

Dirac Medallists 2017



Charles H. Bennett



David Deutsch



Peter W. Shor

Citation:

for pioneering work in applying fundamental concepts of quantum mechanics to solving basic problems in computation and communication, and therefore bringing together the fields of quantum mechanics, computer science and information.

When it all began ...

Villa Gualino meeting in Fall, 1994 (ISI Torino)

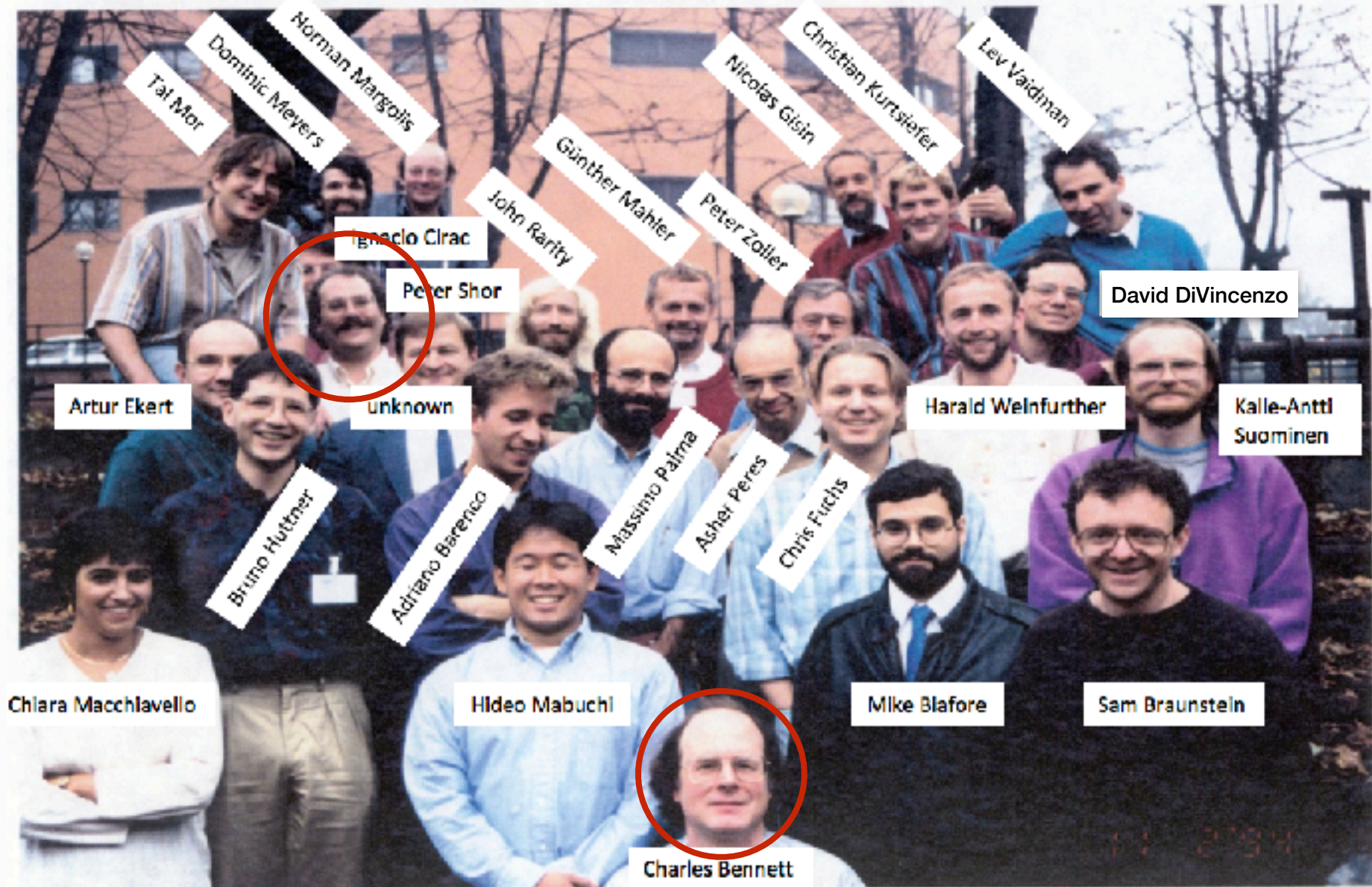


J. I. Cirac

First presentation of ion trap quantum computer

When it all began ...

Villa Gualino meeting in Fall, 1994 (ISI Torino)





Quantum Computing, Quantum Simulation, and Quantum Communications

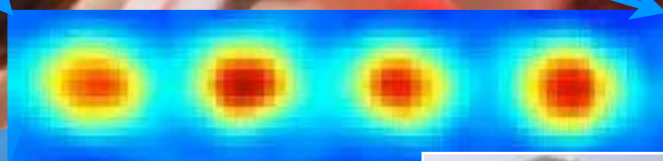
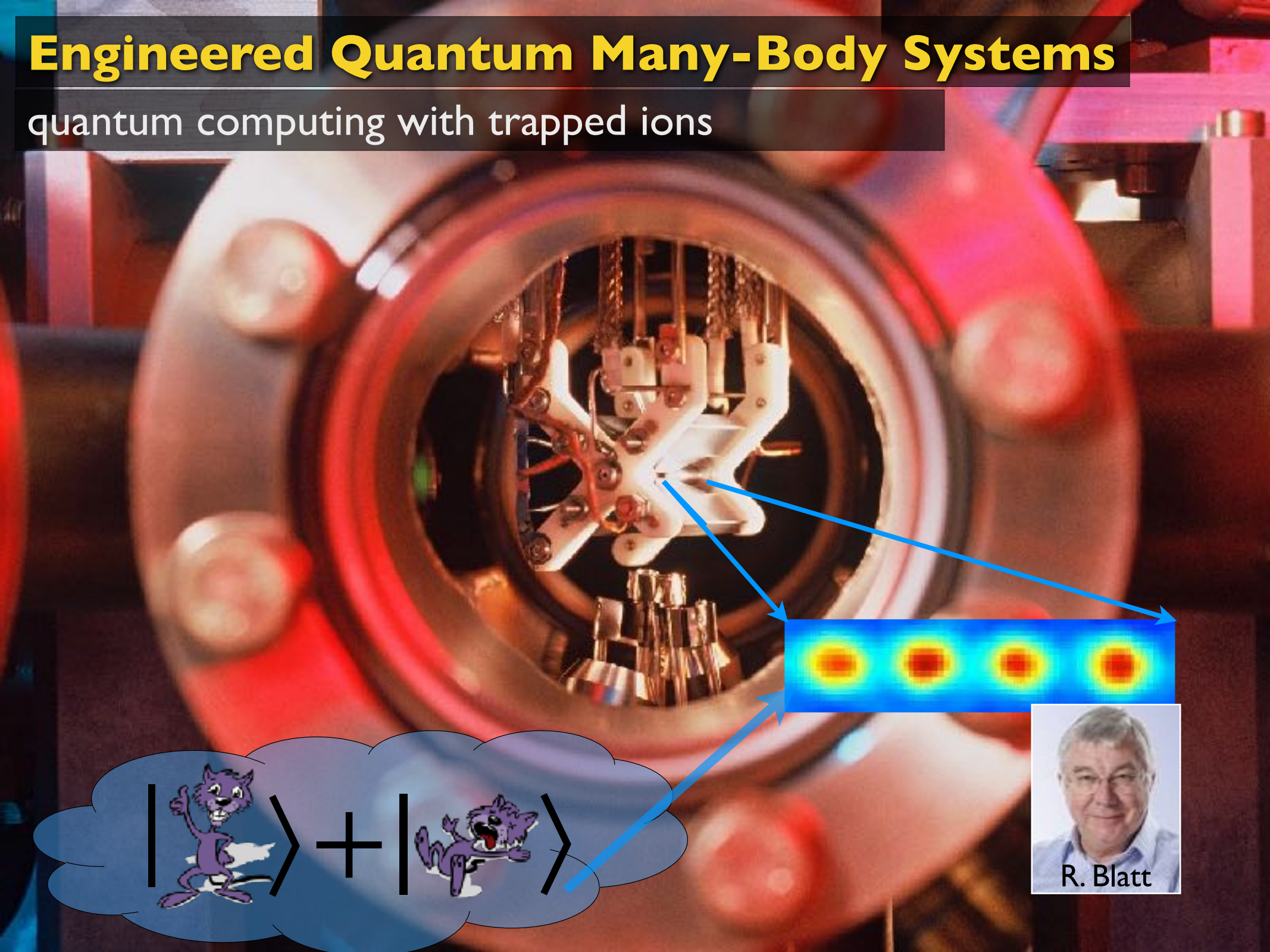
... with Quantum Optics

Peter Zoller
University of Innsbruck &
IQOQI Austrian Academy of Sciences

ICTP Dirac Medal 2017
March 14 2018

Engineered Quantum Many-Body Systems

quantum computing with trapped ions

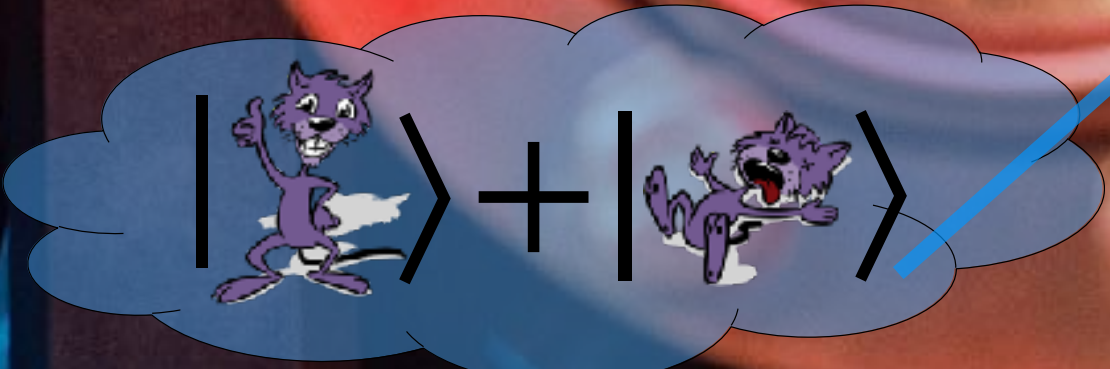
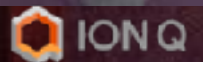


Engineered Quantum Many-Body Systems

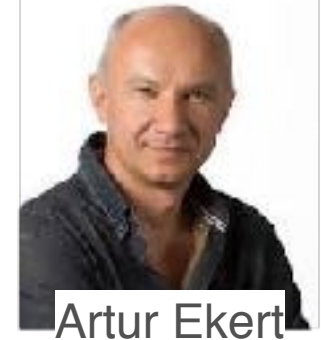
quantum computing with trapped ions



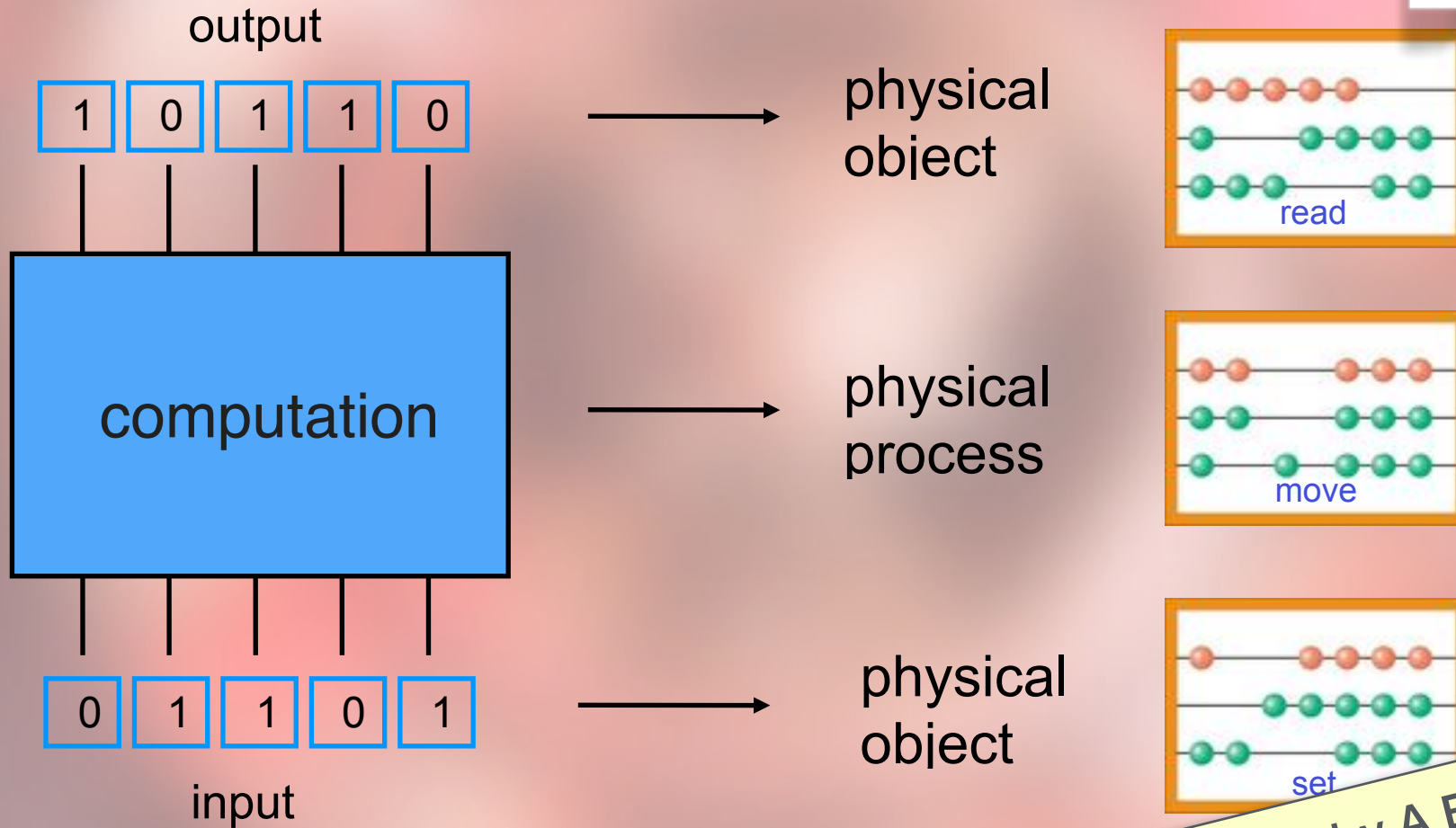
C. Monroe



Information and physics



computing as a physical process

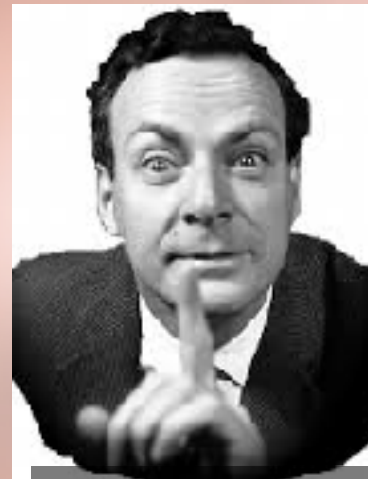
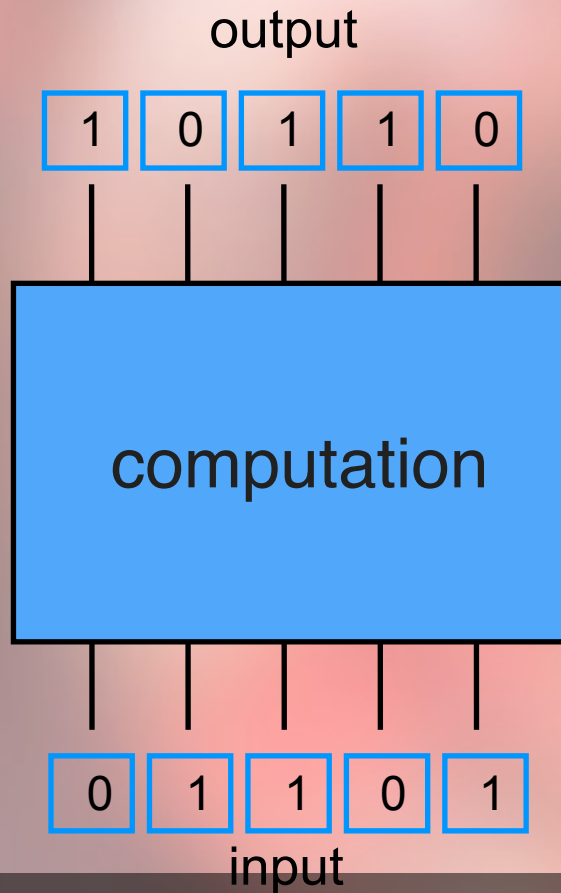


Our present computers process information according to the laws of *classical* physics!

Presentation by A Ekert at ICAP
Aug 1994, Boulder
(C. Wieman & D. Wineland)

Information and physics

computing as a physical process



R. Feynman

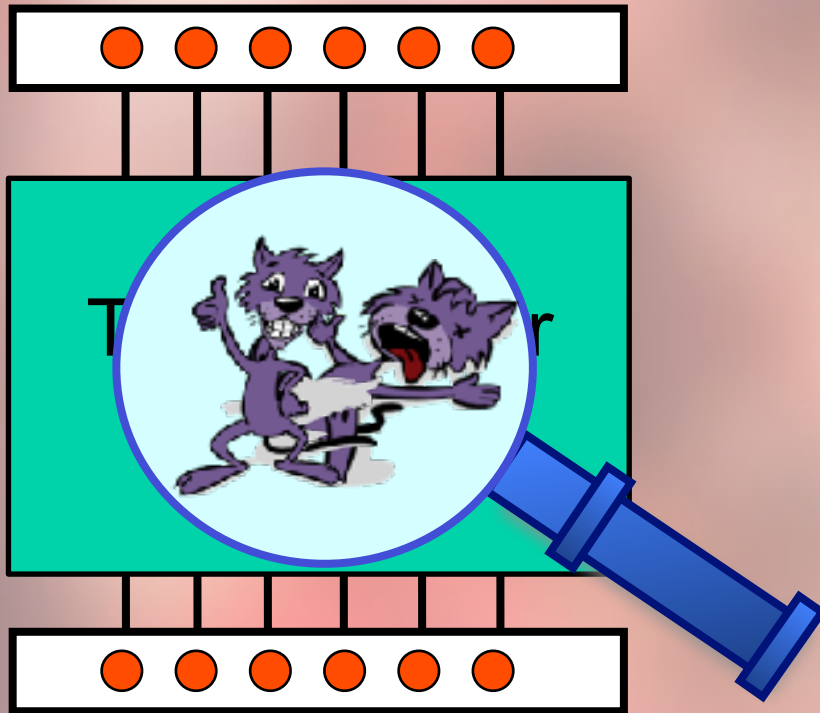


D. Deutsch

“At a fundamental level nature obeys the laws of *quantum physics*. At a fundamental level information science must be a *quantum information science*.”

Information and physics

Quantum Processor



Why Quantum Computing?

Technology: to beat Moore's law

Computer Science: new complexity classes

Physics: to learn about quantum theory

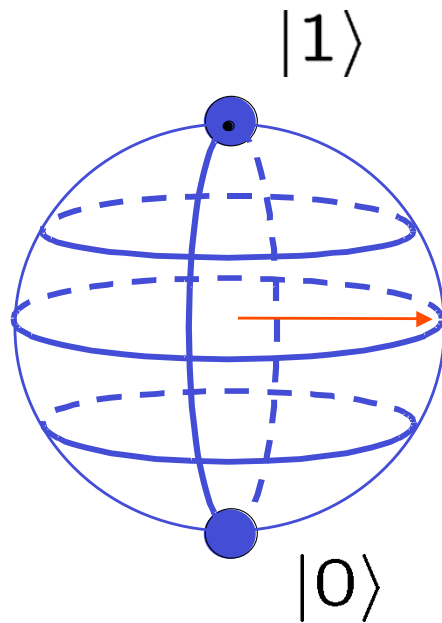
... perform tasks beyond classical computing?

Qubit ...

classical bit
0 or 1

quantum-bit or qubit
0 or 1 or

01



$|0\rangle + |1\rangle$



... and quantum registers

classical bit
0 or 1

quantum-bit or qubit
0 or 1 or 01

classical registers

010

quantum registers

000 001 010 011
100 101 110 111



Erwin Schrödinger:
Entanglement

How big is quantum memory?



n qubits

example: $n \sim 300$

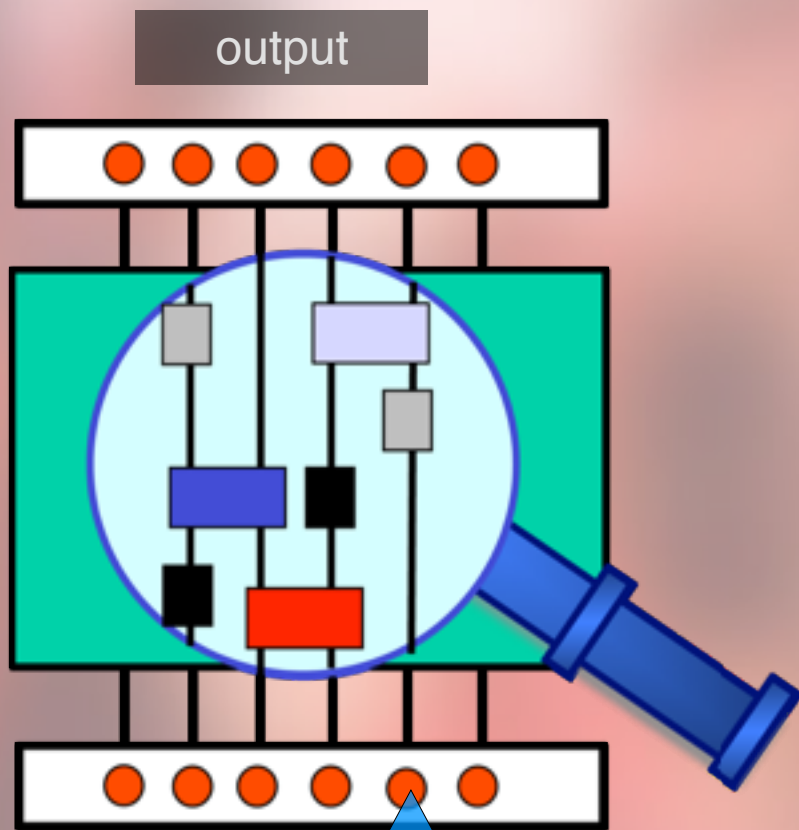
$2^{300} \sim$ atoms in the universe

Hilbert space
is HUGE!



„...it is difficult to simulate quantum mechanics on a classical computer.“ (Richard Feynman, 1986)

How a quantum computer works ...



input

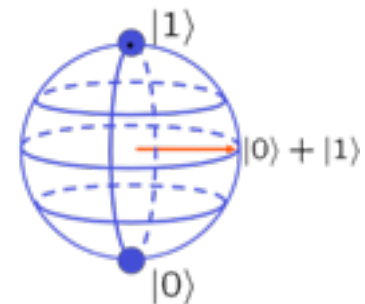
Quantum parallel processing

Single qubit gate

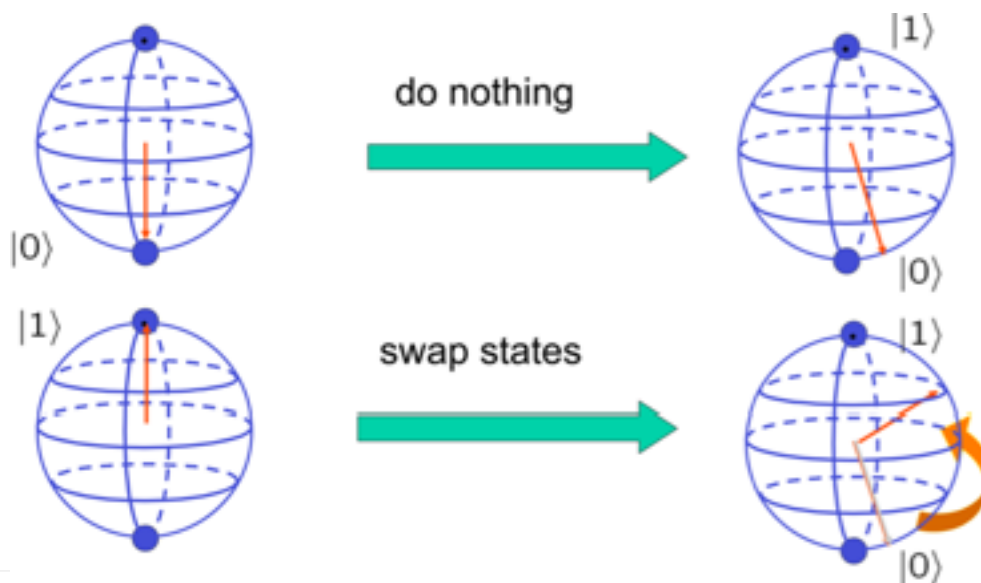
$$|0\rangle + |1\rangle \quad |0\rangle - |1\rangle$$



$$|0\rangle \quad |1\rangle$$

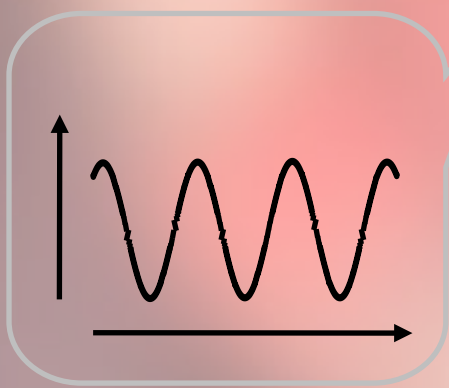
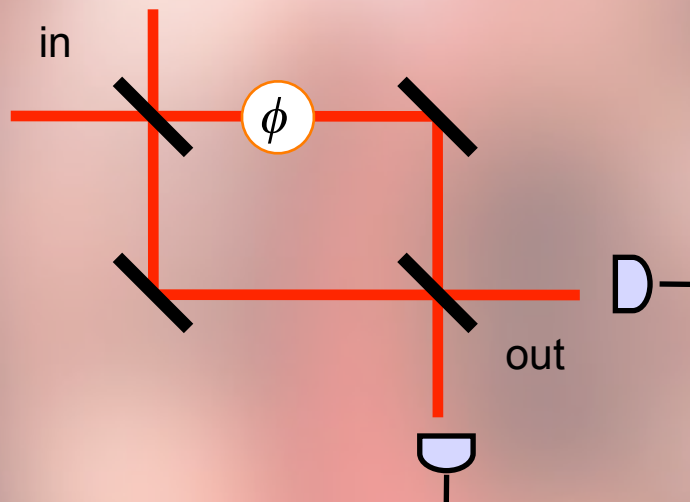


Two-qubit quantum gate



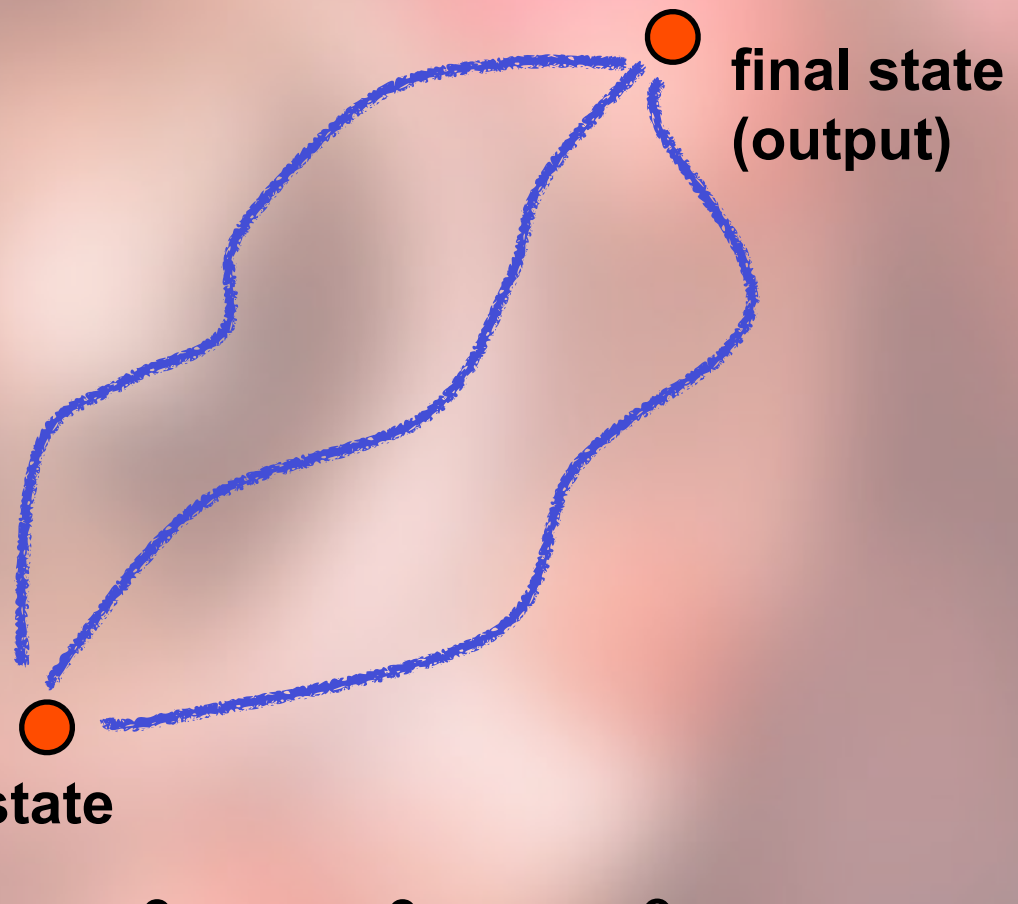
Quantum Parallelism & Algorithms

interferometer



interference pattern

computational paths can interfere in Hilbert space



$$e^{i\theta_1} + e^{i\theta_2} + e^{i\theta_3} + \dots$$

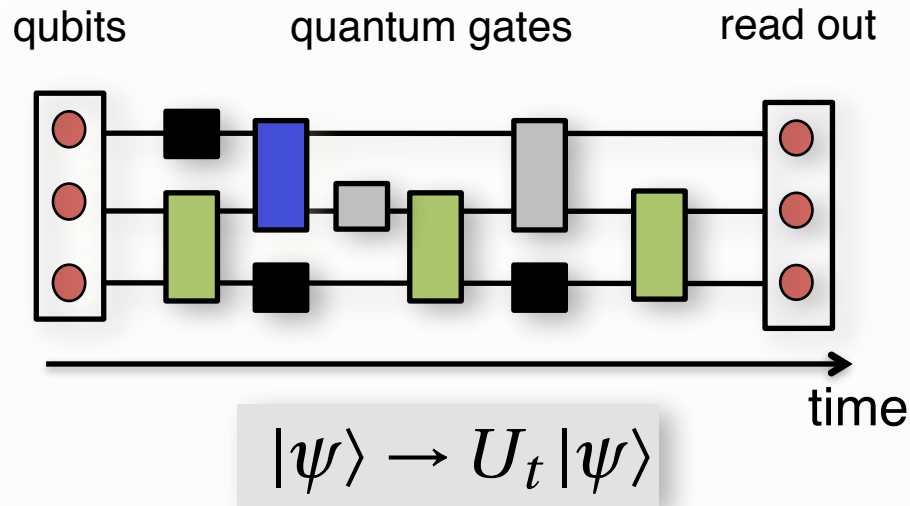
Quantum Computing with Trapped Ions



- **general purpose quantum computing**

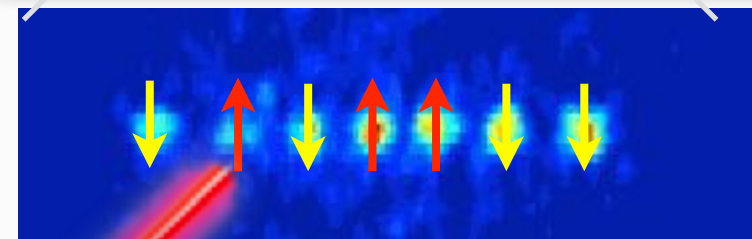
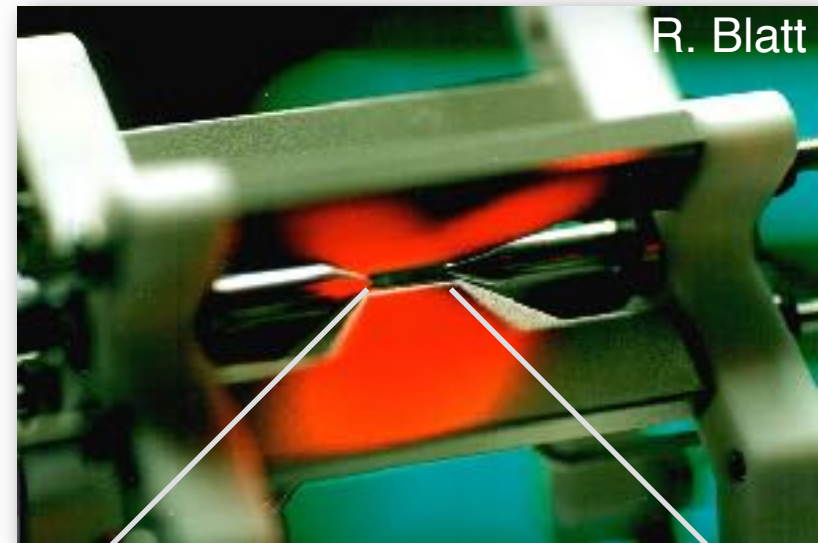
- **atomic physics: trapped ions**

quantum logic network model



coherent Hamiltonian evolution

- quantum gates
- deterministic

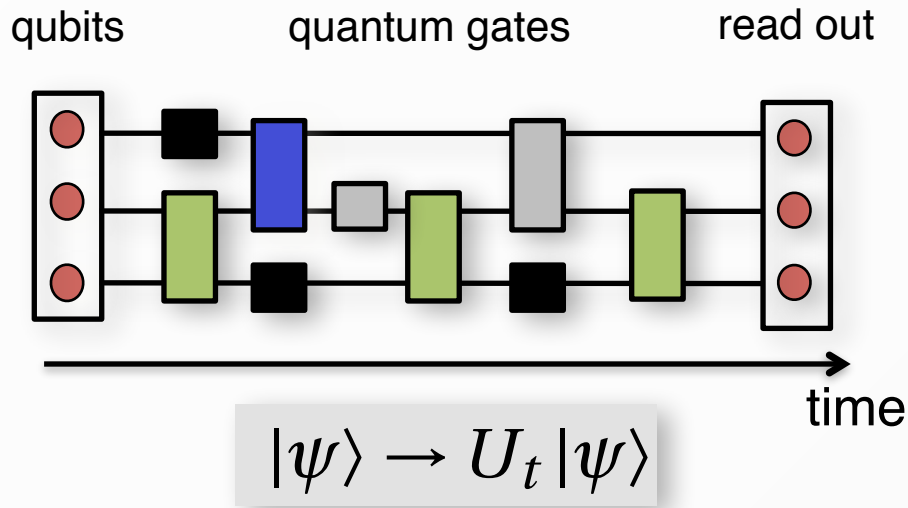


Exp.: Innsbruck, NIST, JQI, MIT, Mainz, MPQ ...

- **general purpose quantum computing**

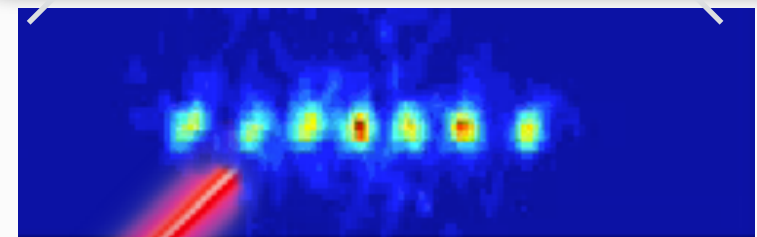
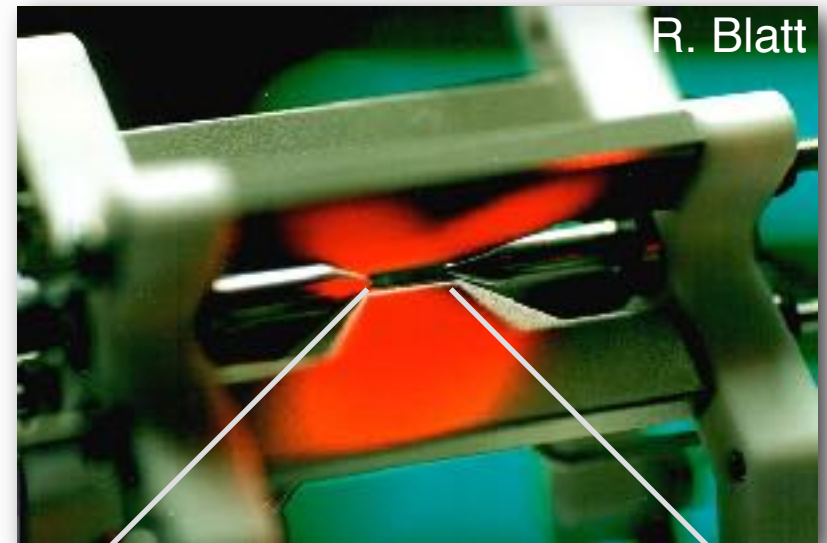
- **atomic physics: trapped ions**

quantum logic network model



coherent Hamiltonian evolution

- quantum gates
- deterministic



phonon bus

laser

Exp.: Innsbruck, NIST, JQI, MIT, Mainz, MPQ ...



Quantum operations & compiler:

Individual light-shift gates

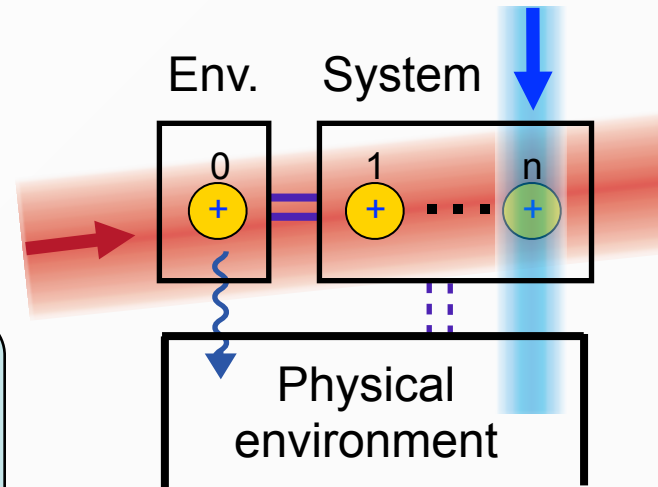
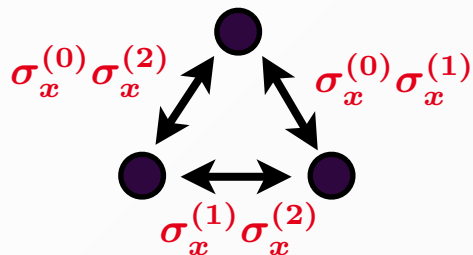
$$\sigma_z^{(0)}, \sigma_z^{(1)}, \sigma_z^{(2)}$$

Collective spin flips

$$S_x, S_y$$

Mølmer-Sørensen gate

$$S_x^2 = \sigma_x^{(0)} \sigma_x^{(1)} + \sigma_x^{(1)} \sigma_x^{(2)} + \sigma_x^{(0)} \sigma_x^{(2)}$$

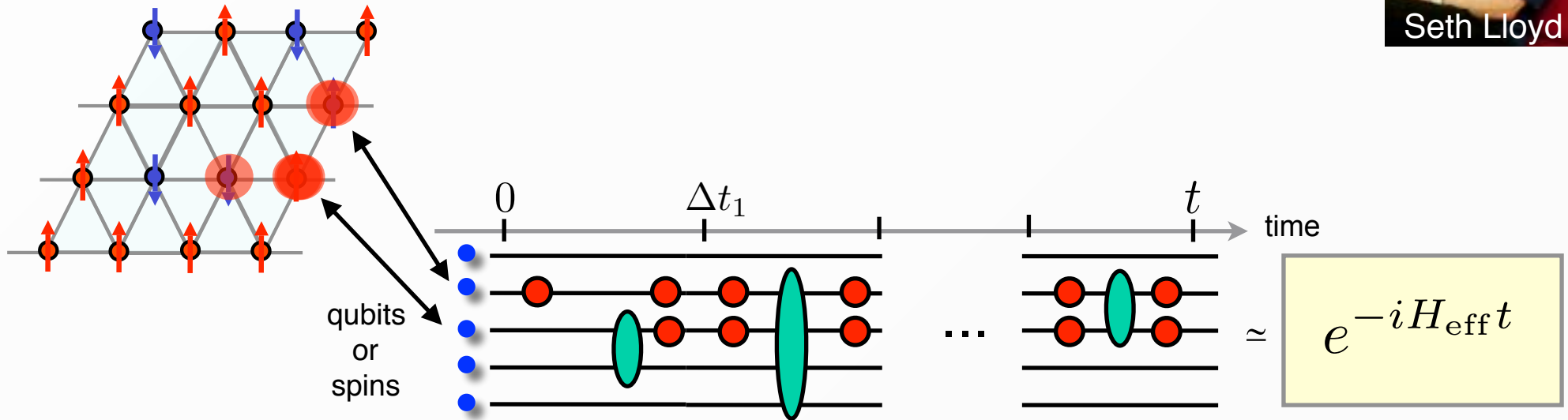


Coupling of environment with physical environment

Optical pumping of „environmental“ ion



Digital Quantum Simulation



idea: approximate time evolution by a stroboscopic sequence of gates

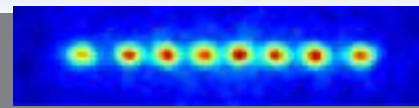
$$U(t) \equiv e^{-iHt/\hbar} = e^{-iH\Delta t_n/\hbar} \dots e^{-iH\Delta t_1/\hbar}$$

Trotter expansion:

$$e^{-iH\Delta t/\hbar} \simeq \underbrace{e^{-iH_1\Delta t/\hbar}}_{\text{first term}} \underbrace{e^{-iH_2\Delta t/\hbar}}_{\text{second term}} \underbrace{e^{\frac{1}{2} \frac{(\Delta t)^2}{\hbar^2} [H_1, H_2]}}_{\text{Trotter errors for non-commuting terms}}$$

$$H = \boxed{-J\sigma_1^z\sigma_2^z} + \boxed{B(\sigma_1^x + \sigma_2^x)}$$

Trotter errors for non-commuting terms



Universal Digital Quantum Simulation with Trapped Ions

B. P. Lanyon,^{1,2*} C. Hempel,^{1,2} D. Nigg,² M. Müller,^{1,3} R. Gerritsma,^{1,2} F. Zähringer,^{1,2}
 P. Schindler,² J. T. Barreiro,² M. Rambach,^{1,2} G. Kirchmair,^{1,2} M. Hennrich,² P. Zoller,^{1,3}
 R. Blatt,^{1,2} C. F. Roos^{1,2}

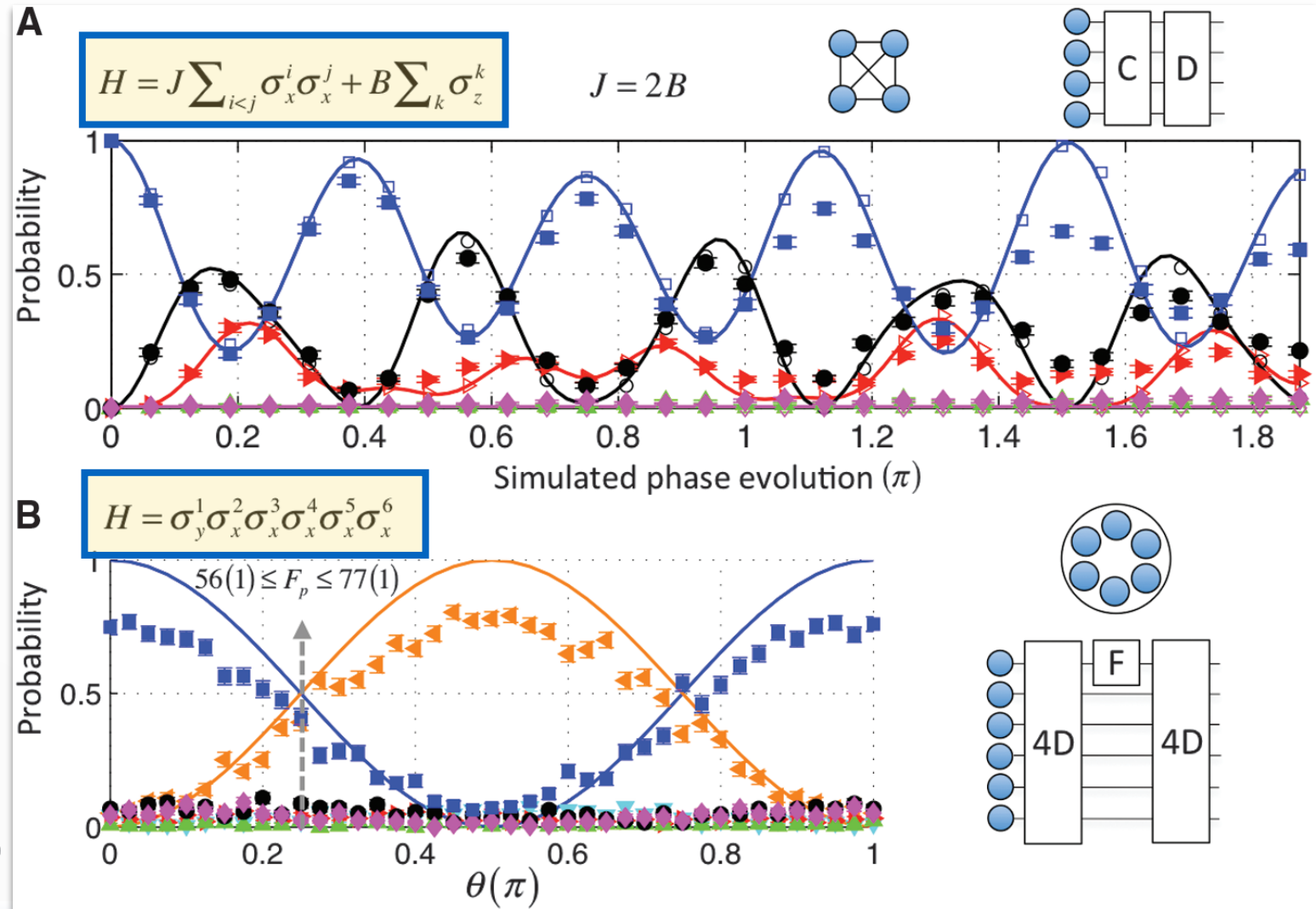


B. Lanyon



C. Roos

4 & 6 Spins



remarks:

- scalability (?)
- error correction (?)

Real-time dynamics of lattice gauge theory on a few-qubit quantum computer

doi:10.1038/nat

Esteban A. Martinez^{1*}, Christine A. Muschik^{2,3*}, Philipp Schindler¹, Daniel Nigg¹, Alexander Erhard¹, Philipp Hauke^{2,3}, Marcello Dalmonte^{2,3}, Thomas Monz¹, Peter Zoller^{2,3} & Rainer Blatt^{1,2}

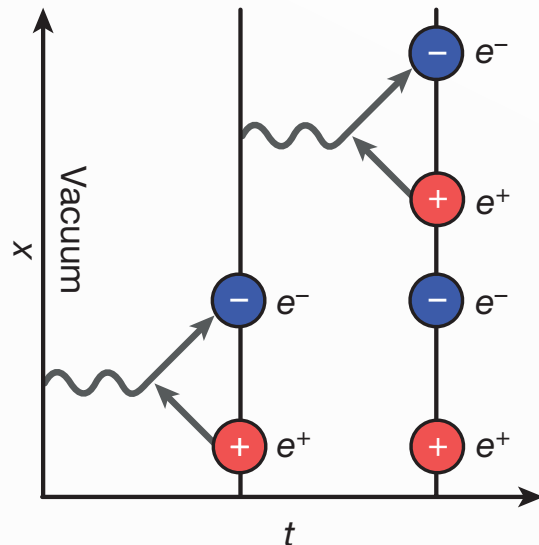


E. Martinez

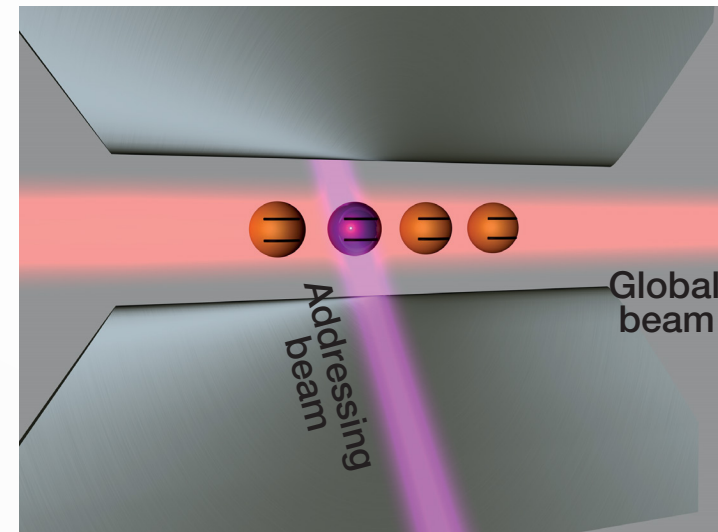


C. Muschik

Schwinger pair production



ion trap quantum computer



Schwinger Model: 1+1D QED

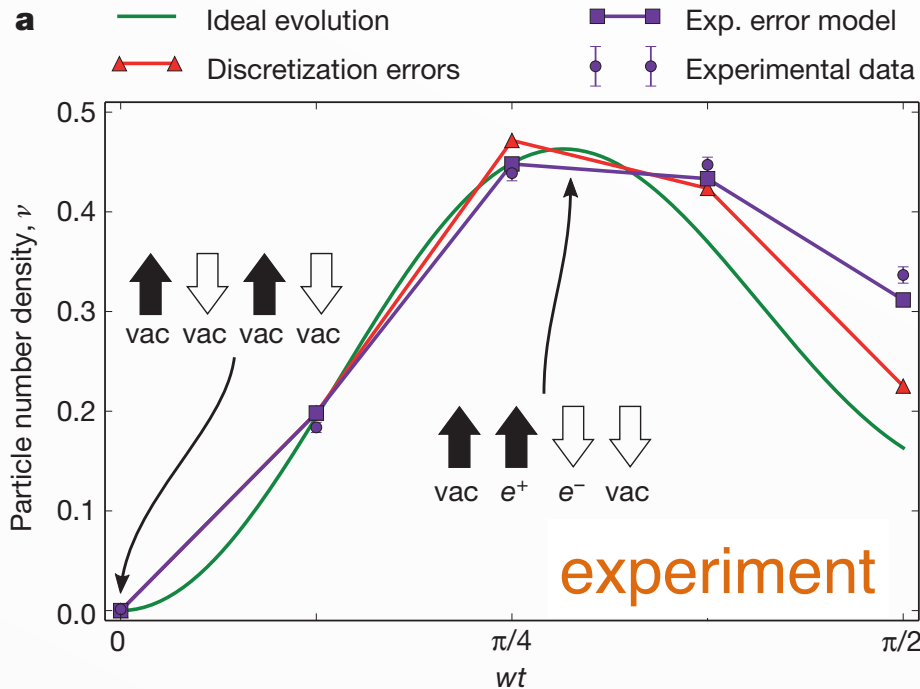


$$\hat{H}_{\text{lat}} = -i w \sum_{n=1}^{N-1} \left[\hat{\Phi}_n^\dagger e^{i\hat{\theta}_n} \hat{\Phi}_{n+1} - \text{h.c.} \right] + J \sum_{n=1}^{N-1} \hat{L}_n^2 + m \sum_{n=1}^N (-1)^n \hat{\Phi}_n^\dagger \hat{\Phi}_n$$

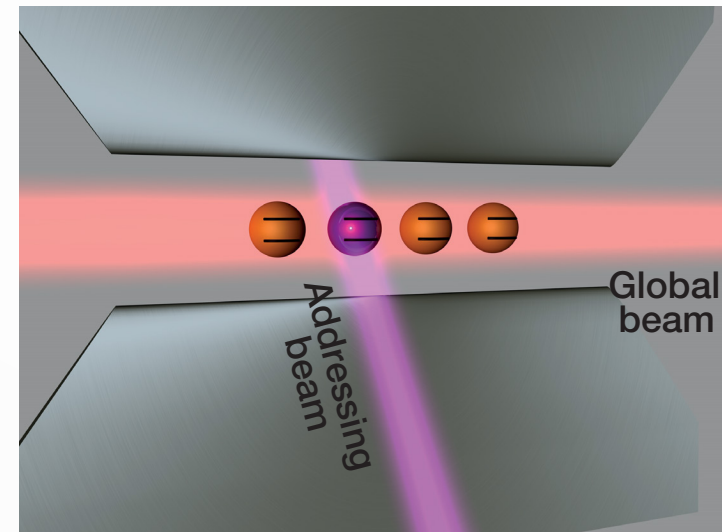
Kogut-Susskind Hamiltonian (Wilson LGT)

Real-time dynamics of lattice gauge theories with a few-qubit quantum computer

Schwinger pair production



ion trap quantum computer



Digital Quantum Simulation of an Exotic Spin Model

- obtained after integrating gauge field

220 quantum gates

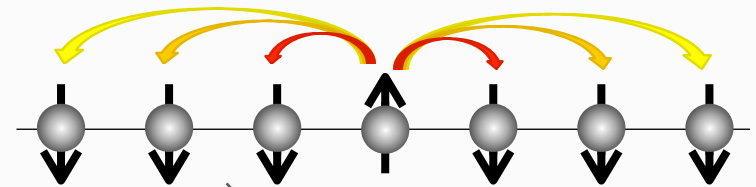
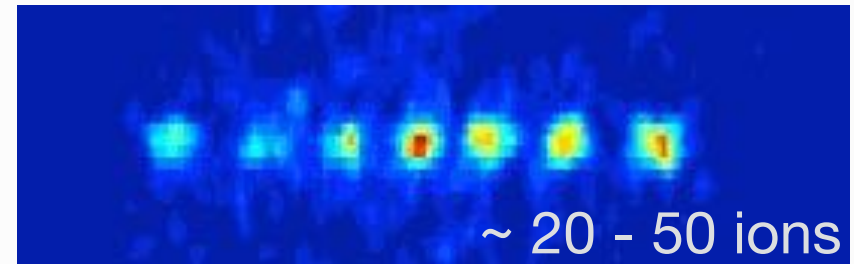
- **spin models**

$$H = \hbar \sum_{i,j} J_{ij} \sigma_x^{(i)} \sigma_x^{(j)} + \hbar B \sum_i \sigma_z^{(i)}$$

K. Kim, C. Monroe et al., Nature (2010)

P Jurcevic, BP Lanyon, R Blatt, C. Roos et al. (2014)

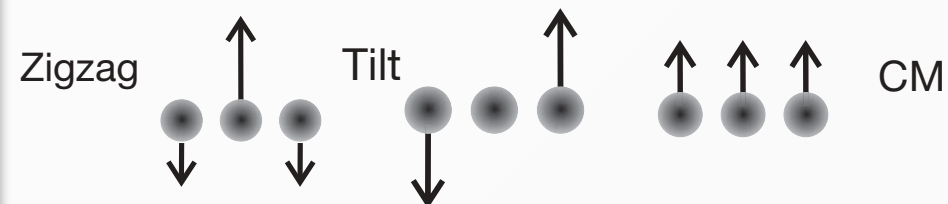
- **trapped ions**



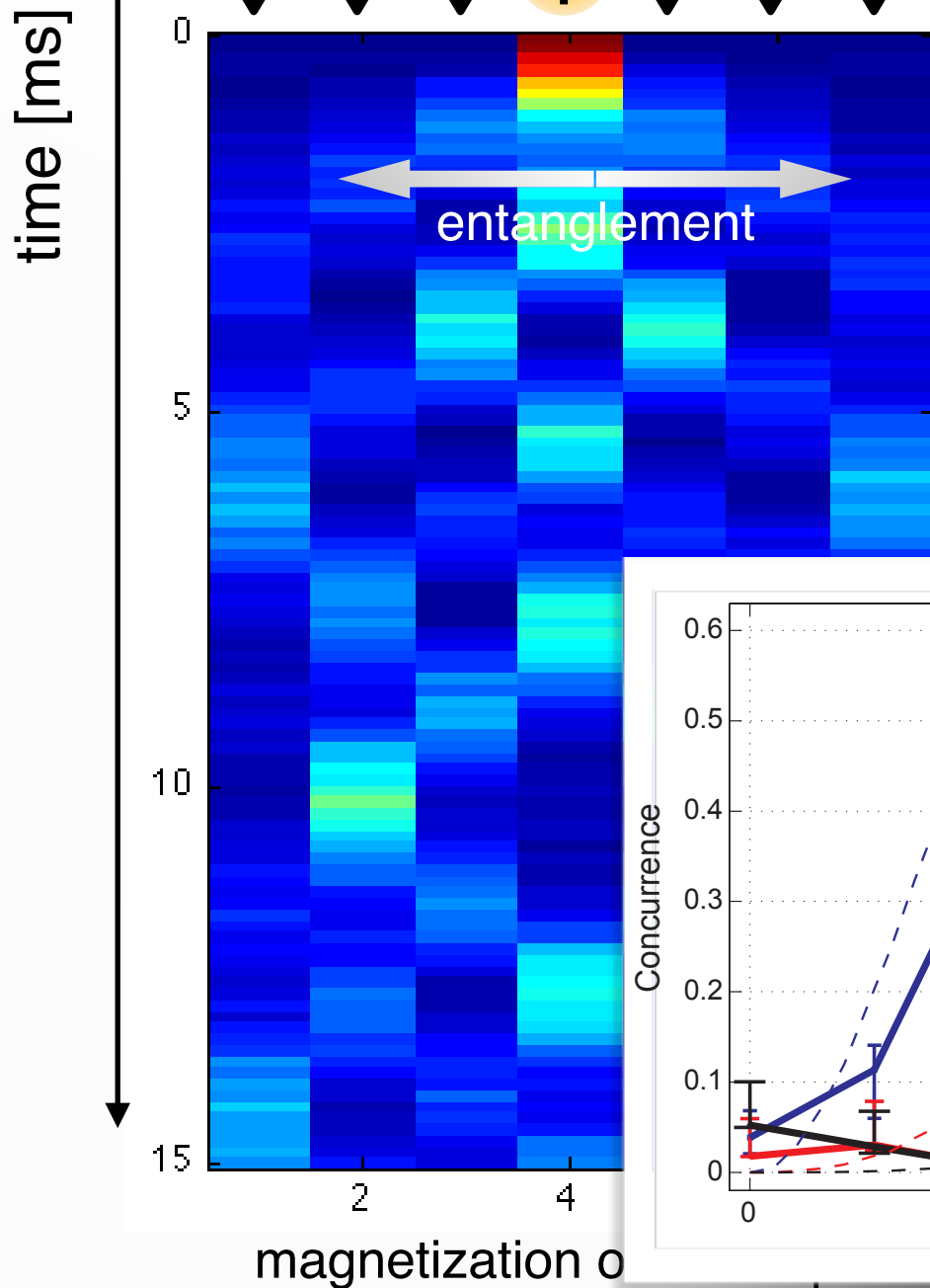
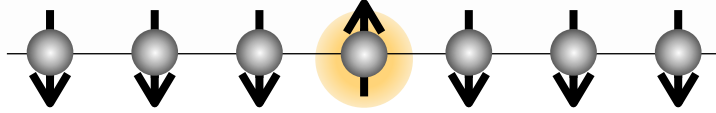
$$J_{ij} \sim \frac{1}{|i-j|^\alpha} \quad 0 \dots 3$$

tunable range interaction

- We “build” a quantum system with desired Hamiltonian & *controllable parameters*, e.g. Hubbard models of atoms in optical lattices

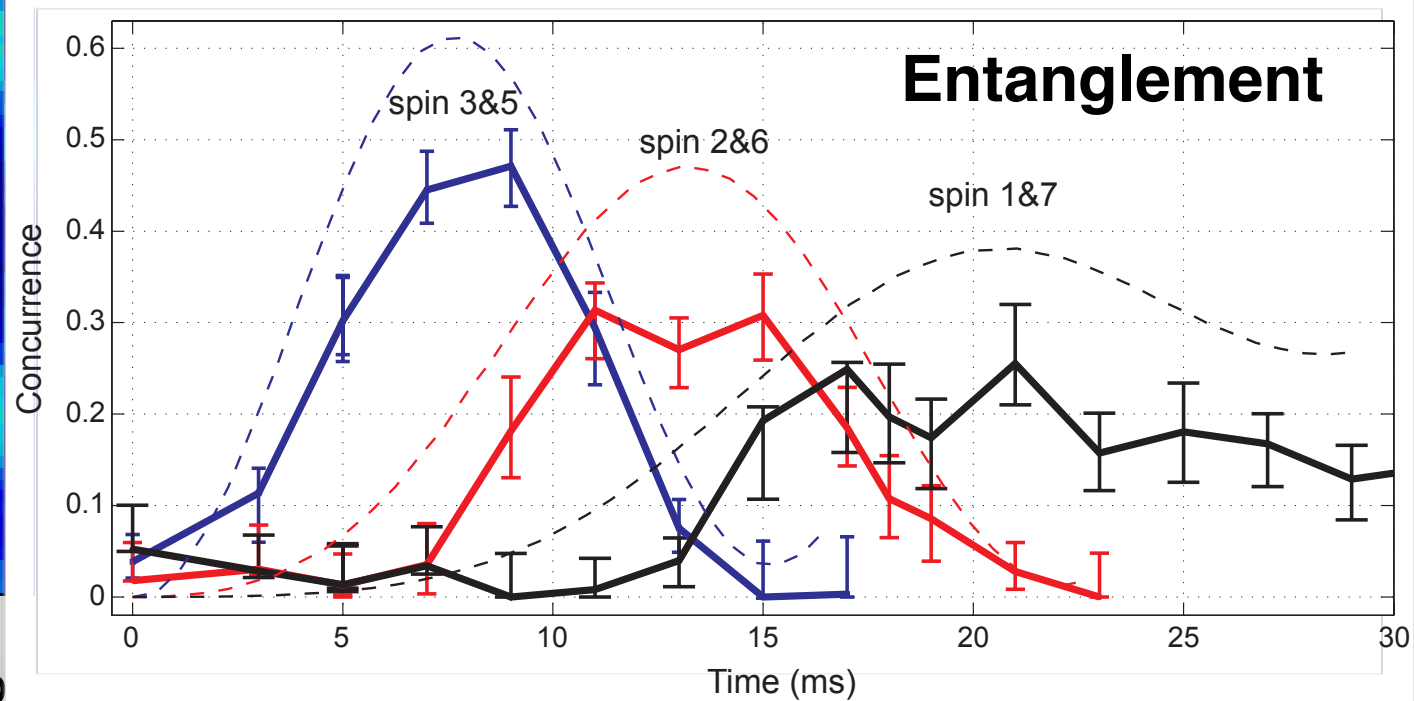


Magnon propagation



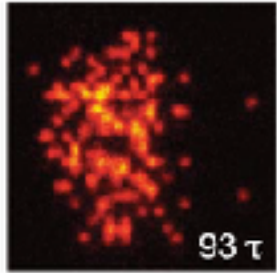
$$\tilde{H}_{XY} = \sum_{i < j} J_{ij} (\sigma^+ \sigma^- + \sigma^- \sigma^+).$$

- ✓ Light-cone-like spreading of entanglement
- ✓ breakdown of the quantum speed-limit due to long range interactions
[exp & theory indistinguishable]

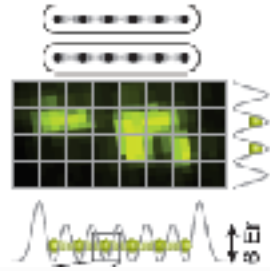


Progress in Engineered Atomic Many-Body Systems

Hubbard models (MPQ, CUA, JQI, ...)

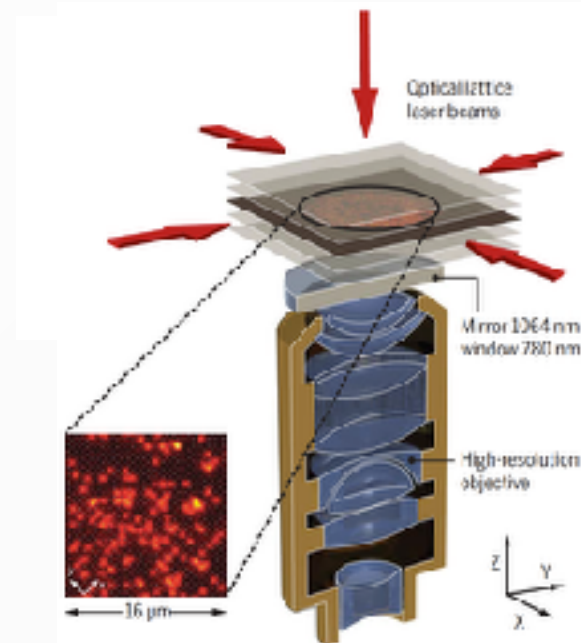


Choi et al., Science (2016)



Kaufman et al., Science (2016)

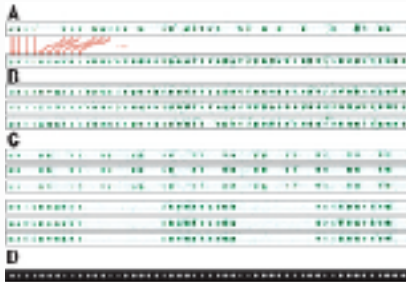
New tools: quantum gas microscope



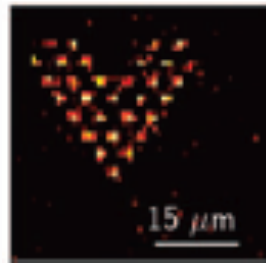
Gross, Bloch, Science (2017)

→ **single site quantum control**
and **measurement**

Rydberg Atoms (MPQ, CUA, IOGS, ...)



Endres et al., Science (2016)



Barredo et al., Science (2016)

Ion Traps (IBK, NIST, JQI, Oxford, ...)

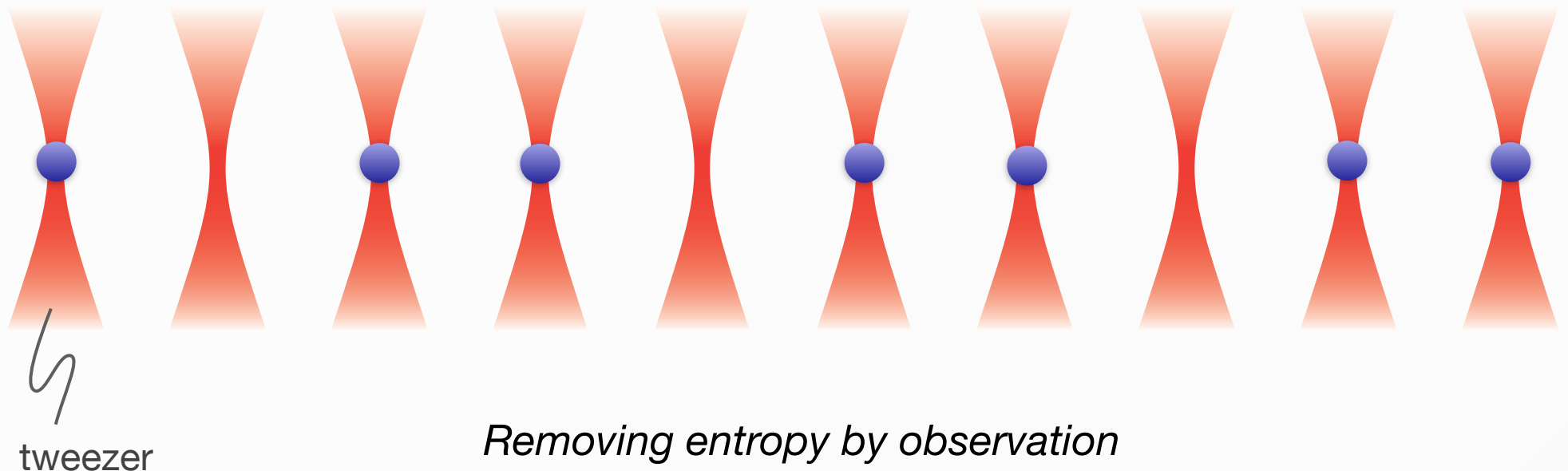
C. Monroe, JQI, Nature 2017

New theory protocols enabled / motivated

How to measure Renyi entropies
demonstrating entanglement?
with repetition rate

Spin Models with Arrays of Rydberg Atoms

Preparation of a *low entropy* state



... or Quantum Computer

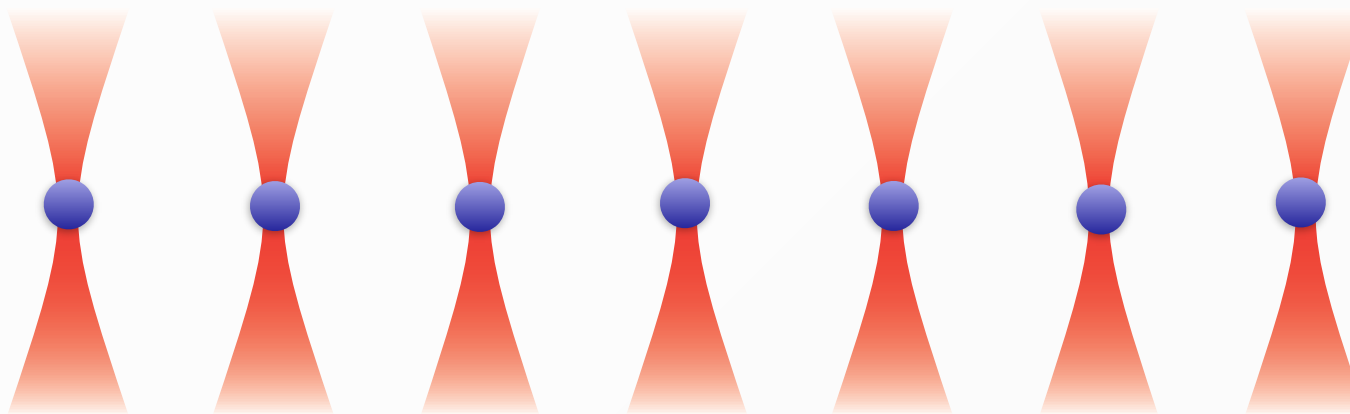
Challenges:

- ✓ controllability
- ✓ scalability

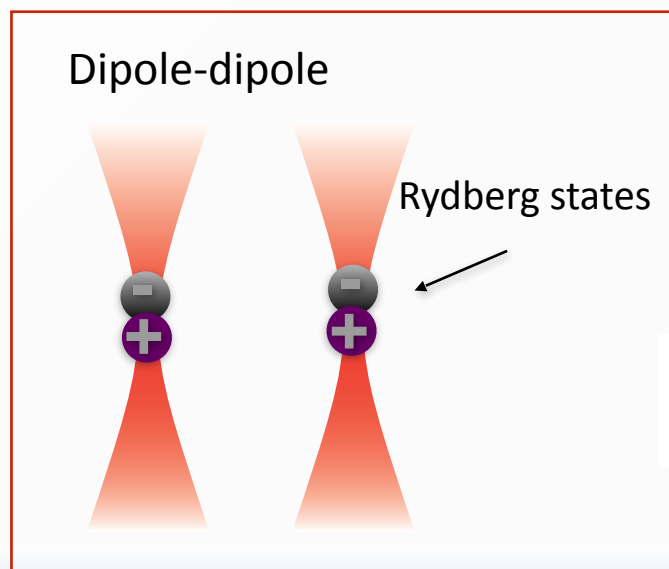
Lukin-Greiner-Vuletic groups (Harvard - MIT), Browaeys (Palaiseau), Saffman (Wisconsin), Biedermann (Sandia)
see also Hubbard-type Models: Regal (JILA), Jochim (Heidelberg), M. Andersen (Otago)

Spin Models with Arrays of Rydberg Atoms

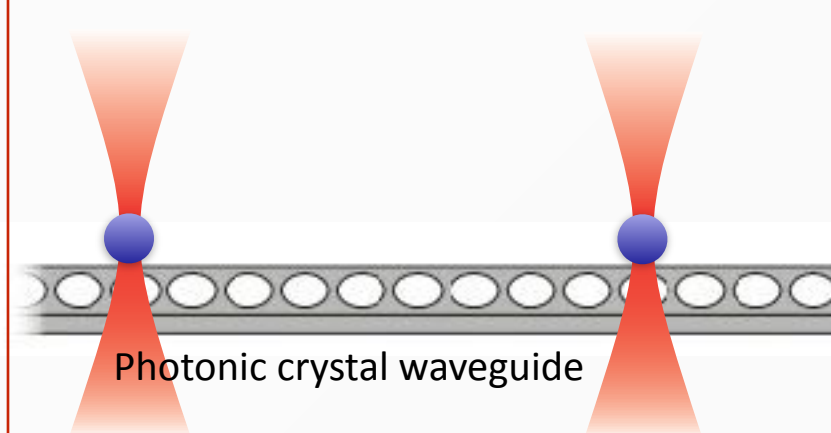
Hamiltonian Engineering



turning on interactions

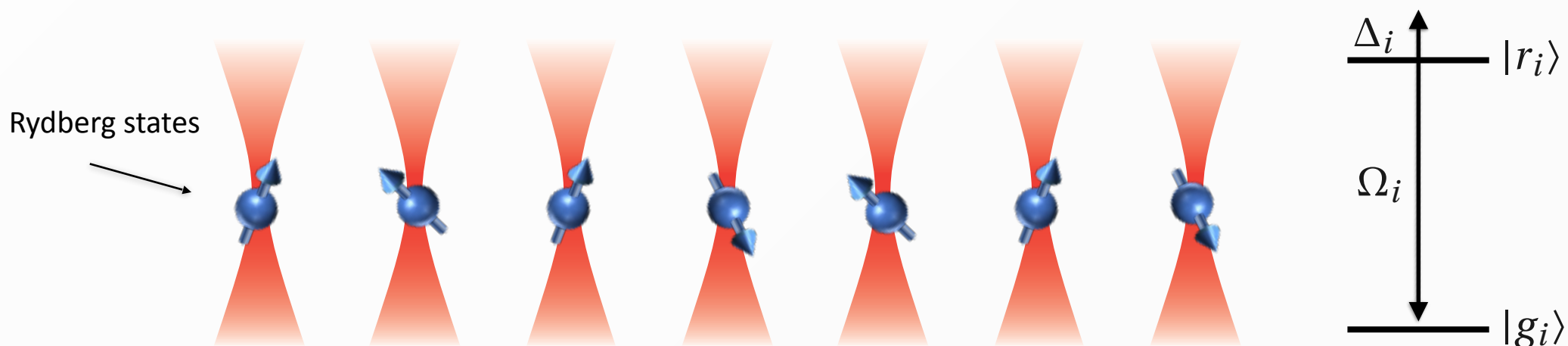


Photon-mediated



Rydberg Spin-Models [or Quantum Computer]

Hamiltonian Engineering



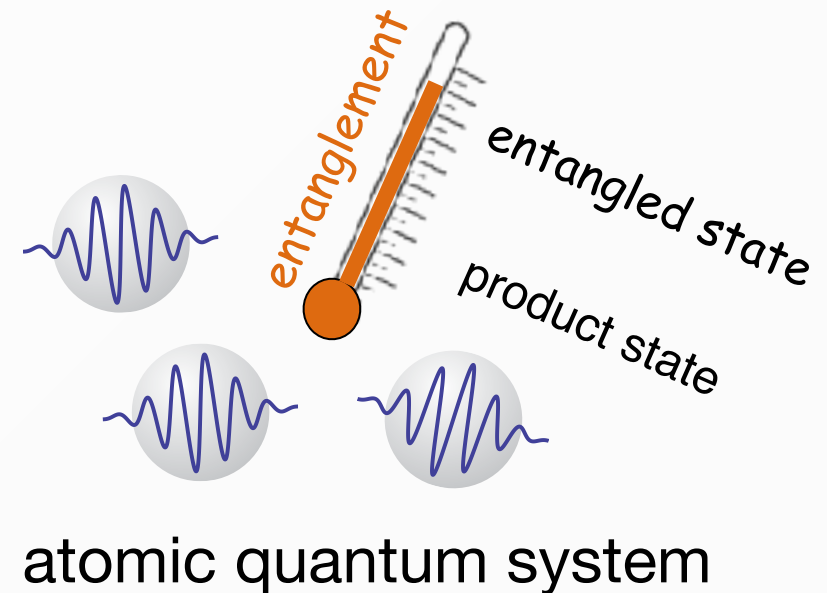
Hamiltonian

$$H = \sum_i \frac{1}{2} \Omega_i \sigma_x^i - \sum_i \Delta_i n_i + \sum_{i < j} V_{ij} n_i n_j$$

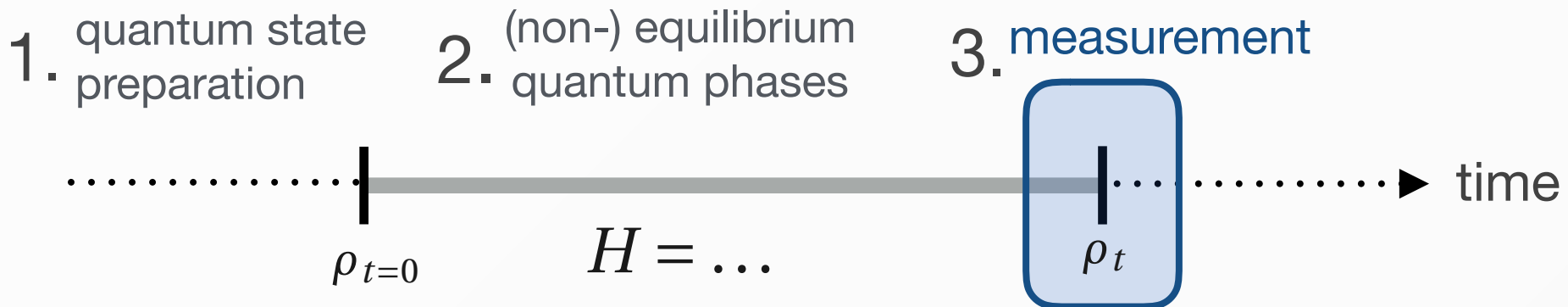
$$\sigma_x^i = |g_i\rangle \langle r_i| + |r_i\rangle \langle g_i|$$

$$V_{ij} = C_6 / r_{ij}^6 \quad n_i = |r_i\rangle \langle r_i| \equiv \frac{1}{2} (1 + \sigma_z^i)$$

Can we measure Entropy/Entanglement?

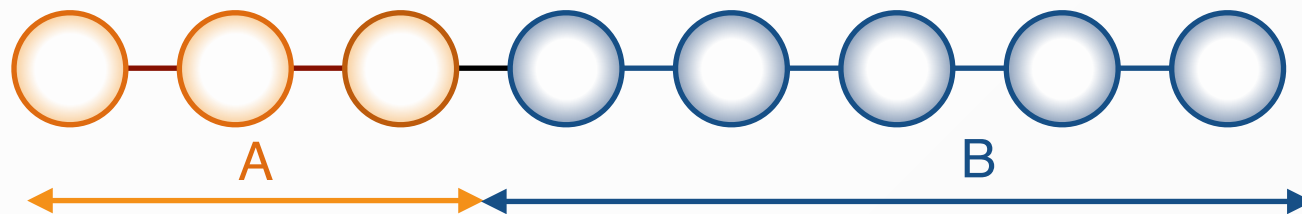


• Cold Atom Experiment



Entanglement Entropy &
Entanglement Spectrum

Quantum many-body system (isolated)



$$H = H_A + H_B + H_{AB}$$

Reduced density matrix & entanglement spectrum

$$\rho_A \equiv \text{Tr}_B |\Psi\rangle \langle \Psi|$$

e.g. ground state

Entanglement entropy

$$S_A = -\text{Tr}_A (\rho_A \log \rho_A)$$

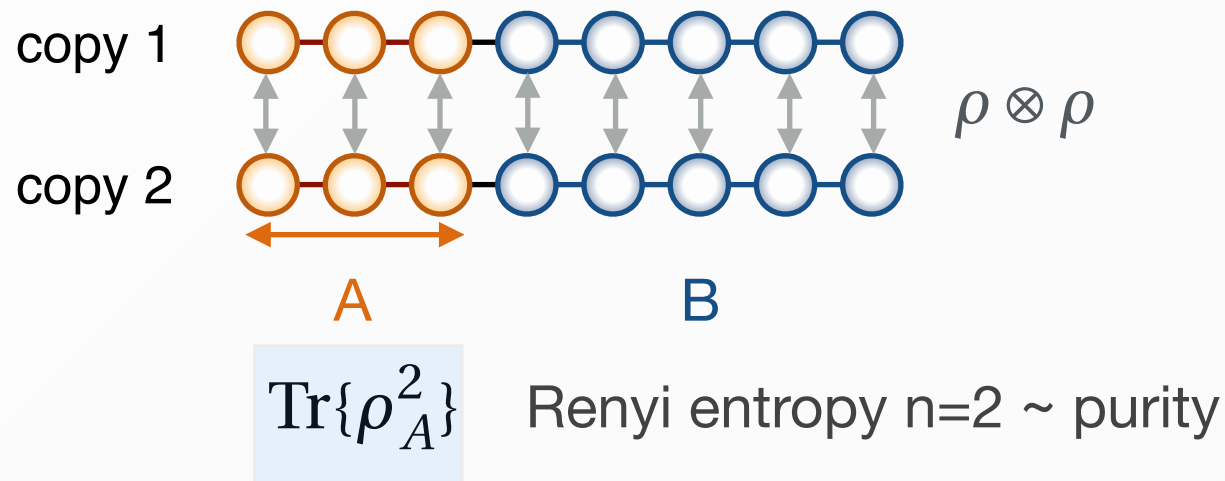
Von Neumann

$$S_A^{(n)} = \frac{1}{1-n} \log \text{Tr}_A \rho_A^n \quad (\leq S_A)$$

Renyi

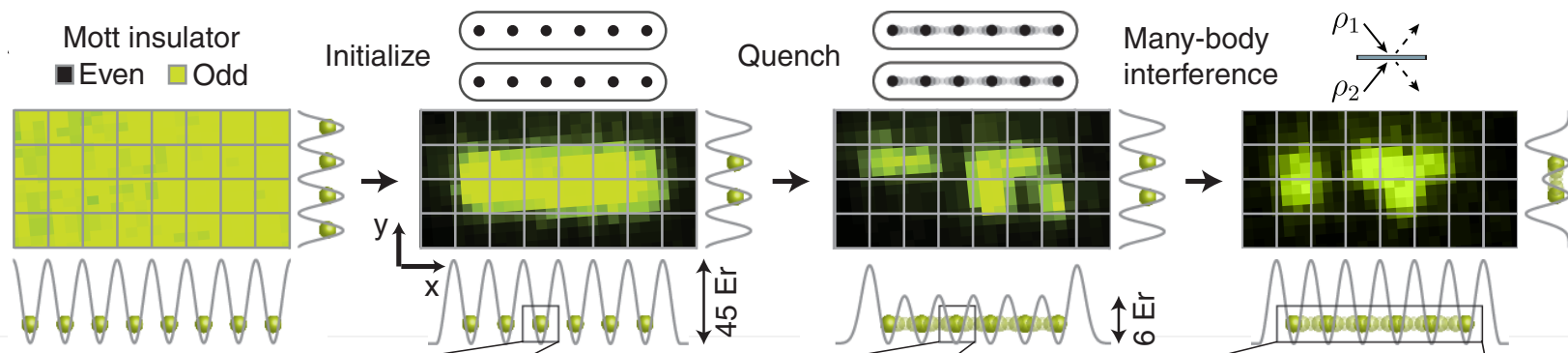
.. Question: Can topological orders strongly correlate entanglement phases

1. Measuring Renyi Entropies *with Copies*



1. Measuring Renyi Entropies *with Copies*

Controlled few-atom systems & quantum gas microscope



R. Islam, M. Greiner et al., Nature (2015)

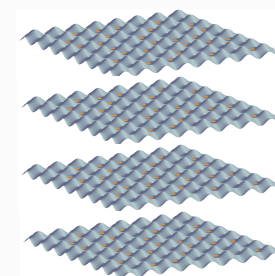
A.M. Kaufmann, M. Greiner et al., Science (2016)

A Quantum Information Perspective

$$\text{Tr}\{\rho^n\}$$

$$\rho \otimes \rho \otimes \rho \otimes \rho \otimes \rho$$

n copies



A Quantum Information Perspective

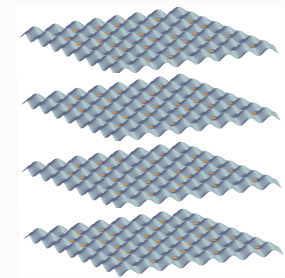
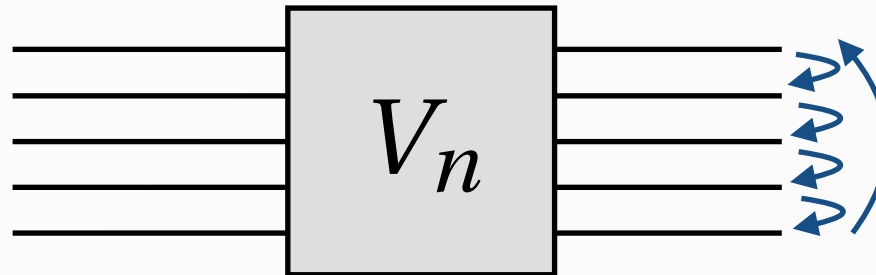
$$\text{Tr}\{\rho^n\} = \text{Tr}\{V^{(n)} \rho \otimes \rho \otimes \dots \otimes \rho \otimes \rho\} \equiv \langle V^{(n)} \rangle \quad \text{expectation value}$$



$$V^{(n)} |\psi_1\rangle \dots |\psi_n\rangle = |\psi_n\rangle |\psi_1\rangle \dots |\psi_{n-1}\rangle$$

shift (or swap) operator

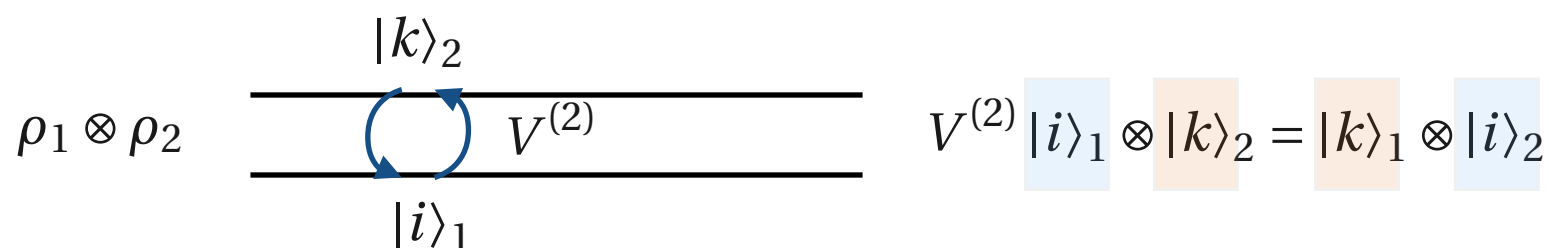
$\rho \otimes \rho \otimes \rho \otimes \rho \otimes \rho$
n copies



A Quantum Information Perspective

$$\text{Tr}\{\rho^n\} = \text{Tr}\{V^{(n)} \rho \otimes \rho \otimes \dots \otimes \rho \otimes \rho\} \equiv \langle V^{(n)} \rangle \quad \text{expectation value}$$

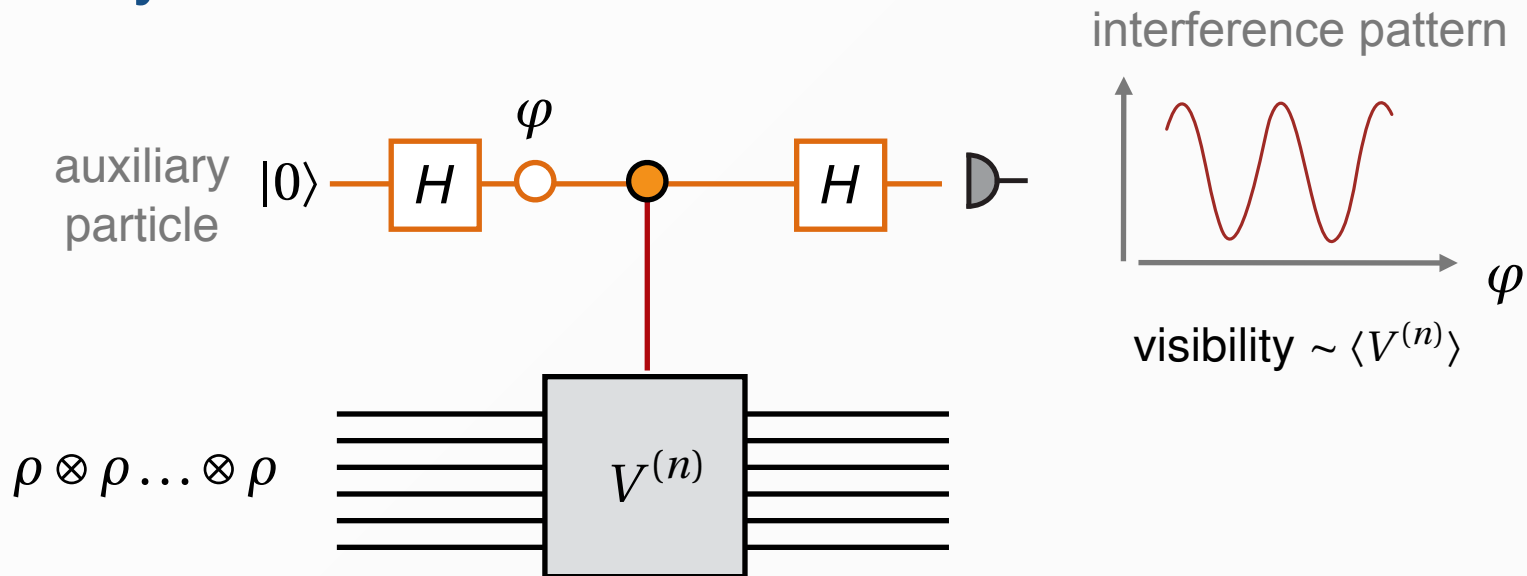
Example: $n=2$



$$\begin{aligned} \text{Tr}\{V^{(2)} \rho_1 \otimes \rho_2\} &= \text{Tr} \left\{ V^{(2)} \sum_{ijkl} \rho_{ij}^{(1)} \rho_{kl}^{(2)} |i\rangle \langle j| \otimes |k\rangle \langle l| \right\} \\ &= \text{Tr}\{\rho_1 \rho_2\} \end{aligned}$$

$$\text{Tr}\{\rho^n\} = \text{Tr}\{V^{(n)} \rho \otimes \rho \otimes \rho \otimes \rho \otimes \rho\} \equiv \langle V^{(n)} \rangle \quad \text{expectation value}$$

Ramsey interferometer:



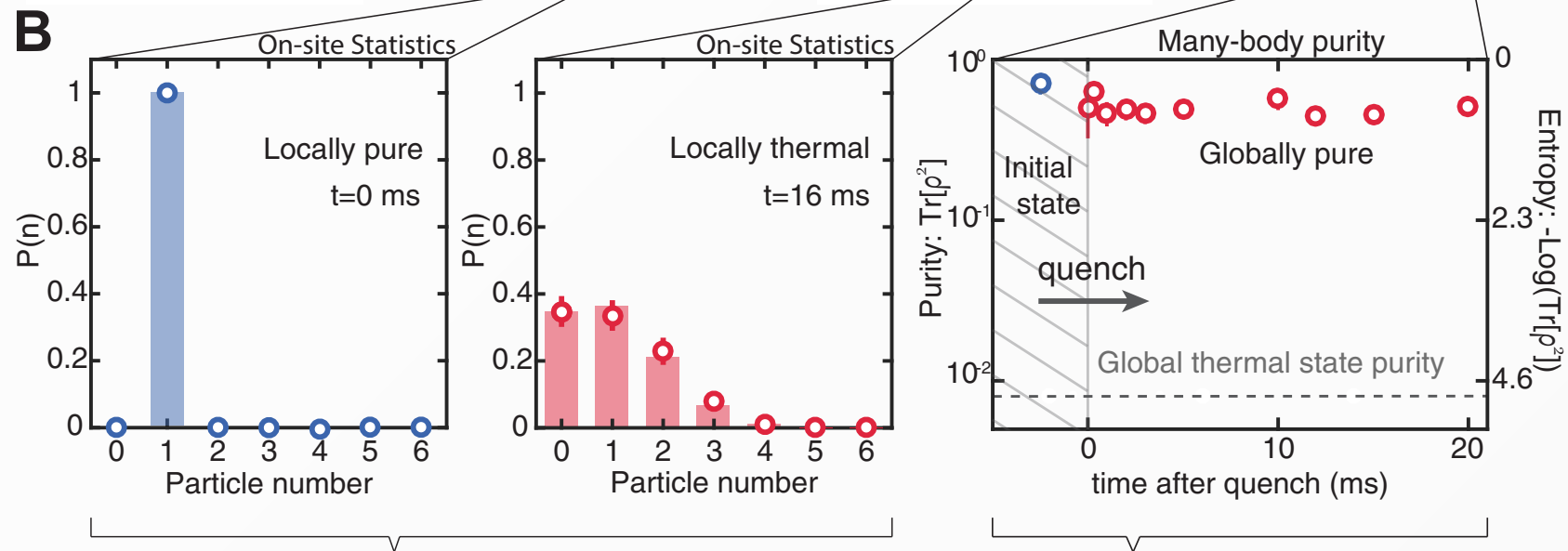
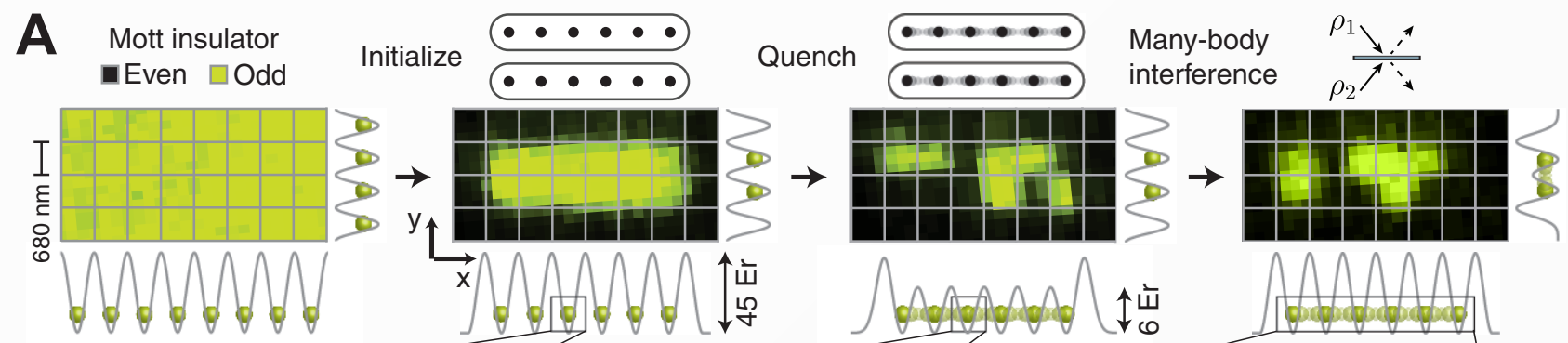
... we need a quantum computer (?)

Quantum thermalization through entanglement in an isolated many-body system

A.J. Daley, H. Pichler, J. Schachenmayer, PZ, PRL (2012).
see also:
C. Moura Alves and D. Jaksch, PRL (2004)

Adam M. Kaufman, M. Eric Tai, Alexander Lukin, Matthew Rispoli, Robert Schittko, Philipp M. Preiss, Markus Greiner*

794 19 AUGUST 2016 • VOL 353 ISSUE 6301



C Expand and Measure Local Occupation Number

Expand and Measure Local and Global Purity





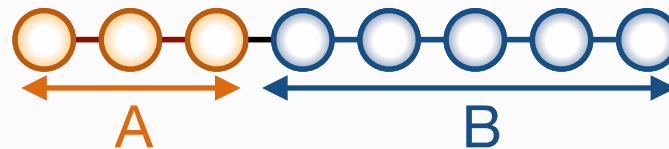
A. Elben



B. Vermersch

2. Measuring Renyi Entropies *via Random Measurements*

single system



$$\rho \otimes \rho$$

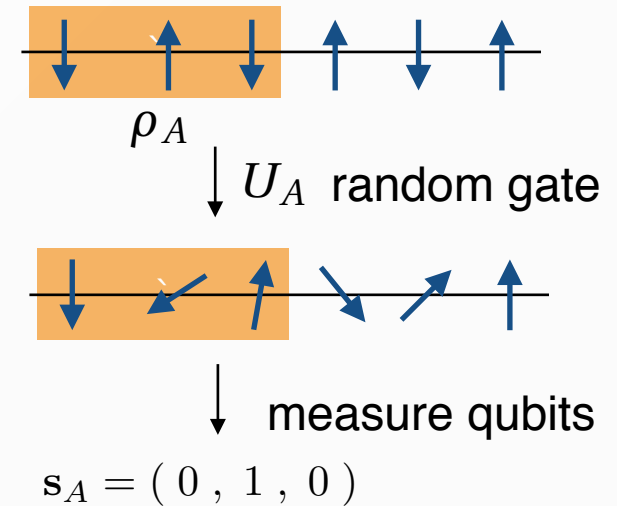
virtual copy
(replica trick)

Random Measurements & Quantum Information

Protocol for chain of qubits:

Random measurement

$$P(\mathbf{s}_A) = \text{Tr}_A \left(|\mathbf{s}_A\rangle \langle \mathbf{s}_A| U_A \rho_A U_A^\dagger \right)$$



Average over Circular Unitary Ensemble (CUE)

$$\langle P(\mathbf{s}_A) \rangle = \frac{1}{N_{H_A}}$$

$$\langle P(\mathbf{s}_A)^2 \rangle = \frac{1 + \text{Tr} \rho_A^2}{N_{H_A} (N_{H_A} + 1)} \quad \approx \quad \begin{matrix} \text{dimension} \\ \text{Hilbert space } A \end{matrix}$$

2-design / t-design

$$\langle P(\mathbf{s}_A)^2 \rangle = \langle \text{Tr}_{1+2} \left[\dots U_A \rho_A U_A^\dagger \otimes U_A \rho_A U_A^\dagger \right] \rangle = \dots$$

2 virtual copies \uparrow CUE

$$\langle U_{ij} U_{kl}^* U_{mn} U_{op}^* \rangle \approx \frac{\delta_{ik} \delta_{jl} \delta_{mo} \delta_{np} + \delta_{io} \delta_{jp} \delta_{mk} \delta_{nl}}{N_H^2}$$

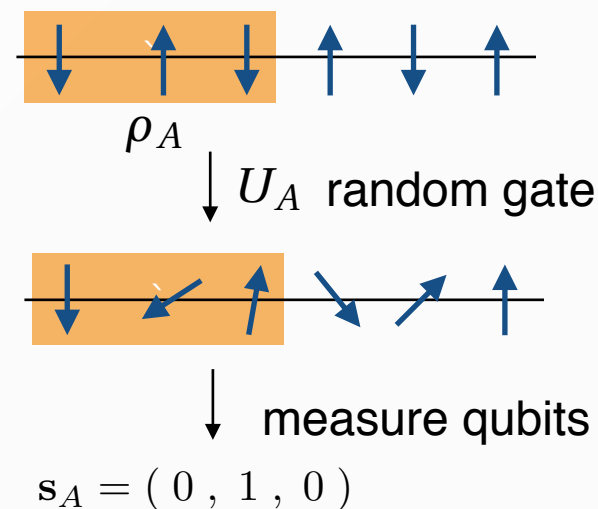
~ Gaussian

Random Measurements & Quantum Information

Protocol for chain of qubits:

Random measurement

$$P(\mathbf{s}_A) = \text{Tr}_A \left(|\mathbf{s}_A\rangle \langle \mathbf{s}_A| U_A \rho_A U_A^\dagger \right)$$



Average over Circular Unitary Ensemble (CUE)

$$\langle P(\mathbf{s}_A) \rangle = \frac{1}{N_{H_A}} \quad \langle P(\mathbf{s}_A)^2 \rangle = \frac{1 + \text{Tr} \rho_A^2}{N_{H_A} (N_{H_A} + 1)}$$

dimension
Hilbert space A

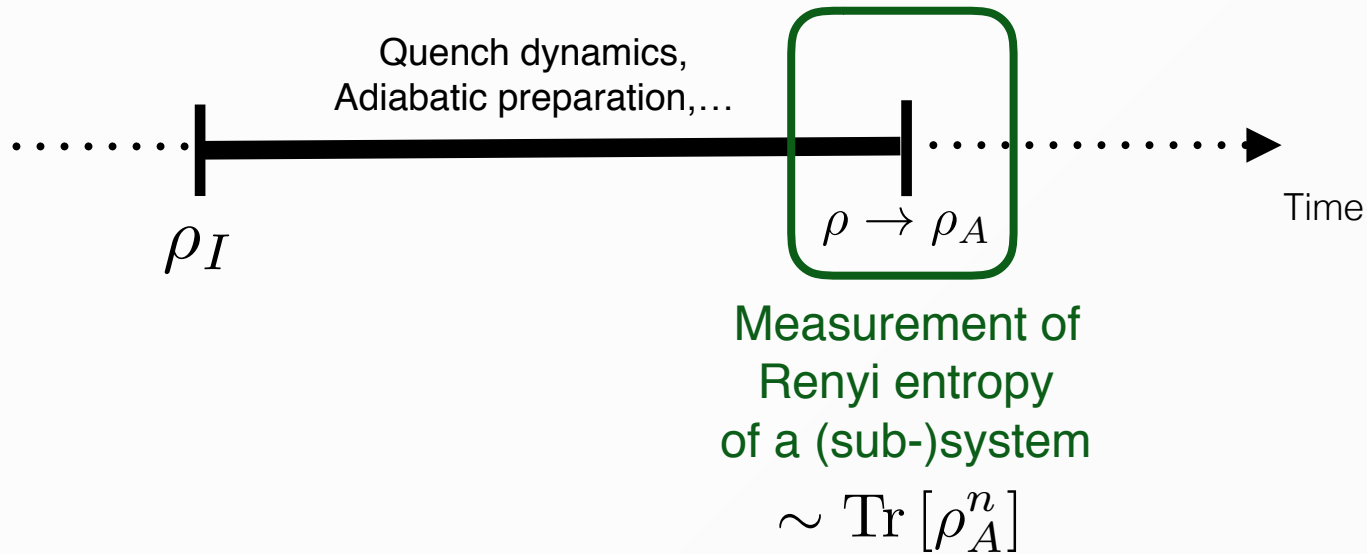
✓ Random measurement gives purity, and higher order Renyi entropies

✓ Required resources:

- how realize random unitaries
- # measurements, # unitaries

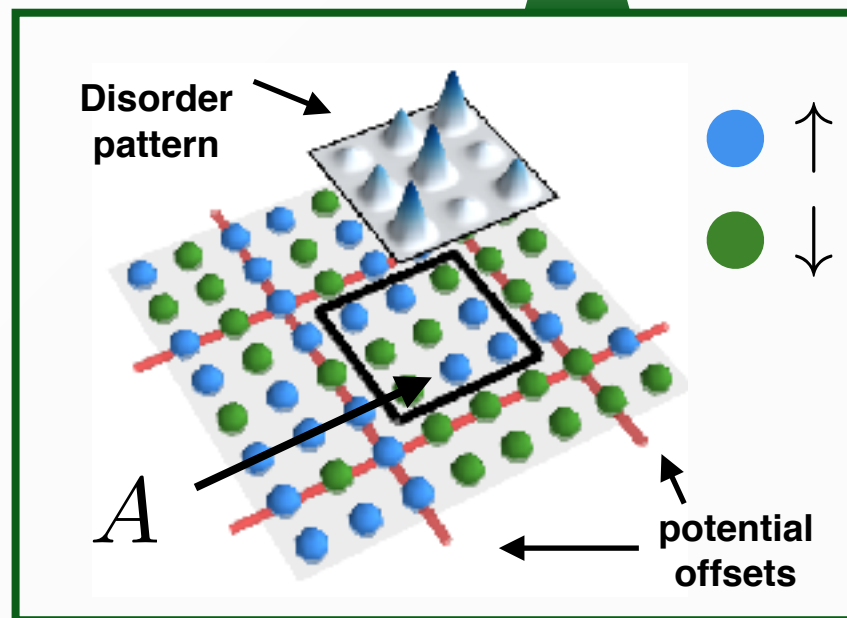
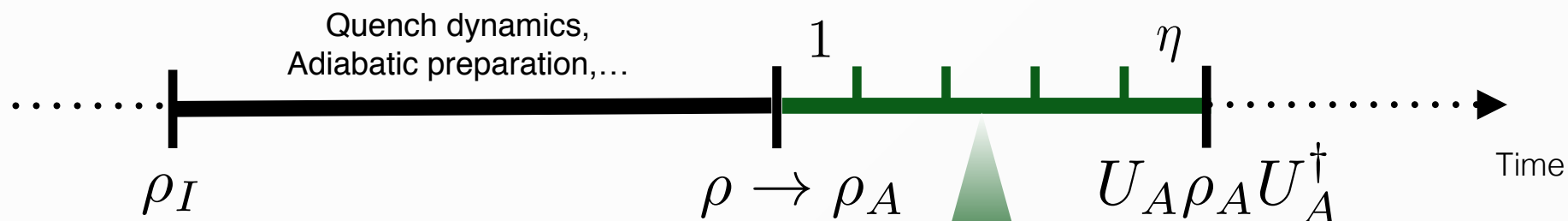
? "the signal is the noise"

Measurement Protocol



Random unitary
as time evolution operator under
random quenches

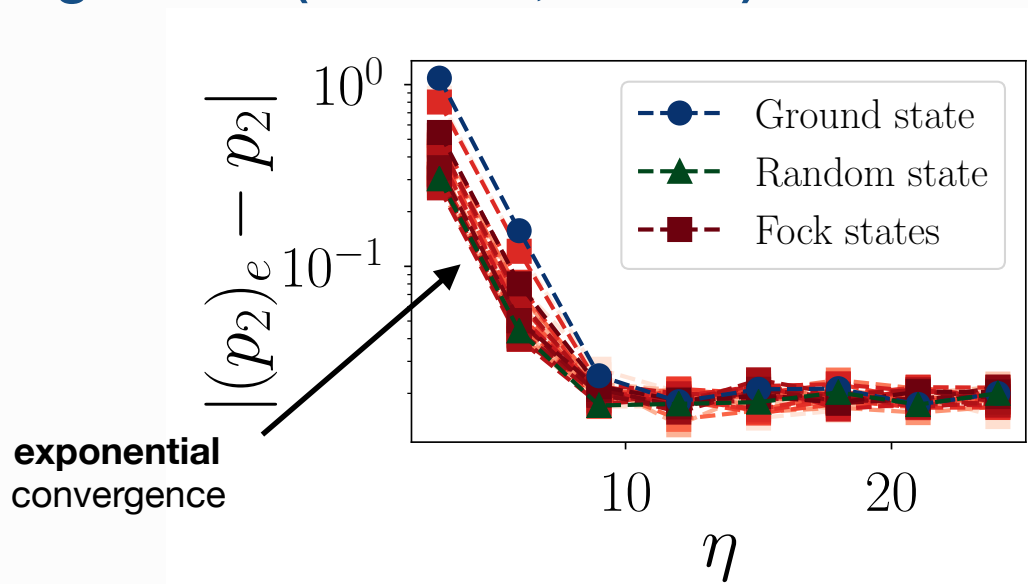
$$U_A = e^{-iH_\eta T} \dots e^{-iH_1 T}$$



Question: Is $U_A = e^{-iH_\eta T} \dots e^{-iH_1 T}$ a random unitary?

→ Apply the protocol to a known input state and compare estimated to true purity to test the ensemble

Ising model (here: 1D, 8 sites)



$$H_j = \Omega \sum_i \sigma_i^x + \sum_{i < j} \frac{C_6}{|r_i - r_j|^6} \sigma_i^z \sigma_j^z + \sum_i \Delta_i^j \sigma_i^z$$

Δ_i^j from gaussian distribution with standard deviation Δ

$$\Omega = C_6/a^6 = \Delta = 1/T$$

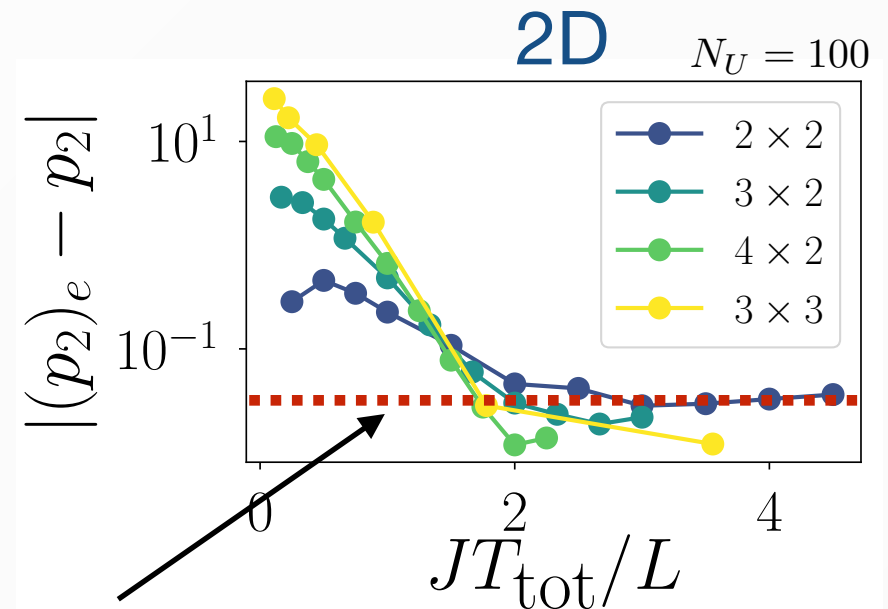
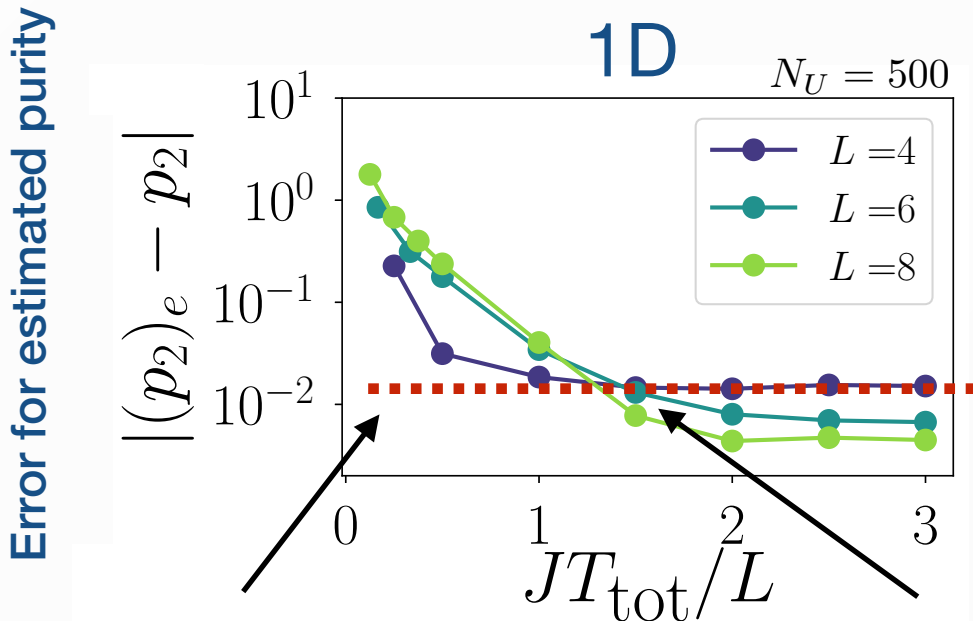
Random unitaries using
 ✓ generic interactions
 ✓ engineered disorder

Ising model with L sites

$$U_A = e^{-iH_\eta T} \dots e^{-iH_1 T}$$

$$H_j = \Omega \sum_i \sigma_i^x + \sum_{i < j} \frac{C_6}{|r_i - r_j|^6} \sigma_i^z \sigma_j^z + \sum_i \Delta_i^j \sigma_i^z$$

Δ_i^j from gaussian distribution
with standard deviation Δ



statistical error threshold
due to finite number (500)
random unitaries

Number of necessary
random quenches
 $\eta \sim L$

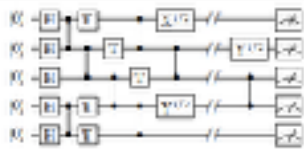
Efficient generation of random
unitaries for purity measurements

S : set of generated random unitaries

$$\langle P(\mathbf{s}_A)^n \rangle_S \stackrel{!}{\sim} \text{Tr}[\rho_A^n] \longleftrightarrow \langle \underbrace{U_{ij}U_{kl}^* \dots U_{mn}U_{op}^*}_{\text{up to } 2n \text{ matrix elements}} \rangle_S \stackrel{!}{=} \langle U_{ij}U_{kl}^* \dots U_{mn}U_{op}^* \rangle_{\text{CUE}}$$

→ S is required to form a **unitary n-design** (not full CUE)

Random quantum circuits: Efficient generation of n-designs



S. Boixo, et al., arXiv:1608.00263

Ohliger et al., New J. Phys., 2014
Brandão et al., Comm. Mat. Phys, 2016

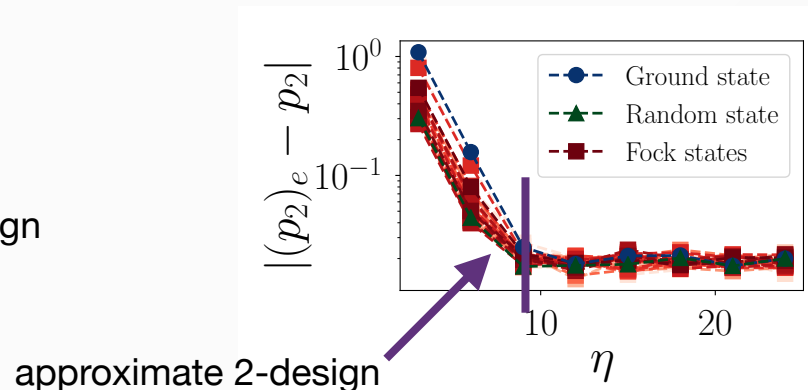
Number of random quenches η ↔ depth of circuit

Certification hierarchy of random unitaries:

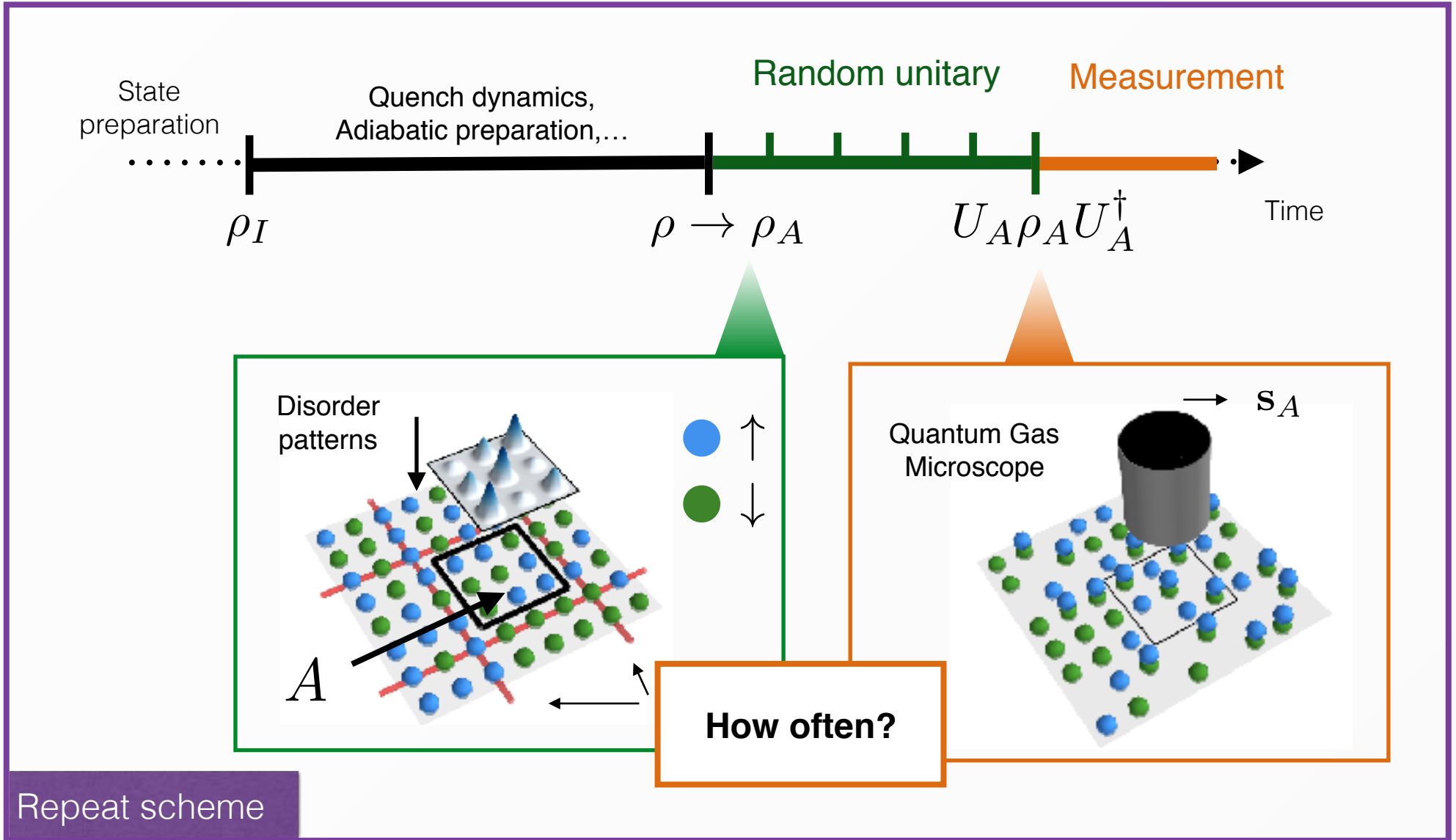
$$\langle P(\mathbf{s}_A) \rangle_S \stackrel{!}{=} \frac{1}{N_{\mathcal{H}_A}} \quad \text{holds if } S \text{ forms 1-design}$$

$$\langle P(\mathbf{s}_A)^2 \rangle_S \stackrel{!}{=} \frac{1 + \text{Tr}[\rho_A^2]}{N_{\mathcal{H}_A}(N_{\mathcal{H}_A} + 1)} \quad \text{holds if } S \text{ forms 2-design}$$

⋮



Measurement Protocol in Hubbard & Spin Models

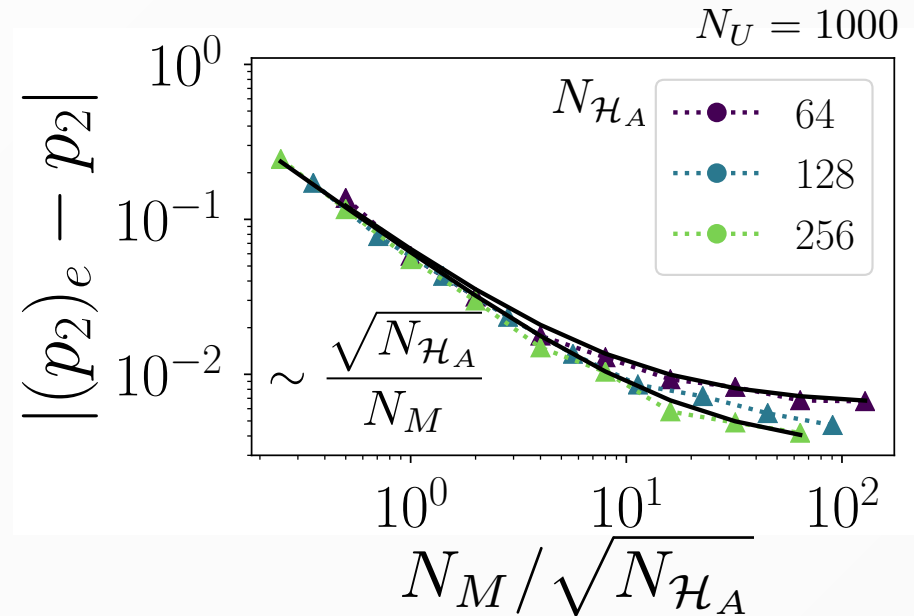
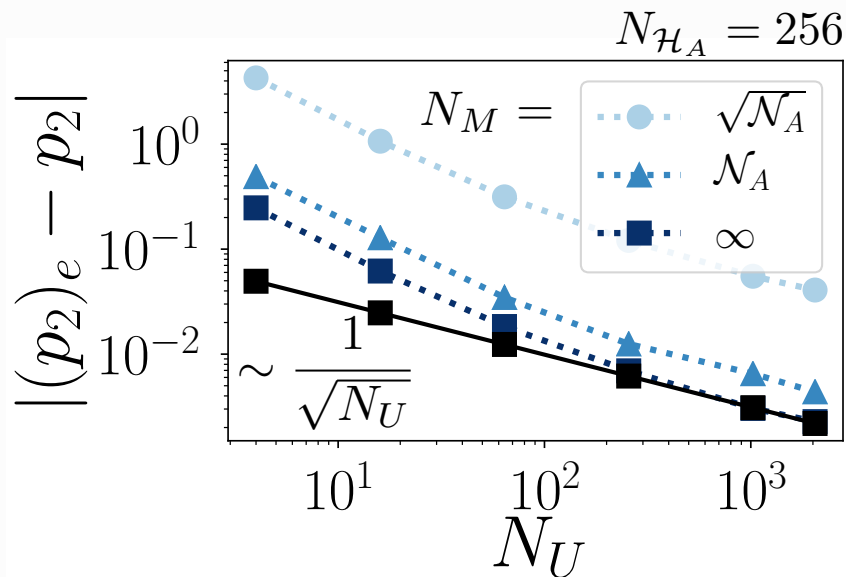


For the same random unitary \rightarrow probabilities $P(s_A)$ for all measurement outcomes s_A

For many random unitaries \rightarrow correlations $\langle P(s_A)^2 \rangle \sim \text{Tr} [\rho_A^2]$

N_M : number of measurements per unitary
 N_U : number of unitaries
 $N_{\mathcal{H}_A}$: Hilbert space dimension of A

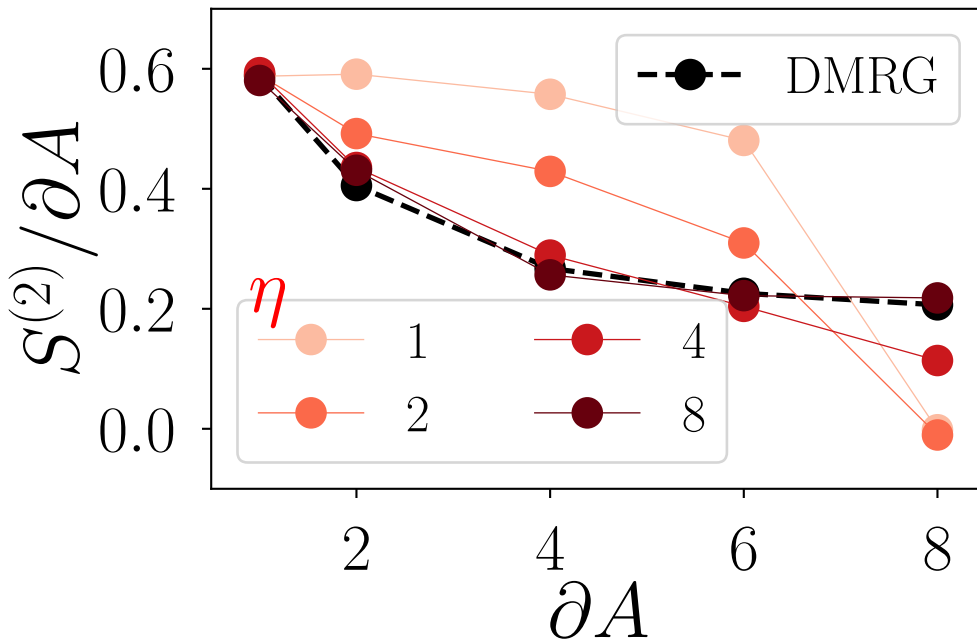
Error for estimated purity (averaged over all outcomes)



Number of measurements to determine p_2
 up to error $\sim 1/\sqrt{N_U}$
 $N_M \sim \sqrt{N_{\mathcal{H}_A}}$

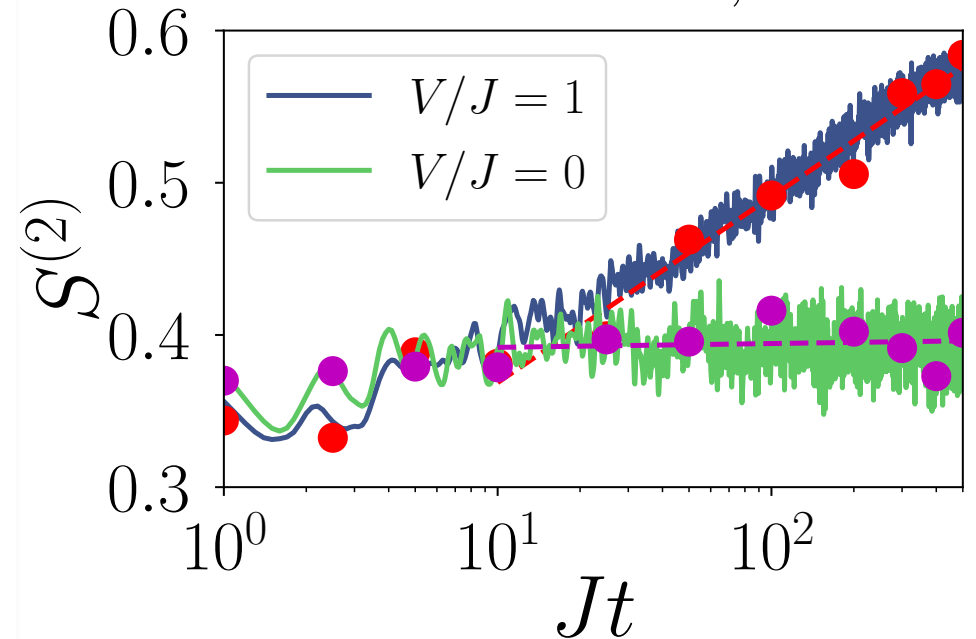
Our protocol allows to *measure* ...

area law in 2D Heisenberg model



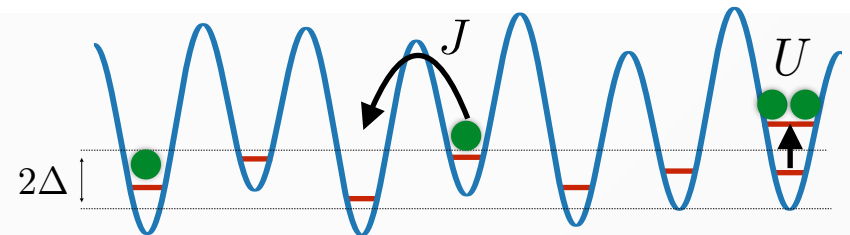
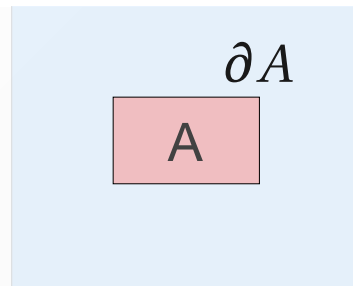
entropy growth in MBL phase

$L = 10, N = 5$



area law

$$S_A \approx \alpha \partial A$$



1D Bose Hubbard with disorder
 $N_M = 100 \quad N_U = 100$

Dirac Medallists 2017



Charles H. Bennett



David Deutsch



Peter W. Shor

Citation:

for pioneering work in applying fundamental concepts of quantum mechanics to solving basic problems in computation and communication, and therefore bringing together the fields of quantum mechanics, computer science and information.

quantum science as fundamental research & quantum technologies