 40.4 \pm 5^\circ \\ \theta_{13} &= 8.42 \pm 0.26^\circ \end{aligned}$$

2 mass
sectors

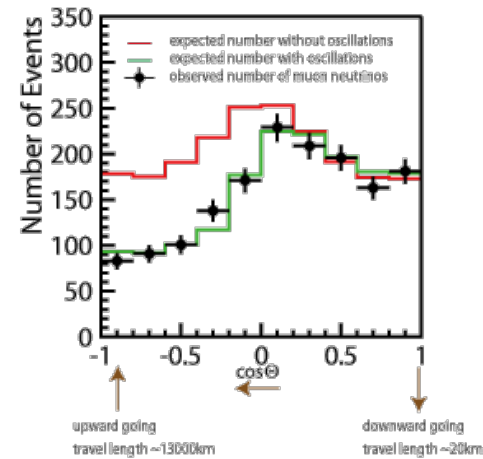
$$\begin{aligned} \Delta m_{21}^2 &= 7.6 \pm 0.2 \cdot 10^{-5} \text{ eV}^2 \\ |\Delta m_{31}^2| &= 2.5 \pm 0.1 \cdot 10^{-3} \text{ eV}^2 \end{aligned}$$

Oscillations

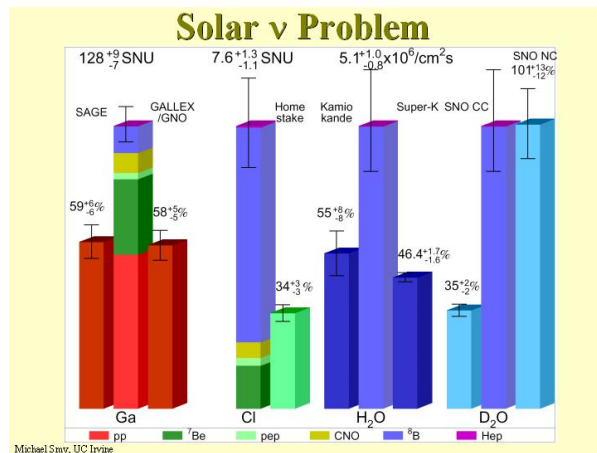
Awarded the Nobel prize 2015



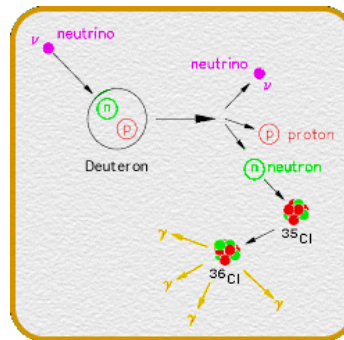
Super K observes the oscillation of atmospheric ν



T. Kajita



SNO solves the solar ν problem



A.B. McDonald

Reactor neutrinos

Neutrinos from fission

Per fission:

- 200 MeV
- 6 neutrinos

→ Intense source
 $2 \cdot 10^{20} \text{ } \nu/\text{s/GW}_{\text{th}}$

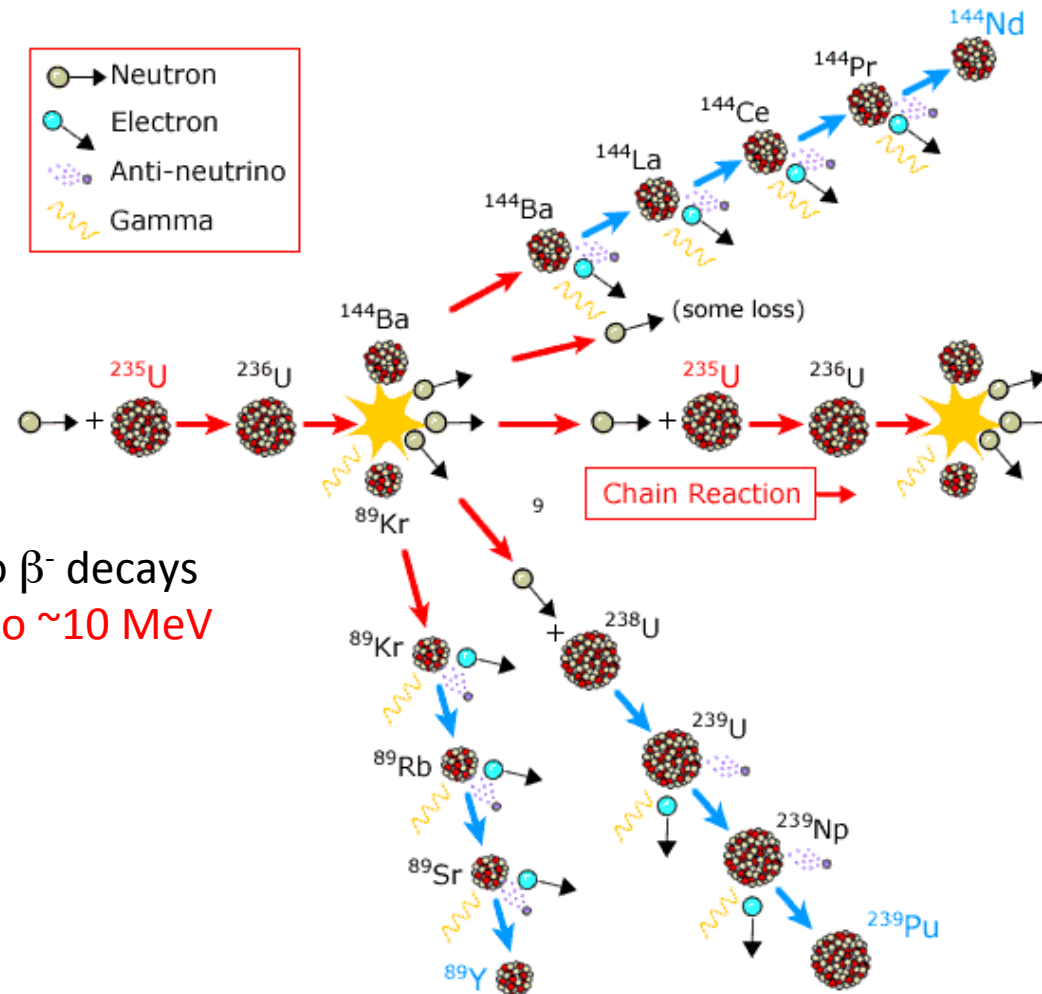
Neutron rich FPs undergo β^- decays
 → Pure source of $\bar{\nu}_e$, up to $\sim 10 \text{ MeV}$

Complex total spectrum

$$S_k(E) = \sum \text{all FPs}$$

~ 800 nuclei

$\sim 10\,000$ β -branches



The guts of $S_k(E)$

Sum of all fission products' activities

$$S_k(E) = \sum_{fp=1}^{N_{fp}} \mathcal{A}_{fp}(T) \times S_{fp}(E)$$

Sum of all β -branches of each fission product

$$S_{fp}(E) = \sum_{b=1}^{N_b} BR_{fp}^b \times S_{fp}^b(Z_{fp}, A_{fp}, E_{0fp}^b, E)$$

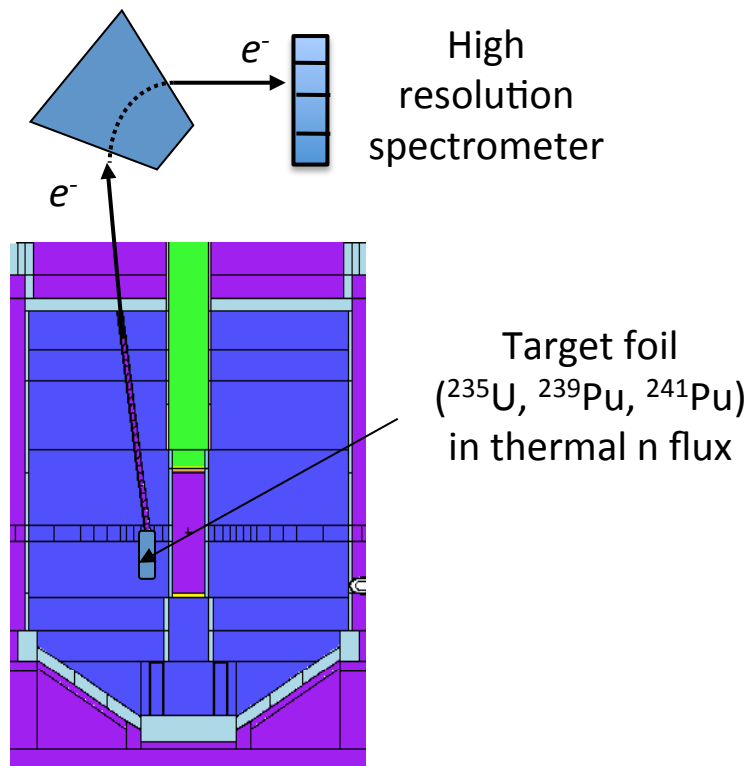
Theory of β -decay

$$S_{fp}^b = \underbrace{K_{fp}^b}_{\text{Norm.}} \times \underbrace{\mathcal{F}(Z_{fp}, A_{fp}, E)}_{\text{Fermi function}} \times \underbrace{pE(E - E_{0fp}^b)^2}_{\text{Phase space}} \\ \times \underbrace{C_{fp}^b(E)}_{\text{Shape factor}} \times \underbrace{\left(1 + \delta_{fp}^b(Z_{fp}, A_{fp}, E)\right)}_{\text{Correction}}$$

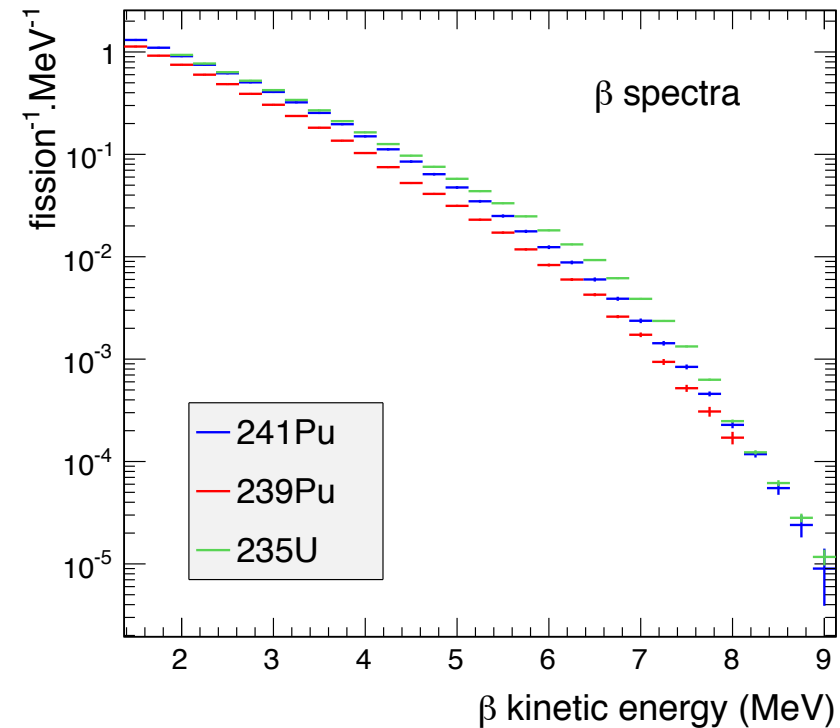
$$\delta_{fp}^b = G_{\nu(QED)} + L_{0(\text{coulomb size})} + C_{(\text{weak size})} + S_{(\text{screening})} + \delta_{WM(\text{weak magnetism})}$$

The ILL electron data

Total electron spectra from the β -decays of ^{235}U , ^{239}Pu and ^{241}Pu fission products.



ILL research reactor
(Grenoble, France)

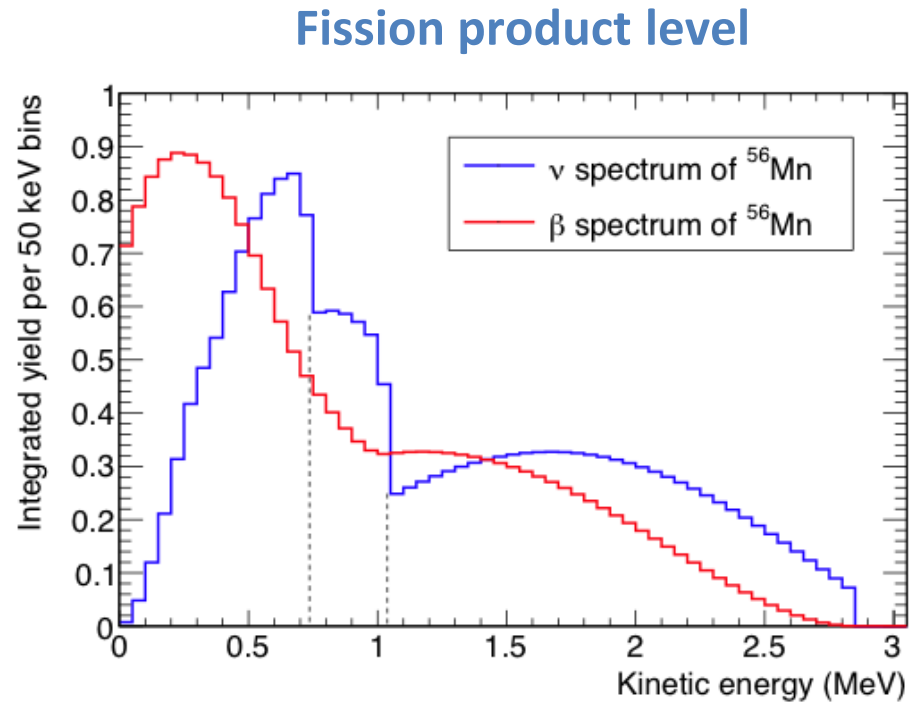
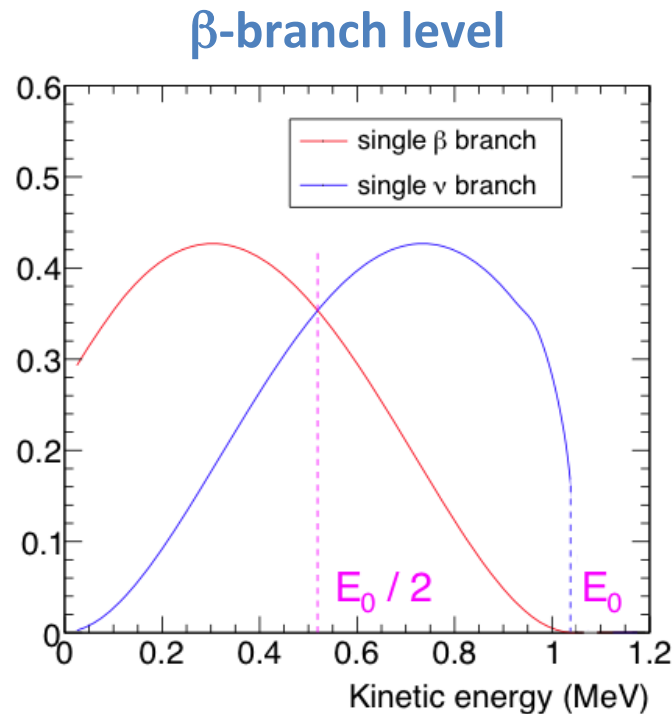


Unique reference of β fission spectra

K. Schreckenbach et al., Phys. Let. B218,365 (1989)+ refs therein

$e^- \rightarrow \nu$ conversion

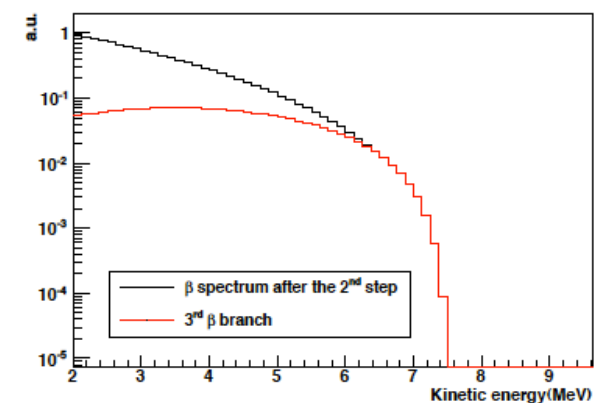
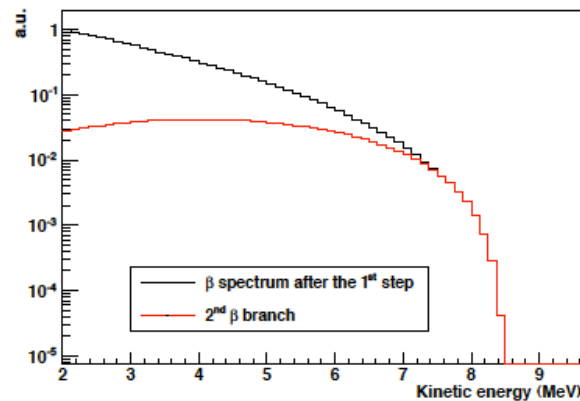
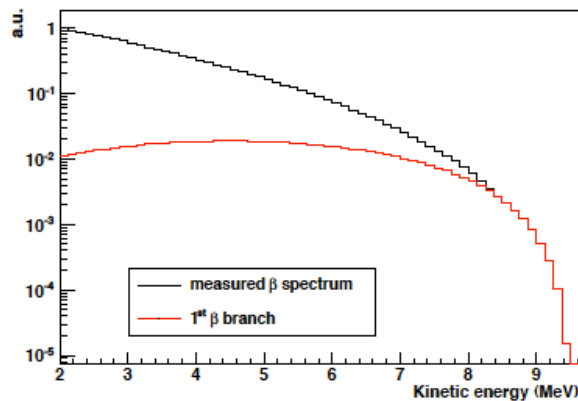
- Exact conversion requires complete knowledge of all contributing β -transitions



- Lot of relevant quantities:** Z , A , End-points, J^π , nuclear matrix elements, branching ratios, fission yields, life time... **but scarce data as E_0 increases**

Conversion of reference e^- spectra

1- Fit total e^- spectrum with a sum of 30 effective β branches determined by iterative method (instead of $\sim 10,000+$ real branches)

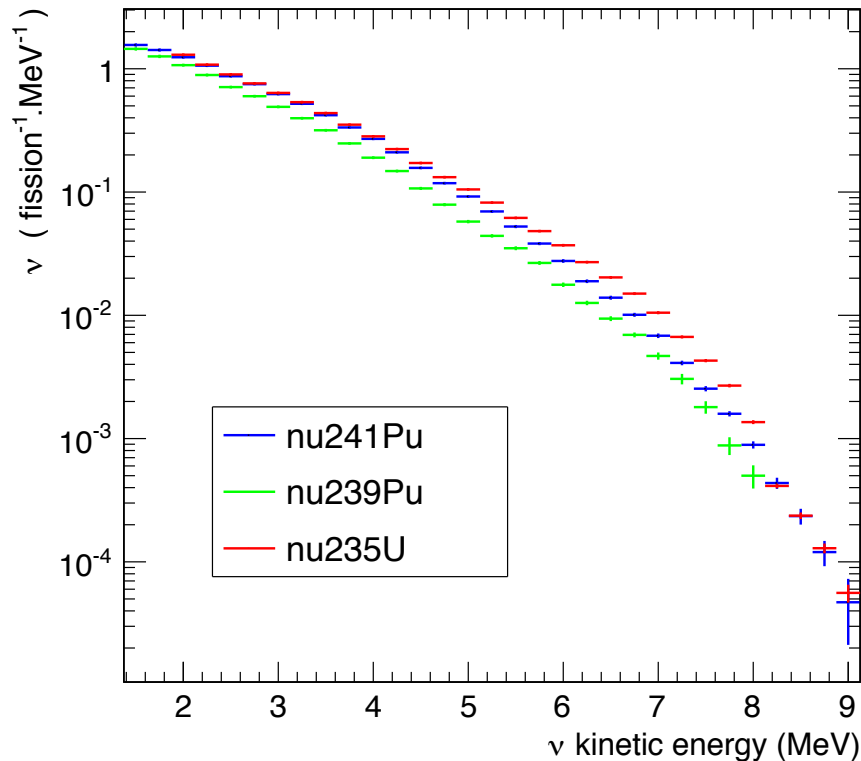


2- Convert each effective e^- branches to ν branches

3- Sum all converted ν branches to get total ν spectrum

K. Schreckenbach et al., Phys. Lett. 99B, 251 (1981).

Reference fission $\bar{\nu}$ spectra

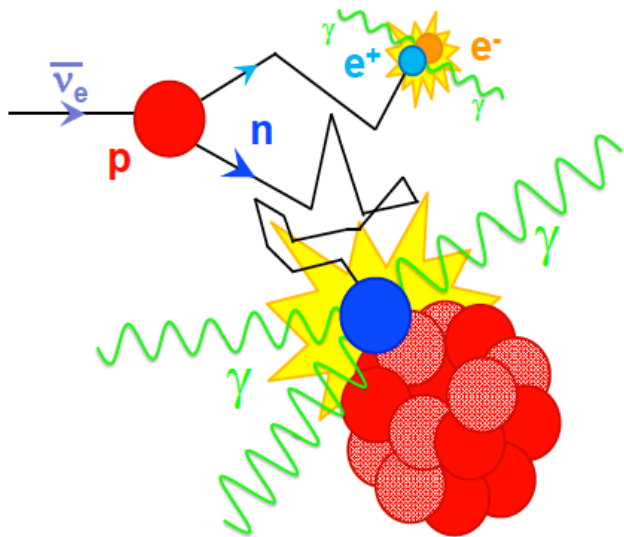


- Antineutrino spectra converted from the β spectra measured at ILL
- Reference spectra over the last 25 years

Detection

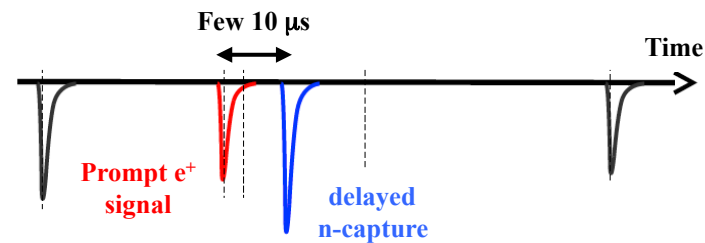
Inverse Beta Decay

$$\bar{\nu}_e + p \rightarrow e^+ + n$$



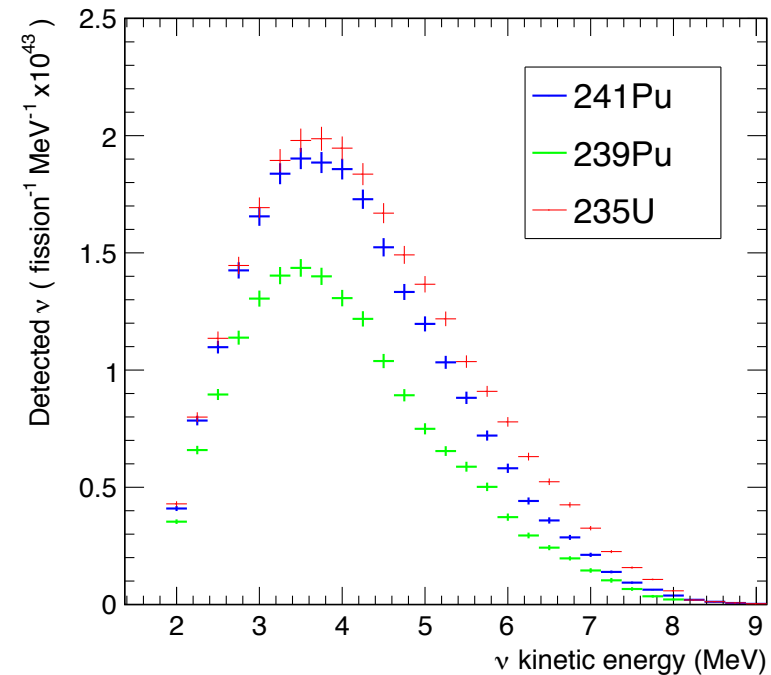
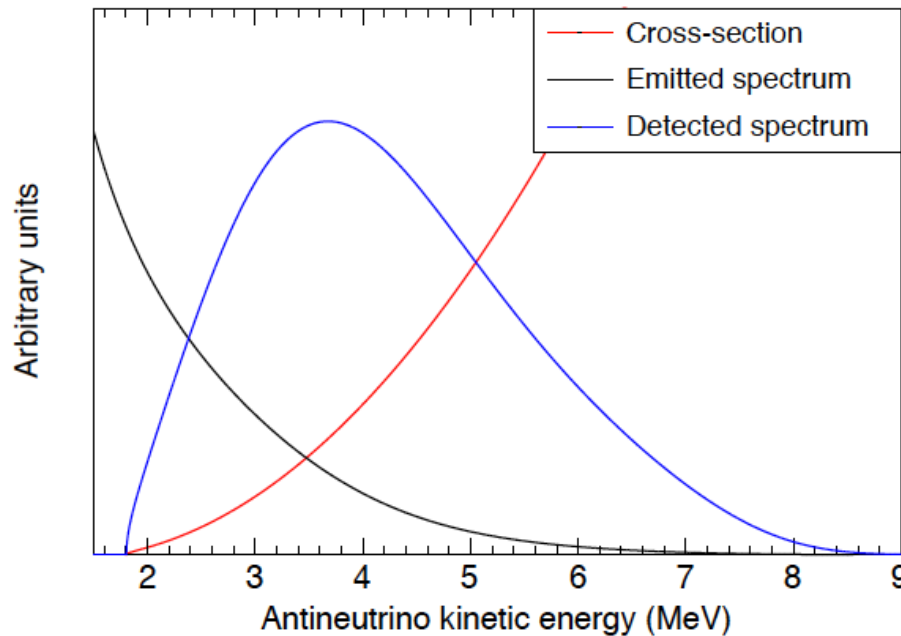
1.8 MeV reaction threshold

- A neutral current process would have to compete with any single e- recoil induced by natural radioactivity
- IBD process provides a selective signal sequence



- An organic scintillator doped with a neutron absorber efficiently combines a proton target with the detection of prompt and delayed signals

Prediction of detected reactor neutrino spectra

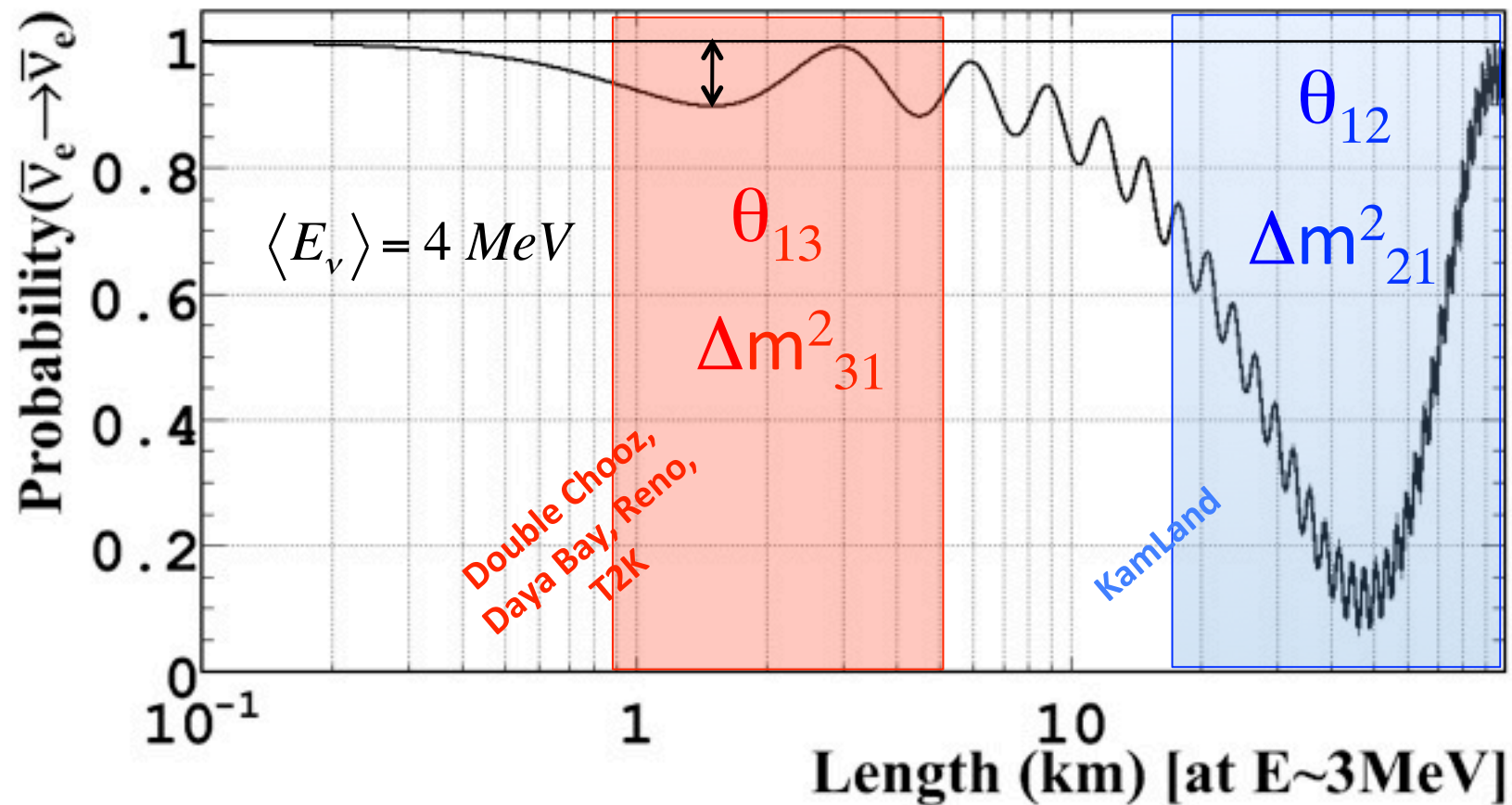


- Error matrix determined from uncertainties of ILL measurements and conversion procedure.
- 239Pu fissions lead to less detected neutrinos than 235U fissions?

ν oscillations at reactors

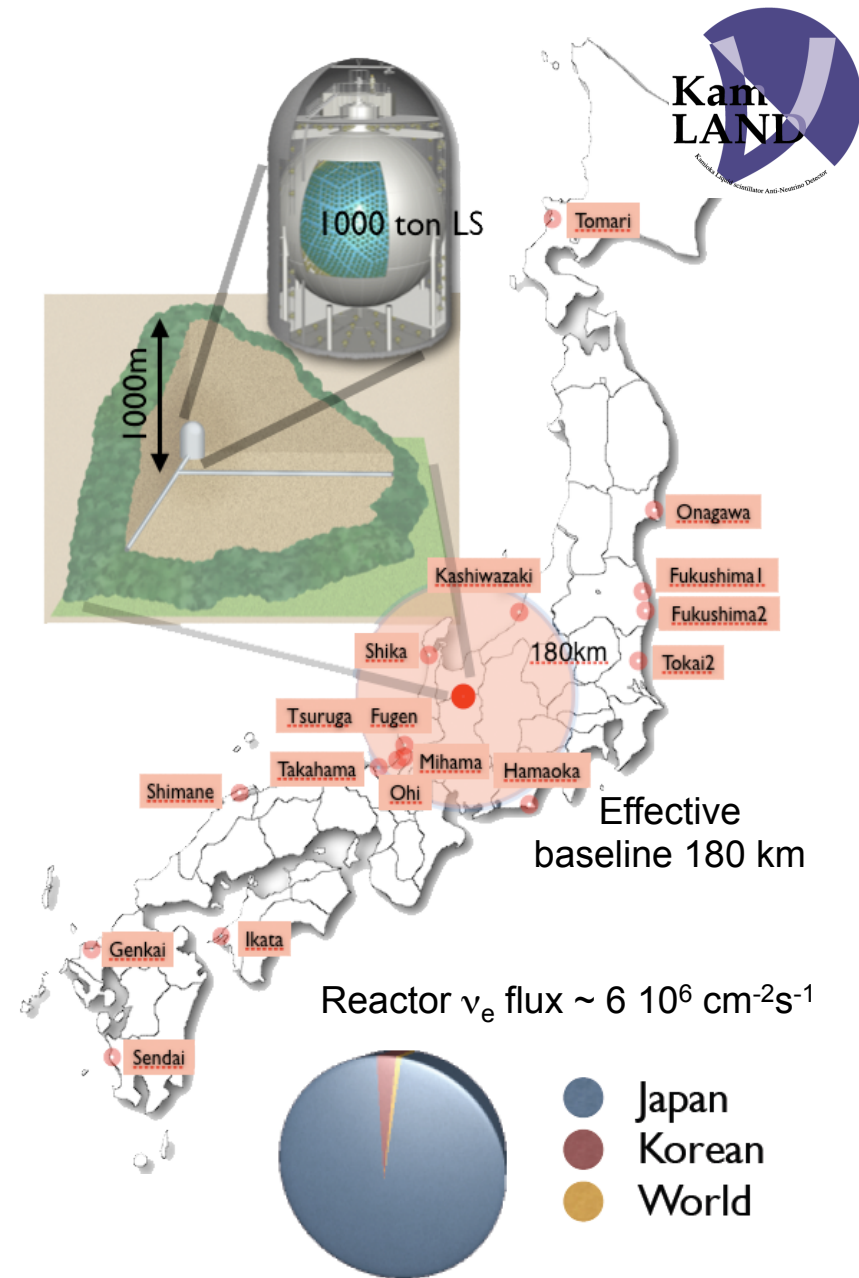
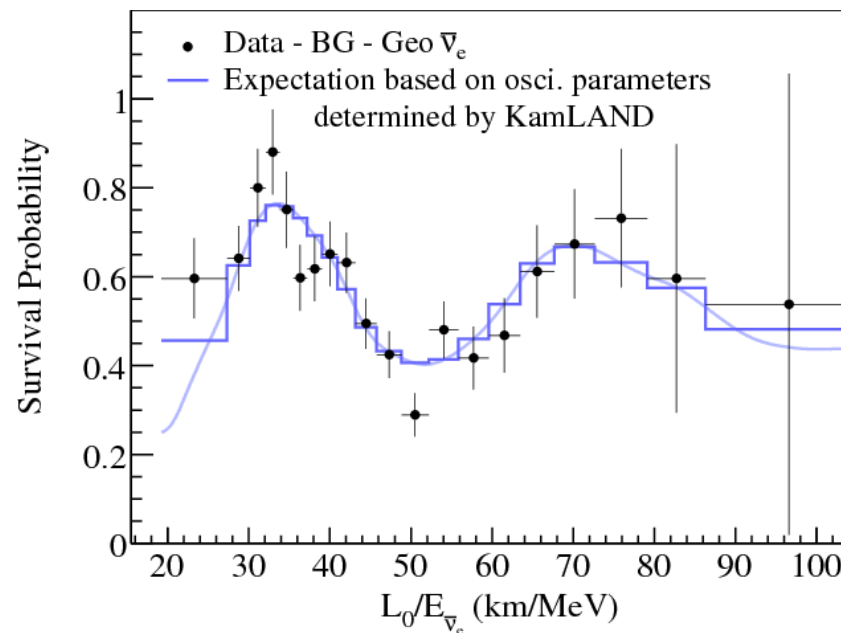
Below μ^+, τ^+ production threshold \rightarrow look for $\bar{\nu}_e$ disappearance

$$P(\nu_e \rightarrow \bar{\nu}_e) = 1 - \sin^2(2\theta) \sin^2(\Delta m^2 L / 4E)$$



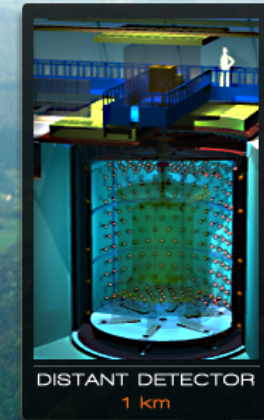
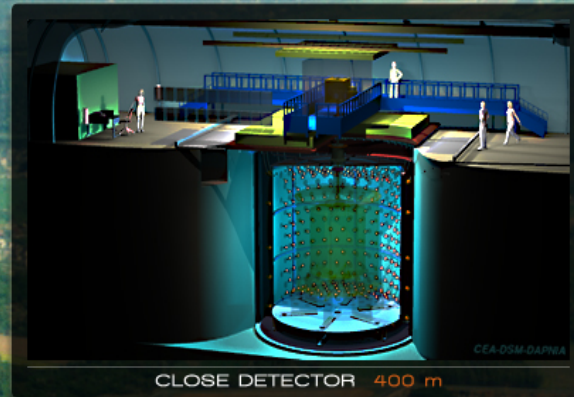
KamLAND

- Large liquid scintillator detector at ~180 km of many Japanese reactors
- Confirmation of the solar oscillation and accurate measurement of Δm^2_{12}



Double Detector concept @ the Chooz Site

Comparison of the ν spectra in 2 identical detectors



WEST REACTOR

EAST REACTOR



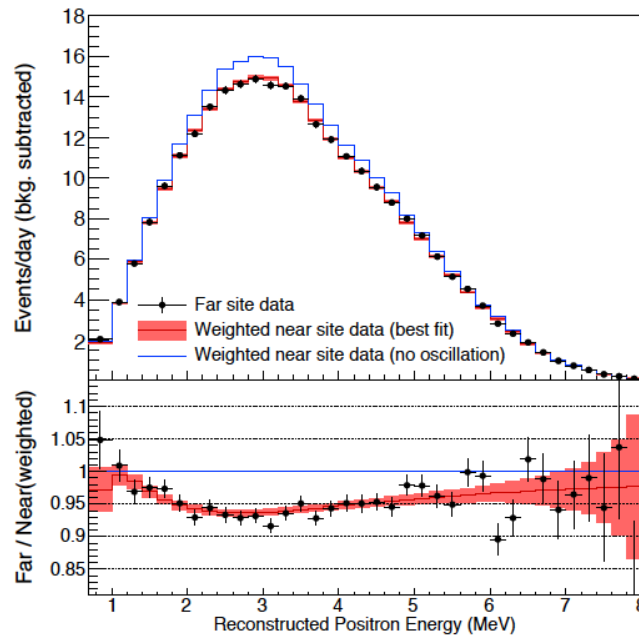
ν_e

2 Chooz reactors
4.25 GW_{th} each

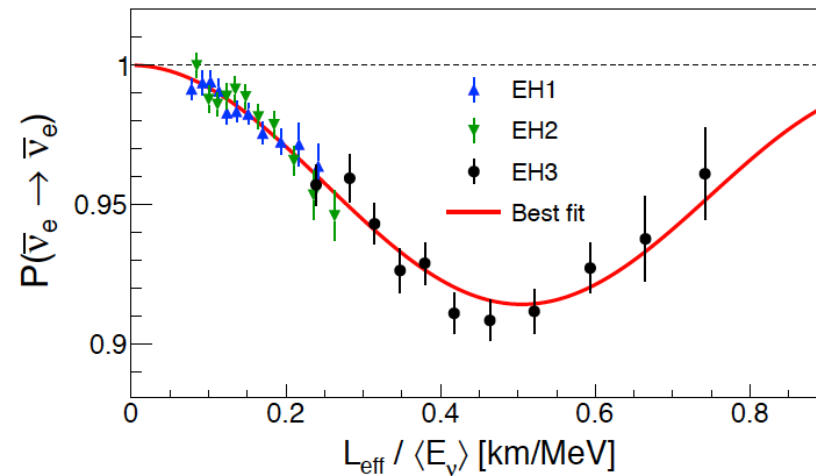
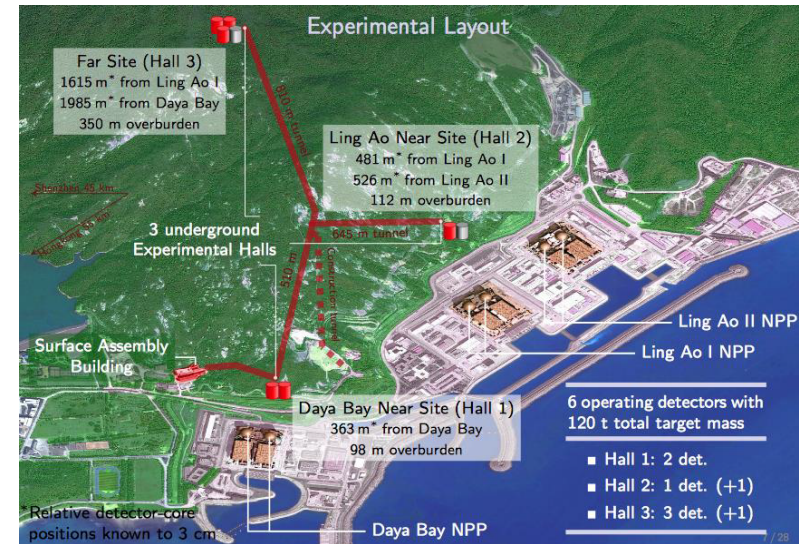
Last mixing angle measured: θ_{13}

$$\sin^2 2\theta_{13} = 0.084 \pm 0.005$$

→ Era of neutrino precision measurements

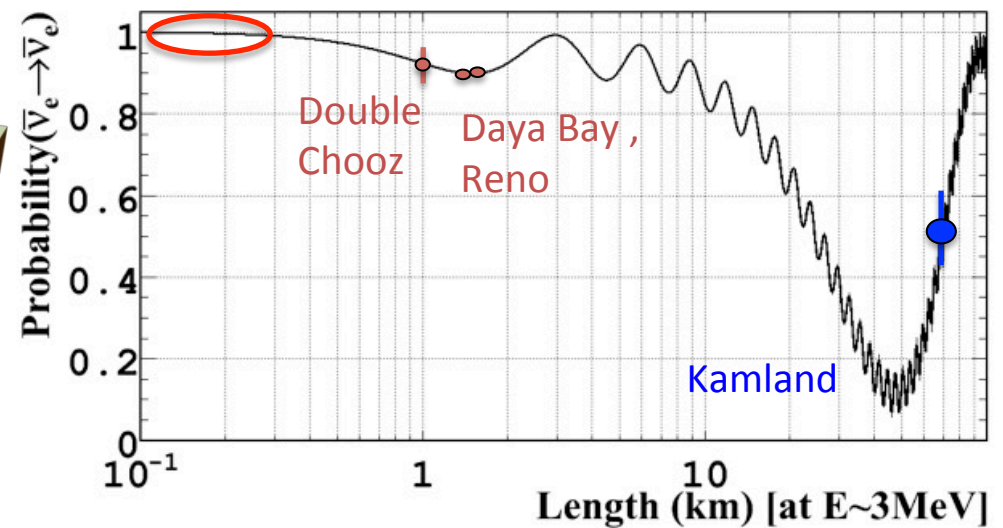
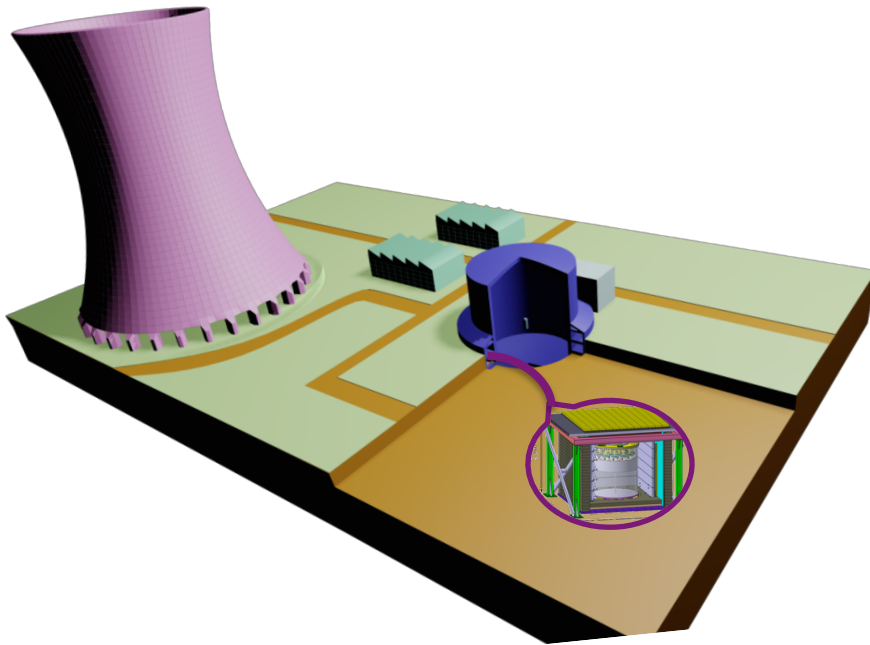


[arXiv:1505.03456](https://arxiv.org/abs/1505.03456)

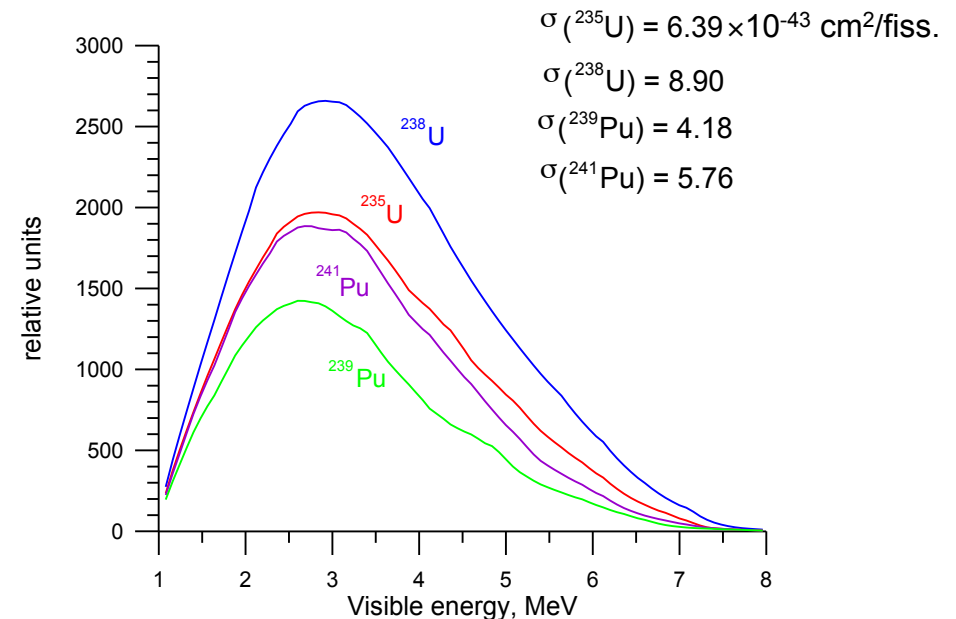
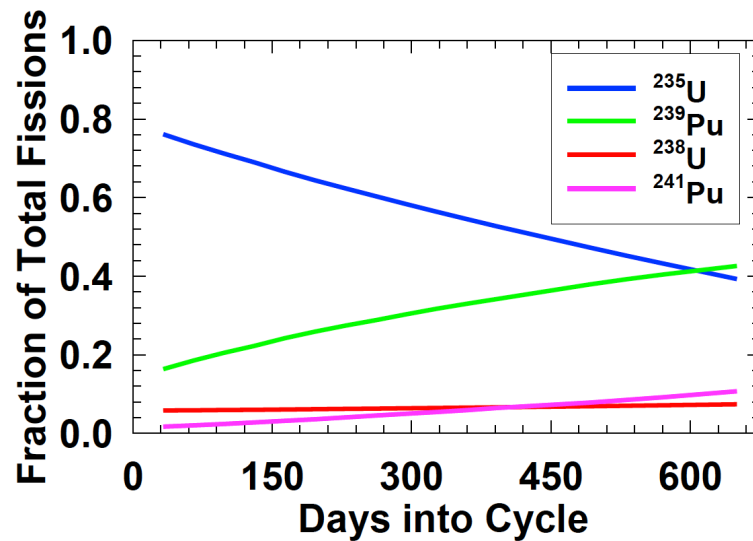


Application of neutrinos to the surveillance of nuclear reactor

- Monitor the operation of a reactor using a small (1m³ scale) neutrino detector at short distance (few 10 m) from the core



Nuclear Fuel Evolution



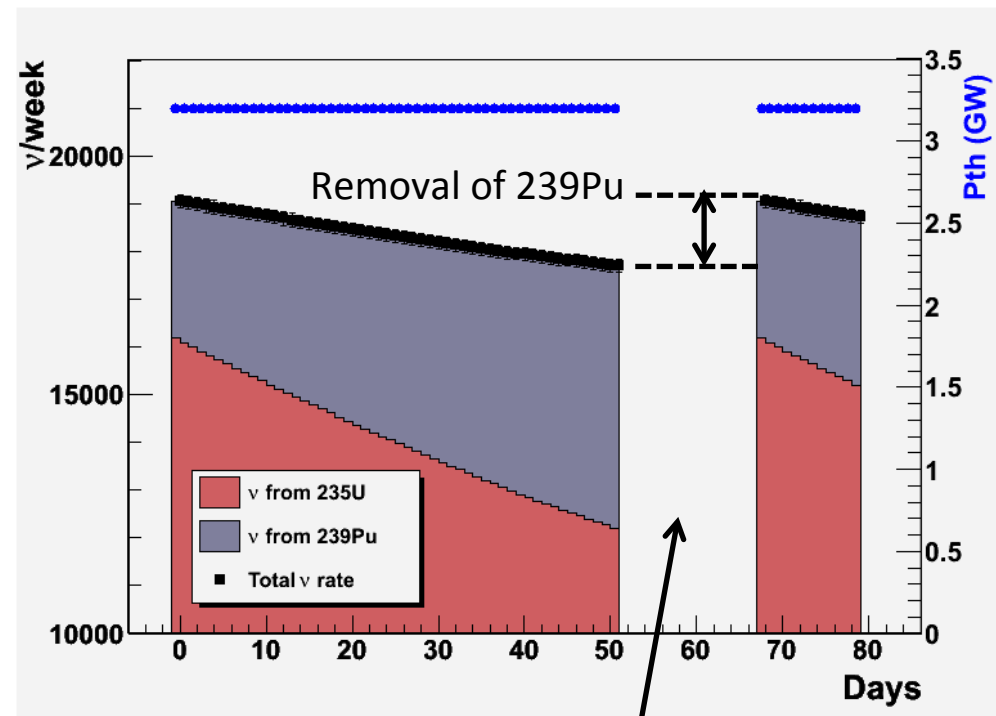
- Antineutrinos encode information about fissile production and consumption in reactors through the burnup effect

Sensitivity to Pu content

1m³ LS target, $\epsilon_{\text{det}}=50\%$
25 m from a 3 GW_{th} core

	²³⁵ U	²³⁹ Pu
E / fission	201.7 MeV	210.0 MeV
$\langle E_\nu \rangle$ ($E_\nu > 1.8$ MeV)	2.94 MeV	2.84 MeV
# ν / fission ($E_\nu > 1.8$ MeV)	1.92	1.45
$\langle \sigma_{\nu \text{ int}} \rangle$	$\approx 3.2 \cdot 10^{-43}$ cm ²	$\approx 2.76 \cdot 10^{-43}$ cm ²

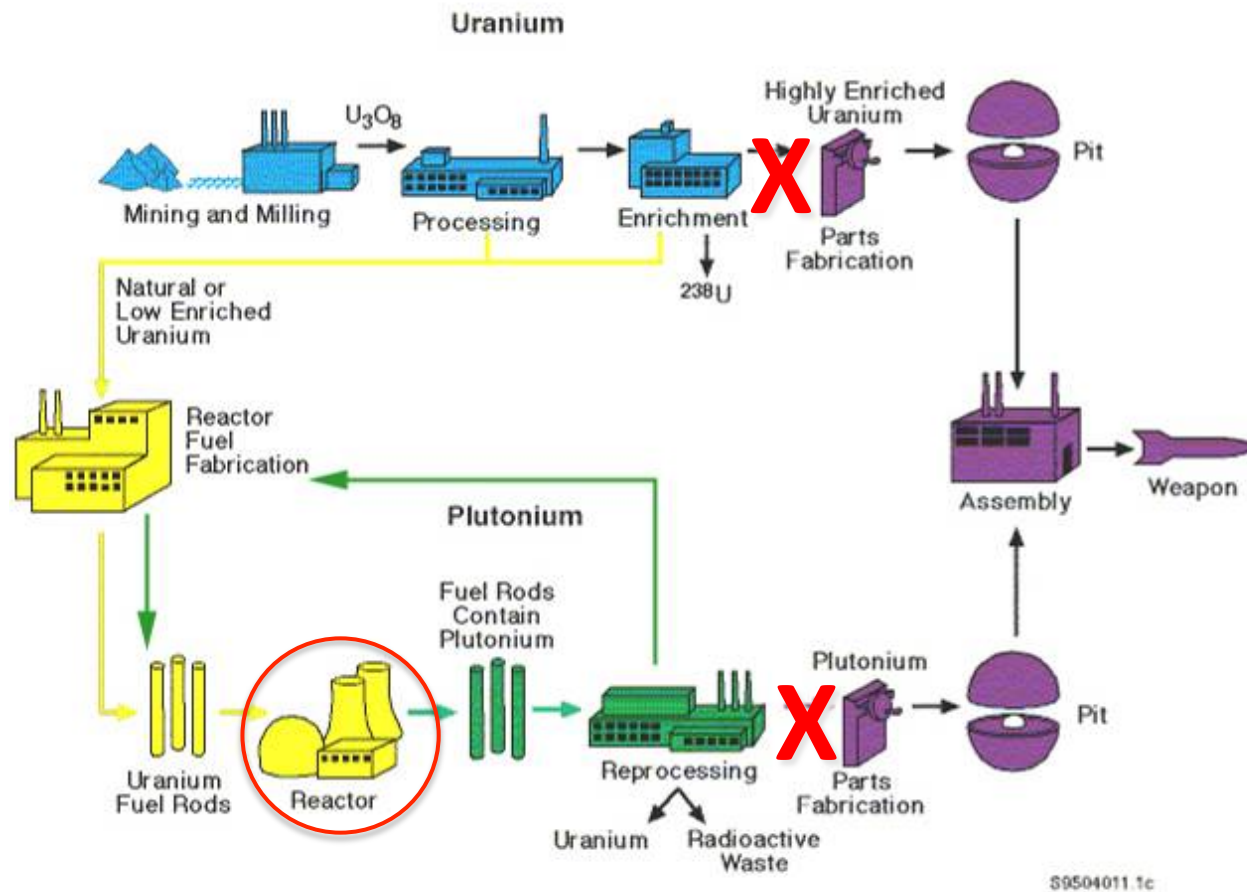
$$\frac{\# \nu \text{ int } (^{235}\text{U})}{\# \nu \text{ int } (^{239}\text{Pu})} = \frac{210.0}{201.7} \times \frac{1.92}{1.45} \times \frac{3.2}{2.76} = 1.60$$



Control of Fissile Material



IAEA goal: detect diversion of fissile material from peaceful to military programs



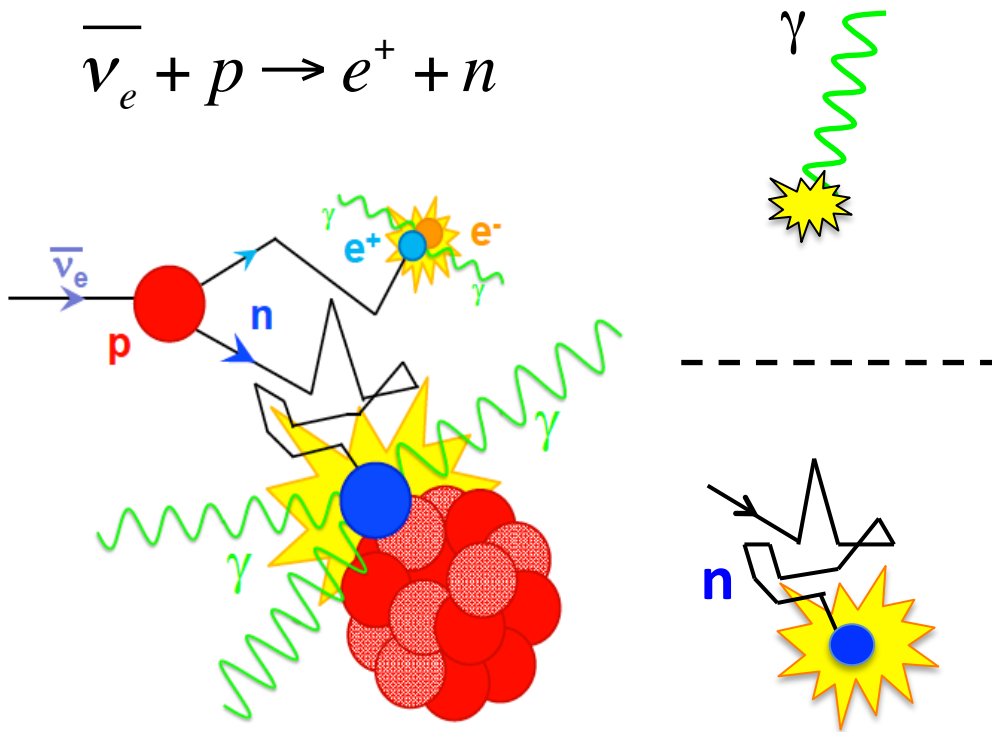
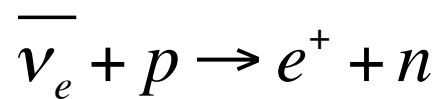
$\bar{\nu}_e$ measurements could monitor the fissile inventories in operating reactors

Specifications of a neutrino monitor

- Compact, portable (*easy installation, operate close to surface*)
 - Safe (*no connection with the core or the primary circuit, no impact on reactor operation and safety*)
 - Simple (*operated by non-experts*)
 - Remotely controlled (*no on site inspection, little maintenance*)
-
- Challenging simplification w.r.t state of the art neutrino detectors
 - Challenging mitigation of external background

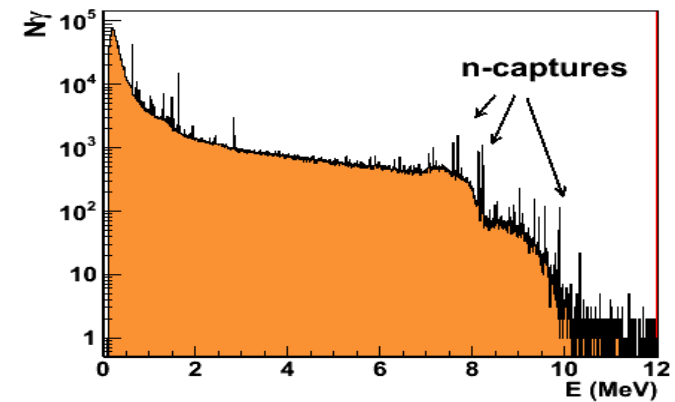
Background

Inverse Beta Decay



Accidental $[\gamma-n_{th}]$ coinc

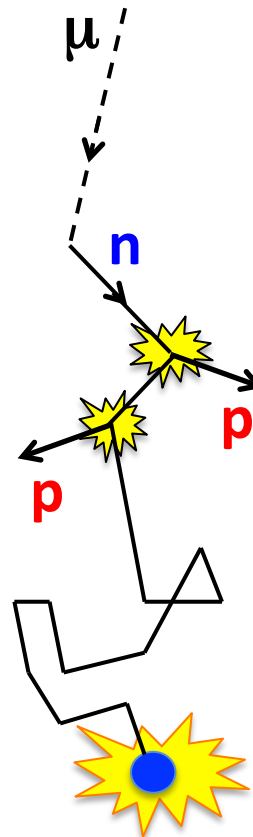
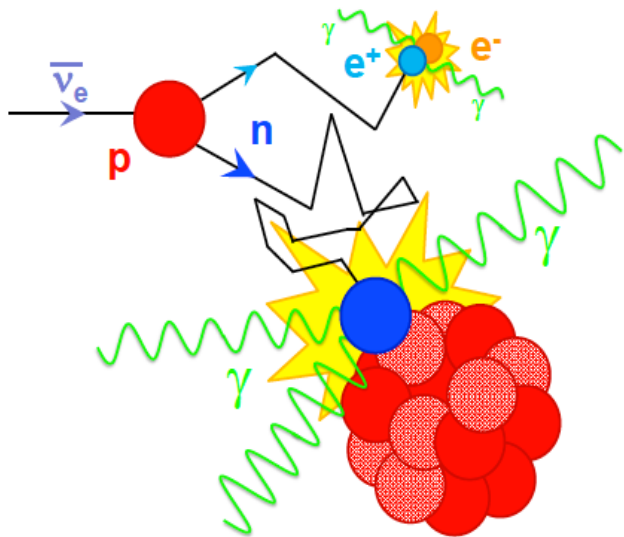
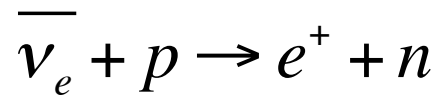
- Natural radioactivity + reactor leakage
- High E γ 's from n-capture on metals, ^{16}N production in primary circuit, ...



- On site measurements + heavy shielding design
- Online measurement and subtraction

Background

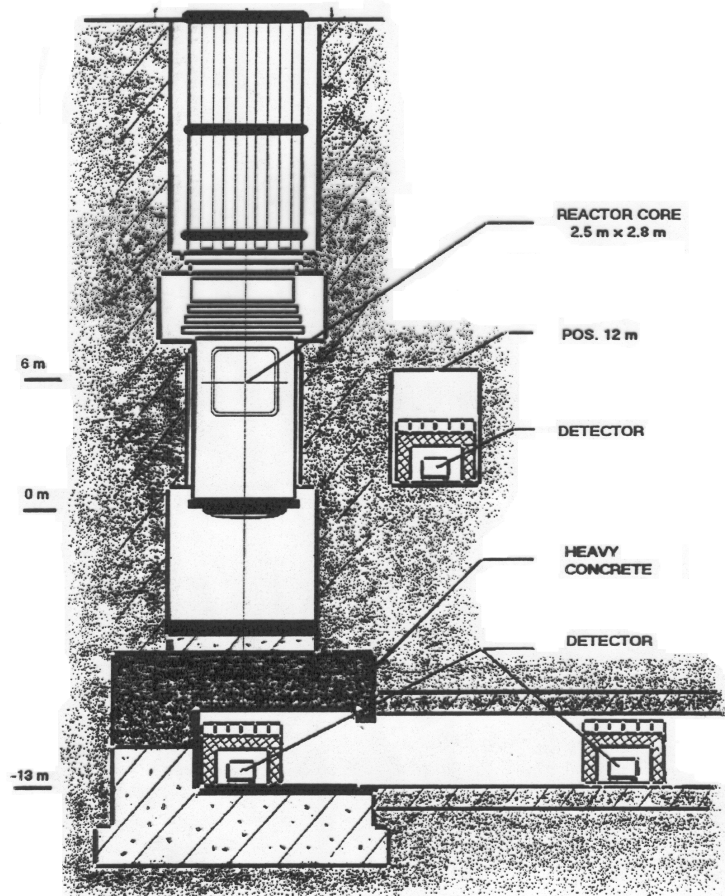
Inverse Beta Decay



Correlated background from fast n

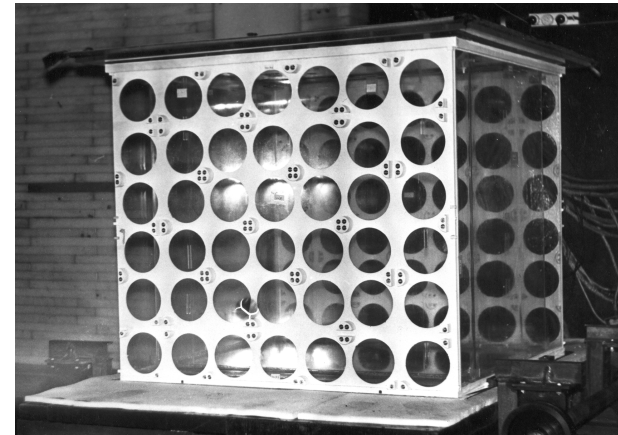
- Cosmic rays - induced
- **Online rejection:**
 - Minimal overburden
 - Active μ veto around target
 - PSD
- **Offline:**
 - Subtraction of reactor OFF data
- Must suppress all fast n from reactor

Rovno experiment



3GW Rovno NPP – Ukraine

Detector vessel



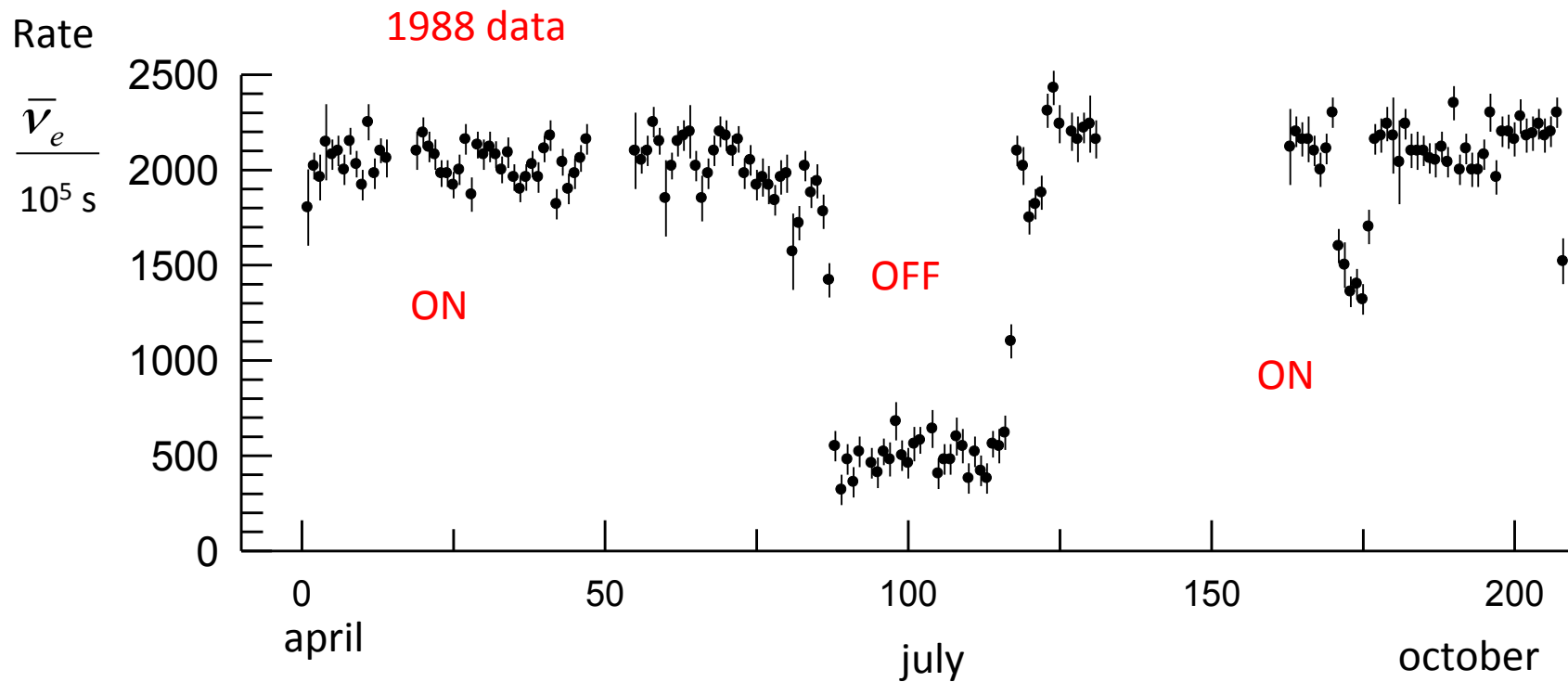
1m³ of liquid scintillator + 0.5 g/l Gd
84 PMT

$\epsilon_{\text{det}} \sim 50\%$

Favorable signal/noise configuration:

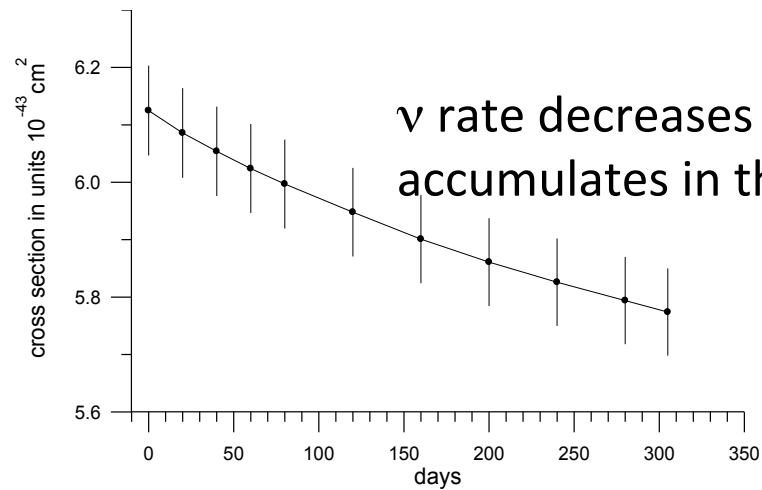
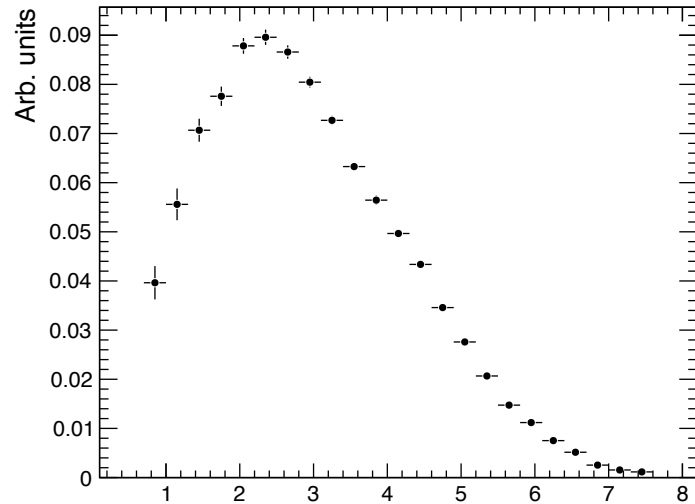
- 18 m from the core
- 30 m.w.e. oberburden

Rovno experiment

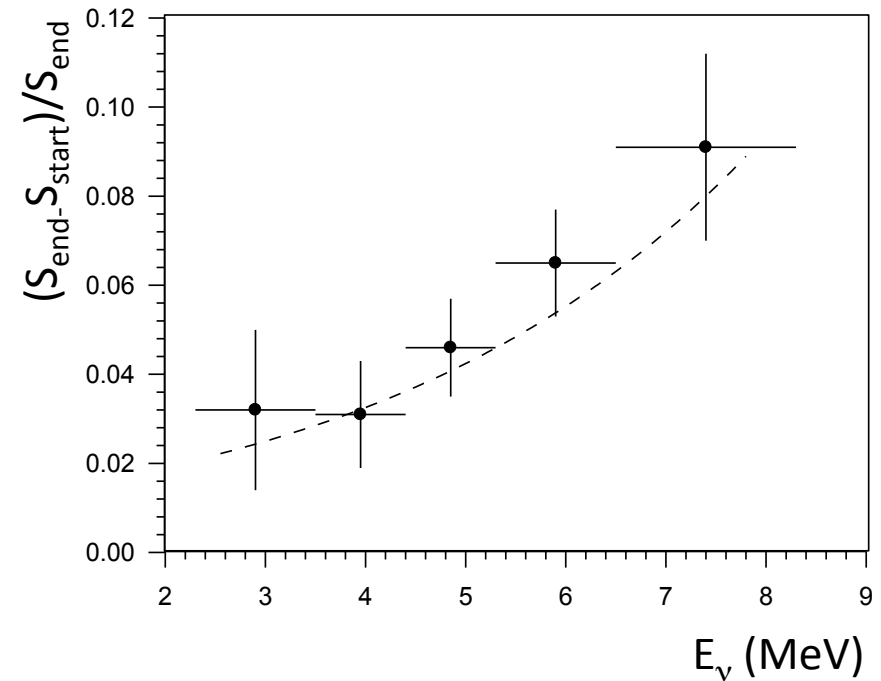


Rovno experiment

174000 ν events



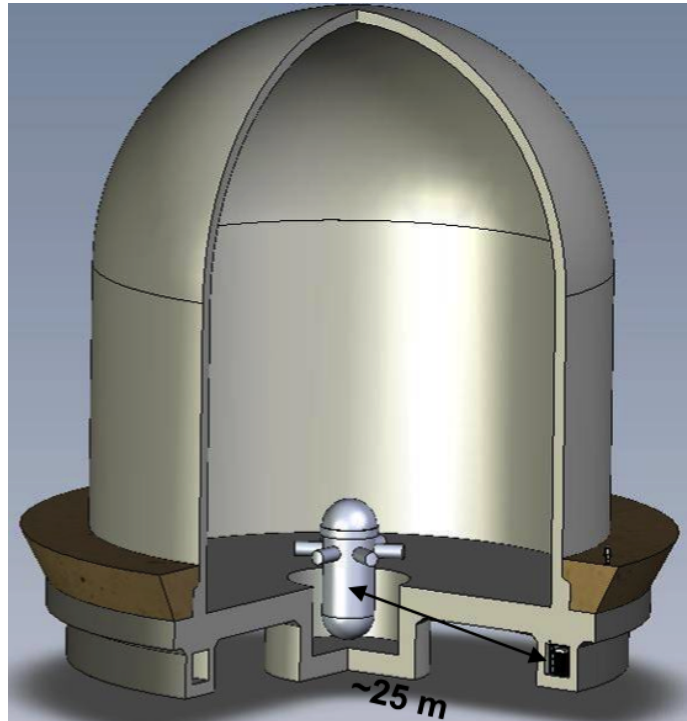
ν rate decreases as Pu accumulates in the core



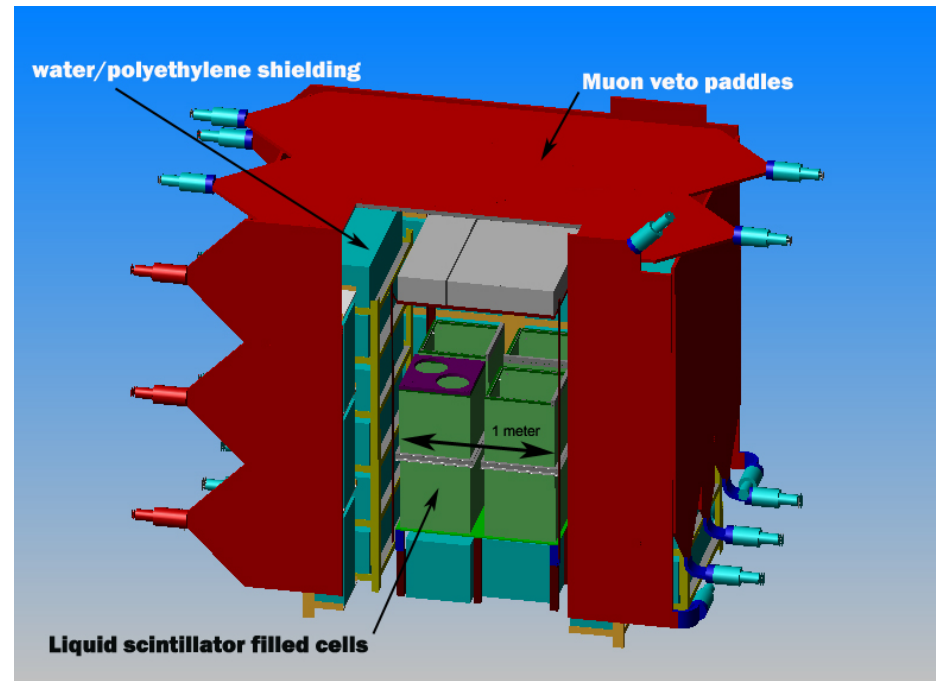
ν spectrum gets harder

SONGS

SONGS detector deployed at the San Onofre Nuclear Generating Station



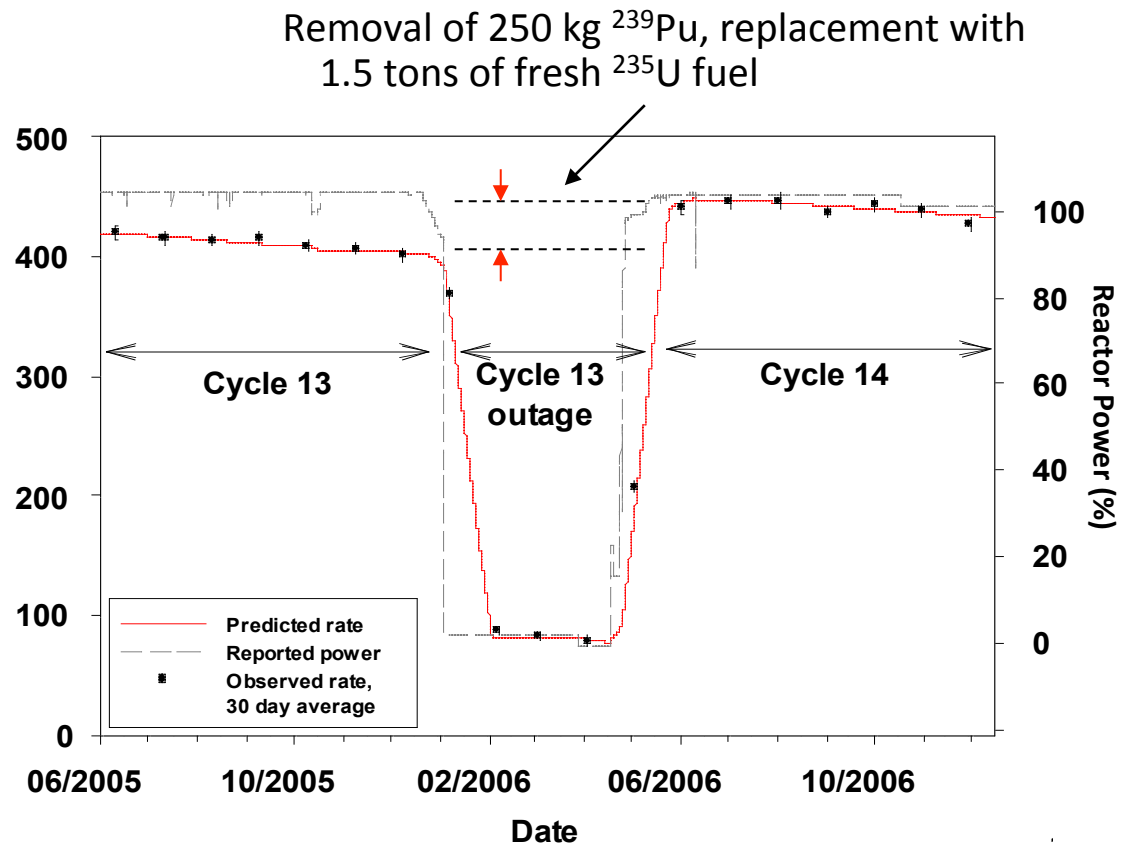
- $3.4 \text{ GW}_{\text{th}} \Rightarrow \sim 10^{21} \text{ v/s}$
- 3800 int. expected per day in 1m^3 liq. scint. target



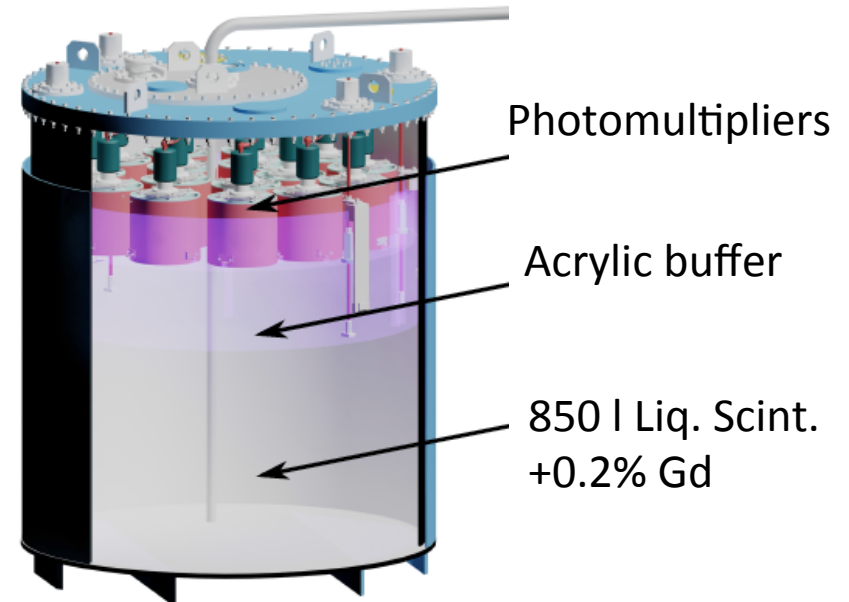
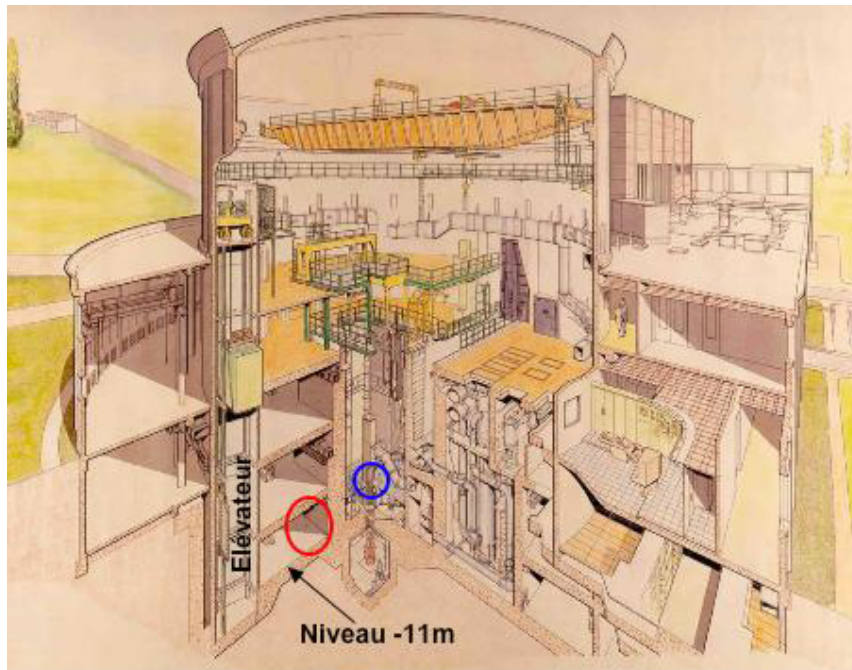
- **Low cost and robust detector**
- **Very simple design and light shielding**
- **Automated, non intrusive measurement**

SONGS

- Remarkable monitoring of the reactor operation
- Good signal/noise ratio
- Sensitivity limited by the low detection efficiency



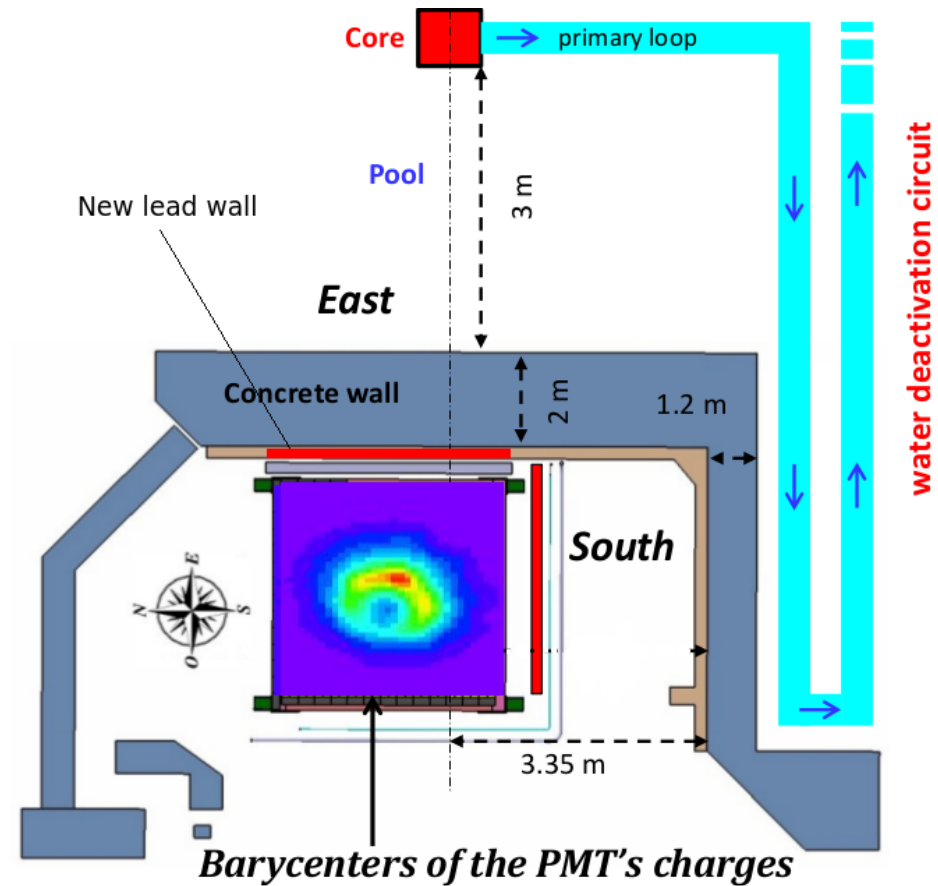
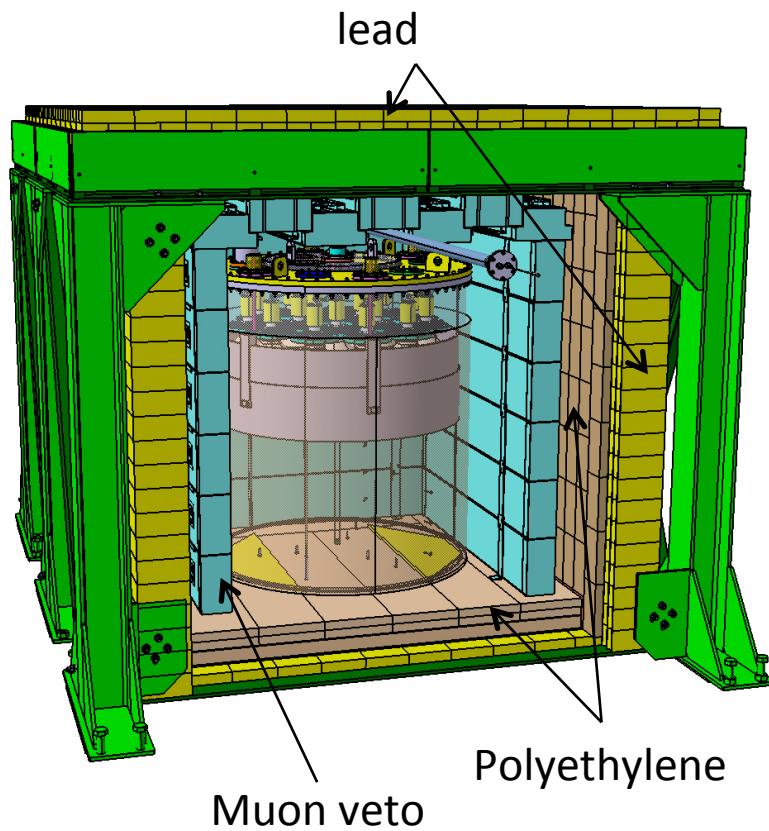
NUCIFER



- OSIRIS: 70 MW research reactor at CEA-Saclay, France
- Nucifer is only 7.2 m away from the core!
- Aiming at demonstrating the concept of “neutrino-metry” at the pre-industrialized stage.

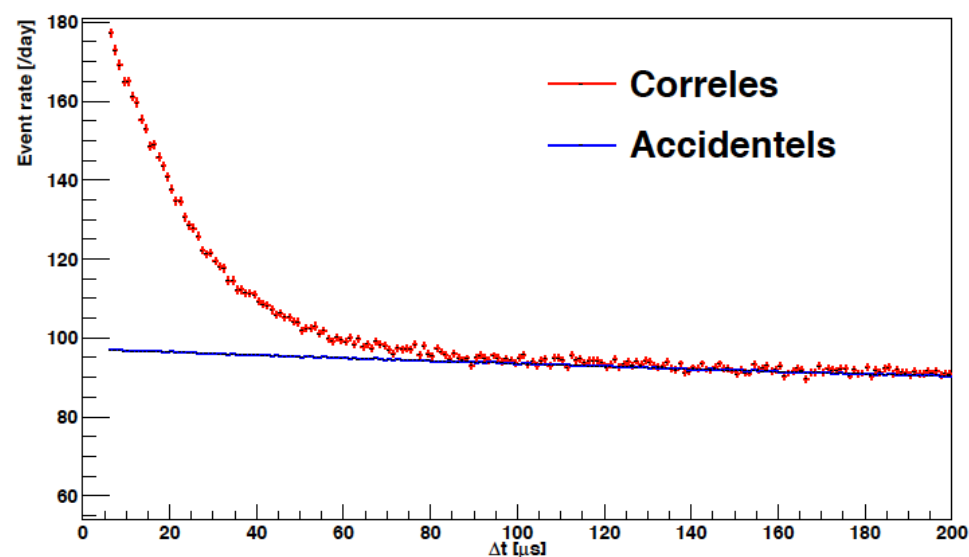
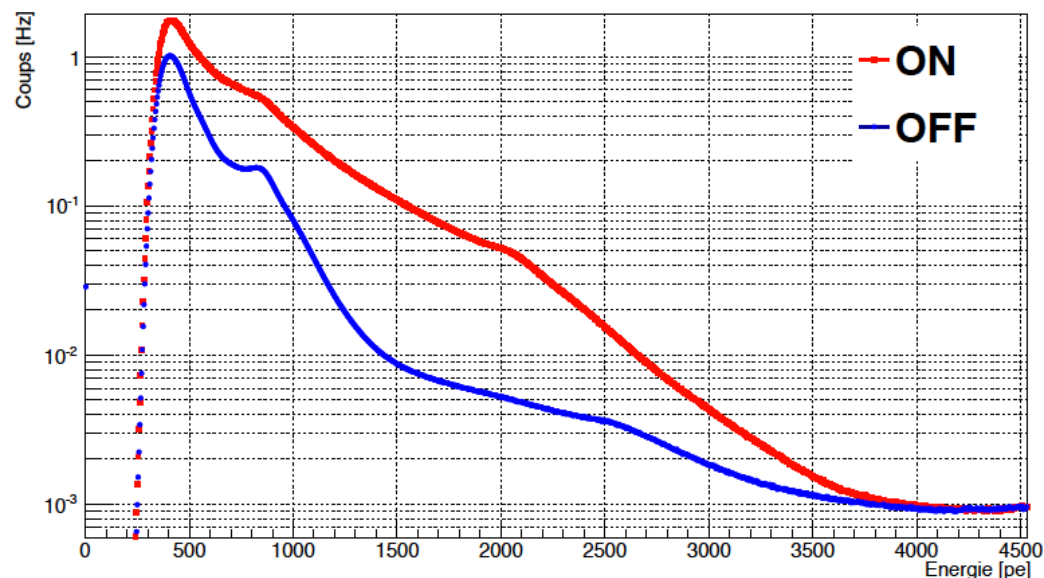
Nucifer shielding

Very challenging background induced by the reactor and the cosmic rays



Background

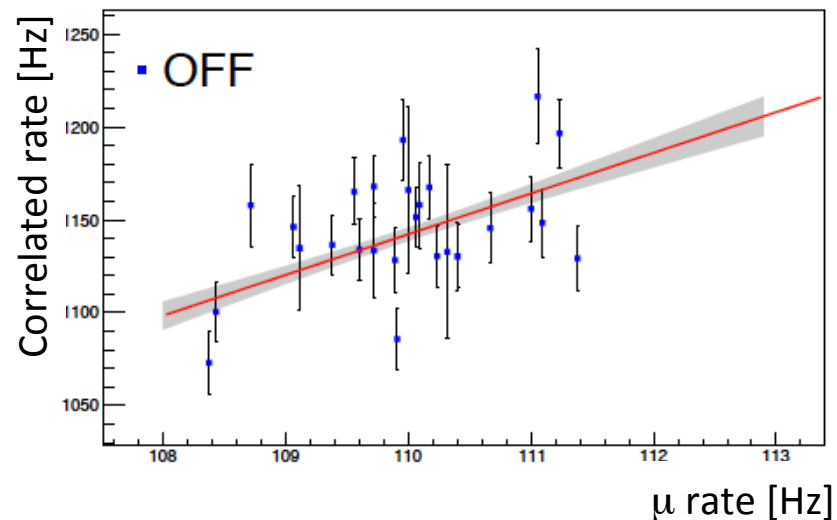
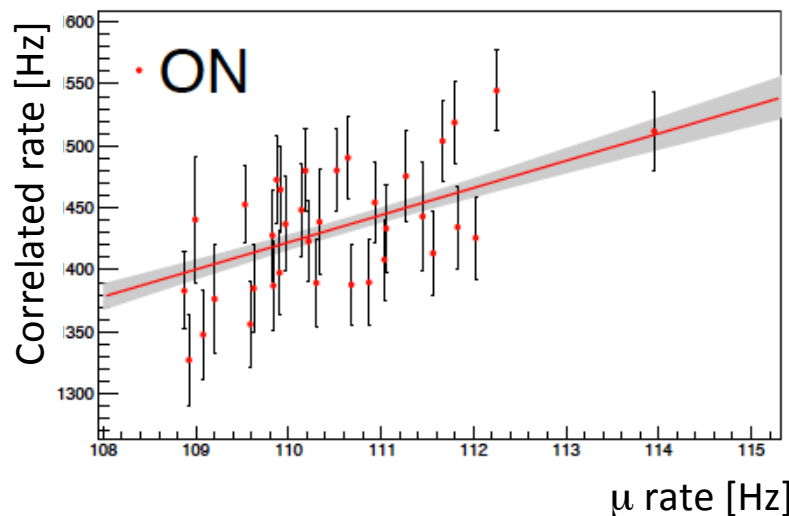
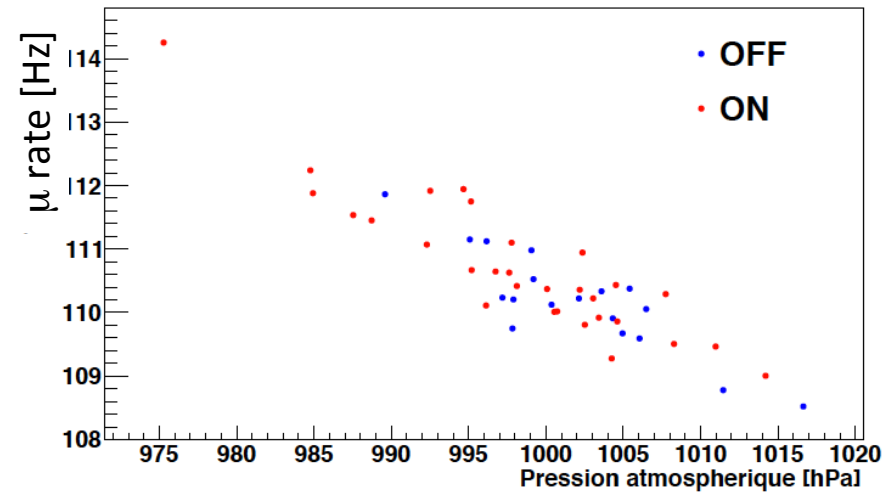
- Contamination of high gammas from the reactor in the prompt and delayed energy window
- Large accidental background scaling with $(P_{\text{reactor}})^2$
- 4733 ± 6 candidate pairs /day
 3332.8 ± 0.5 accidentals/days,
accurately measured online



Cosmic rays induced background

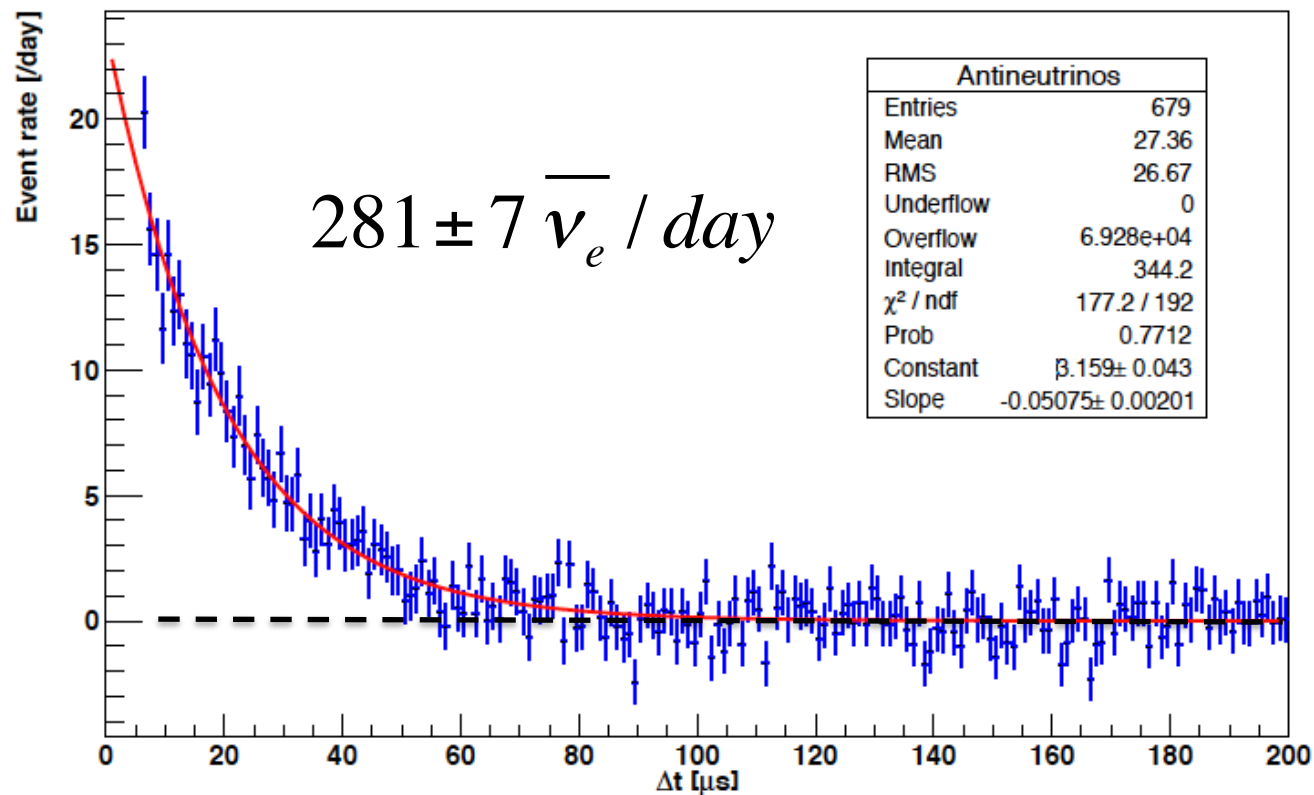
A clear correlation between P_{atmos} and the μ rate in Nucifer is observed

→ The subtraction of reactor OFF background has to be corrected for the \neq of μ rate between ON and OFF periods.

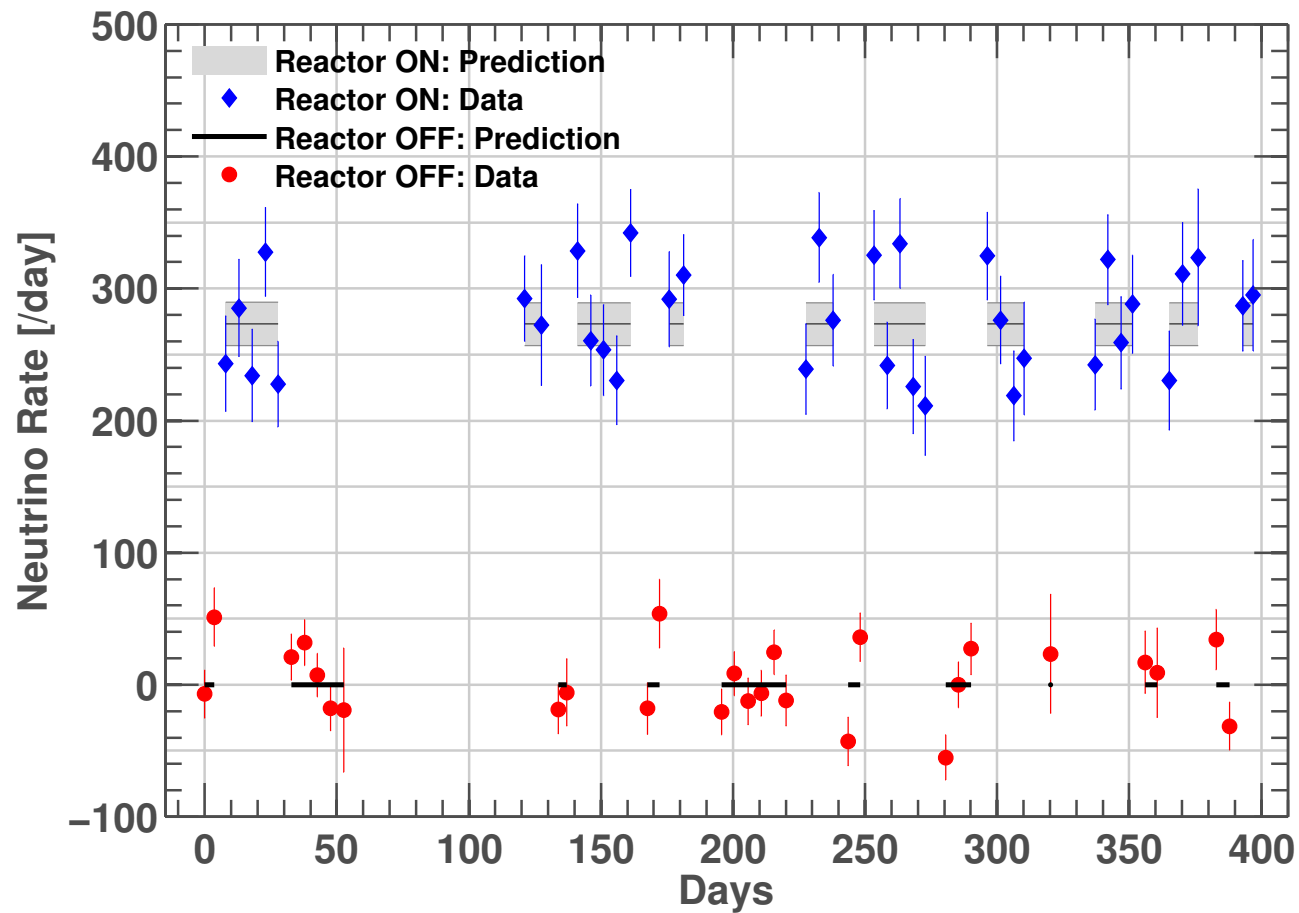


Nucifer results

$$R_{\bar{\nu}_e}^{ON}(R_\mu) = R_{paires}^{ON} - R_{acc}^{ON} - R_{corr}^{OFF}(R_\mu^{ON})$$



Nucifer results

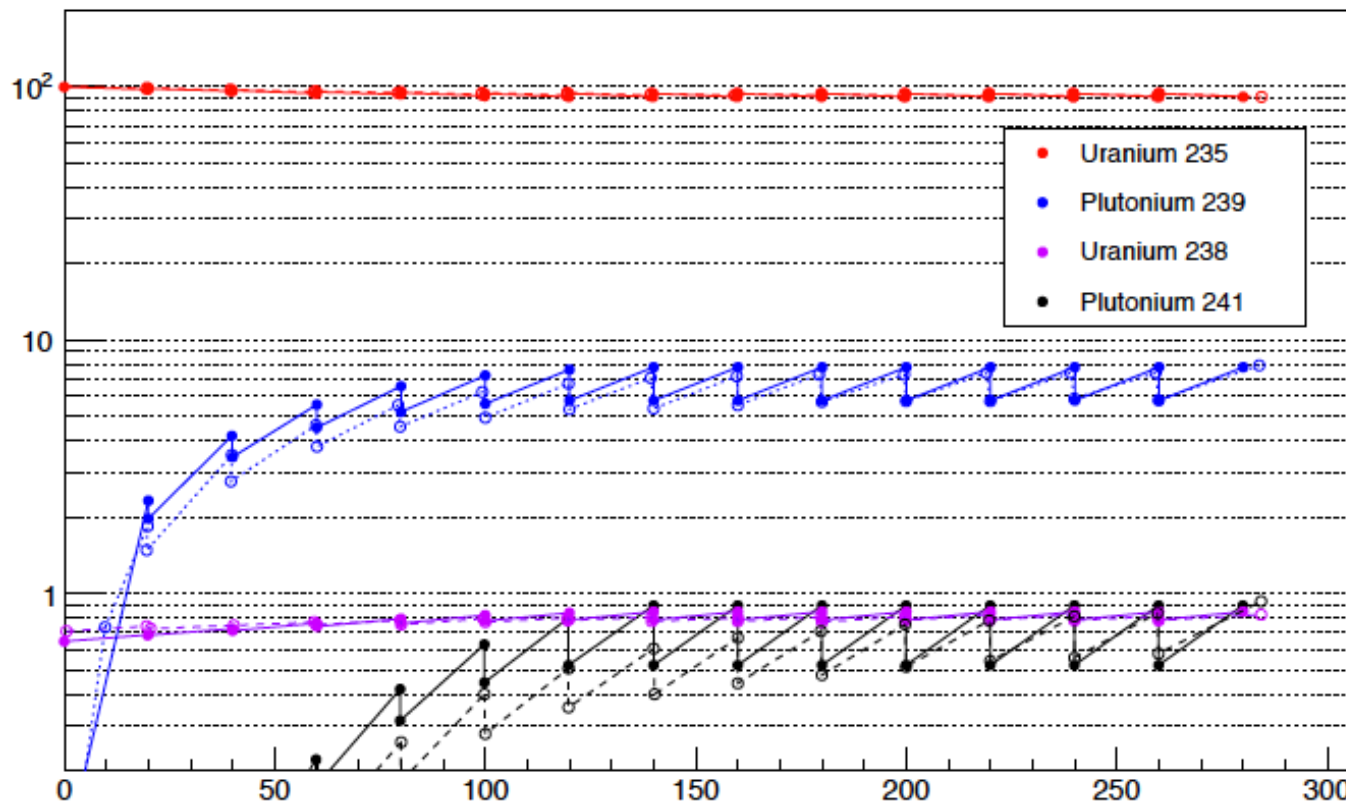


[arXiv:1509.05610](https://arxiv.org/abs/1509.05610)

Sensitivity study

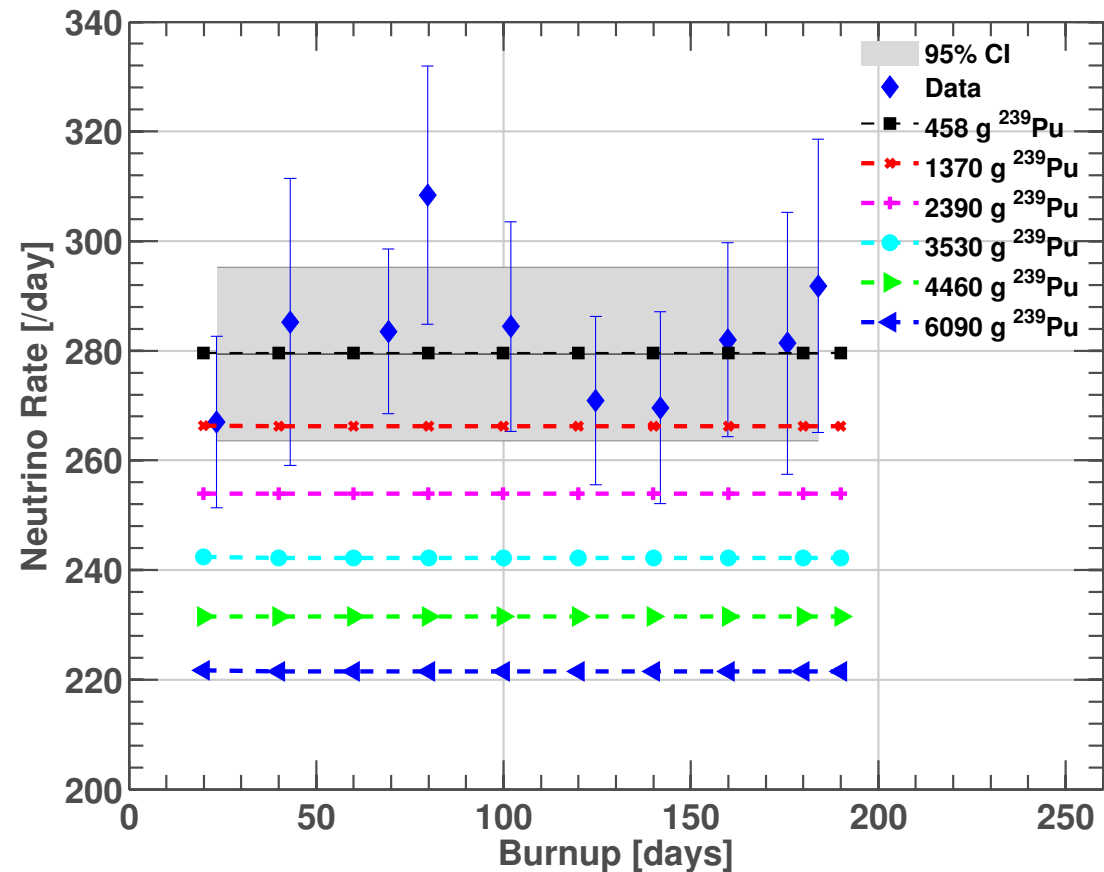
Osiris uses highly enriched ^{235}U fuel with short cycles

→ The evolution of the fuel composition is undetectable (<1% change in the neutrino rate)



Sensitivity study

- Various levels of Pu were simulated in the core
- Within the current background conditions Nucifer would have seen the effect of ~1.5 kg of Pu in Osiris at 95% C.L., that is ~10% of the total fissile mass.

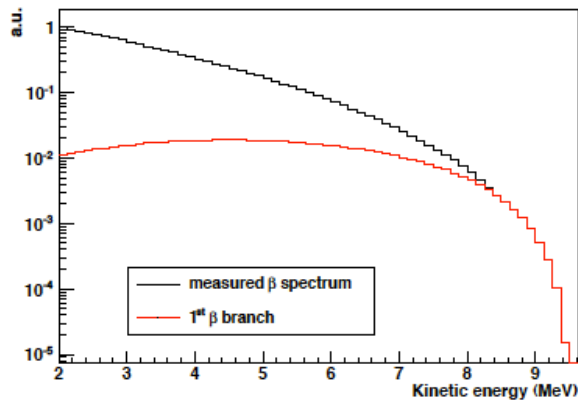


Reactor surveillance with neutrinos

- Antineutrino monitoring can assess the plutonium content in water-cooled nuclear reactors for nonproliferation applications.
- Need to demonstrate further sensitivity and deployment capacities before competing with the current tools of reactor surveillance.
- Direct information on the fission processes in the core is an asset for the surveillance of online refueling/reprocessing reactors. But the lever arm on the neutrino signal is small.
- Disposal of weapon-grade plutonium by irradiation in fast reactors (Plutonium Management and Disposition Agreement) could be monitored by neutrinos with simple observables like total burnup and $\langle N_{\bar{\nu}} \rangle / \text{GW}$
- New detection technics are being developed in synergy with fundamental research activities.

Accurate prediction of neutrino fission spectra

Conversion of the ILL spectra with virtual β -Branches



Fitted on ILL β -spectrum

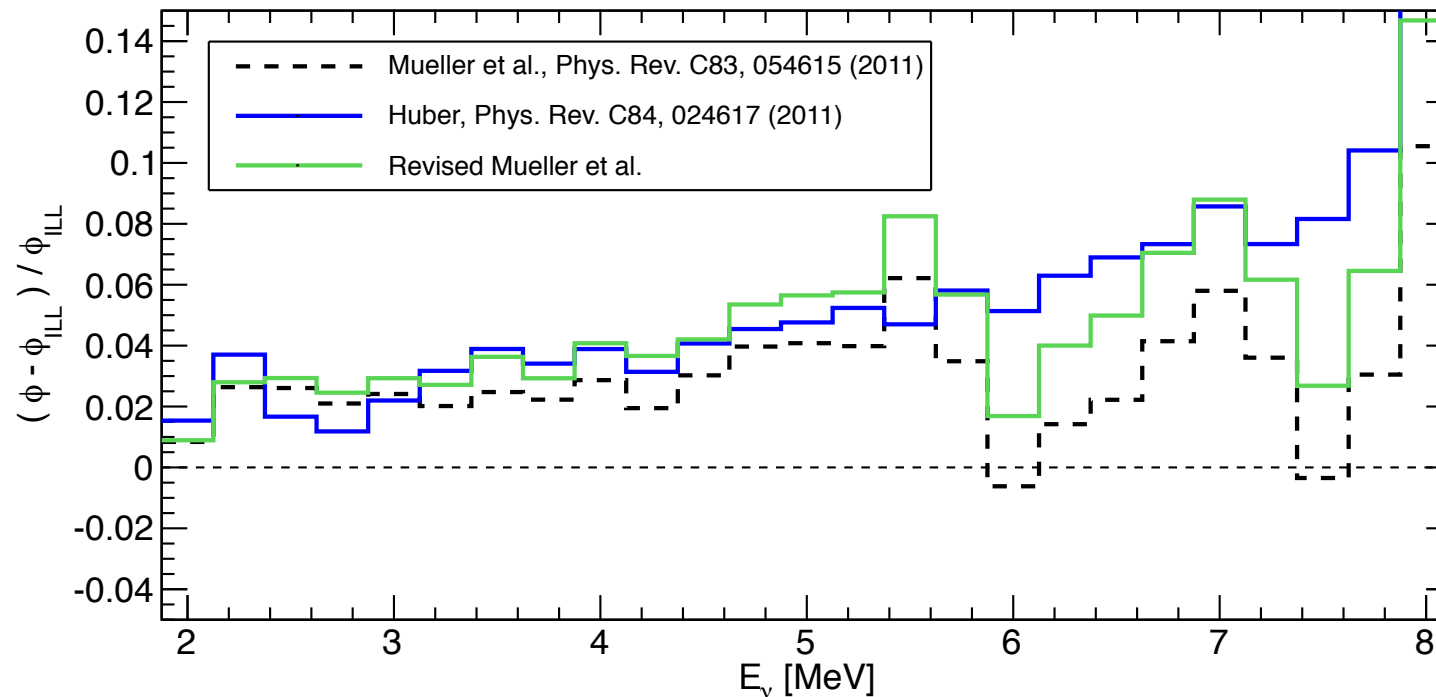
$$S_{fp}^b = \underbrace{K_{fp}^b}_{\text{Norm.}} \times \underbrace{\mathcal{F}(Z_{fp}, A_{fp}, E)}_{\text{Fermi function}} \times \underbrace{pE(E - E_{0fp}^b)^2}_{\text{Phase space}} \times \underbrace{C_{fp}^b(E)}_{\text{Shape factor}} \times \underbrace{\left(1 + \delta_{fp}^b(Z_{fp}, A_{fp}, E)\right)}_{\text{Correction}}$$

The mean $Z(E0)$
from nuclear
databases

Set to 1

Effective δ_{WM} and δ_{CW} corrections
applied after the conversion process

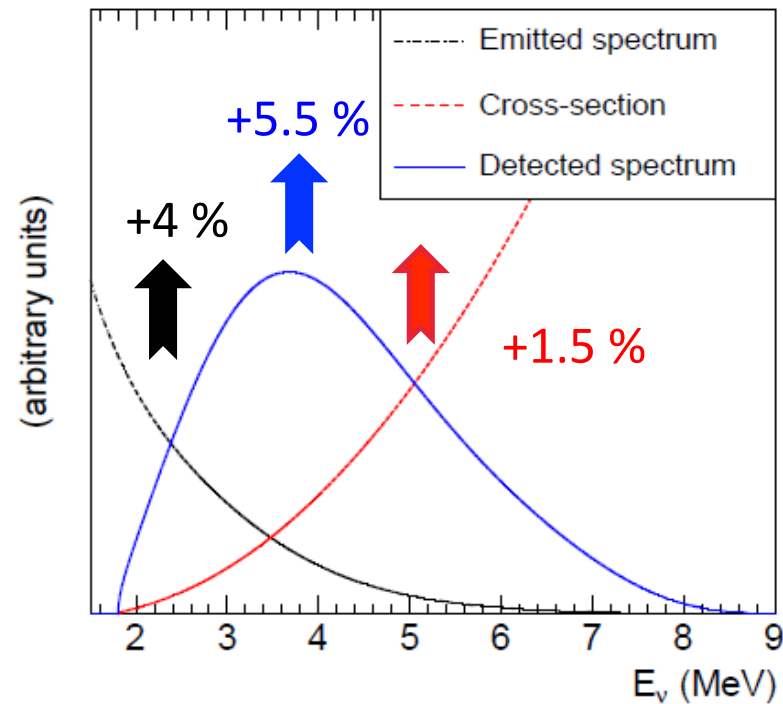
Conversion of ILL Spectra Revisited



- Confirms global increase of predicted spectrum
- Fixes remaining oscillations of mixed-approach prediction
- Extra deviation at high energy from more complete correction to Fermi theory (weak interaction in the finite volume of the parent nucleus).

New Prediction

Improved
converted
spectra

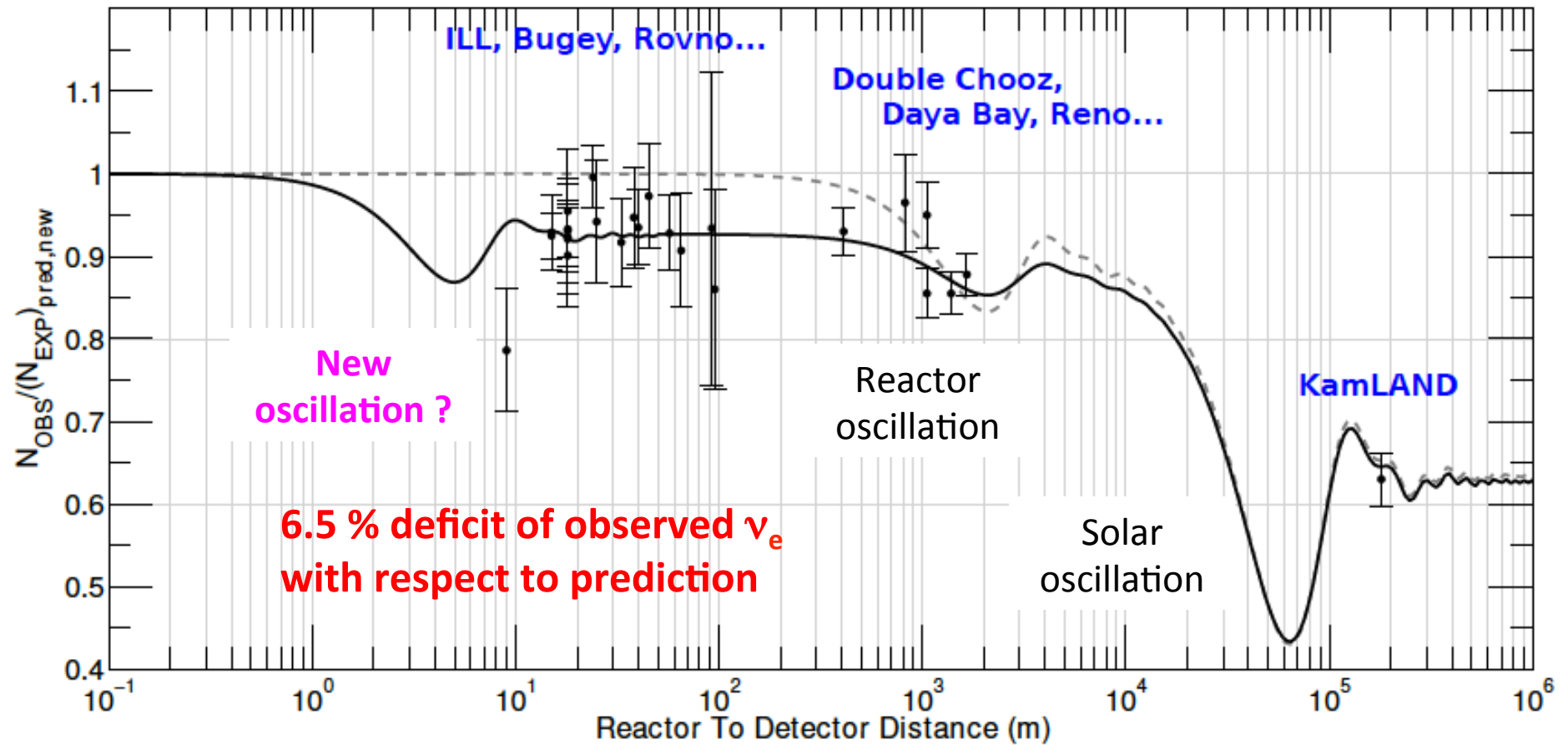


Measured value of
neutron lifetime has
evolved in time and
 $\sigma_{\text{int}} \propto 1/\tau_n$

*G. Mention et al,
Phys. Rev. D83, 073006, 2011*

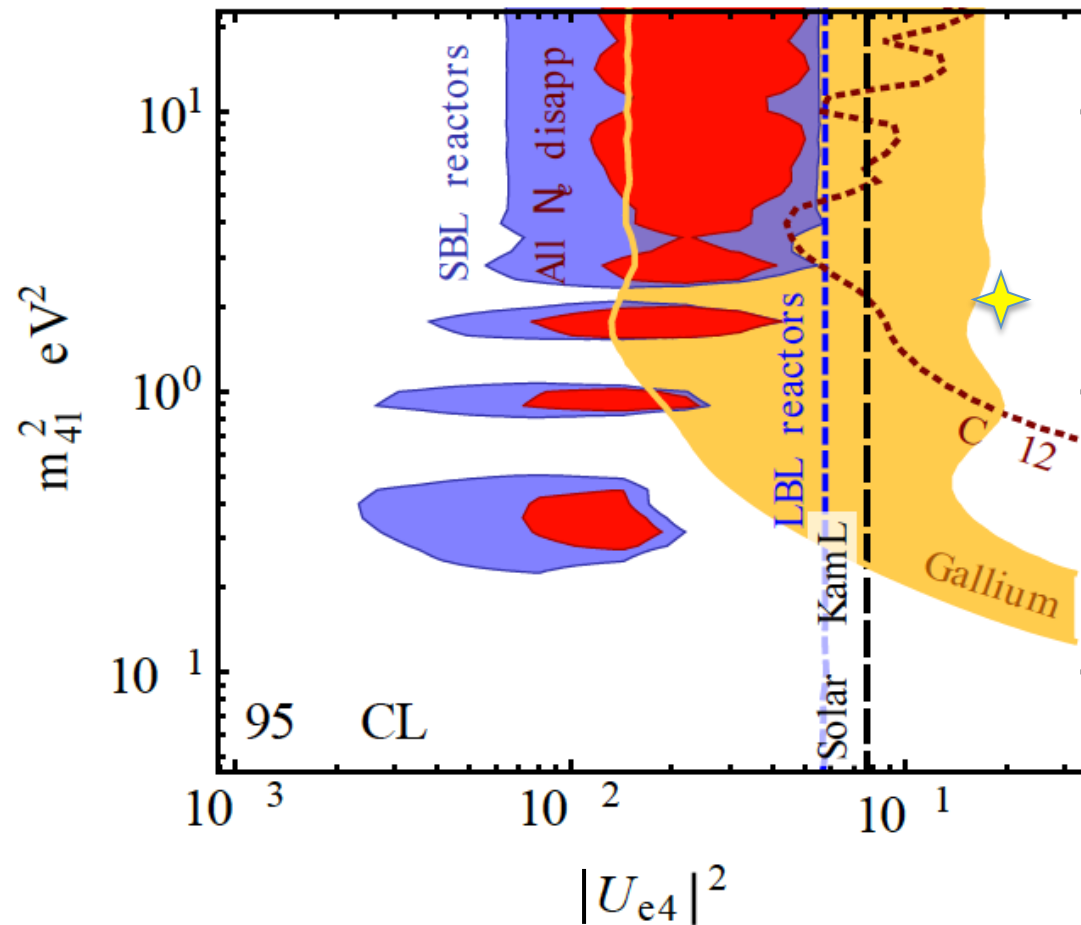
→ Predicted flux increases by ~5.5 %

Reactor Antineutrino Anomaly



Sterile Neutrino ?

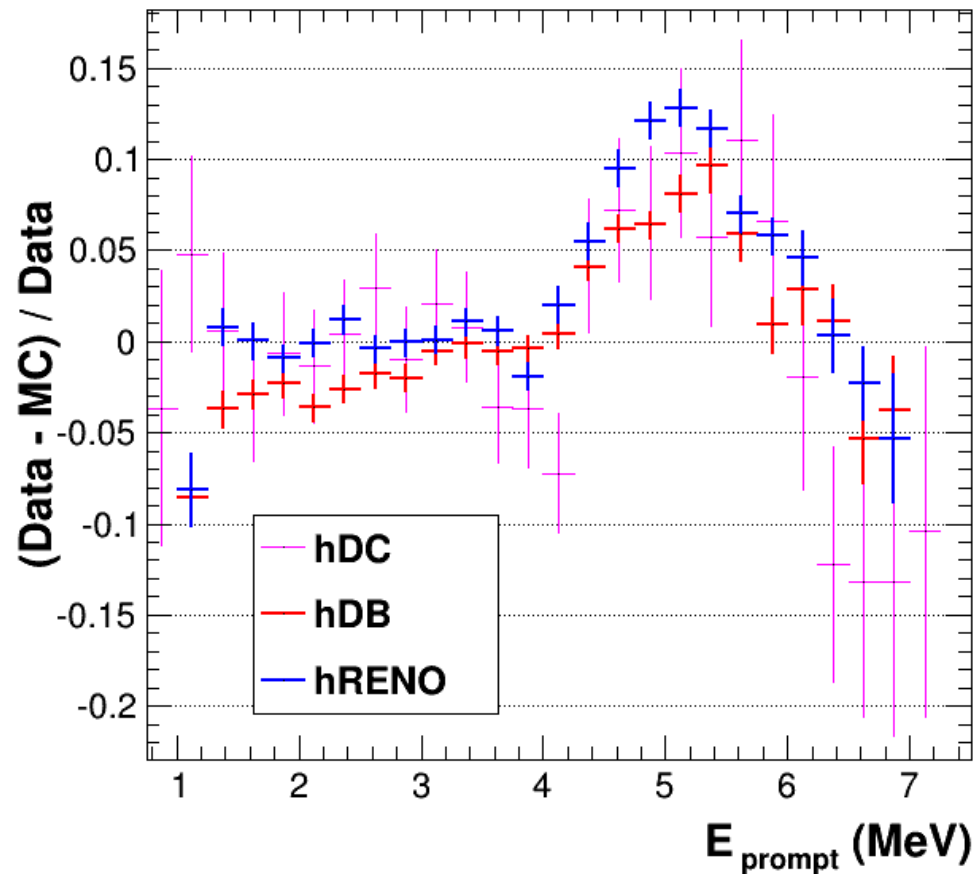
J. Kopp et al, JHEP 1305 (2013) 050



- Puzzling convergence of several ν_e disappearances toward the same oscillation parameters.
- Tension with nm disappearance data and cosmology fits
- Triggered a worldwide experimental program for a direct search of a new oscillation pattern.

... or bias in the converted spectra?

Double Chooz, Daya Bay, Reno



3 experiments point to the same deviation

→ Biased prediction?

→ Biased reactor anomaly?

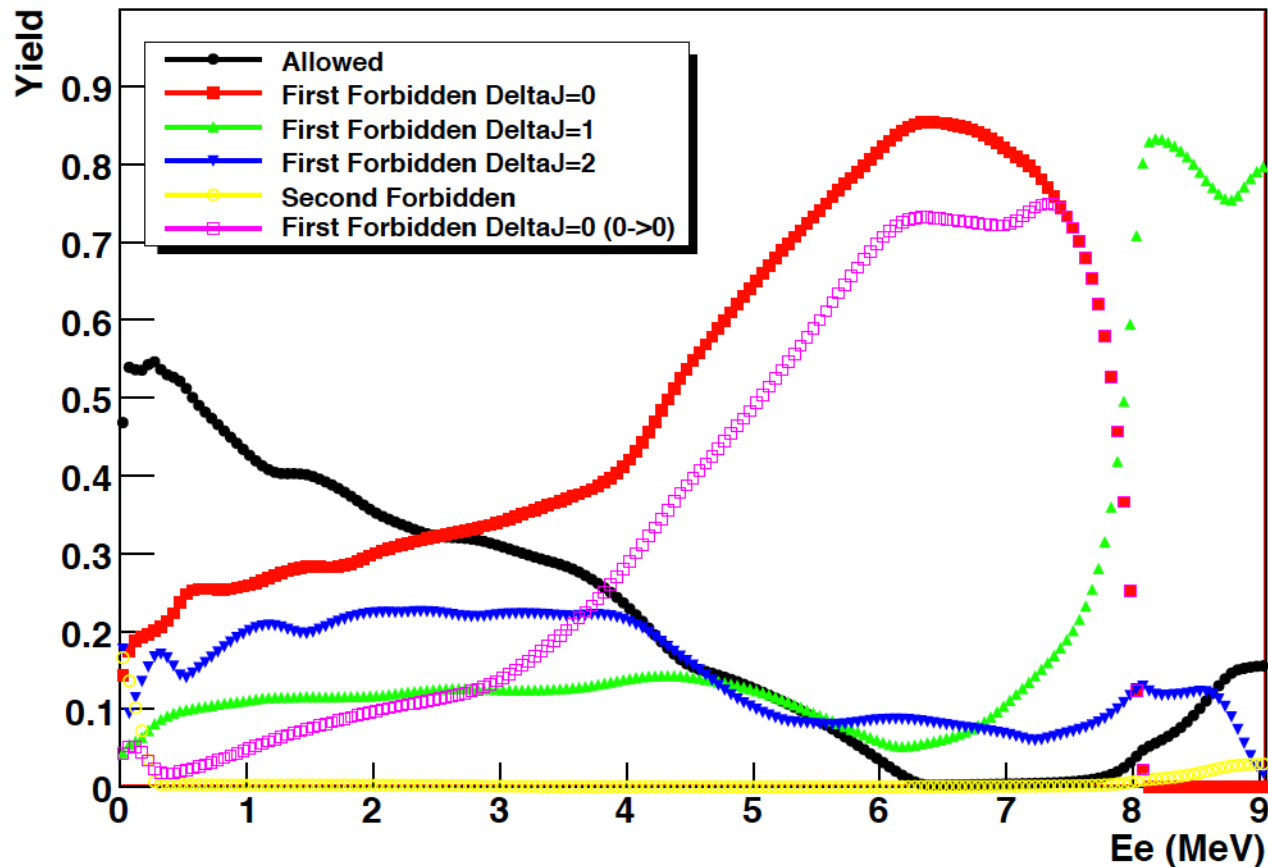
Potential Biases

A.C. Hayes et al, Phys. Rev. D92 (2015) 3, 033015

- Differences between ILL and commercial reactors:
 - neutron energy spectrum and fission yields
 - Off-equilibrium effects (under control, impact at low energy only, <3 MeV)
 - non-fission sources of antineutrinos (not likely)
- Error in the ILL electron spectra (?)
 - Would need a complementary experiment with similar accuracy
- Forbidden transitions
- Corrections to Fermi theory

Contribution of Forbidden Decays

- Relative contribution of different transition types in each energy bins of the total fission spectrum
- Only well known decays computed, quite few after 6 MeV.

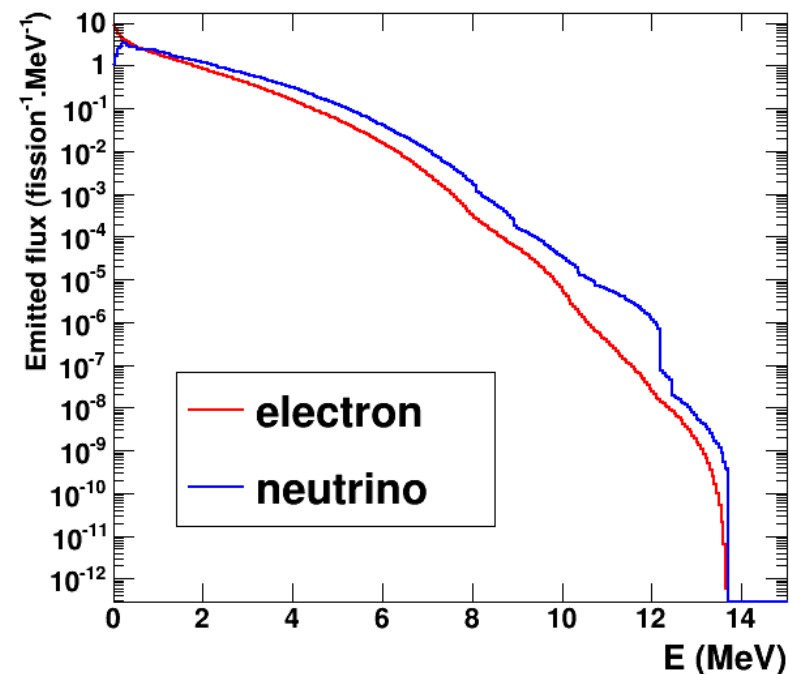


A. Letourneau

Tool for numerical studies of the conversion process

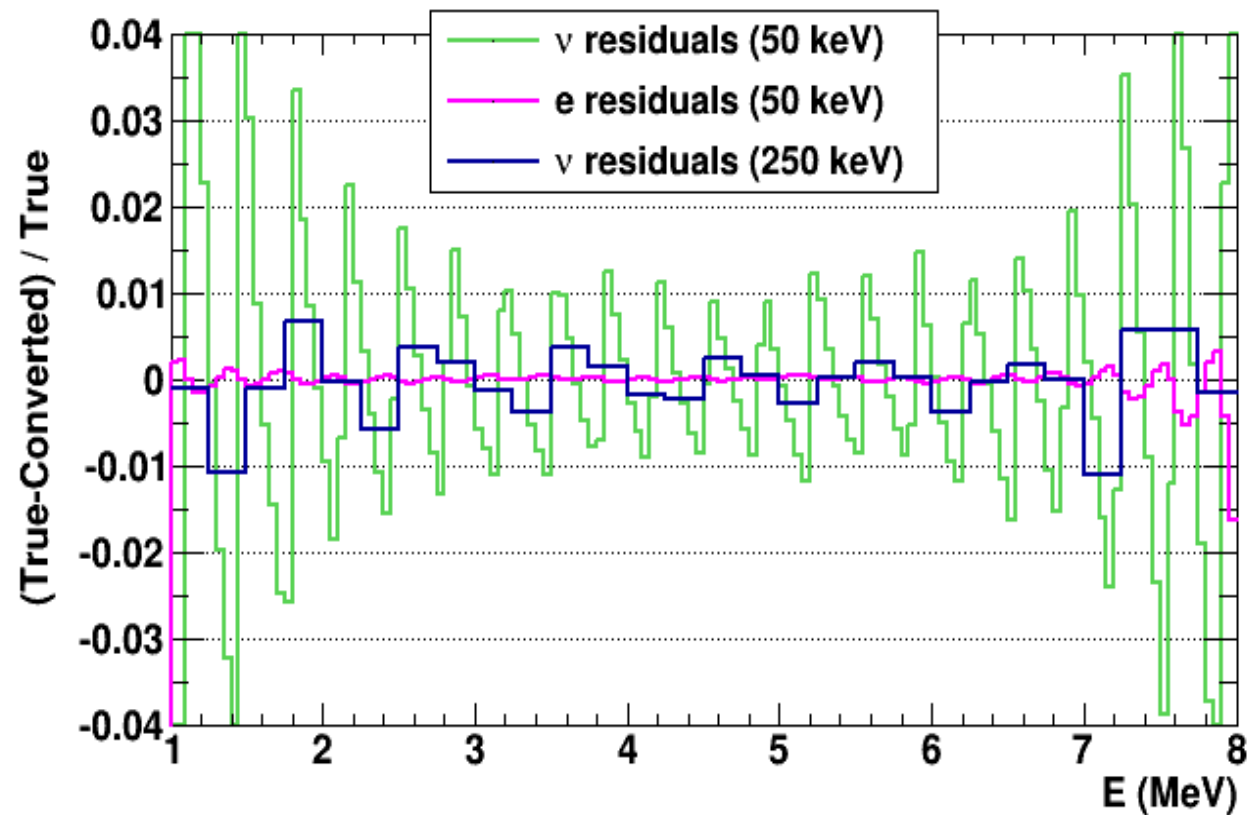
- Numerical studies of the impact of forbidden decays using calculated fission spectra based on all **decays indexed in the ENSDF nuclear database**.
- Each decay shape is determined from spin-parity data.
- Despite all systematics of such calculation, **electron and neutrino spectra are true images of each other**.
- Although they do not match perfectly the ILL spectra they contain a **good approximation of the physical distribution of β -decays with full control of the forbidden shapes we put in**.

ENSDF based ^{235}U fission spectra



Validation with allowed β -branches

- Consistency cross-check: when all ENSDF branches are forced to be of allowed type and the effective branches are allowed as well, both electron and neutrino residuals are zeroed.



Forbidden Decays

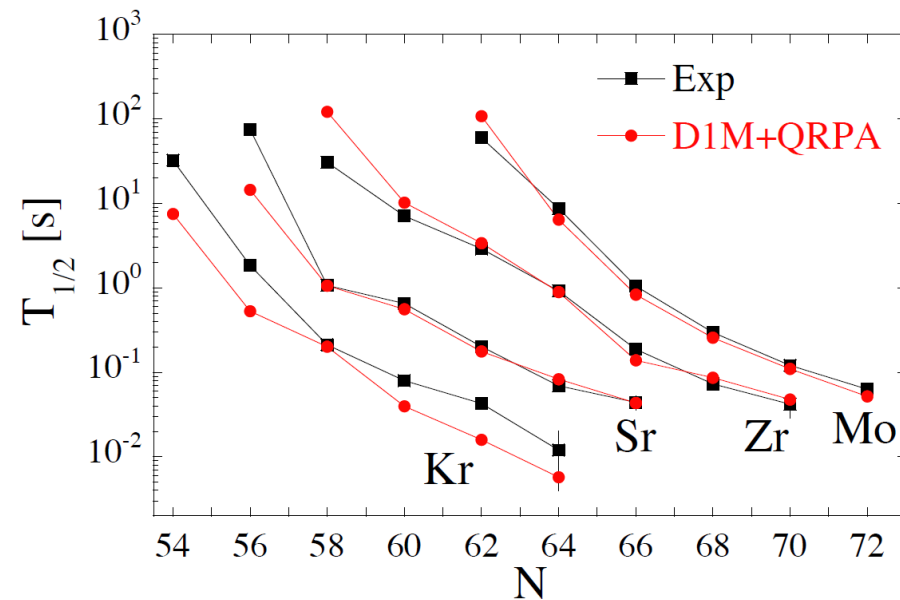
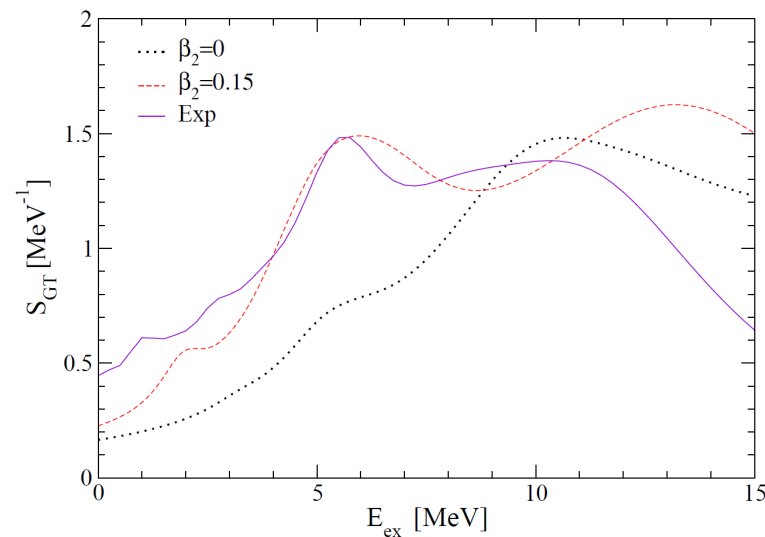
Nuclear operators

		ALLOWED	FIRST FORBIDDEN	SECOND FORBIDDEN
Polar V.	Matrix ΔJ Parity Change	$\int 1$ 0 no	$\int \mathbf{r}, \int \boldsymbol{\alpha}$ 0, ± 1 (no $0 \rightarrow 0$) yes	R_{ij}, A_{ij} $\pm 1, \pm 2$ (no $1 \leftrightarrow 0$) no $\int \boldsymbol{\alpha} \times \mathbf{r}$ ± 1 no
Axial V.	Matrix ΔJ Parity Change	$\int \boldsymbol{\sigma}$ 0, ± 1 (no $0 \rightarrow 0$) no	$\int \boldsymbol{\sigma} \cdot \mathbf{r}, \int \gamma_5$ 0 yes $\int \boldsymbol{\sigma} \times \mathbf{r}$ 0, ± 1 (no $0 \rightarrow 0$) yes B_{ij} 0, $\pm 1, \pm 2$ (no $0 \rightarrow 0$) $\frac{1}{2} \rightarrow \frac{1}{2}, 1 \leftrightarrow 0$ yes	T_{ij} ± 2 no S_{ijk} $\pm 2, \pm 3$ (no $2 \leftrightarrow 0$) no

- One single nuclear operator.
- No E dependence.
- Several nuclear operators with basically unknown relative contributions.
- Complex E dependence.

pnQRPA Strength Distributions with Gogny (D1M and D1S) force

- Calculation of all pair-pair nuclei available
- Deformation of the nuclei included in the prediction of β -strengths
- Current effort to extend it to all fission products
- Should bring valuable new inputs to study the shape of 1st forbidden decays and the weak magnetism corrections

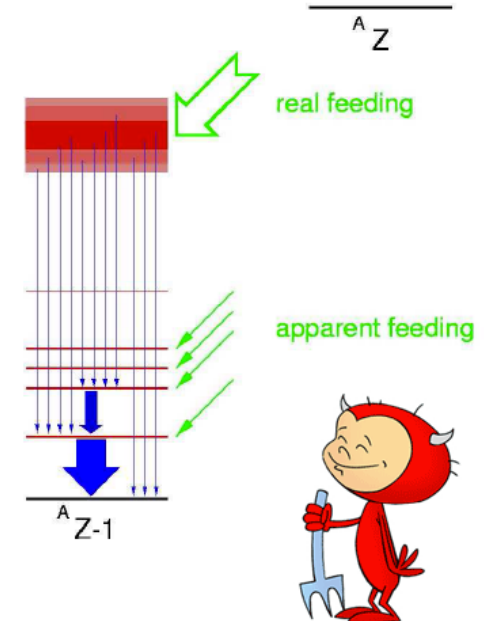


S. Peru et al.

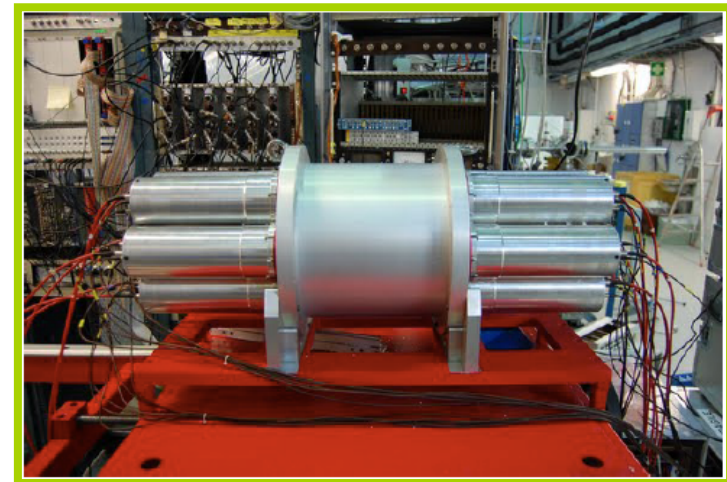
Summation Method

TAGS measurements

- Pandemonium effect: experimental bias giving too much weight to high energy transitions → Total Absorption γ -ray Spectrometry.
- Short list of main contributor fission products



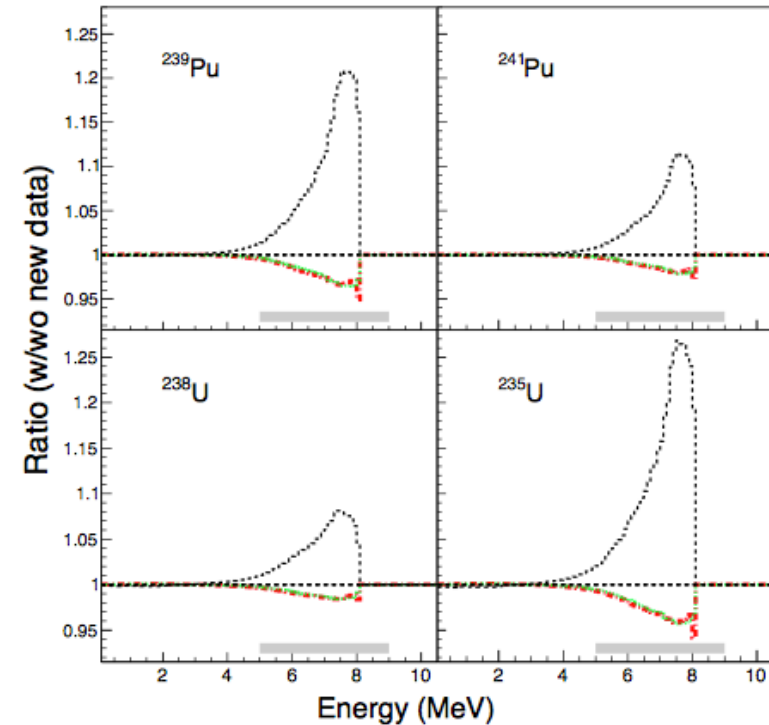
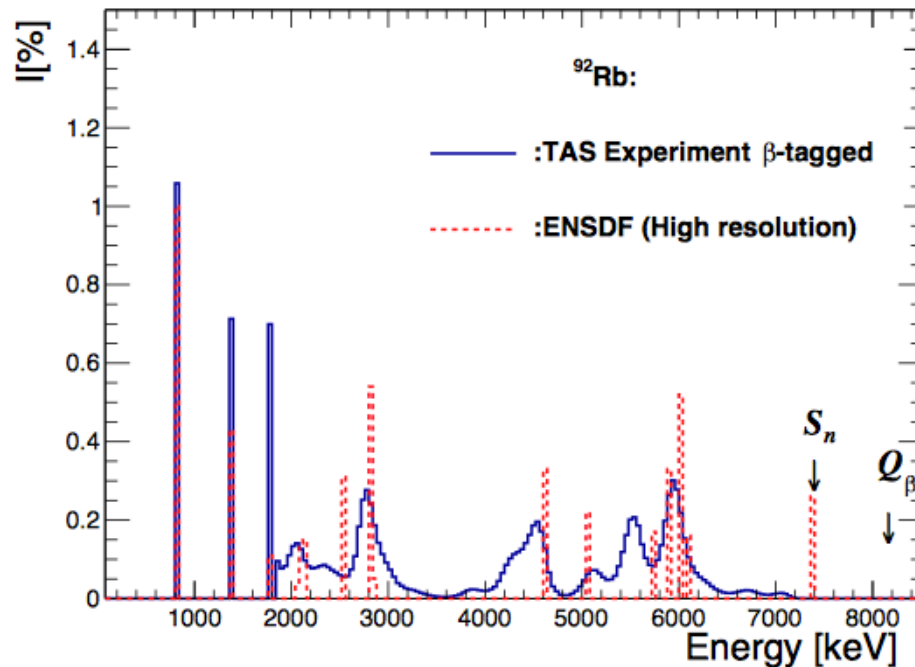
	4 - 5 MeV	5 - 6 MeV	6 - 7 MeV	7 - 8 MeV
^{92}Rb	4.74%	11.49%	24.27%	37.98%
^{96}Y	5.56%	10.75%	14.10%	-
^{142}Cs	3.35%	6.02%	7.93%	3.52%
^{100}Nb	5.52%	6.03%	-	-
^{93}Rb	2.34%	4.17%	6.78%	4.21%
^{98m}Y	2.43%	3.16%	4.57%	4.95%
^{135}Te	4.01%	3.58%	-	-
^{104m}Nb	0.72%	1.82%	4.15%	7.76%
^{90}Rb	1.90%	2.59%	1.40%	-
^{95}Sr	2.65%	2.96%	-	-
^{94}Rb	1.32%	2.06%	2.84%	3.96%



A. Algora et al., Phys. Rev. Lett. 105, 202501 (2010)

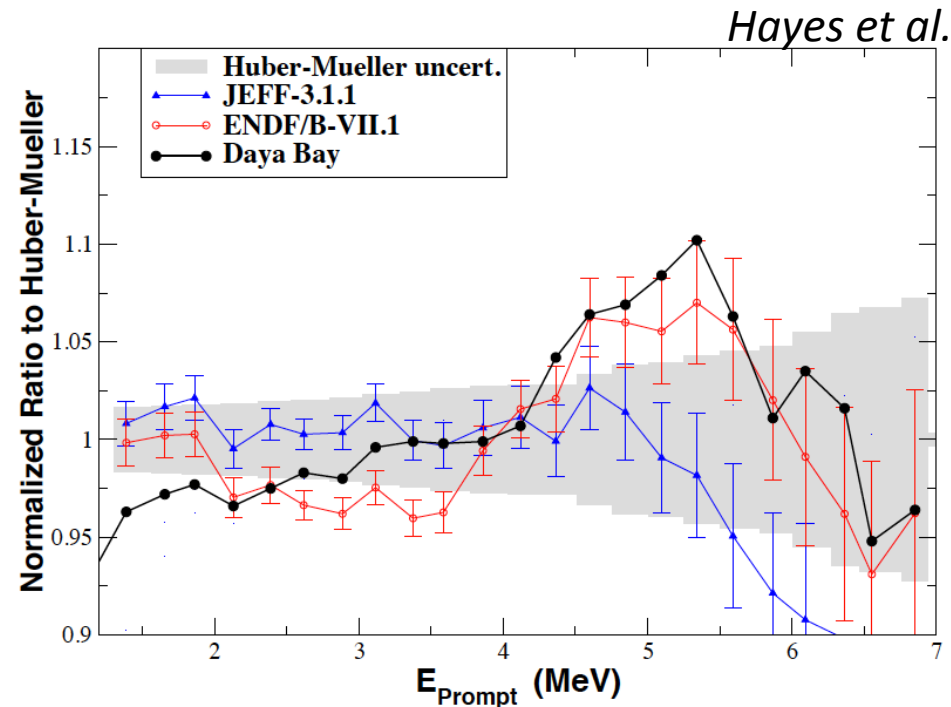
TAGS MEASUREMENTS

[arXiv:1504.05812](https://arxiv.org/abs/1504.05812)



- More and more complete data set reducing the impact of pandemonium
- Improvement of the shape uncertainty and impact of unknown nuclei
- Estimation of the final uncertainty at the few % level remains difficult

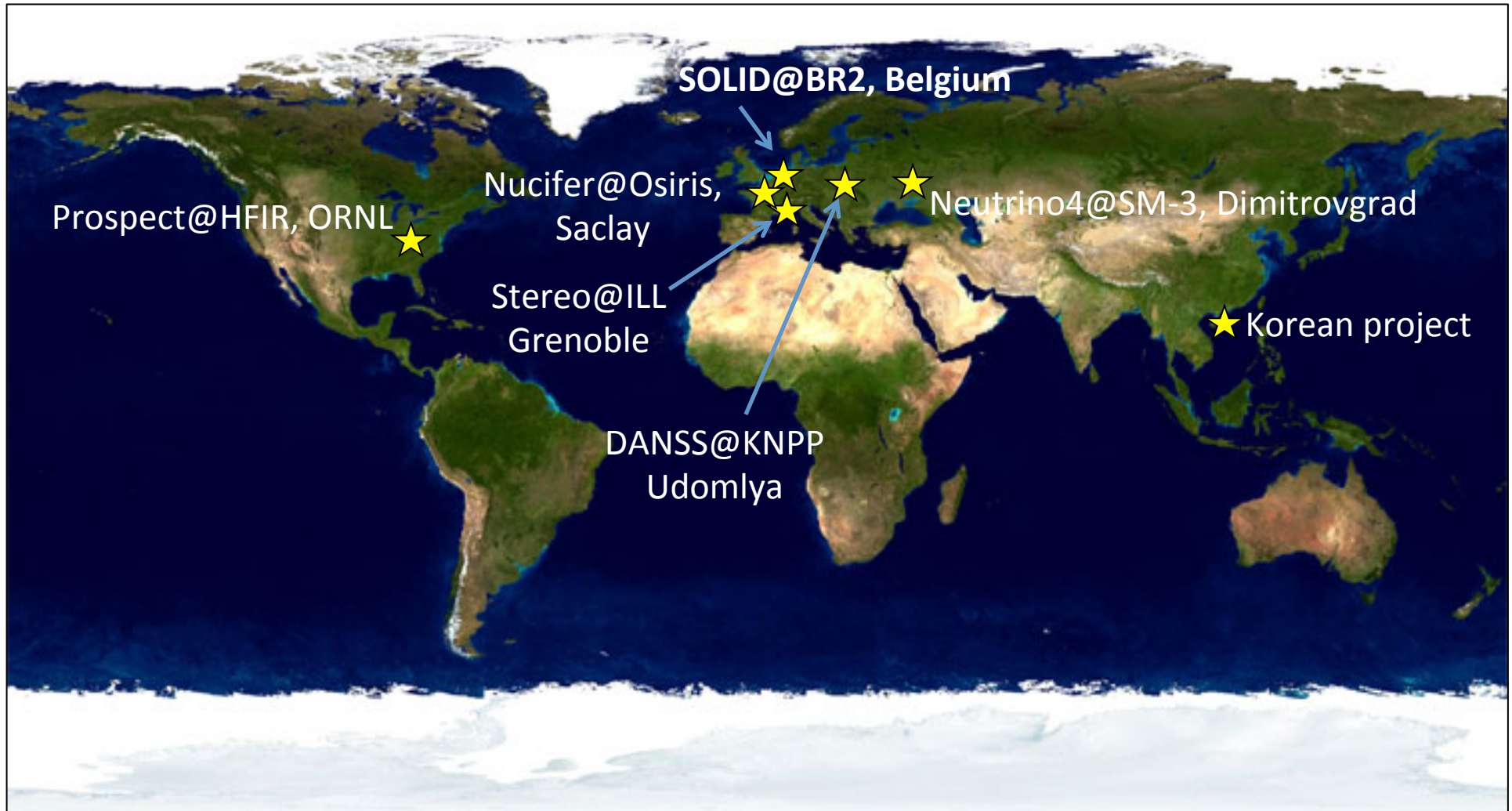
Sensitivity to Evaluated Fission Yield



- The predicted neutrino spectrum is quite sensitive to the choice of the nuclear database of evaluated fission yields.
- TAGS data don't provide information on the shape of forbidden decays
→ synergy with the ongoing studies on the conversion method

Search for a sterile neutrino

Worldwide Overview



Research Reactors

- **Compact sources**

- No oscill. smearing.

- **High statistics, typically few 100 evts/day/t**

- Intense source
- Very short baselines available (5-50 m)

- **Alternation of reactor ON/OFF periods**

- moderate overburden compensated by accurate measurement of the cosmogenic component.

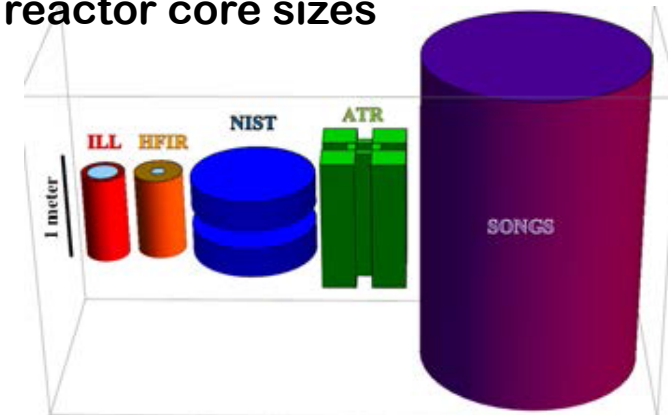
- **Highly enriched fuel**

- Well known ^{235}U fission spectrum.

- **But challenging reactor-induced backgrounds (γ and n)**

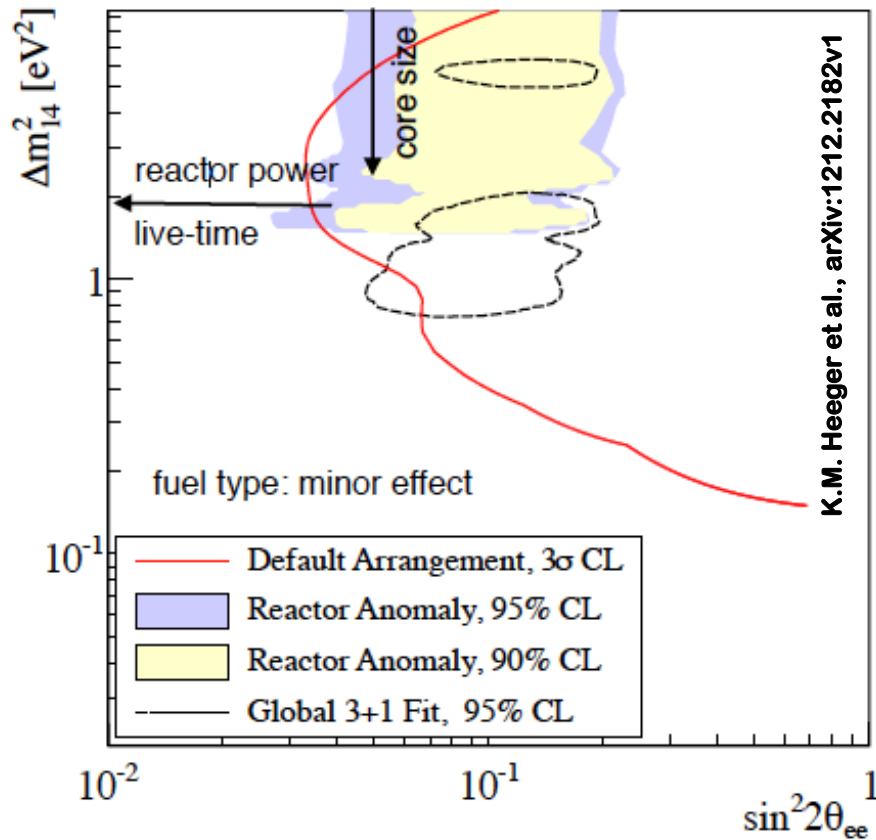
→ Requires comprehensive site characterization.

Typical reactor core sizes

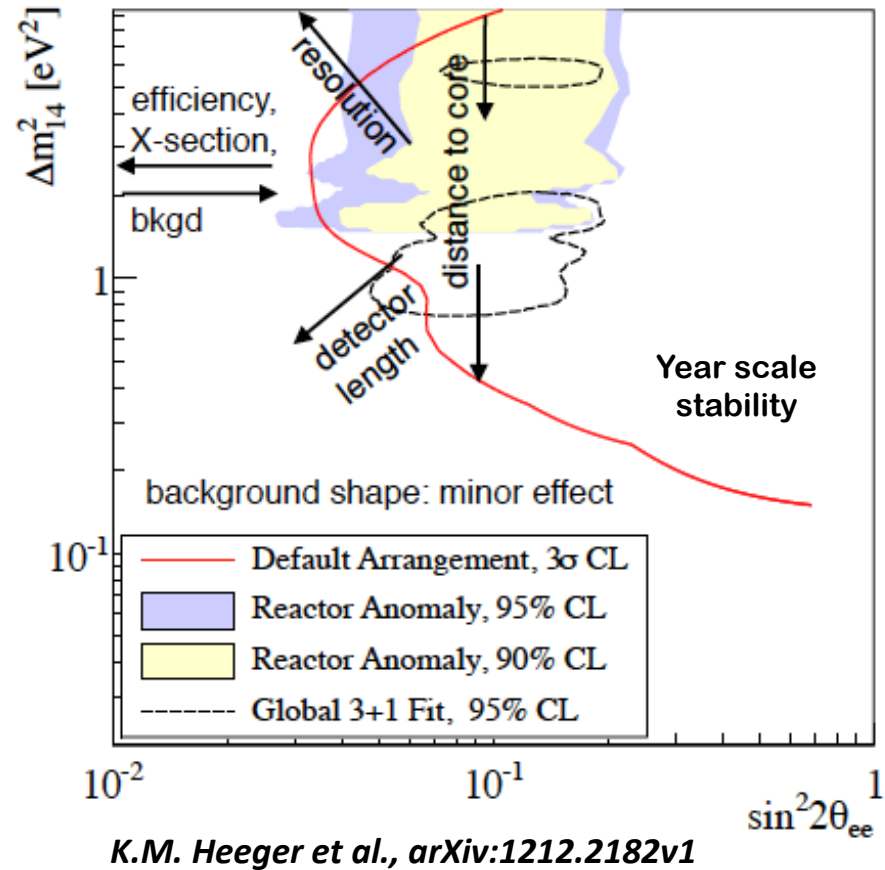


Key experimental parameters




















Core



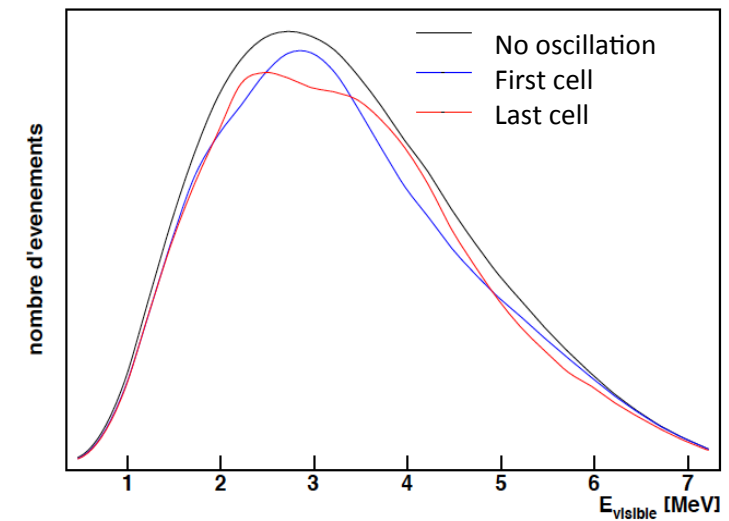
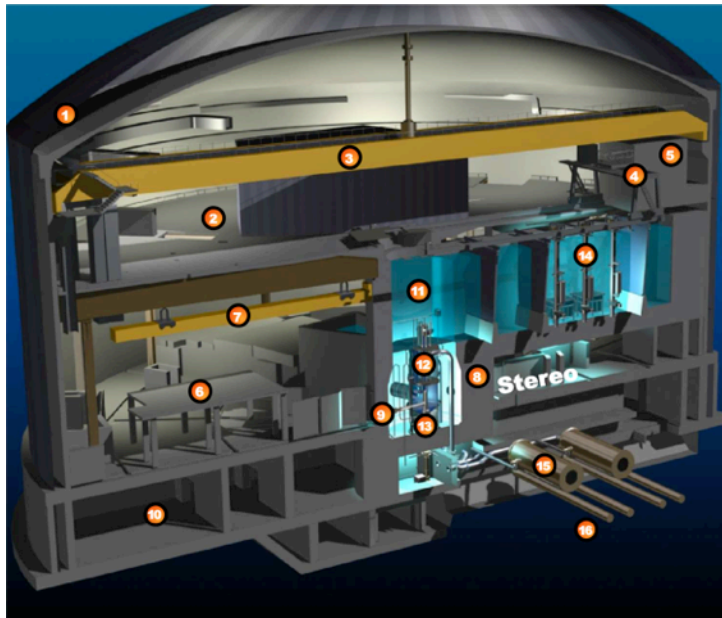
Detector



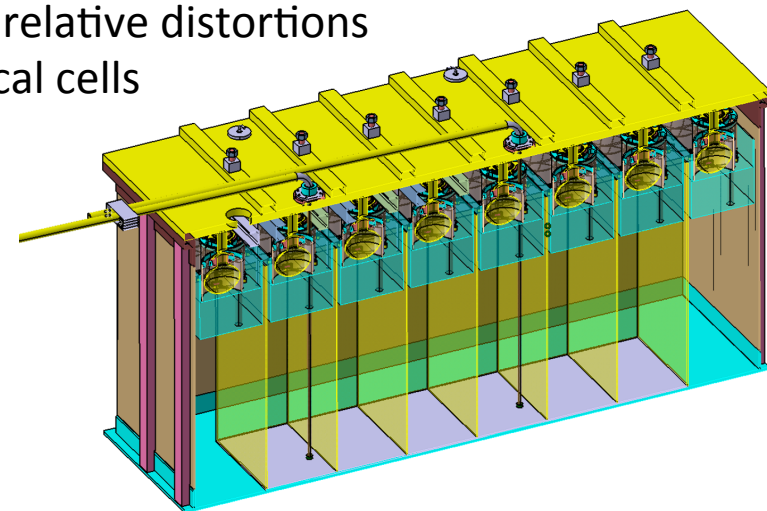
Reactor Proposals

	Gd	^6Li	Highly Segmented	Moving detector	2 det.
Nucifer (FRA)					
Poseidon (RU)					
Stéreo (FRA)					
Neutrino 4 (RU)					
Hanaro (KO)					
DANSS (RU)					
Prospect (USA)					
SoLid (UK)					

STEREO @ ILL



Look for relative distortions
in identical cells

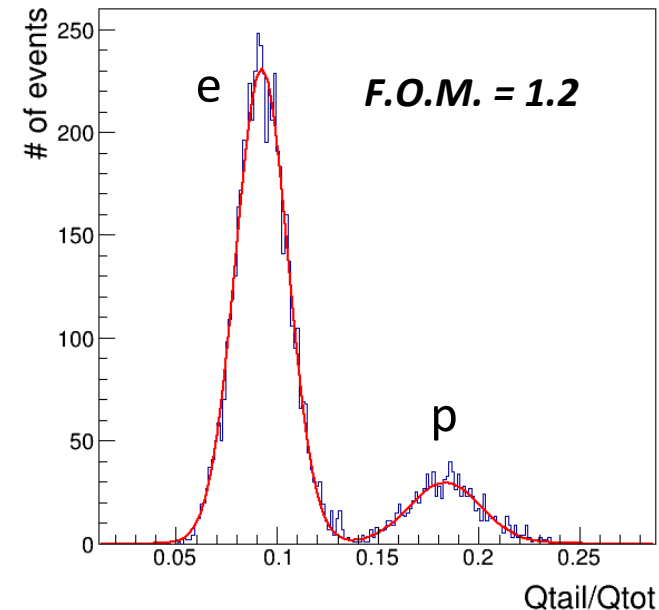
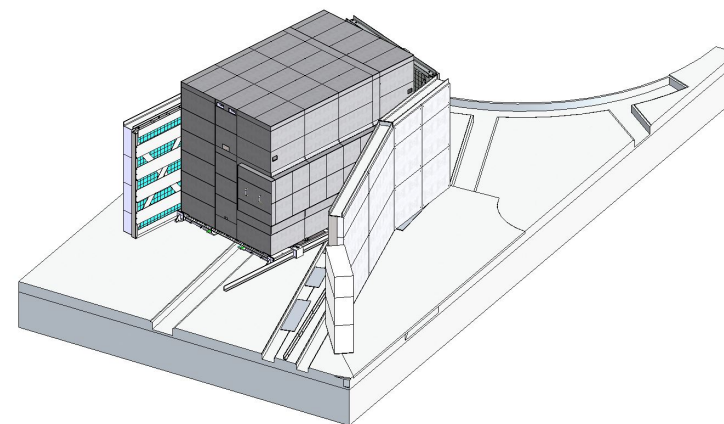


ILL site:

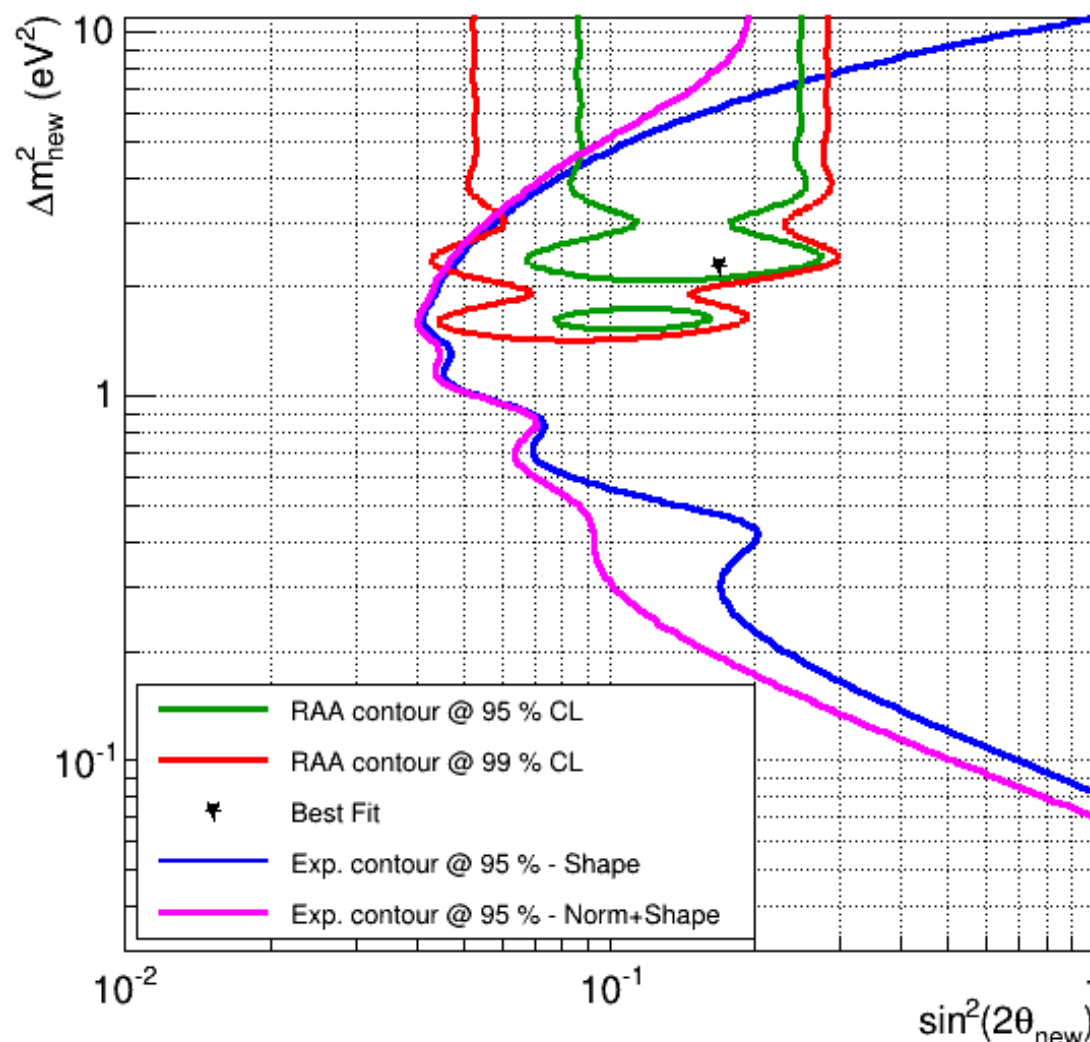
- 57 MW, compact core < 1m
- [8.9–11.1] m from core,
possible extension to 12.3 m.
- 15 mwe overburden
- High level of reactor background

Background rejection

- Comprehensive on site measurements performed last year: μ , n and γ backgrounds.
- Sequential installation/validation of external shielding before detector installation early 2015.
- Muon induced:
 - Overburden
 - Active outer-crown and μ -veto
 - Liquid scintillator optimized for PSD capability and light yield

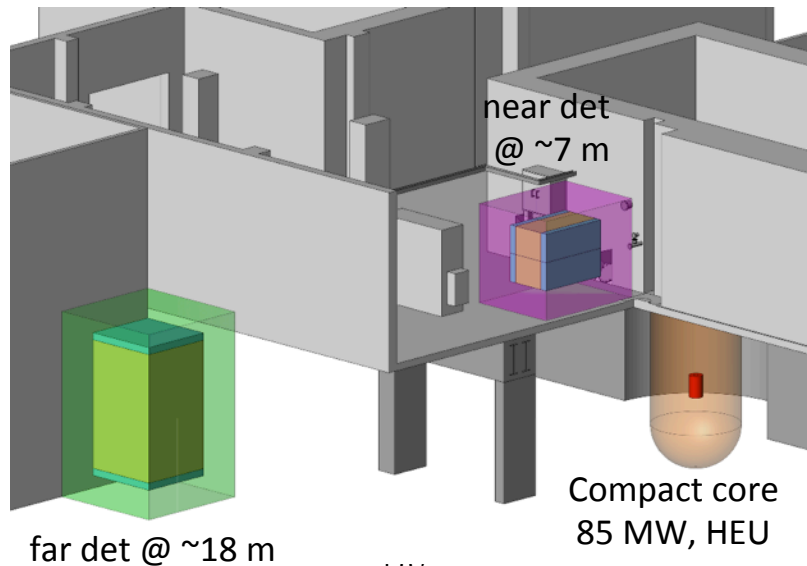


STEREO Sensitivity



- 300 days, $L_0 = 10 \text{ m}$
- $E_{\text{prompt}} > 2 \text{ MeV}$, $E_{\text{delayed}} > 5 \text{ MeV}$
- $\sim 410 \nu_e / \text{day}$
- $\delta E_{\text{scale}} = 2\%$
- All syst. of predicted spectra
- $S/B = 1.5$, $1/E + \text{flat model}$
- Norm 4%
- Start data taking mid-2016

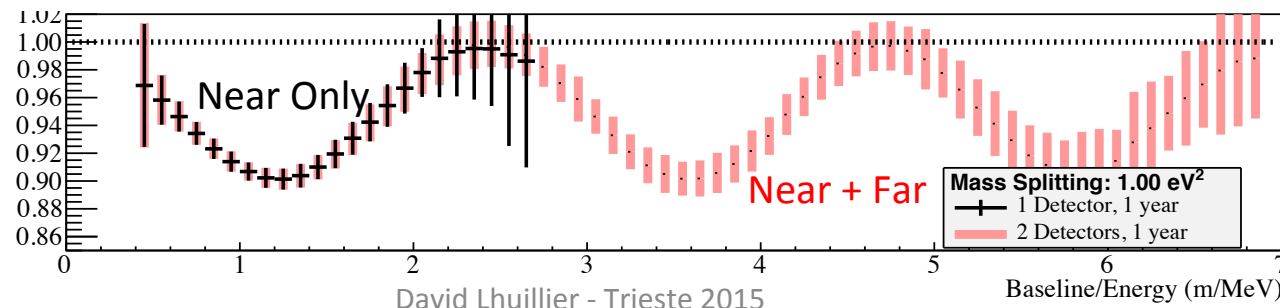
PROSPECT Experiment



Physics Goals

- Search for sterile ν_e oscillations at short-baseline.
- Probe and resolve “reactor anomaly”.
- Precision measurement of reactor ν_e spectrum for physics and safeguards.

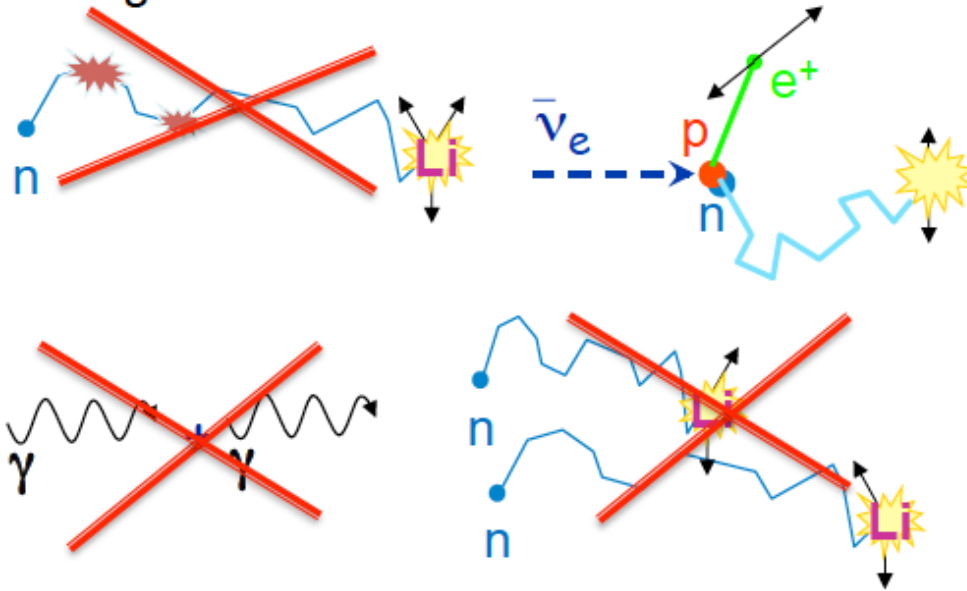
Map out L/E Oscillations



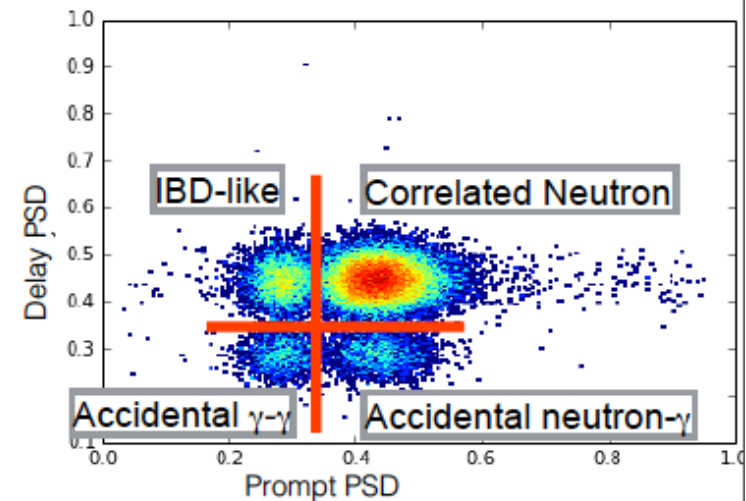
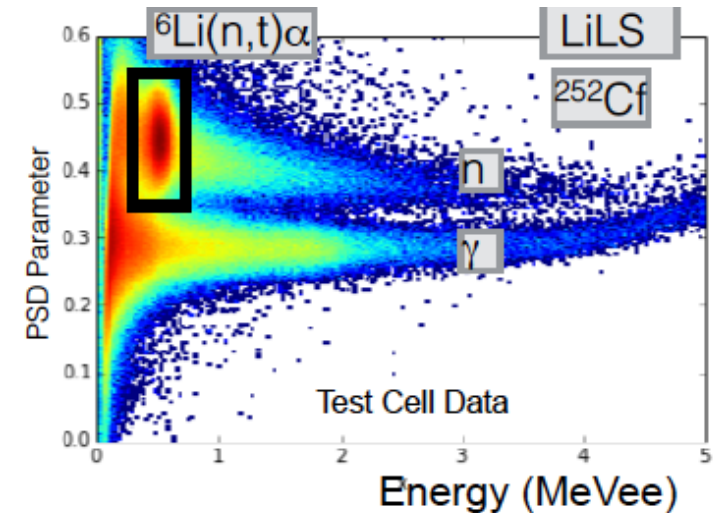
6Li-doped liquid scintillator

- ^6Li -capture, Pulse Shape Discrimination, and topology from segmentation

- Strong rejection of accidental and correlated backgrounds



- Using simulation/deployment data to understand and mitigate electromagnetic-neutron capture correlated backgrounds



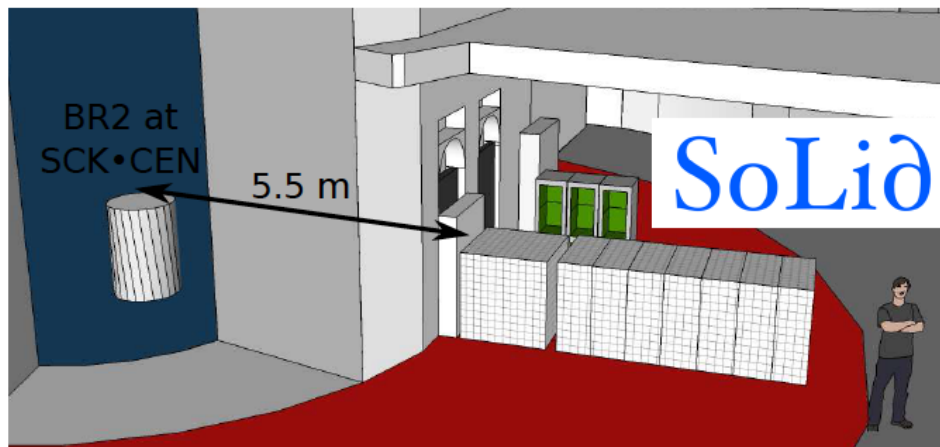
SoLi θ Experiment

BR2 REACTOR, Mol, Belgium

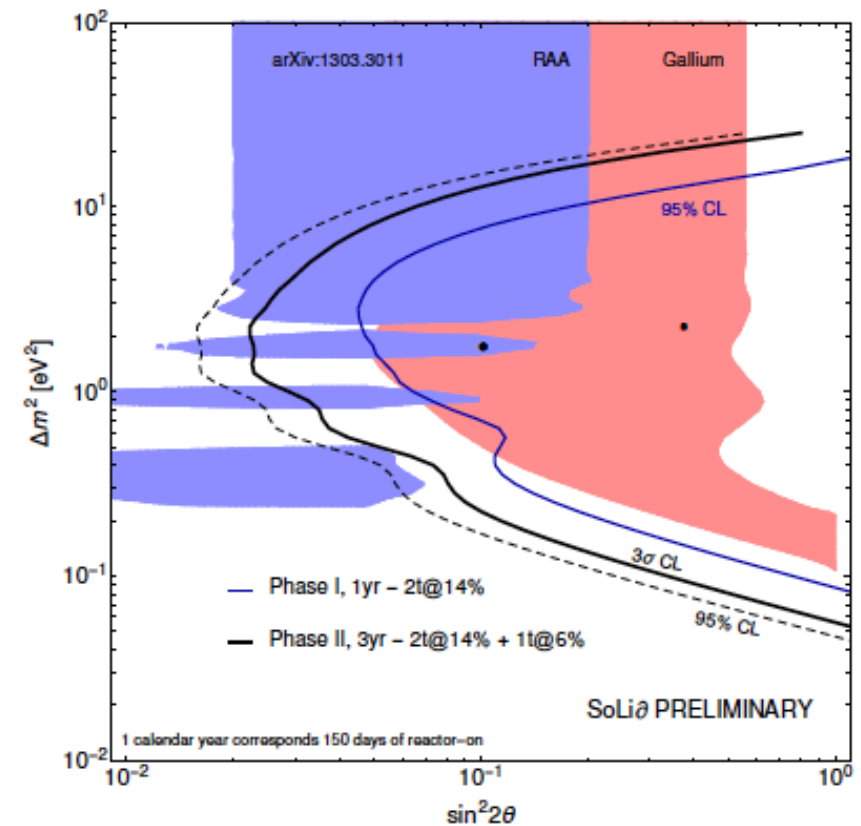
- Core: 45-80 MW, HEU fuel
- Favorable reactor background level

DETECTOR

- **Novel type of composite solid scintillator detector** (PVT + $^6\text{LiF}:\text{ZnS}$)
- 2.88 t fiducial volume, highly segmented.
- 288 kg prototype under test

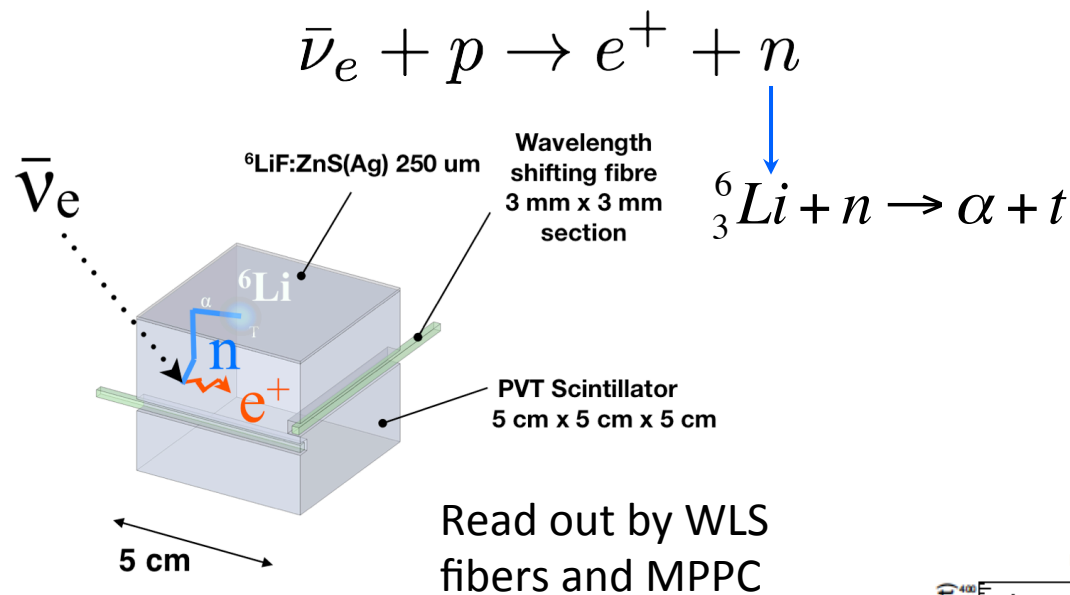


Expected sensitivity

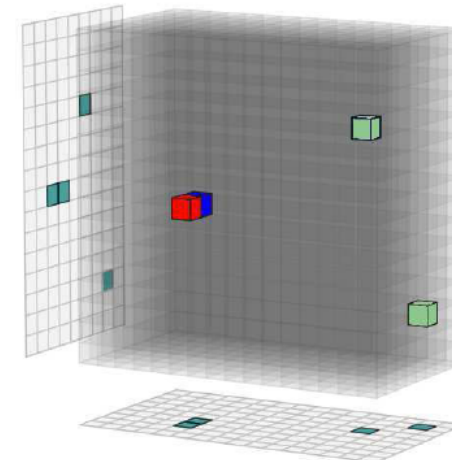


<http://arxiv.org/abs/1510.07835>

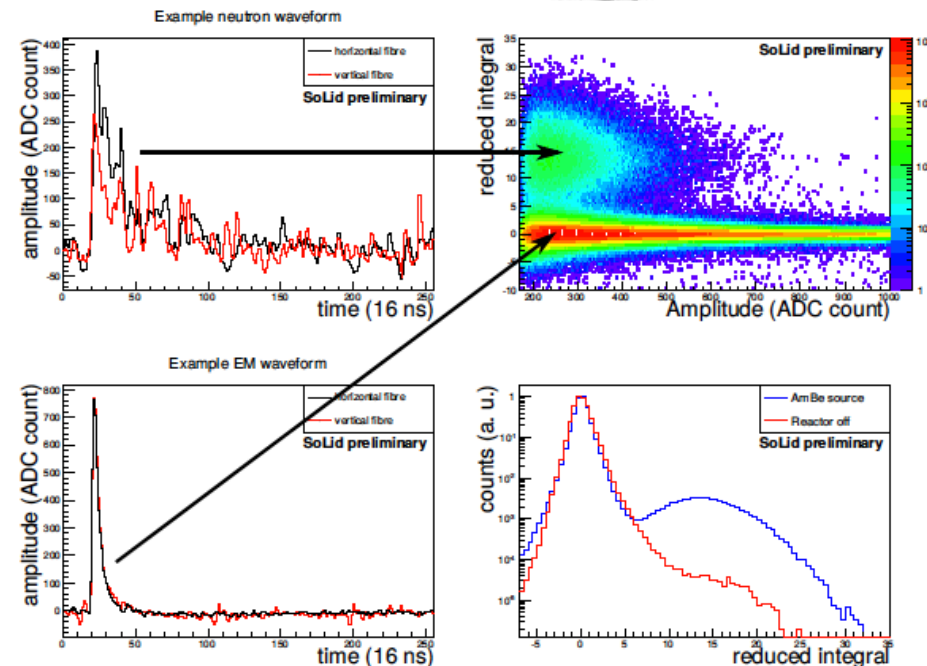
SoLiD detection



IBD candidate: **positron** + **neutron**
(+ accidental gammas)



- Very **discriminant neutron signal** in ${}^6\text{LiF:ZnS}$. High neutron- γ rejection factor
- 3D reconstruction close to interaction point : **high background rejection capability using topological information of IBD.**
- Limited resolution due to light collection

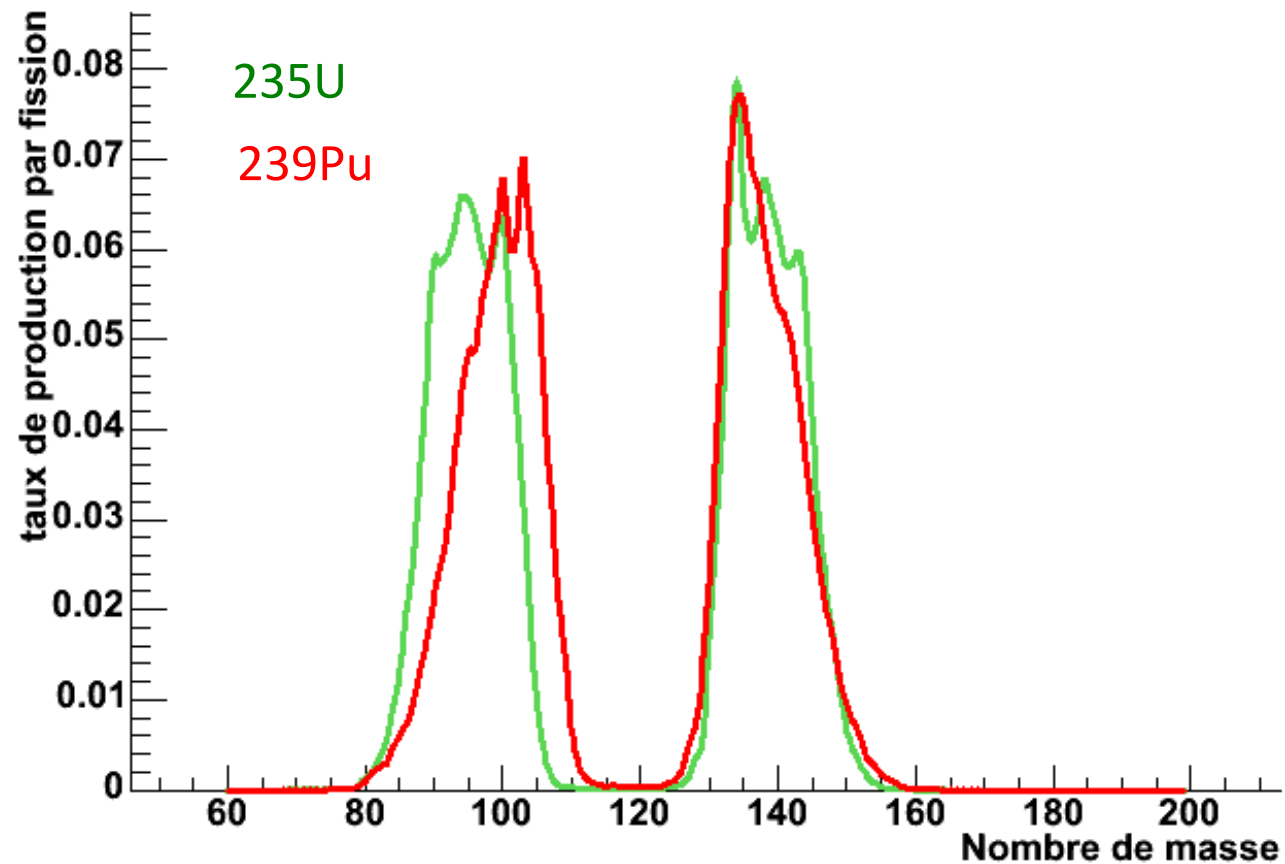


Conclusion

- Neutrinos connect a wide range of physics, from nuclear decays to cosmology.
- Modern experiments at reactors perform precision measurements, becoming sensitive to few percent effects on the neutrino flux and spectrum shape. The last mixing angle θ_{13} has been measured and a new “anomaly” has been revealed, pointing to a possible new sterile neutrino state.
- A mature detection technology allows the development of the first application of neutrino to non-proliferation. Common effort of all short baseline experiment is ongoing to improve the background rejection in small detectors. Lithium doped scintillator should allow an efficient online rejection of the cosmic-rays induced background, which is the ultimate limitation of the sensitivity of neutrino measurements close to surface.
- Reactor neutrino benefit from the huge information available in nuclear databases. The comparison of measurements and absolute predictions is now reaching a highlevel of accuracy providing a sensitive probe of the quality of the nuclear data.

Backup Slides

Mass distribution



Forbidden Decays

Konopinsky and Ulenbeck *Phys. Rev. vol. 60, June 1941.*

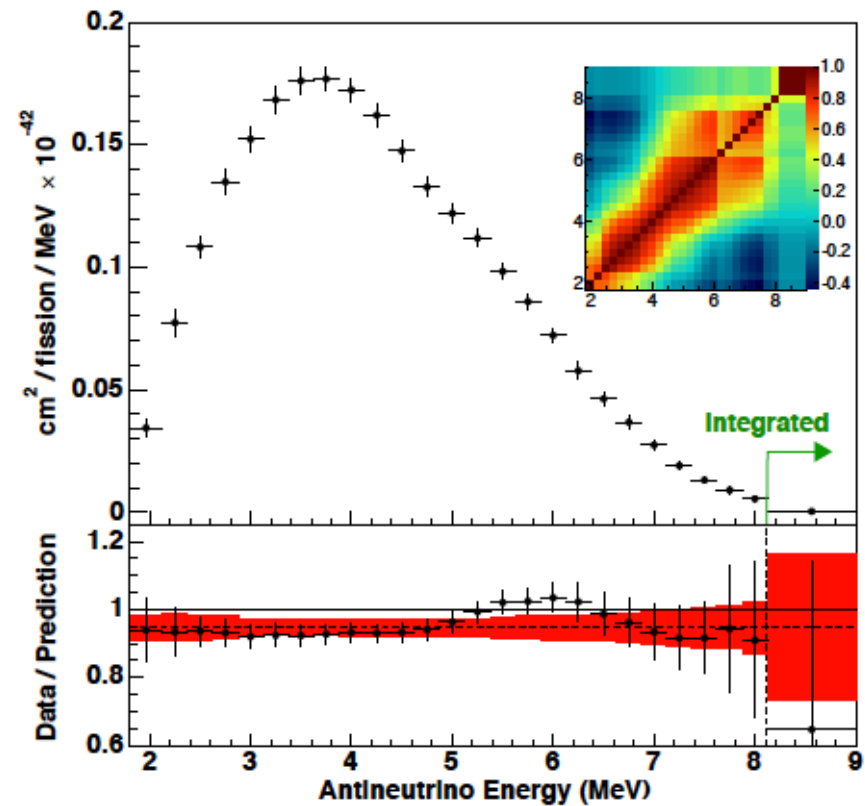
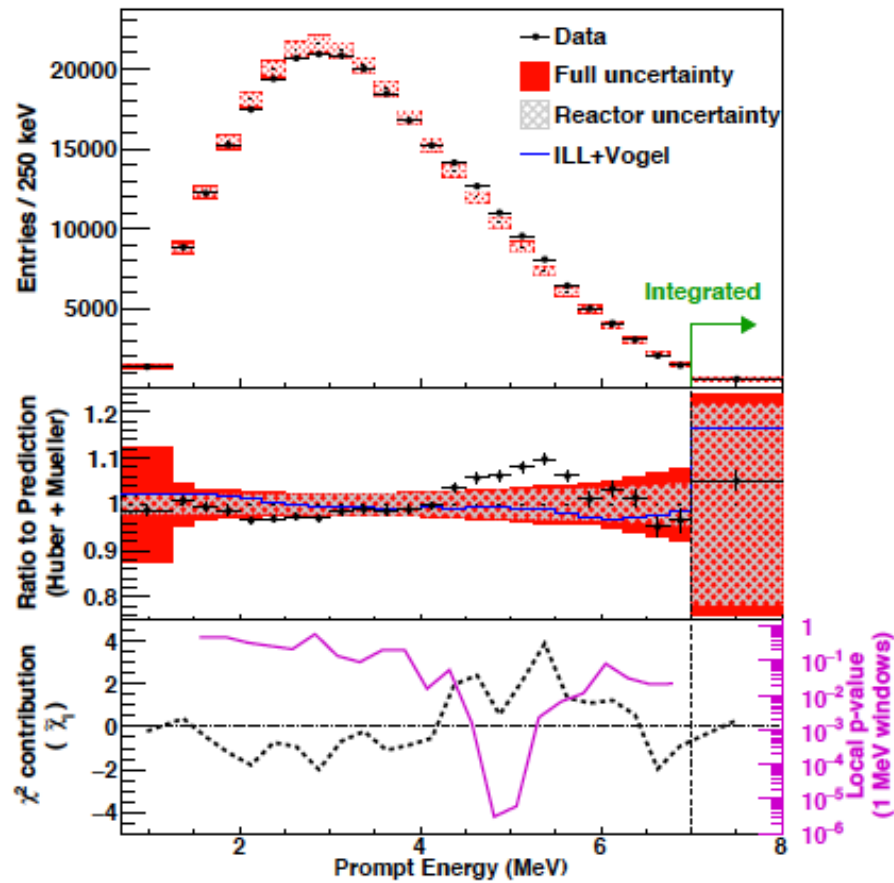
Shape factor $C(E)$ = Nuclear matrix element \times Lepton matrix element

ΔJ^π	Nuclear matrix	Lepton matrix
Vector		
$0^-, 1^-$	$ \int \mathbf{r} ^2$	$\frac{1}{3}q^2 L_0 + 2L_1 + M_0 - \frac{2}{3}qN_0$
$0^-, 1^-$	$ \int \boldsymbol{\alpha} ^2$	L_0
$0^-, 1^-$	$(\int \boldsymbol{\alpha}) \cdot (\int \mathbf{r})$	$\frac{1}{3}qL_0 - N_0$
Axial		
0^-	$ \int \boldsymbol{\sigma} \cdot \mathbf{r} ^2$	$\frac{1}{9}q^2 L_0 + M_0 - \frac{2}{3}qN_0$
0^-	$ \int \gamma_5 ^2$	L_0
0^-	$(\int \boldsymbol{\sigma} \cdot \mathbf{r}) \cdot (\int \gamma_5)$	$\frac{1}{3}qL_0 - N_0$
$0^-, 1^-$	$ \int \boldsymbol{\sigma} \times \mathbf{r} ^2$	$\frac{1}{6}q^2 L_0 + \frac{1}{2}L_1 + M_0 + \frac{2}{3}qN_0$
$0^-, 1^-, 2^-$	$ B_{ij} ^2$	$\frac{1}{12}q^2 L_0 + \frac{3}{4}L_1$

Adds dependence on lepton momenta



Daya Bay

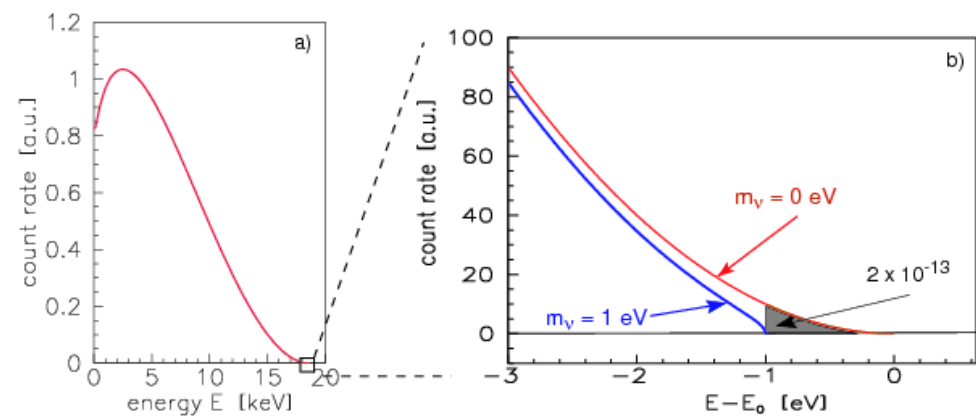
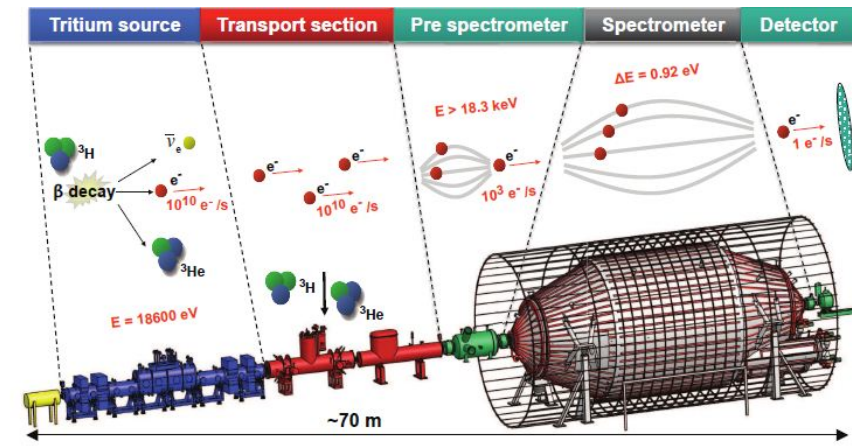


Open questions

Absolute Mass

Katrin experiment

- Direct measurement
- Looking at the end of the tritium β -spectrum

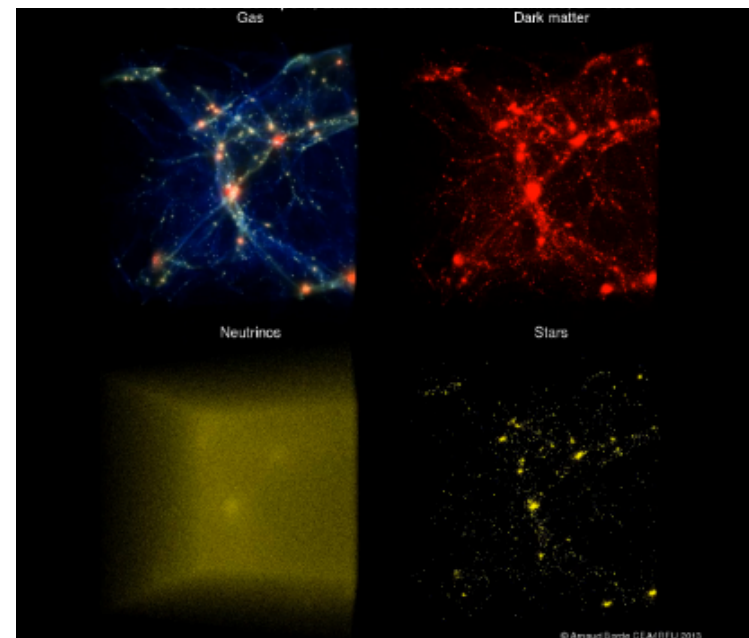


Open questions

Absolute Mass

All combined cosmology data

$$\Sigma M_\nu < 0.12 \text{ eV (95\% C.L.)}$$

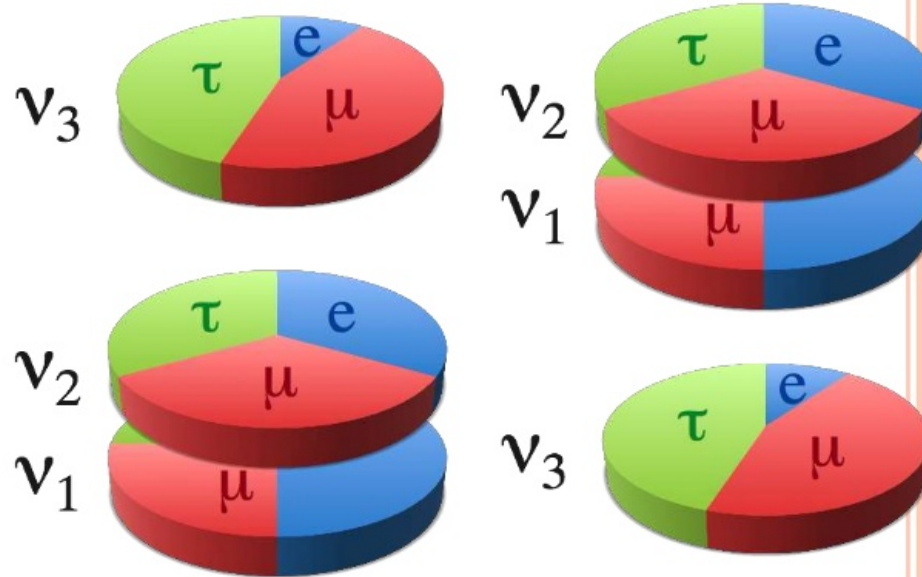


<http://arxiv.org/abs/1506.05976>

Open questions

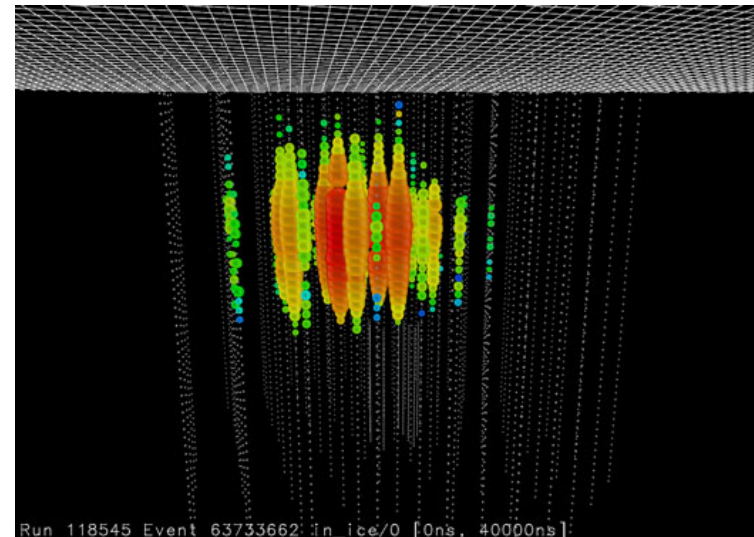
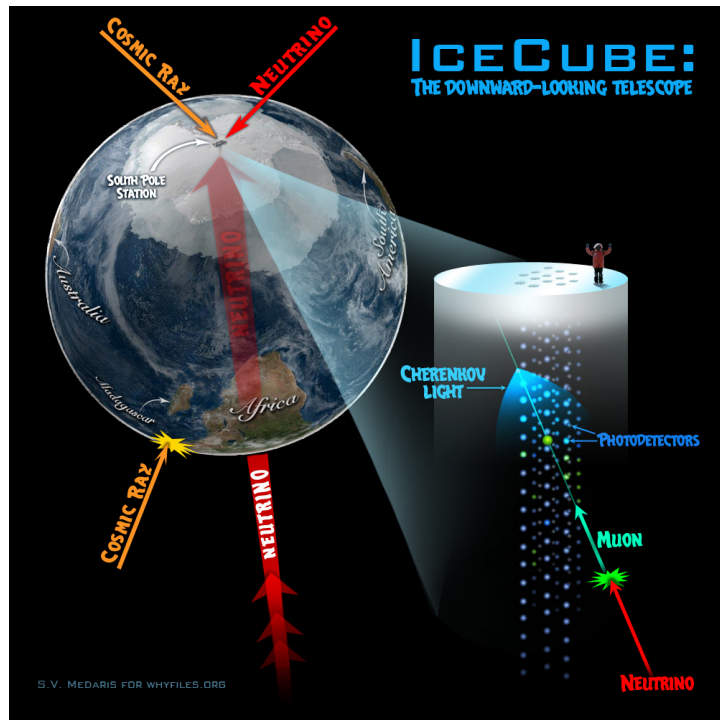
Mass hierarchy

CURRENT UNDERSTANDING OF NEUTRINO PARAMETERS



S. K. Verma

Neutrino Astronomy



PeV neutrino event

Geo-Neutrinos

